Math 337 - Elementary Differential Equations Lecture Notes - Laplace Transforms: Part B

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Outline

- Inverse Laplace Transforms
 - Solving Differential Equations
 - Laplace for Systems of DEs
 - Laplace Transforms and Maple
- 2 Special Functions
 - Heaviside or Step function
 - Periodic functions
 - Impulse or δ Function



Inverse Laplace Transforms

Theorem (Inverse Laplace Transform)

If f(t) and g(t) are piecewise continuous and have exponential order with exponent a on $[0,\infty)$ and F=G, where $F=\mathcal{L}[f]$ and $G=\mathcal{L}[g]$, then f(t)=g(t) at all points where both f and g are continuous. In particular, f and g are continuous on $[0,\infty)$, then f(t)=g(t) for all $t\in [0,\infty)$.

The functions may disagree at points of discontinuity

Definition (Inverse Laplace Transform)

If f(t) is piecewise continuous and has exponential order with exponent a on $[0, \infty)$ and $\mathcal{L}[f(t)] = F(s)$, then we call f the **inverse Laplace transform** of F, and denote it by

$$f(t) = \mathcal{L}^{-1}[F(s)].$$



Linearity of Inverse Laplace Transforms

Theorem (Linearity of Inverse Laplace Transform)

Assume that $f_1 = \mathcal{L}^{-1}[F_1]$ and $f_2 = \mathcal{L}^{-1}[F_2]$ are piecewise continuous and has exponential of order with exponent a on $[0, \infty)$. Then for any constants c_1 and c_2 ,

$$\mathcal{L}^{-1}[c_1F_1 + c_2F_2] = c_1\mathcal{L}^{-1}[F_1] + c_2\mathcal{L}^{-1}[F_2] = c_1f_1 + c_2f_2.$$

Example: Find
$$\mathcal{L}^{-1}\left[\frac{2}{(s-3)^4} + \frac{12}{s^2+16} + \frac{5(s+2)}{s^2+4s+5}\right]$$
.

Rewrite as

$$\frac{1}{3}\mathcal{L}^{-1}\left[\frac{3!}{(s-3)^4}\right] + 3\mathcal{L}^{-1}\left[\frac{4}{s^2 + 16}\right] + 5\mathcal{L}^{-1}\left[\frac{(s+2)}{(s+2)^2 + 1}\right]$$

With Exponential Shift Theorem

$$\frac{1}{3}e^{3t}t^3 + 3\sin(4t) + 5e^{-2t}\cos(t)$$



Example: Consider the initial value problem:

$$y'' + y = e^{-t}\cos(2t)$$
 with $y(0) = 2$, $y'(0) = 1$

Let $Y(s) = \mathcal{L}[y(t)]$, then taking Laplace transforms gives

$$s^{2}Y(s) - sy(0) - y'(0) + Y(s) = \frac{s+1}{(s+1)^{2} + 4}$$

or

$$(s^2 + 1)Y(s) = 2s + 1 + \frac{s+1}{(s+1)^2 + 4}$$

Equivalently,

$$Y(s) = \frac{2s+1}{s^2+1} + \frac{s+1}{(s^2+1)((s+1)^2+4)}$$



Example: Since

$$Y(s) = \frac{2s+1}{s^2+1} + \frac{s+1}{(s^2+1)((s+1)^2+4)},$$

and the first term is already in simplest form, partial fractions decomposition gives

$$\frac{s+1}{(s^2+1)((s+1)^2+4)} = \frac{As+B}{s^2+1} + \frac{C(s+1)+D\cdot 2}{(s+1)^2+4}.$$

It follows that

$$s+1 = (As+B)(s^2+2s+5) + (C(s+1)+2D)(s^2+1)$$

Let s = i, then

$$1 + i = (B + Ai)(4 + 2i) = (4B - 2A) + i(4A + 2B)$$



Example: Equating the real and imaginary parts of the previous equation give:

$$-2A + 4B = 1$$
 and $4A + 2B = 1$

Solving the **linear equations** gives $A = \frac{1}{10}$ and $B = \frac{3}{10}$

Since

$$s+1 = (As+B)(s^2+2s+5) + (C(s+1)+2D)(s^2+1),$$

the cubic (s^3) terms give 0 = A + C or $C = -\frac{1}{10}$.

