Partial Differential Equations - Class Notes

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1 Chapter 1

Sidenotes

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What is a PDE?
A PDE is an equation which contains partial derivatives of an unknown function and we want to find that unknown function.
Example: $F(t, x, y, z, u, \frac{\partial u}{\partial t}, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z}, \frac{\partial^2 u}{\partial t^2}, \frac{\partial^2 u}{\partial x \partial y}, \ldots) = 0$.
Note, the first partial derivatives are considered 1^{st} ordered partials
whereas the second ordered partials are considered 2^{nd} ordered partials.
The variables that are not u are considered independent variables and u is considered a dependent variable.
What PDEs do we study?
Generally, we restrict our attention to equations that model some phenomenom from physics, engineering, economics, geology, ... etc. We can use physical intuition to help guide the math.
Classification of PDEs

1. Order of PDE: Highest derivative.
   Example: $\frac{\partial^3 u}{\partial x^3} - \sin(y)u^7 = 3$ is a third order PDE.
   Example: $(\frac{\partial y}{\partial t})^5 - \frac{\partial^2 y}{\partial x \partial t} = e^x$ is a second order PDE.
2. Number of independent variables.
   Example: $\frac{du}{dt} = \frac{\partial^2 u}{\partial x^2}$ has two independent variables: t, x.
   This is the $1-D$ heat equation.
   Example: $\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} = \Delta u$ has 4 independent variables.
   This is the $3-D$ heat equation. Δu is Laplacian of u.
   Notation
   $\Delta u = \nabla^2 u = \nabla \cdot \nabla u = (\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}) \cdot (\frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z}) = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2}$
   $\Delta u = 0$ is considered Laplace's equation.
3. Linear vs non-linear
   A linear PDE is any equation of the form $L[u(x)] = f(x)$ where $f(x)$ is a known function is a linear partial differential operator.

Definition: A differential operator is any rule that takes a function as its input and returns an expression that involves the derivatives of that function.
Example:

$$\begin{aligned} u(x,t) \qquad & v(x,t) && (1)\\ O[u] &= \frac{\partial^2 u}{\partial x^2} + \sin x + \pi - 7e^{tu} && (2)\\ O[u+3v] &= \frac{\partial^2}{\partial x^2}(u+3v) + \sin x + \pi - 7e^{tu+3tv} && (3)\\ &= \frac{\partial^2 u}{\partial x^2} + 3\frac{\partial^2 v}{\partial x^2} + \sin x + \pi - 7e^{tu+3tv} && (4) \end{aligned}$$

Definition: A linear operator, L, is an operator that has the property:

$$L[au + bv] = aL[u] + bL[v] \tag{5}$$

Where a and b are constants.
Theorem: If u and v are vectors and L is linear, then L can be represented by a matrix.
Theorem: If L is linear ordinary operator, it must take the form:

$$L[u] = f_0(x)u + f_1(x)u' + f_2(x)u'' + \ldots + f_n(x)u^{(n)} \tag{6}$$

Where the f_i's are known functions.
Definition: A linear ODE is any ODE of the form where $f(x)$ is known is the following:

$$L[u] = f(x) \tag{7}$$

If $f(x) = 0$, then the equation is homogeneous. Otherwise, the equation is non-homogeneous.
Ex: $(u')^2 = 0 \Rightarrow u' = 0 \rightarrow$ linear, homogeneous.
Theorem: If L is a linear partial differential operator, it must take the form (x is a vector with n unknowns)

$$L[u(x)] = f_0(x)u + \sum_{i=1}^{n} f_i(x)\frac{\partial u}{\partial x_i} + \sum_{i=1}^{n}\sum_{j=1}^{n} f_{ij}(x)\frac{\partial^2 u}{\partial x_i \partial x_j} + \ldots \tag{8}$$

Definition: A linear PDE is any PDE of the form

$$L[u(x)] = f(x) \tag{9}$$

If $f(x) = 0$, the equation is homogeneous, else it is non-homogeneous.
Ex: $u_t = 4u_x$ - Linear, homogeneous.

