

Lecture 27

November 16, 2021

3:10 PM

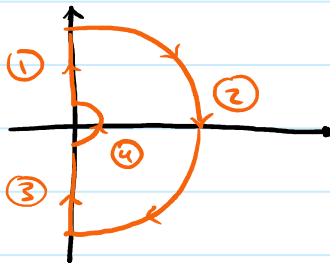
Review:

↳ Nyquist Stability

↳ Care when $G(s)$ has poles on imag axis

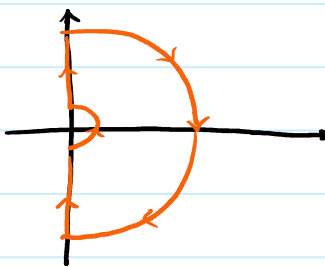
Idea:

↳ Use contours to exclude the poles on the imag axis. The Nyquist criterion then remains the same.



①, ②, ③ as before

④: $s = re^{i\theta}$ $\theta \in [-\pi/2, \pi/2]$



Suppose $G(s)$ has k poles at 0

↳ What is the effect of part 4 of the contour on the Nyquist plot

$$G(s) = \frac{N(s)}{s^k D(s)}$$

On part ④:

$$G(re^{i\theta}) = \frac{N(re^{i\theta})}{(re^{i\theta})^k D(re^{i\theta})}$$

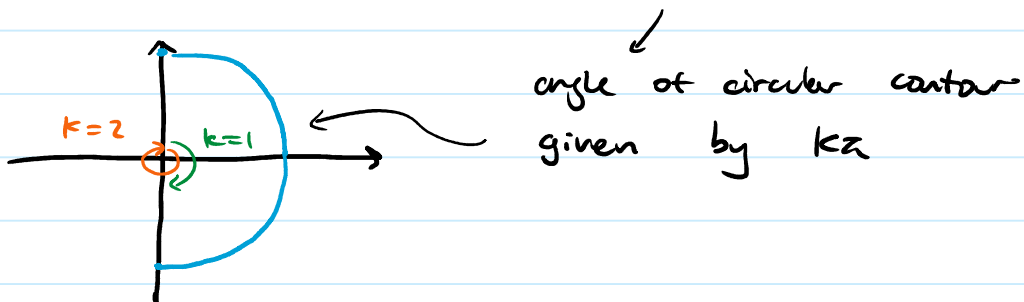
N.D.S

circle of radius $\frac{1}{r^k}$
 $\rightarrow 0$
 $1 - ik\theta$

case 1: $(re^{i\theta})^k P_c(re^{i\theta})$

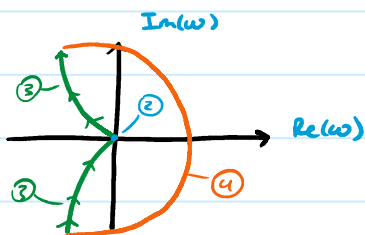
As $r \rightarrow 0$, $\frac{N_c(0)}{(re^{i\theta})^k P_c(0)} = \frac{1}{r} e^{-ik\theta} \rightarrow 0$

$\frac{N_c(0)}{P_c(0)}$ an integer describing angle of $k\pi$ oriented CW



e.g. $L(s) = \frac{10}{s(s+1)}$

(last time we drew parts of ①, ②, ③ of γ)



L has $k=1$ pole at 0

↳ part ④ will give a semicircle of ∞ radius, oriented CW

Nyquist Criterion

↳ L has 0 poles in DRHP

↳ roots of $1+L$ are in OLHP iff γ encircles -1 0 times CCW.

↳ This is the case in above example

Frequency Response

↳ In order to plot the Nyquist contour

of a TF $L(s)$ we need to plot the

complex # $L(i\omega)$ as ω varies in $(-\infty, +\infty)$

Suppose $L(s)$ is BIBO stable (i.e. all poles in OLHP).

↳ any bounded input yields bounded output

$L(i\omega)$ is the Frequency Response

$$u(t) = \sin(\omega t) \rightarrow \boxed{L(s)} \rightarrow y(t)$$

$$U(s) = \frac{\omega}{s^2 + \omega^2}$$

$$Y(s) = \underbrace{L(s)}_{\text{poles in OLHP}} \cdot \underbrace{\frac{\omega}{s^2 + \omega^2}}_{\text{poles at } \pm i\omega}$$
$$y(t) = \sum_i \text{Res}(Y(s)e^{st}, p_i)$$

$$= \sum \text{Res}(Ye^{st}, \text{poles of } L \text{ in OLHP}) \leftarrow \tilde{y}(t)$$
$$+ \underbrace{\sum \text{Res}(Ye^{st}, s = i\omega) + (\bar{*})}_{*} \leftarrow \bar{y}(t)$$

Important Property

↳ each residue in $\tilde{y}(t)$ gives a function that tends to 0 exponentially. Say that a general pole of $L(s)$ is at $-a \pm ib$ ($a > 0$).

$$\text{Res}(Y(s)e^{st}, -a \pm ib) = \dots e^{(-a \pm ib)t} = e^{-at} e^{\pm ibt} \rightarrow 0$$

$\tilde{y}(t)$ is the transient component of $y(t)$

$\bar{y}(t)$ is the steady-state behaviour

$$y(t) \rightarrow \bar{y}(t) = \text{Res}(Y(s)e^{st}, s = +i\omega) + \text{Res}(Y(s)e^{st}, s = -i\omega)$$

$$\bar{y}(t) = \text{Res}\left(\frac{L(s)\omega}{s^2 + \omega^2} e^{st}, s = i\omega\right) + (\dots)$$

$$= \frac{L(i\omega)\cancel{\omega}e^{i\omega t}}{2i\omega} + \frac{L(-i\omega)\cancel{\omega}e^{-i\omega t}}{-2i\omega}$$

$$\angle(i\omega) = |\angle(i\omega)| e^{i\angle(i\omega)}$$

$$\bar{y}(t) = |\angle(i\omega)| \frac{e^{i(\omega t + \angle(i\omega))} - e^{i(\omega t + \angle(i\omega))}}{2i}$$

$$= |\angle(i\omega)| \sin(\omega t + \angle(i\omega))$$

More generally

$$u(t) = A \sin(\omega t + \phi)$$

$y(t)$ will converge exponentially to

$$\bar{y}(t) = A |\angle(i\omega)| \sin(\omega t + \phi + \angle(i\omega))$$

RMK:

1. If $u(t)$ sinusoidal, then $y(t)$ is too (in S.S.)
↳ same frequency
2. $|\angle(i\omega)|$ scales the input magnitude
3. $\angle(i\omega)$ shifts the phase of the input signal

∴ $\angle(i\omega)$ is the frequency response

Plotting N is hard since we don't know how $\angle(i\omega)$ moves in ω -plane

Idea:

- ↳ produce reasonably accurate plots of $|\angle(i\omega)|$ and $\angle(i\omega)$ for $\omega \geq 0$ in logarithmic scale.

↳ Bode Plots

Bode Mag Plot : $20 \log |G(j\omega)|$ vs $\log(\omega)$

Bode Phase Plot : $\angle G(j\omega)$ vs $\log(\omega)$