The constant terms give

$$1 = 5B + C + 2D$$
 or $D = -\frac{1}{5}$



Example: Thus,

$$\begin{array}{lcl} Y(s) & = & \frac{2s+1}{s^2+1} + \frac{s+1}{(s^2+1)((s+1)^2+4)}, \\ Y(s) & = & \frac{2s+1}{s^2+1} + \frac{s/10}{s^2+1} + \frac{3/10}{s^2+1} - \frac{(s+1)/10}{(s+1)^2+4} - \frac{2/5}{(s+1)^2+4} \\ Y(s) & = & \frac{\left(\frac{21}{10}\right)s}{s^2+1} + \frac{\left(\frac{13}{10}\right)\cdot 1}{s^2+1} - \frac{\left(\frac{1}{10}\right)(s+1)}{(s+1)^2+4} - \frac{\left(\frac{1}{5}\right)\cdot 2}{(s+1)^2+4} \end{array}$$

This last line allows easy application of the Inverse Laplace Transform to obtain the solution

$$y(t) = \frac{21}{10}\cos(t) + \frac{13}{10}\sin(t) - \frac{1}{10}e^{-t}\cos(2t) - \frac{1}{5}e^{-t}\sin(2t).$$



Laplace for Systems of DEs

Laplace for Systems of Differential Equations: Consider

$$\dot{y}_1 = a_{11}y_1 + a_{12}y_2 + f_1(t), \qquad y_1(0) = y_{10}$$

 $\dot{y}_2 = a_{21}y_1 + a_{22}y_2 + f_2(t), \qquad y_2(0) = y_{20}$

Taking Laplace transforms gives

$$sY_1 - y_1(0) = a_{11}Y_1 + a_{12}Y_2 + F_1(s),$$

 $sY_2 - y_2(0) = a_{21}Y_1 + a_{22}Y_2 + F_2(s),$

which can be written

$$(s - a_{11})Y_1 - a_{12}Y_2 = y_1(0) + F_1(s),$$

$$-a_{21}Y_1 + (s - a_{22})Y_2 = y_2(0) + F_2(s).$$

Or in matrix form

$$(s\mathbf{I} - \mathbf{A})\mathbf{Y} = \mathbf{v}_0 + \mathbf{F}(s)$$





Laplace for Systems of DEs

Laplace for Systems of Differential Equations: The matrix form is

$$(s\mathbf{I} - \mathbf{A})\mathbf{Y} = \mathbf{y}_0 + \mathbf{F}(s),$$

which is readily solved as

$$\mathbf{Y} = (s\mathbf{I} - \mathbf{A})^{-1}\mathbf{y}_0 + (s\mathbf{I} - \mathbf{A})^{-1}\mathbf{F}(s)$$

The inverse satisfies

$$(s\mathbf{I} - \mathbf{A})^{-1} = \frac{1}{|s\mathbf{I} - \mathbf{A}|} \begin{pmatrix} s - a_{22} & a_{12} \\ a_{21} & s - a_{11} \end{pmatrix},$$

where

$$|s\mathbf{I} - \mathbf{A}| = s^2 - (a_{11} + a_{22})s + a_{11}a_{22} - a_{12}a_{21}$$

is the characteristic polynomial

This is not hard to solve algebraically, but the **inverse Laplace** transform may be messy



Example: Consider the nonhomogeneous system

$$\dot{\mathbf{y}} = \begin{pmatrix} -4 & -1 \\ 1 & -2 \end{pmatrix} \mathbf{y} + \begin{pmatrix} 2e^t \\ \sin(2t) \end{pmatrix}, \qquad \mathbf{y}_0 = \begin{pmatrix} 1 \\ 2 \end{pmatrix}$$

