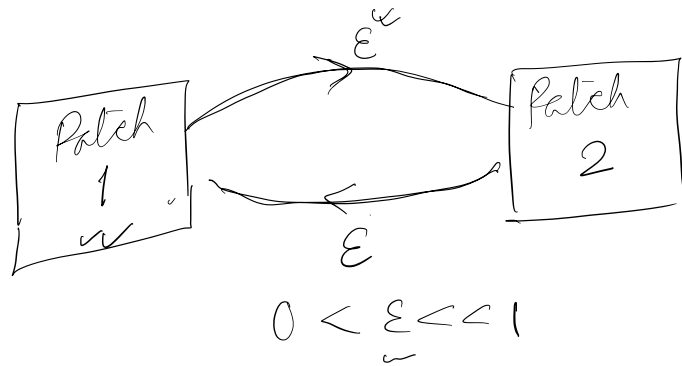


Oct 07, 2024

Goal: To study the impact of spatial dispersal on the dynamics of the whole environment.

Special case:



Model:

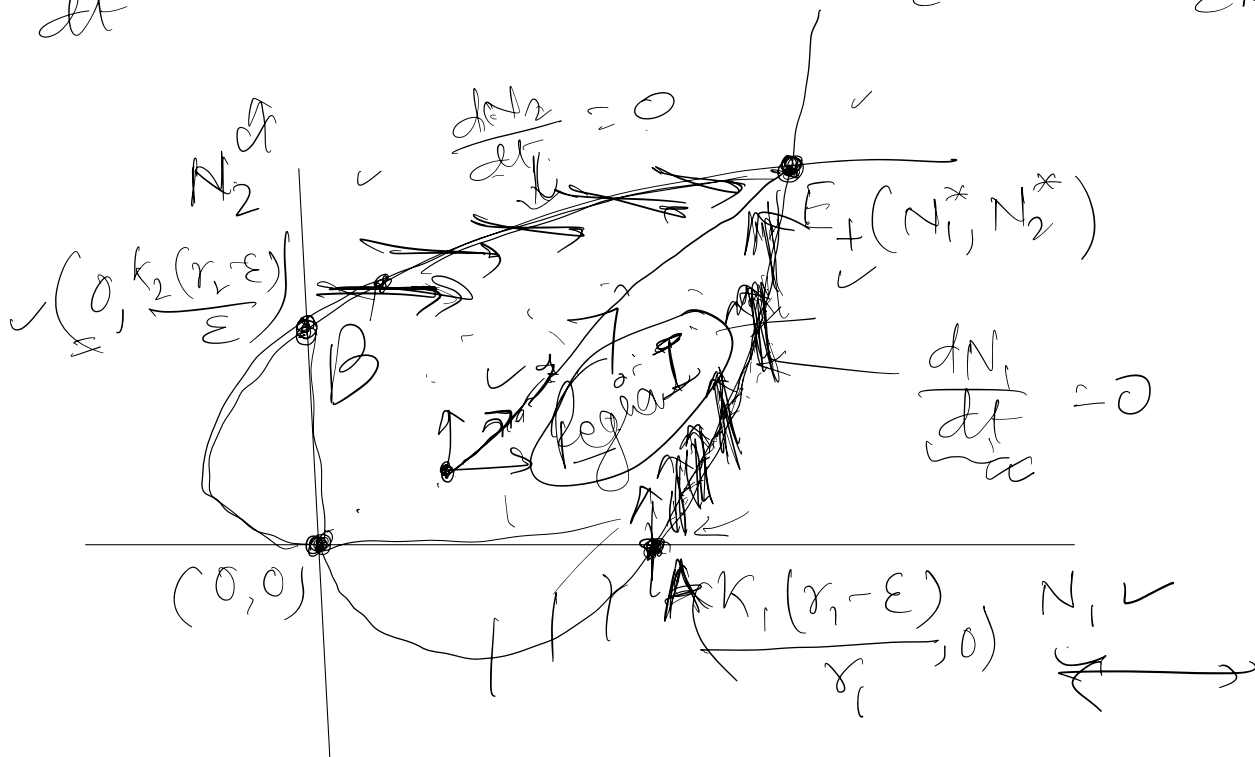
$$\checkmark \frac{dN_1}{dt} = r_1 N_1 \left(1 - \frac{N_1}{K_1}\right) + \varepsilon (N_2 - N_1) \quad \text{--- (I)}$$

$$\checkmark \frac{dN_2}{dt} = r_2 N_2 \left(1 - \frac{N_2}{K_2}\right) + \varepsilon (N_1 - N_2) \quad \text{--- (II)}$$