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Example:

$$u_{tt} = u_{xx} + uyy \quad \text{Linear, homogeneous} \tag{10}$$
$$\cos(xt) = u + u_t + u_{xyz} \quad \text{Linear, non-homogeneous} \tag{11}$$
$$u_t u_{xt} = 0 \quad \text{non-linear} \tag{12}$$
$$u_{xt} + e^x \cos t \; u_t = 0 \quad \text{linear, homogeneous} \tag{13}$$
$$u_t + u_{xx} + ue^u = 0 \quad \text{non-linear} \tag{14}$$

Note: You can add linear combinations of solutions to linear homogeneous equations and still get a solution. Example: $u_x = u_t$. Some solutions to this are:

1. $u_1(x,t) = 3$
2. $u_2(x,t) = x + t$
3. $u_3(x,t) = e^{x+t}\cos(x+t)$
4. $\vdots$
   $Au_1 + Bu_2 + Cu_3$ is also a solution.

How do we solve an ODE?

1. Use some technique to find an explicit solution.
2. Use power series and determine the coefficients

$$y(x) = \sum_{n=0}^{\infty} a_n x^n \tag{15}$$

3. Laplace Transforms

How do we solve PDEs?

1. Try to locate an explicit solution
2. We don't use power series, instead, we use a trigonometric series $\Rightarrow$ Fourier Series.

$$y(x) = \sum_{n=0}^{\infty} a_n \sin(nx) + b_n \cos(nx) \tag{16}$$

3. Laplace Transforms are good if the domain is $[0, \infty)$.
   Fourier Transforms are good if the domain is $(-\infty, \infty)$.
4. Reduce the PDE to a system of ODEs.

Initial Condiction

1. For ODEs, to solve a 1^{st} order equation, you need $y(0)$.
   2^{nd} order $\to y(0), y'(0)$
   3^{rd} order $\to y(0), y'(0), y''(0)$
   $\vdots$
   n^{th} order $\to y(0), y'(0), y''(0), \ldots, y^{(n-1)}(0)$
2. For PDEs, it's more complicated $\Rightarrow$ it depends on the PDE.
   Example: $u(x,t), x \in [a,b], t \in [0,\infty)$
   If $u_t = u_{xx}$
3. Boundary conditions:

$$u(a,t) = g_1(t) \tag{17}$$
$$u(b,t) = g_2(t) \tag{18}$$

   If $u_{tt} = u_{xx}$, we must specify:

   (a) Initial Conditions

$$u(x,0) = f_1(x) \tag{19}$$
$$u_t(x,0) = f_2(x) \tag{20}$$

   (b) Boundary Conditions

$$u(a,t) = g_1(t) \tag{21}$$
$$u(b,t) = g_2(t) \tag{22}$$

1-D Heat Equation
Assume cross sections are uniform Imagine a cross section:

O o==========o L

Assume cross sections are uniform and the lateral sides are well insulated $\Rightarrow$ heat only flows in the x-direction.
We need the following:

- $u(x,t)$: Temperature of rod at position x and time t.
- $u(x,0)$: Initial temperature

- $u(0,t)$ and $u(L,t)$: Boundary Conditions

<u>Definition:</u>

- $g(x,t)$: heat flux (energy / area time)
- $Q(x,t)$: heat energy density (energy / volume)
- A : Cross sectional area
- C_P : Heat capacity or specific heat
- ρ : Density
- K : Thermal conductivity

We want to find an equation for the temperature evolution. We will use conservation of energy : Look at a little Δx section of the rod starting at x_0.

Δx
o=====|o|=====o
x_0 $x_0\Delta x$

Conservation of energy : heat in - heat out = heat accumulated
Heat in $=' qA\Delta t' = A\int_{t_0}^{t_0+\Delta t} q(x_0,t)$ dt
Heat out $= A\int_{t_0}^{t_0+\Delta t} q(x_0+\Delta x,t)$ dt
Heat Accumulated $= QA\Delta x|_{t_0+\Delta t} - QA\Delta x|_{t_0}$

$$= A\int_{x_0}^{x_0+\Delta x} Q(x,t_0+\Delta t)\ \mathrm{dx} - A\int_{x_0}^{x_0+\Delta x} Q(x,t_0)\ \mathrm{dx} \tag{23}$$

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Heat Equation
Conservation of energy
Heat in - heat out = heat accumulated

$$A \int_{t_0}^{t_0 \to \Delta t} g(x_0, t)\ \mathrm{dt}\ - A \int_{t_0}^{t_0 \to \Delta t} q(x_0 + \Delta x, t)\ \mathrm{dt}\ = A \int_{t_0}^{t_0 \to \Delta t} Q(x, t_0 + \Delta t)\ \mathrm{dx}\ - A \int_{t_0}^{t_0 \to \Delta t} Q(x, t_0)\ \mathrm{dx} \tag{24}$$