Taking Laplace transforms gives

$$sY_1 - 1 = -4Y_1 - Y_2 + \frac{2}{s-1}$$

 $sY_2 - 2 = Y_1 - 2Y_2 + \frac{2}{s^2 + 4}$

Equivalently,

$$\left(\begin{array}{cc} s+4 & 1 \\ -1 & s+2 \end{array}\right) \left(\begin{array}{c} Y_1 \\ Y_2 \end{array}\right) = \left(\begin{array}{c} 1+\frac{2}{s-1} \\ 2+\frac{2}{s^2+4} \end{array}\right)$$



Example: If

$$(s\mathbf{I} - \mathbf{A}) = \begin{pmatrix} s+4 & 1 \\ -1 & s+2 \end{pmatrix}, \quad \mathbf{y}_0 + \mathbf{F}(s) = \begin{pmatrix} 1 + \frac{2}{s-1} \\ 2 + \frac{2}{s^2+4} \end{pmatrix},$$

then the solution is

$$\begin{pmatrix} Y_1(s) \\ Y_2(s) \end{pmatrix} = (s\mathbf{I} - \mathbf{A})^{-1}(\mathbf{y}_0 + \mathbf{F}(s)),$$

where

$$(s\mathbf{I} - \mathbf{A})^{-1} = \frac{1}{(s+3)^2} \begin{pmatrix} s+2 & -1\\ 1 & s+4 \end{pmatrix}$$

The expressions for $Y_1(s)$ and $Y_2(s)$ are fairly complex, so we show how Maple can help solve these expressions into a form, which readily has an inverse Laplace transform



Example: From the previous page, it is easy to see that

$$Y_1(s) = \frac{s}{(s+3)^2} + \frac{2(s+2)}{(s-1)(s+3)^2} - \frac{2}{(s^2+4)(s+3)^2}$$

$$Y_2(s) = \frac{2s+9}{(s+3)^2} + \frac{2}{(s-1)(s+3)^2} + \frac{2(s+4)}{(s^2+4)(s+3)^2}$$

One can perform **Partial Fractions Decomposition** on these expressions, which is a very messy process.

We demonstrate how this can be done with Maple, which can readily perform both Partial Fractions Decompositions and inverse Laplace transforms

A **Special Maple Sheet** is provided along with a complete solution using **Maple**



Example: Maple gives the Laplace transform

$$Y_1(s) = \frac{\frac{3}{8}}{s-1} + \frac{\frac{12s-10}{169}}{s^2+4} - \frac{\frac{69}{26}}{(s+3)^2} + \frac{\frac{749}{1352}}{s+3}$$

$$Y_2(s) = \frac{\frac{1}{8}}{s-1} + \frac{\frac{88-38s}{169}}{s^2+4} + \frac{\frac{69}{26}}{(s+3)^2} + \frac{\frac{2839}{1352}}{s+3}$$

The inverse Laplace transform gives the solution

$$y_1(t) = \frac{3}{8}e^t + \frac{12}{169}\cos(2t) - \frac{5}{169}\sin(2t) - \frac{69}{26}te^{-3t} + \frac{749}{1352}e^{-3t}$$
$$y_2(t) = \frac{1}{8}e^t + \frac{44}{169}\sin(2t) - \frac{38}{169}\cos(2t) + \frac{69}{26}te^{-3t} + \frac{2839}{1352}e^{-3t}$$



Discontinuous Functions

Unit Step function or Heaviside function satisfies

$$u(t) = \begin{cases} 0, & t < 0, \\ 1, & t \ge 0. \end{cases}$$

The translated version of the **Unit Step function** by c units is

$$u_c(t) = \begin{cases} 0, & t < c, \\ 1, & t \ge c, \end{cases}$$

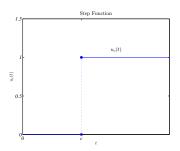
which represents a switch turning on at t = c

