

**Homework 9**  
**Partial Differential Equations**  
**Math 531**  
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**Excercise 8.2.2:** Consider the heat equation with time-dependent sources and boundary conditions

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + Q(x, t) \quad u(x, 0) = f(x)$$

Reduce the problem to one with homogeneous boundary conditions if

(b)

$$u(0, t) = A(t) \quad \text{and} \quad \frac{\partial u}{\partial x}(L, t) = B(t)$$

Let  $u_E(x, t)$  be the linear PDE that satisfies the steady state problem:

$$u_E(0, t) = A(t) \quad \frac{\partial u_E}{\partial x}(L, t) = B(t) \quad \rightarrow \quad u_E(x, t) = A(t) + xB(t)$$

We can now define  $v(x, t) = u(x, t) - u_E(x, t)$  with the following PDE's:

$$\frac{\partial v}{\partial t} = \frac{\partial u}{\partial t} - \frac{\partial u_E}{\partial t} \quad \rightarrow \quad \frac{\partial u}{\partial t} = \frac{\partial v}{\partial t} + \frac{\partial u_E}{\partial t} \quad \frac{\partial^2 v}{\partial x^2} = \frac{\partial^2 u}{\partial x^2} - \frac{\partial^2 u_E}{\partial x^2} = \frac{\partial^2 u}{\partial x^2}$$

Now we can convert our original PDE:

$$\frac{\partial v}{\partial t} + \frac{\partial u_E}{\partial t} = k \frac{\partial^2 v}{\partial x^2} + Q(x, t)$$

Now notice this is a PDE with the following homogeneous boundary conditions:

$$v(0, t) = u(0, t) - u_E(0, t) = 0 \quad \frac{\partial v}{\partial x}(L, t) = \frac{\partial u}{\partial x}(L, t) - \frac{\partial u_E}{\partial x}(L, t) = 0$$

and the following initial condition:

$$v(x, 0) = u(x, 0) - u_E(x, 0) = f(x) - (A(0) + xB(0))$$

**Excercise 8.2.5:** Solve the initial value problem for a two-dimensional heat equation inside a circle (of radius  $a$ ) with time-independent boundary conditions

$$\frac{\partial u}{\partial t} = k \nabla^2 u \quad u(a, \theta, t) = g(\theta) \quad u(r, \theta, 0) = f(r, \theta)$$

Let  $u_E(r, \theta)$  be the equilibrium temperature distribution:

$$\nabla^2 u_E = 0 \quad \rightarrow \quad \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_E}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_E}{\partial \theta^2} = 0$$

Using separation of variables, we get:  $u_E(r, \theta) = F(r)G(\theta)$ :

$$\frac{r}{F} \frac{d}{dr} \left( r \frac{dF}{dr} \right) + \frac{1}{G} \frac{d^2 G}{d\theta^2} = 0 \quad \rightarrow \quad \frac{r}{F} \frac{d}{dr} \left( r \frac{dF}{dr} \right) = -\frac{1}{G} \frac{d^2 G}{d\theta^2} = \lambda$$

Thus we get the following ODE's:

$$r^2 F'' + r F' - \lambda F = 0 \quad G'' + \lambda G = 0$$

Using the fact of the geometry of the circle, and that  $F(r)$  is finite at  $r = 0$ , we get the following eigenvalues and eigenfunction:

$$\lambda = 0 \quad \rightarrow \quad G(\theta) = 1 \quad F(r) = 1$$

$$\lambda = m^2 \quad \rightarrow \quad G(\theta) = a \cos m\theta + b \sin m\theta \quad F(r) = c_1 r^m$$

From this, we get the following:

$$u_E(r, \theta) = \sum_{m=0}^{\infty} (A_m \cos m\theta + B_m \sin m\theta) r^m$$

Using the boundary conditions, we get:

$$u_E(a, \theta) = A_0 + \sum_{m=1}^{\infty} (A_m \cos m\theta + B_m \sin m\theta) a^m$$