Let us simplify and divide by A. Then, let us combine the integrals:

$$\int_{t_0}^{t_0 \to \Delta t} [q(x_0, t) - q(x_0 + \Delta x, t)]\ \mathrm{dt}\ = \int_{t_0}^{t_0 \to \Delta t} [Q(x, t_0 + \Delta t) - Q(x, t_0)]\ \mathrm{dx} \tag{25}$$

Divide by $\Delta x \Delta t$ and take limit as $\Delta x, \Delta t \to 0$

$$\lim_{\Delta t, \Delta x \to 0} \frac{1}{\Delta x \Delta t} \int_{t_0}^{t_0 \to \Delta t} [q(x_0, t) - q(x_0 + \Delta x, t)]\ \mathrm{dt}\ = \lim_{\Delta t, \Delta x \to 0} \frac{1}{\Delta x \Delta t} \int_{t_0}^{t_0 \to \Delta t} [Q(x, t_0 + \Delta t) - Q(x, t_0)]\ \mathrm{dx} \tag{26}$$

$$\lim_{\Delta t} \frac{1}{\Delta t} \int_{t_0}^{t_0 \to \Delta t} [\lim_{\Delta x \to 0} \frac{q(x_0, t) - q(x_0 + \Delta x, t)}{\Delta x}]\ \mathrm{dt}\ = \lim_{\Delta x \to 0} \frac{1}{\Delta x} \int_{t_0}^{t_0 \to \Delta t} \lim_{\Delta t \to 0} \frac{Q(x, t_0 + \Delta t) - Q(x, t_0)}{\Delta t}\ \mathrm{dx} \tag{27}$$

On the left side, we see the order is a bit difference. We want the delta to come first, such as in the difference quotient. The eft is now $-q_x(x_0, t)$ and the right is $Q_t(x, t_0)$.

$$\lim_{\Delta t \to} \frac{1}{\Delta t} \int_{t_0}^{t_0 + \Delta t} -q_x(x_0 t)\ \mathrm{dt}\ = \lim_{\Delta x \to 0} \frac{1}{\Delta x} \int_{x_0}^{x_0 + \Delta x} Q_t(x, t_0)\ \mathrm{dx} \tag{28}$$

$$\lim_{\Delta t \to 0} -q_x(x_0, t_0 + \Delta t) = \lim_{\Delta x \to 0} Q_t(x_0 + \Delta x, t_0) \tag{29}$$

At step 28, we used the fundamental theorem of calculus and derived both sides.

$$-q_x(x_0, t_0) = Q_t(x_0, t_0) \tag{30}$$

Since x_0 and t_0 are arbitrary, $-q_x(x, t) = Q_t(x, t)$
q and Q are related to u:

$$Q = \rho c_p u \qquad q = -K u_x \tag{31}$$

$$-q_x = Q_t \Rightarrow K u_{xx} = \rho c_p u_t \tag{32}$$

$$\Rightarrow u_t = \frac{k}{\rho c_p} u_{xx} \tag{33}$$

$$\Rightarrow u_t = \alpha^2 u_{xx} \tag{34}$$

$$\alpha = \sqrt{\frac{K}{\rho c_p}} \tag{35}$$

α is thermal diffusivity
$u_t = \alpha^2 u_{xx} \leftarrow$ 1-D heat equation (diffusivity equation)
We have a steady-state: $(t \to \infty)$, $u_t = 0 \Rightarrow u_{xx} = 0 \Rightarrow$ straight line
1-D: $-q_x = Q_t \Rightarrow -\nabla \cdot \vec{q} = Q_t$, $\vec{q}$ is a vector.

$$q = -K \nabla u \Rightarrow -\nabla \cdot (-K \nabla u) = \rho c_p u_t \tag{36}$$

$$\Rightarrow K \Delta u = \rho c_p u_t \tag{37}$$

$$\Rightarrow u_t = \alpha^2 \Delta u \tag{38}$$

What about a steady-state? $u_t = 0$

$$\Delta u = 0 \tag{39}$$

Here, we have Laplace's equation.
Note: It is not dependent on time.
The Wave Equation $u(x, t)$ is the height of the rope. We use Newton's 2^{nd} law on small segments of rope.