An indicator function, which is on for $c \le t < d$, satisfies

$$u_{cd}(t) = u_c(t) - u_d(t) = \begin{cases} 0, & t < c \text{ or } t \ge d, \\ 1, & c \le t < d, \end{cases}$$



Laplace Transform of Step Function



Laplace Transform of Step Function;

$$\mathcal{L}[u_c(t)] = \int_0^\infty e^{-st} u_c(t) dt = \int_c^\infty e^{-st} dt = \lim_{A \to \infty} \int_c^A e^{-st} dt$$
$$= \lim_{A \to \infty} \left(\frac{e^{-cs}}{s} - \frac{e^{-sA}}{s} \right) = \frac{e^{-cs}}{s}, \quad s > 0$$



Laplace Transform of Step Function

Theorem

If $F(s) = \mathcal{L}[f(t)]$ exists for $s > a \ge 0$, and if c is a nonnegative constant, then

$$\mathcal{L}[u_c(t)f(t-c)] = e^{-cs}\mathcal{L}[f(t)] = e^{-cs}F(s), \qquad s > a.$$

Conversely, if $f(t) = \mathcal{L}^{-1}[F(s)]$, then

$$u_c(t)f(t-c) = \mathcal{L}^{-1}[e^{-cs}F(s)].$$

This theorem states that the translation of f(t) a distance c in the positive t direction corresponds to the multiplication of F(s) by e^{-cs}



Example with Step Function

Example with Step Function: Consider the following initial value problem:

$$y'' + 2y' + 5y = u_2(t) - u_5(t),$$
 $y(0) = 0,$ $y'(0) = 0.$

Take Laplace transforms and obtain

$$(s^{2} + 2s + 5)Y(s) = \frac{e^{-2s}}{s} - \frac{e^{-5s}}{s}$$

This rearranges to

$$Y(s) = \frac{e^{-2s} - e^{-5s}}{s(s^2 + 2s + 5)}$$

Partial fraction decomposition gives

$$\frac{1}{s(s^2+2s+5)} = \frac{A}{s} + \frac{B(s+1)+2C}{(s+1)^2+4}$$



Example with Step Function

Example with Step Function: Partial fraction decomposition gives

$$1 = A((s+1)^2 + 4) + (B(s+1) + 2C) s$$

With s = 0, 1 = 5A or $A = \frac{1}{5}$

The s^2 coefficient gives 0 = A + B, so $B = -\frac{1}{5}$

The s^1 coefficient gives 0 = 2A + B + 2C, so $C = -\frac{1}{10}$

Thus,

$$Y(s) = \left(\frac{\frac{1}{5}}{s} - \frac{\frac{1}{5}(s+1) + \frac{2}{10}}{(s+1)^2 + 4}\right) (e^{-2s} - e^{-5s})$$



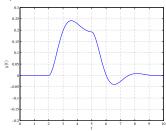
Example with Step Function

Example with Step Function: With the Laplace transform

$$Y(s) = \left(\frac{\frac{1}{5}}{s} - \frac{\frac{1}{5}(s+1) + \frac{2}{10}}{(s+1)^2 + 4}\right)(e^{-2s} - e^{-5s})$$

The theorem for step functions allows the inverse Laplace transform yielding

$$y(t) = \frac{u_2(t)}{10} \left(2 - 2e^{-(t-2)} \cos(2(t-2)) - e^{-(t-2)} \sin(2(t-2)) \right) - \frac{u_5(t)}{10} \left(2 - 2e^{-(t-5)} \cos(2(t-5)) - e^{-(t-5)} \sin(2(t-5)) \right)$$





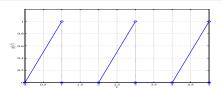
Periodic Functions

Definition

A function f is said to be **periodic with period** T > 0 if

$$f(t+T) = f(t)$$

for all t in the domain of f.