Analysis:

$$\frac{dN_1}{dt} = 0 \Rightarrow N_2 = -\frac{(r_1 - \varepsilon)}{\varepsilon} N_1 + \frac{r_1}{\varepsilon K_1} N_1^2$$

$$\frac{dN_2}{dt} = 0 \Rightarrow N_1 = -\frac{(r_2 - \varepsilon)}{\varepsilon} N_2 + \frac{r_2}{\varepsilon K_2} N_2^2$$



Equilibria : intersection of the curve
 $\frac{dN_1}{dt} = 0$ and $\frac{dN_2}{dt} = 0$.

\Rightarrow exactly two equilibria:

$$E_0 = (0, 0), \quad E_+ = (N_1^*, N_2^*)$$

$\frac{dN_2}{dt} > 0$ on the curve $\frac{dN_1}{dt} = 0$
between A and E_+

✓

$$\left[\because \text{put } N_2 = 0 \text{ in (1)}, \right. \\ \left. \frac{dN_2}{dt} = \varepsilon N_1 > 0 \right]$$

Similarly,

$$\frac{dN_1}{dt} > 0 \text{ on the curve } \frac{dN_2}{dt} = 0$$

between B and E_+ .

If $(N_1(0), N_2(0)) \in \text{Region (I)}$,
where $\frac{dN_1}{dt} > 0, \frac{dN_2}{dt} > 0, (N_1(t), N_2(t))$
can not escape from Region (I).

$$\Rightarrow N_1 \uparrow, N_2 \uparrow$$

$$N_1 \leq N_1^*, N_2 \leq N_2^*$$

$$\Rightarrow \lim_{t \rightarrow \infty} N_1(t) = N_1^*, \quad \lim_{t \rightarrow \infty} N_2(t) = N_2^*.$$

Conclusions:

I. All solutions (except the equilibrium E_0) converge to E_+ as $t \rightarrow \infty$.

II. E_0 is unstable and E_+ is asymptotically stable.

Question: What is E_+ and what is the impact of ε on E_+ ?

One way:

$$\begin{aligned} \checkmark \quad & r_1 N_1 \left(1 - \frac{N_1}{K}\right) + \varepsilon (N_2 - N_1) = 0 \\ \checkmark \quad & r_2 N_2 \left(1 - \frac{N_2}{K}\right) + \varepsilon (N_2 - N_1) = 0 \end{aligned} \quad \left. \begin{array}{l} \text{Solve} \\ \text{Solve} \end{array} \right\}$$

— $N_2 = \text{quadratic in } N_1 \text{ from 1st eq.}$

— substitute into the 2nd eq.

→ polynomial of N_1 of order 4 = 0

→ polynomial of N_1 of order 3 = 0 as $N_1 = 0$

⇒ solve for $N_1 \Rightarrow N_2$.

Alternate: Asymptotic Analysis

Try to find:

$$N_1 = K_1 + \varepsilon x_1 + \dots$$

$$N_2 = K_2 + \varepsilon x_2 + \dots$$

$$\begin{cases} \gamma_1 [K_1 + \varepsilon x_1] \left[1 - \frac{K_1 + \varepsilon x_1}{K_1} \right] + \varepsilon [K_2 - K_1 + \varepsilon(x_1 - x_2)] = 0 \\ \gamma_2 [K_2 + \varepsilon x_2] \left[1 - \frac{K_2 + \varepsilon x_2}{K_2} \right] + \varepsilon [K_1 - K_2 + \varepsilon(x_1 - x_2)] = 0 \end{cases}$$