- ρ = density of rope.
- $\mathrm{dm} = \rho\ \mathrm{dx}$

$$F = ma \tag{40}$$

$$T \sin(\theta(x + \Delta x)) - T \sin(\theta(x)) = \int_x^{x + \Delta x} u_{tt}\ \mathrm{dm} \tag{41}$$

$$T[\sin(\theta(x + \Delta x)) - \sin(\theta(x))] = \rho \int_x^{x + \Delta x} u_{tt}\ \mathrm{dx} \tag{42}$$

Let us assume θ is small, $\sin \theta \approx \tan \theta$

$$T[\tan(\theta(x + \Delta x)) - \tan(\theta(x))] = \rho \int_x^{x + \Delta x} u_{tt}\ \mathrm{dx} \tag{43}$$

Also, $\tan(\theta(x)) = u_x(x, t)$.

$$T[u_x(x + \Delta x, t) - u_x(x, t)] = \rho \int_x^{x + \Delta x} u_{tt}\ \mathrm{dx} \tag{44}$$

Now, let us divide both sides by Δx and take the limit as $\Delta x \to 0$

$$\lim_{\Delta x \to 0} T \left[\frac{u_x(x + \Delta x, t) - u_x(x, t)}{\Delta x} \right] = \rho \lim_{\Delta x \to 0} \frac{\int_x^{x+\Delta x} u_{tt} \text{ dx}}{\Delta x} \tag{45}$$

On the left side, we have $u_x x$ and the right side we have $u_{tt}(x + \Delta x, t)$.

$$Tu_{xx}(x, t) = \rho u_{tt}(x, t) \tag{46}$$

$$u_{tt} = \frac{T}{\rho} u_{xx} = c^2 u_{xx} \tag{47}$$

$$c = \sqrt{\frac{T}{\rho}} = \text{ wave speed} \tag{48}$$

On the left, we have the $1 - D$ wave equation which is used for light, sound, rope, etc.
In 2-D, it corresponds to a vibrating membrane (drum)

$$u_{tt} = c^2 \Delta u \tag{49}$$

<u>Remark</u>:

$$u_t = u_{xx} \quad \text{Heat Equation} \tag{50}$$

$$u_{xx} + u_{yy} = 0 \quad \text{Laplace Equation} \tag{51}$$

$$u_{tt} = u_{xx} \quad \text{wave} \tag{52}$$

Here, we can replace:
u_t with t
u_x with x
u_{xx} with x^2

1. $t = x^2$ parabola
2. $x^2 + y^2 = 0$ ellipse
3. $t^2 = x^2$ hyperbolas

So, the equations behave like the following:

1. The Heat Equation is called parabolic
2. The Laplace Equation is called elliptic
3. The Wave Equation is called hyperbolic

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Approximating Functions with other Functions

1. Prove Series

$$f(x) = \sum_{n=0}^{M} a_n x^n \quad \text{Finite Power Series} \tag{53}$$

This is not the best way to approximate a function.

We choose the a_n's so that the power series is "close" to $f(x)$ which means we want to minimize the error.

We increase M to get a better approximation.

The problem begins when you change M, the values of a_n's change as well. Therefore, recalculating is a lot of work.

If we let $M \to \infty$ and if $f \in C^\infty$, so then $a_n = \frac{f^{(n)}(0)}{n!}$ and we get the Taylor series.

Note: C^∞: C means Continuous and the ∞ indicates the number of derivatives that are continuous.

Problem: This is only good inside the radius of convergence.

A Fourier Series is a trigonometric polynomial

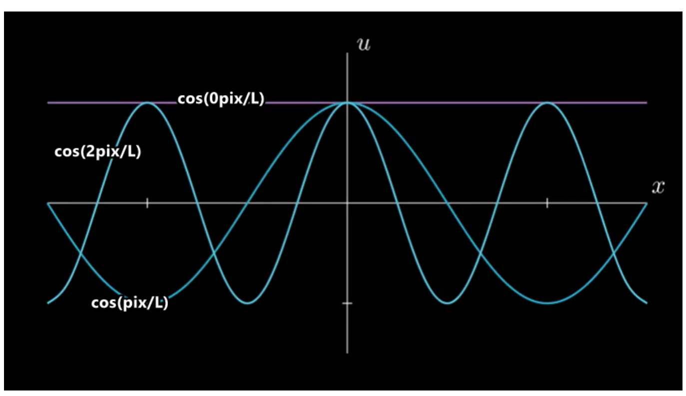
$$\sum_{n=0}^{M} a_n \sin\left(\frac{n\pi x}{L}\right) + b_n \cos\left(\frac{n\pi x}{L}\right) \longleftarrow \text{period} = 2L \tag{54}$$