A sawtooth waveform

$$f(t) = \begin{cases} t, & 0 \le t < 1, \\ 0, & 1 < t < 2. \end{cases}$$

and f(t) has period 2



Periodic and Window functions

Consider a **periodic function** f(t). Define the **window function**, $f_T(t)$, as follows:

$$f_T(t) = f(t) [1 - u_T(t)] = \begin{cases} f(t), & 0 \le t \le T \\ 0, & \text{otherwise} \end{cases}$$

The Laplace transform $F_T(s)$ satisfies:

$$F_T(s) = \int_0^\infty e^{-st} f_T(t) dt = \int_0^T e^{-st} f_T(t) dt.$$

The **window function** specifies values of f(t) over a single period

This can be replicated k periods to the right as

$$f_T(t - kT)u_{kT}(t) = \begin{cases} f(t - kT), & kT \le t \le (k+1)T \\ 0, & \text{otherwise} \end{cases}$$



Laplace for Periodic Functions

By summing n time shifted replications of the **window function**, $f_T(t - kT)u_{kT}(t)$, k = 0,...,n - 1, gives $f_{nT}(t)$, the periodic extension of $f_T(t)$ to the interval [0, nT],

$$f_{nT}(t) = \sum_{k=0}^{n-1} f_T(t - kT)u_{kT}(t)$$

Theorem

If f is periodic with period T and is piecewise continuous on [0,T], then

$$\mathcal{L}[f(t)] = \frac{F_T(s)}{1 - e^{-sT}} = \frac{\int_0^T e^{-st} f(t) dt}{1 - e^{-sT}}.$$



Laplace for Periodic Functions

Proof: From our earlier theorem, we have for each $k \geq 0$,

$$\mathcal{L}[f_T(t-kT)u_{kT}(t)] = e^{-kTs}\mathcal{L}[f_T(t)] = e^{-kTs}F_T(s).$$

By linearity of \mathcal{L} , the **Laplace transform** of f_{nT} is

$$F_{nT}(s) = \int_0^{nT} e^{-st} f(t) dt = \sum_{k=0}^{n-1} \mathcal{L}[f_T(t-kT)u_{kT}(t)]$$

$$= \sum_{k=0}^{n-1} e^{-kTs} F_T(s) = F_T(s) \sum_{k=0}^{n-1} \left(e^{-Ts}\right)^k = F_T(s) \frac{1 - \left(e^{-Ts}\right)^n}{1 - e^{-sT}}.$$

The last term comes from summing a geometric series. With $e^{-sT} < 1$.

$$F(s) = \lim_{n \to \infty} \int_0^{nT} e^{-st} dt = \lim_{n \to \infty} F_T(s) \frac{1 - \left(e^{-Ts}\right)^n}{1 - e^{-sT}} = \frac{F_T(s)}{1 - e^{-sT}}$$



Sawtooth Function

Return to **sawtooth** waveform

$$f(t) = \begin{cases} t, & 0 \le t < 1, \\ 0, & 1 \le t < 2. \end{cases}$$
 and $f(t)$ has period 2

The **theorem** for the **Laplace transform** of **periodic function** gives

$$\mathcal{L}[f(t)] = \frac{\int_0^2 e^{-st} f(t)dt}{1 - e^{-2s}}$$

But

$$\int_0^2 e^{-st} f(t)dt = \int_0^1 t e^{-st} dt = \frac{1 - se^{-s} - e^{-s}}{s^2},$$

SO

$$\mathcal{L}[f(t)] = \frac{1 - se^{-s} - e^{-s}}{s^2 (1 - e^{-2s})}$$



IVP with Periodic Forcing Function

Example: Consider the following initial value problem:

$$y'' + 4y = f(t),$$
 $y(0) = 0,$ $y'(0) = 0,$

with the **square** waveform as the **periodic forcing function**:

$$f(t) = \begin{cases} 1, & 0 \le t < 1, \\ 0, & 1 \le t < 2. \end{cases}$$
 and $f(t)$ has period 2