From here, we get the following coefficients:

$$A_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} g(\theta) d\theta \quad A_m = \frac{1}{a^m \pi} \int_{-\pi}^{\pi} g(\theta) \cos m\theta d\theta \quad B_m = \frac{1}{a^m \pi} \int_{-\pi}^{\pi} g(\theta) \sin m\theta d\theta$$

We let the following be true now:

$$v(r, \theta, t) = u(r, \theta, t) - u_E(r, \theta) \quad \text{with} \quad v(a, \theta, t) = 0 \quad v(r, \theta, 0) = f(r, \theta) - u_E(r, \theta)$$

Similar to previous homeworks, we get the following for  $v(r, \theta, t)$ :

$$v(r, \theta, t) = \sum_{n=1}^{\infty} A_{0n} J_0(\sqrt{\lambda_{0n}} r) e^{-\lambda_{0n} k t} + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} J_m(\sqrt{\lambda_{mn}} r) (A_{mn} \cos m\theta + B_{mn} \sin m\theta) e^{-\lambda_{mn} k t}$$

From here, we use our initial conditions:

$$\begin{aligned} v(r, \theta, 0) &= f(r, \theta) - u_E(r, \theta) \\ &= \sum_{n=1}^{\infty} A_{0n} J_0(\sqrt{\lambda_{0n}} r) + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} J_m(\sqrt{\lambda_{mn}} r) (A_{mn} \cos m\theta + B_{mn} \sin m\theta) \end{aligned}$$

From here, we get the following coefficients:

$$\begin{aligned} A_{0n} &= \frac{\int_0^a \int_{-\pi}^{\pi} (f(r, \theta) - u_E(r, \theta)) J_0(\sqrt{\lambda_{0n}} r) r dr d\theta}{2\pi \int_0^a J_0^2(\sqrt{\lambda_{0n}} r) dr} \\ A_{mn} &= \frac{\int_0^a \int_{-\pi}^{\pi} (f(r, \theta) - u_E(r, \theta)) J_m(\sqrt{\lambda_{mn}} r) \cos m\theta r dr d\theta}{2\pi \int_0^a J_m^2(\sqrt{\lambda_{mn}} r) dr} \\ B_{mn} &= \frac{\int_0^a \int_{-\pi}^{\pi} (f(r, \theta) - u_E(r, \theta)) J_m(\sqrt{\lambda_{mn}} r) \sin m\theta r dr d\theta}{2\pi \int_0^a J_m^2(\sqrt{\lambda_{mn}} r) dr} \end{aligned}$$

**Excercise 8.4.2:** Use the method of eigenfunction expansions to solve, without reducing to homogeneous boundary conditions,

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} \quad u(0, t) = A \quad u(L, t) = B \quad u(x, 0) = f(x)$$

Using an eigenfunction expansion, we choose get the Sturm-Liouville Problem in the spatial domain:

$$\frac{d^2 \phi}{dx^2} + \lambda \phi = 0 \quad \text{with} \quad \phi(0) = 0 \quad \phi(L) = 0$$

Notice the eigenvalues and eigenfunctions:

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \quad \phi_n(x) = \sin \frac{n\pi x}{L}$$

Thus we can let the following be true:

$$u(x, t) = \sum_{n=1}^{\infty} B_n(t) \phi_n(x) \quad \text{with} \quad u(x, 0) = f(x) = \sum_{n=1}^{\infty} B_n(0) \phi_n(x) \quad \rightarrow \quad B_n(0) = \frac{2}{L} \int_0^L f(x) \phi_n(x) dx$$

Notice the following from that and orthogonality of sines:

$$\frac{\partial u}{\partial t} = \sum_{n=1}^{\infty} \frac{dB_n(t)}{dt} \phi_n(x) = k \frac{\partial^2 u}{\partial x^2} \quad \frac{dB_n(t)}{dt} = \frac{2k}{L} \int_0^L \frac{\partial^2 u}{\partial x^2} \phi_n(x)$$