We use Fourier Series for a function on a bounded interval and we will use $x \in [-L, L]$

Advantages of Fourier Series

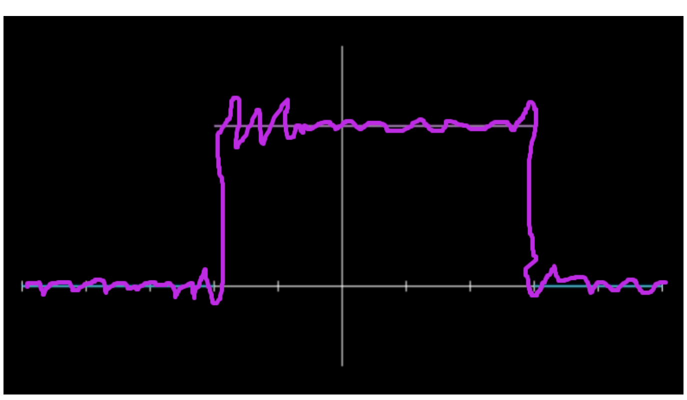
1. If M increases, we only need to calculate the new a_n's and b_n's. This property is due to the fact that the basis functions are orthogonal.
2. If $M = \infty$ and f is continuous, then the Fourier Series $= f(x) \forall x \in (-L, L)$. Our interval must be open for the case that $f(-L) \neq f(L)$.

Basis Functions : $\sin\left(\frac{n\pi x}{L}\right), \cos\left(\frac{n\pi x}{L}\right)$

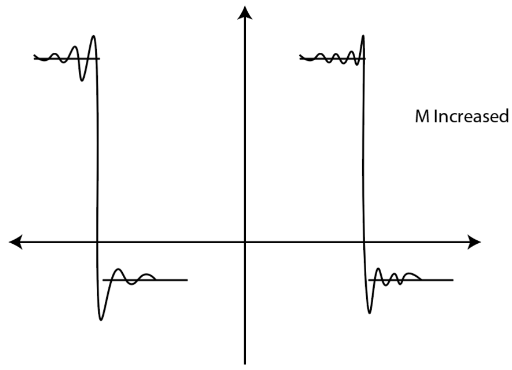


What happens if you use a Fourier Series on a discontinuous function?

$$f(x) = \begin{cases} 1 & x \in (-4, 6) \\ 0 & x \in [-10, -4] \bigcup [6, 10] \end{cases} \tag{55}$$

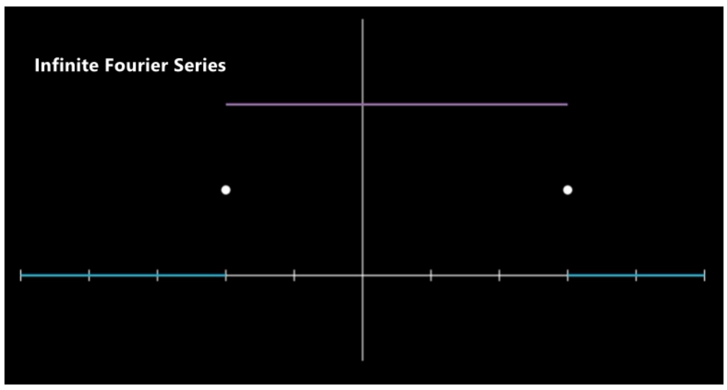


The Oscillations around the discontinuities are called Gibbs phenomenon. As M increases, the oscillation's amplitude does not change. However, the oscillations do get progressively closer to the discontinuities.