The **theorem** for the **Laplace transform** of **square** waveform gives

$$\mathcal{L}[f(t)] = \frac{\int_0^2 e^{-st} f(t)dt}{1 - e^{-2s}}$$

But

$$\int_0^2 e^{-st} f(t)dt = \int_0^1 e^{-st} dt = \frac{1 - e^{-s}}{s},$$

SO

$$\mathcal{L}[f(t)] = \frac{1 - e^{-s}}{s(1 - e^{-2s})} = \frac{1}{s(1 + e^{-s})}$$



IVP with Periodic Forcing Function

Example: Taking the Laplace transform of the **IVP** with $\mathcal{L}[y(t)] = Y(s)$, we have:

$$s^{2}Y(s) - sy(0) - y'(0) + 4Y(s) = \frac{1}{s(1 + e^{-s})}$$

Thus,

$$Y(s) = \frac{1}{s(s^2 + 4)(1 + e^{-s})}$$

Partial fractions decomposition gives

$$\frac{1}{s(s^2+4)} = \frac{1/4}{s} - \frac{s/4}{s^2+4},$$

while

$$\frac{1}{1+e^{-s}} = \frac{1}{1-(-e^{-s})} = 1 - e^{-s} + e^{-2s} - \dots + (-1)^n e^{-ns} +$$



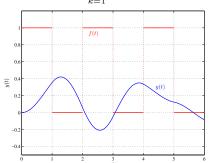
IVP with Periodic Forcing Function

Example: So

$$Y(s) = \frac{1}{4} \left(\frac{1}{s} - \frac{s}{s^2 + 4} \right) \sum_{k=0}^{\infty} (-1)^k e^{-ks}$$

Taking the inverse Laplace transform gives:

$$y(t) = \frac{1}{4} (1 - \cos(2t)) + \frac{1}{4} \sum_{k=1}^{\infty} (-1)^k u_k(t) (1 - \cos(2(t-k)))$$





Impulse Function

Impulse Function: Some applications have phenomena of an **impulsive nature**, *e.g.*, a large magnitude force over a very short time

$$ay'' + by' + cy = g(t),$$

where g(t) is very large for $t \in [t_0, t_0 + \varepsilon)$ and is otherwise zero

Example: Let $t_0 = 0$ be a real number and ε be a small positive constant

Suppose $t_0 = 0$ and $g(t) = I_0 \delta_{\varepsilon}(t)$, where

$$\delta_{\varepsilon}(t) = \frac{u_0(t) - u_{\varepsilon}(t)}{\varepsilon} = \begin{cases} \frac{1}{\varepsilon}, & 0 \le t < \varepsilon, \\ 0, & t < 0 \text{ or } t \ge \varepsilon. \end{cases}$$

Consider the **mass-spring** system m = 1, $\gamma = 0$, and k = 1

$$y'' + y = I_0 \delta_{\varepsilon}(t), \quad y(0) = 0, \quad y'(0) = 0.$$



Impulse Function - Example

Example: With $\delta_{\varepsilon}(t) = \frac{u_0(t) - u_{\varepsilon}}{\varepsilon}$, the **Laplace transform** is easy for

$$y'' + y = I_0 \delta_{\varepsilon}(t), \qquad y(0) = 0, \quad y'(0) = 0.$$

It satisfies

$$(s^2+1)Y(s) = \frac{I_0}{\varepsilon} \left(\frac{1-e^{-\varepsilon s}}{s}\right),$$

SO

$$Y(s) = \frac{I_0}{\varepsilon} \left(\frac{1}{s} - \frac{s}{s^2 + 1} \right) \left(1 - e^{-\varepsilon s} \right)$$