Notice the following from Green's formula:

$$\begin{aligned} \int_0^L \left( u \frac{d^2 \phi_n}{dx^2} - \phi_n \frac{\partial^2 u}{\partial x^2} \right) dx &= u \frac{d\phi_n}{dx} - \phi_n \frac{\partial u}{\partial x} \Big|_0^L \\ &= \left( u(L, t) \frac{d\phi_n(L)}{dx} - \phi_n(L) \frac{\partial u}{\partial x}(L, t) \right) - \left( u(0, t) \frac{d\phi_n(0)}{dx} - \phi_n(0) \frac{\partial u}{\partial x}(0, t) \right) \\ &= B \frac{d\phi_n(L)}{dx} - A \frac{d\phi_n(0)}{dx} \\ &= \frac{n\pi}{L} \left( (-1)^n B - A \right) \end{aligned}$$

Notice the following equality:

$$\begin{aligned} \int_0^L u \frac{d^2 \phi_n}{dx^2} dx - \int_0^L \phi_n \frac{\partial^2 u}{\partial x^2} dx &= \int_0^L \left( u \frac{d^2 \phi_n}{dx^2} - \phi_n \frac{\partial^2 u}{\partial x^2} \right) dx \\ \int_0^L u \frac{d^2 \phi_n}{dx^2} dx - \int_0^L \left( u \frac{d^2 \phi_n}{dx^2} - \phi_n \frac{\partial^2 u}{\partial x^2} \right) dx &= \int_0^L \phi_n \frac{\partial^2 u}{\partial x^2} dx \end{aligned}$$

Thus we get the following:

$$\begin{aligned}
\int_0^L \phi_n \frac{\partial^2 u}{\partial x^2} dx &= \int_0^L u \frac{d^2 \phi_n}{dx^2} dx - \int_0^L \left( u \frac{d^2 \phi_n}{dx^2} - \phi_n \frac{\partial^2 u}{\partial x^2} \right) dx \\
&= \int_0^L u \frac{d^2 \phi_n}{dx^2} dx - \frac{n\pi}{L} \left( (-1)^n B - A \right) \\
&= \int_0^L u (-\lambda_n \phi_n) dx - \frac{n\pi}{L} \left( (-1)^n B - A \right) \\
&= -\lambda_n \int_0^L \left( \sum_{m=1}^{\infty} B_m(t) \phi_m(x) \right) \phi_n dx - \frac{n\pi}{L} \left( (-1)^n B - A \right) \\
&= -\lambda_n \left( \frac{L}{2} \right) B_n(t) - \frac{n\pi}{L} \left( (-1)^n B - A \right)
\end{aligned}$$

From here, we get the following:

$$\begin{aligned}
\frac{dB_n(t)}{dt} &= \frac{2k}{L} \int_0^L \frac{\partial^2 u}{\partial x^2} \phi(x) \\
&= \frac{2k}{L} \left( -\lambda_n \left( \frac{L}{2} \right) B_n(t) - \frac{n\pi}{L} \left( (-1)^n B - A \right) \right) \\
&= -\lambda_n k B_n(t) - \frac{2kn\pi}{L^2} \left( (-1)^n B - A \right)
\end{aligned}$$

Now we multiply by  $e^{-\lambda_n kt}$  and solve for  $B_n$ , we get:

$$B_n(t) = B_n(0)e^{-\lambda_n kt} - \frac{2}{n\pi} \left( (-1)^n B - A \right) \left( 1 - e^{-\lambda_n kt} \right)$$

**Excercise 8.4.3:** Consider

$$c(x)\rho(x)\frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left[ K_0(x)\frac{\partial u}{\partial x} \right] + q(x)u + f(x, t)$$

$$u(x, 0) = g(x) \quad u(0, t) = \alpha(t) \quad u(L, t) = \beta(t)$$

Assume that the eigenfunction  $\phi_n(x)$  of the related homogeneous problem are known.

(a) Solve without reducing to a problem with homogeneous boundary conditions.