$$\Rightarrow -\frac{\varepsilon x_1}{K_1} \gamma_1 [K_1 + \varepsilon x_1] + \varepsilon [K_2 - K_1 + \varepsilon(x_1 - x_2)] = 0$$

$$\gamma_1 - \frac{x_1}{K_1} \gamma_1 [K_1 + \varepsilon x_1] + K_2 - K_1 + \varepsilon(x_1 - x_2) = 0$$

$$\varepsilon \ll 1 \Rightarrow -\gamma_1 x_1 + K_2 - K_1 \quad [\text{ignoring } \varepsilon]$$

$$\Rightarrow x_1 = \frac{K_2 - K_1}{\gamma_1}$$

Similarly, $x_2 = \frac{K_1 - K_2}{\gamma_2}$

$$\left. \begin{aligned} N_1 &\sim K_1 + \varepsilon \frac{K_2 - K_1}{\gamma_1} \\ N_2 &\sim K_2 + \varepsilon \frac{K_1 - K_2}{\gamma_2} \end{aligned} \right\} \xrightarrow{\text{if } K_1 > K_2} \begin{aligned} N_1 &\sim < K_1, N_1 \downarrow \text{ as } \varepsilon \uparrow \\ N_2 &\sim > K_2, N_2 \uparrow \text{ as } \varepsilon \uparrow \end{aligned}$$

Higher-order ODEs: Simple Pendulum (Harmonic)

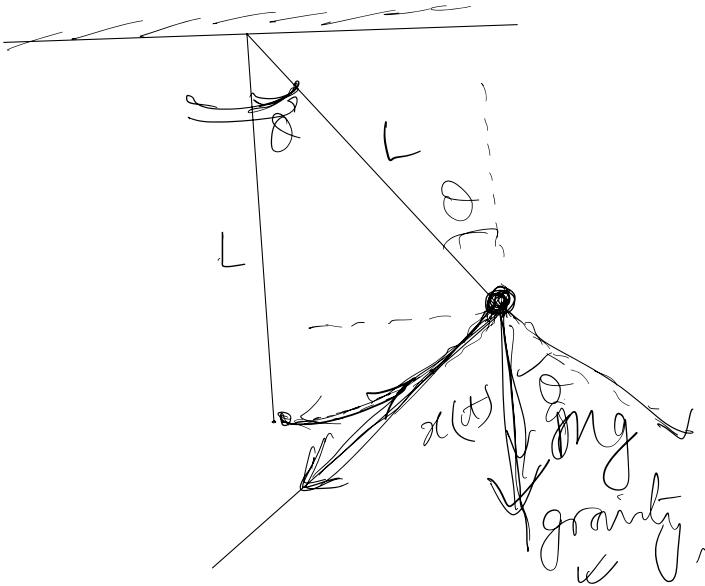
Objective: - model motion of pendulum (periodic)
- identify period of the motion.

Assumption:

- Friction is negligible
- swings in a perfect plane
- arm can not bend or stretch/

compress

- arm is massless
- gravity is constant ($g = 9.8 \text{ m/s}^2$)



m = mass

g = acceleration due to gravity

L = length

θ = angle between string position with vertical line (rest)

t = time

T = period of pendulum.

$$x(t) = L\theta(t)$$

$$\Rightarrow \text{acceleration, } a = \frac{d^2 x}{dt^2} = L \frac{d^2 \theta}{dt^2}$$

$$\text{Force, } F = -mg \sin \theta$$

$$\text{Newton's 2nd law: } F = ma$$

$$\Rightarrow -mg \sin \theta = m \cdot L \frac{d^2 \theta}{dt^2}$$

$$\Rightarrow \boxed{\frac{d^2 \theta}{dt^2} = -\frac{g}{L} \sin \theta}$$

$$\left\{ \begin{array}{l} \frac{d^2 \theta}{dt^2} + \frac{g}{L} \sin \theta = 0 \quad \left(\text{2nd order ODE} \right) \\ \theta(0) = \theta_0 \quad \checkmark \\ \theta'(0) = 0 \quad \checkmark \end{array} \right.$$