The inverse Laplace transform gives

$$y_{\varepsilon}(t) = \frac{I_0}{\varepsilon} \left(u_0(t)(1 - \cos(t)) - u_{\varepsilon}(t)(1 - \cos(t - \varepsilon)) \right)$$



Impulse Function - Example

Example: Since

$$y_{\varepsilon}(t) = \frac{I_0}{\varepsilon} \left(u_0(t)(1 - \cos(t)) - u_{\varepsilon}(t)(1 - \cos(t - \varepsilon)) \right),$$

equivalently:

$$y_{\varepsilon}(t) = \begin{cases} 0, & t < 0, \\ \frac{I_0}{\varepsilon} (1 - \cos(t)) & 0 \le t < \varepsilon, \\ \frac{I_0}{\varepsilon} (\cos(t - \varepsilon) - \cos(t)) & t \ge \varepsilon. \end{cases}$$

The limiting case is

$$y_0(t) = \lim_{\varepsilon \to 0} y_{\varepsilon}(t) = u_0(t)I_0 \sin(t) = \begin{cases} 0, & t < 0, \\ I_0 \sin(t), & t \ge 0. \end{cases}$$



Unit Impulse Function

Unit Impulse Function: Rather than using the definition of $\delta_{\varepsilon}(t-t_0)$ to model an impulse, then take the limit as $\varepsilon \to 0$, we define an idealized **unit Impulse Function**, δ

- The "function" δ imparts an impulse of magnitude **1** at $t = t_0$, but is **zero** for all other values of t
- Properties of $\delta(t-t_0)$
 - Limiting behavior:

$$\delta(t - t_0) = \lim_{\varepsilon \to 0} \delta_{\varepsilon}(t - t_0) = 0$$

2 If f is continuous for $t \in [a, b)$ and $t_0 \in [a, b)$, then

$$\int_{a}^{b} f(t)\delta(t-t_0)dt = \lim_{\varepsilon \to 0} \int_{a}^{b} f(t)\delta_{\varepsilon}(t-t_0)dt = f(t_0).$$



$$\delta(t-t_0)$$

Dirac delta function, $\delta(t-t_0)$: This is not an ordinary function in elementary calculus, and it satisfies:

$$\int_{a}^{b} \delta(t - t_0) dt = \begin{cases} 1, & \text{if} \quad t_0 \in [a, b), \\ 0, & \text{if} \quad t_0 \notin [a, b). \end{cases}$$

The Laplace transform of $\delta(t-t_0)$ follows easily:

$$\mathcal{L}[\delta(t-t_0)] = \int_0^\infty e^{-st} \delta(t-t_0) dt = e^{-st_0}$$

Note: $\mathcal{L}[\delta(t)] = 1$.

The delta function is the symbolic derivative of the Heaviside function, so

$$\delta(t - t_0) = u'(t - t_0)$$

This is rigorously true in the theory of **generalized functions** or **distributions**



Example for $\delta(t-t_0)$

Example: Consider the initial value problem:

$$y'' + 2y' + 2y = \frac{t}{\pi}\delta(t - \pi), \qquad y(0) = 0, \quad y'(0) = 1$$

The Laplace transform of the forcing function is

$$F(s) = \int_0^\infty e^{-st} \left(\frac{t}{\pi} \delta(t - \pi) \right) dt = e^{-\pi s}$$

It follows that the **Laplace transform** of the IVP is

$$s^{2}Y(s) - 1 + 2sY(s) + 2Y(s) = e^{-\pi s},$$

so

$$Y(s) = \frac{1 + e^{-\pi s}}{(s+1)^2 + 1}$$



Example for $\delta(t-t_0)$

Example: Since $Y(s) = \frac{1 + e^{-\pi s}}{(s+1)^2 + 1}$, the inverse Laplace transform satisfies:

$$y(t) = e^{-t}\sin(t) + u_{\pi}(t)e^{-(t-\pi)}\sin(t-\pi)$$