Let the following be true:

$$\sigma = c\rho \quad u(x, t) = \sum_{n=0}^{\infty} B_n(t)\phi_n(x)$$

Substituting, we get:

$$\sigma \sum_{n=0}^{\infty} B'_n(t)\phi_n(x) = \frac{\partial}{\partial x} \left[ K_0(x)\frac{\partial u}{\partial x} \right] + q(x)u + f(x, t)$$

Similar to our last problem, we can use Green's formula to get:

$$\sigma \frac{dB_n(t)}{dt} = -\lambda_n k B_n(t) + f_n(t) - \frac{k\sqrt{\lambda_n}((-1)^n\beta(t) - \alpha(t))}{\int_0^L \phi_n^2(x)\sigma dx}$$

with the following:

$$f_n(t) = \frac{\int_0^L f(x, t)\phi_n(x)\sigma dx}{\int_0^L \phi_n^2(x)\sigma dx}$$

Multiplying by  $e^{\frac{\lambda_n}{\sigma}t}$  and solving for  $B_n$ , we get:

$$B_n(t) = e^{-\frac{\lambda_n}{\sigma}t} \left( \frac{1}{\sigma} \int e^{\frac{\lambda_n}{\sigma}t} f_n(t) dt + \frac{\frac{k}{n}\sqrt{\lambda_n}}{\int_0^L \phi_n^2(x)\sigma, dx} \int -e^{\frac{\lambda_n}{\sigma}t} ((-1)^n\beta(t) - \alpha(t)) dt \right) + B_n(0)e^{-\frac{\lambda_n}{\sigma}t}$$

with the following coefficients found from the initial condition:

$$B_n(0) = \frac{\int_0^L g(x)\phi_n(x)\sigma dx}{\int_0^L \phi_n^2(x)\sigma, dx}$$

(b) Solve by first reducing to a problem with homogeneous boundary conditions

Let the following be true:

$$u_E(x, t) = \alpha(t) + \frac{x}{L} (\beta(t) - \alpha(t))$$

Now we set the following and substitute:

$$v(x, t) = u(x, t) - u_E(x, t) \quad c\rho \frac{\partial v}{\partial t} = \frac{\partial}{\partial x} \left( K_0 \frac{\partial v}{\partial x} \right) + vq(x) + f(x, t)$$

with the following conditions:

$$v(0, t) = v(\pi, t) = 0 \quad v(x, 0) = g(x) - \alpha(0) - \frac{x}{L} (\beta(0) - \alpha(0))$$

We now let the following:

$$v(x, t) = \sum_{n=1}^{\infty} B_n(t) \phi_n(x) \quad \sigma = c\rho$$

Using Green's formula, we get:

$$\sigma \frac{dB_n(t)}{dt} = -\lambda_n B_n(t) + f_n(t)$$

where

$$f_n(t) = \frac{\int_0^L f(x, t) \phi_n(x) \sigma dx}{\int_0^L \phi_n^2(x) \sigma dx}$$

Multiplying by  $e^{\frac{\lambda_n}{\sigma} t}$  and solving for  $B_n$ , we get:

$$B_n(t) = e^{-\frac{\lambda_n}{\sigma} t} \frac{1}{\sigma} \int e^{\frac{\lambda_n}{\sigma} t} f_n(t) dt + B_n(0) e^{-\frac{\lambda_n}{\sigma} t}$$

with the following coefficients found from the initial condition:

$$B_n(0) = \frac{\int_0^L g(x) - \alpha(0) - \frac{x}{L} (\beta(0) - \alpha(0)) \phi_n(x) \sigma dx}{\int_0^L \phi_n^2(x) \sigma dx}$$

**Excercise 9.2.1:** Consider

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + Q(x, t) \quad u(x, 0) = g(x)$$

In all cases, obtain formulas similar to (2.20) by introducing a Green's function.

Notice equation (2.20):

$$u(x, t) = \int_0^L g(x_0) G(x, t; x_0, 0) dx_0 + \int_0^L \int_0^t Q(x_0, t_0) G(x, t; x_0, 0) dt_0 dx_0 \quad (2.20)$$

(c) Solve using any method if

$$\frac{\partial u}{\partial t}(0, t) = 0 \quad \text{and} \quad \frac{\partial u}{\partial t}(L, t) = 0$$

(d) Use Green's formula instead of term-by-term differentiation if

$$\frac{\partial u}{\partial x}(0, t) = A(t) \quad \text{and} \quad \frac{\partial u}{\partial x}(L, t) = \beta(t)$$