$$\equiv \left\{ \begin{array}{l} x_1(t) = \theta(t) \\ x_2(t) = \frac{d\theta}{dt} \end{array} \right.$$

$$\Rightarrow \left\{ \begin{array}{l} \frac{dx_1(t)}{dt} = x_2(t), \quad x_1(0) = \theta_0 \\ \frac{dx_2(t)}{dt} = -\frac{g}{L} \sin x_1(t), \quad x_2(0) = 0 \end{array} \right.$$

Solve
 $\xrightarrow{?}$

$$x_1(t) \equiv \theta(t)$$

• small angle approximation:

$$\theta \approx \sin(\theta), \quad \forall \theta, \quad x_1 \approx \sin x_1$$

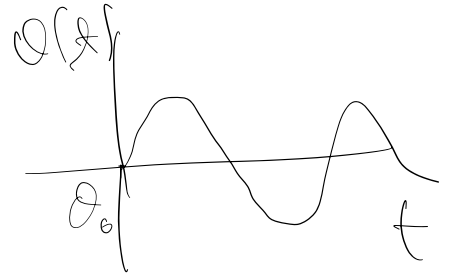
$$\Rightarrow \frac{d^2 \theta}{dt^2} + \frac{g}{L} \sin \theta = 0$$

$$\Rightarrow \frac{d^2 \theta}{dt^2} + \frac{g}{L} \theta = 0$$

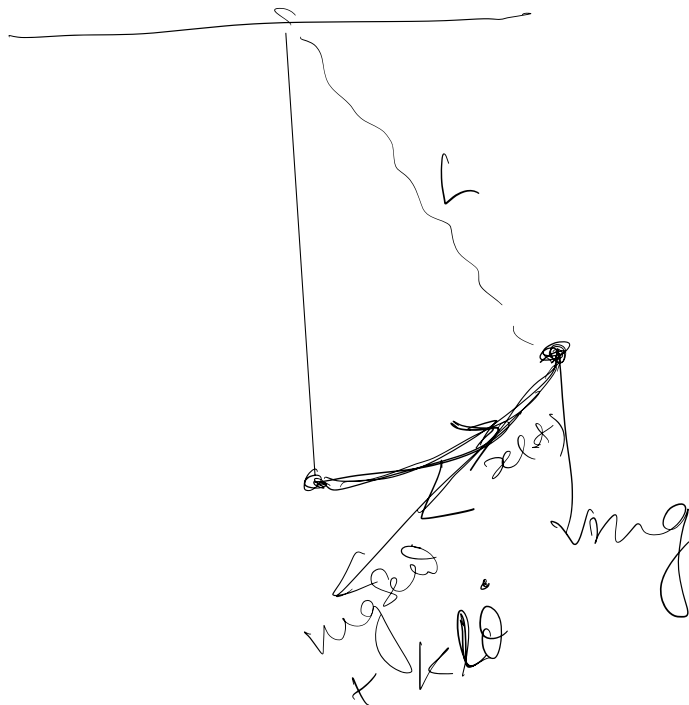
Solve (Lin Diff. Eq)

$$\theta(t) \approx \theta_0 \cos\left(\underbrace{\sqrt{\frac{g}{L}} t}_{\omega} \right) \quad \omega = \frac{2\pi}{T}$$

$$\Rightarrow T \approx 2\pi \underbrace{\sqrt{\frac{L}{g}}}$$



□ Damped pendulum (oscillation)



$$\therefore \underbrace{\frac{d^2\theta}{dt^2}} + \underbrace{\frac{g}{L} \sin\theta} + \frac{k}{m} \frac{d\theta}{dt} = 0$$

$$\Rightarrow \begin{cases} \frac{dx_1}{dt} = x_2 \\ \frac{dx_2}{dt} = -\frac{g}{L} \sin x_1 - \frac{k}{m} x_2 \end{cases}$$