If $M = \infty$, then we have:

$$\text{Fourier Series } = \begin{cases} f(x) & = \text{ if } x \text{ is a point of continuity} \\ \lim_{c \to 0^+} \frac{f(x+c)+f(x-c)}{2} & \text{ if x is a point of discontinuity} \end{cases} \tag{56}$$



Orthogonality
Recall: The vectors

$$u = \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_n \end{bmatrix} \quad \text{and} \quad v = \begin{bmatrix} v_1 \\ v_2 \\ \vdots \\ v_n \end{bmatrix} \tag{57}$$

are orthogonal if the dot product is zero.

$$u \circ v = \sum_{i=1}^{n} u_i v_i = 0 \tag{58}$$

We want to generalize this to function $x \in [-L, L]$.

Definition: Two functions $f(x)$ and $g(x)$ are orthogonal on $[a, b]$ if

$$\int_a^b f(x)g(x) \text{ dx } = 0 \tag{59}$$

Theorem: All basis functions in the Fourier Series are mutually orthogonal

$$\int_{-L}^{L} \sin\left(\frac{m\pi x}{L}\right) \sin\left(\frac{n\pi x}{L}\right) \text{ dx } = 0 \quad n \neq m \tag{60}$$

$$\int_{-L}^{L} \cos\left(\frac{m\pi x}{L}\right) \cos\left(\frac{n\pi x}{L}\right) \text{ dx } = 0 \quad n \neq m \tag{61}$$

What happens if $m = n$?

$$\int_{-L}^{L} \sin^2\left(\frac{m\pi x}{L}\right) \text{ dx} \tag{62}$$

Here, we want to use the double angle formula: $\cos(2\theta) = 1 - 2\sin^2\theta$.

$$\begin{aligned} \int_{-L}^{L} \sin^2\left(\frac{m\pi x}{L}\right) \text{ dx } &= \frac{1}{2}\int_{-L}^{L} 1 - \cos\left(\frac{2m\pi x}{L}\right) \text{ dx} && (63) \\ &= \frac{1}{2}\Big[x - \frac{L}{2m\pi}\sin\left(\frac{2m\pi x}{L}\right)\Big]_{-L}^{L} && (64) \\ &= \frac{1}{2}\Big[L - \frac{L}{2m\pi}\sin(2m\pi) - \Big(-L - \frac{2}{2m\pi}\sin(-2m\pi)\Big)\Big] && (65) \\ &= L && (66) \end{aligned}$$

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Similarly,

$$\int_{-L}^{L} \cos^2(\frac{n\pi x}{L}) \text{ dx } = L \tag{67}$$

If $n = 0$,

$$\int_{-L}^{L} 1 \text{ dx } = 2L \tag{68}$$

<u>Note:</u> You cannot differentiate the Fourier Series term-by-term $f'(x)$ like you can with Taylor series.

Let's show $\cos(\frac{n\pi x}{L})$ and $\sin(\frac{m\pi x}{L})$ are orthogonal on $[-L, L]$.

$$\int_{-L}^{L} \sin\left(\frac{m\pi x}{L}\right)\cos\left(\frac{n\pi x}{L}\right) \text{ dx } = \frac{1}{2}\int_{-L}^{L} \sin\left(\frac{(m+n)\pi x}{L}\right) + \sin\left(\frac{(m-n)\pi x}{L}\right) \text{ dx} \tag{69}$$

$$= -\frac{1}{2}\Big[\frac{L}{(m+n)\pi}\cos\left(\frac{(m+n)\pi x}{L}\right) + \frac{L}{(m-n)\pi}\cos\left(\frac{(m-n)\pi x}{L}\right)\Big]_{-L}^{L} \tag{70}$$

Here, we expand our difference and notice we have even and odd functions.

In general, the coefficients are:

$$a_n = \frac{1}{L}\int_{-L}^{L} f(x)\sin\left(\frac{n\pi x}{L}\right) \text{ dx} \tag{71}$$

$$b_n = \frac{1}{L}\int_{-L}^{L} f(x)\cos\left(\frac{n\pi x}{L}\right) \text{ dx} \tag{72}$$

$$b_0 = \frac{1}{2L}\int_{-L}^{L} f(x) \text{ dx} \tag{73}$$

<u>Example:</u> $f(x) = x, \quad x \in [-3, 3]$.

Find the Fourier Series for f.

$$a_n = \frac{1}{3}\int_{-3}^{3} x\sin\left(\frac{n\pi x}{L}\right) \text{ dx} \tag{74}$$

Here, we want to integrate by parts:

x	$\sin\left(\frac{n\pi x}{L}\right)$
1	$-\frac{3}{n\pi}\cos\left(\frac{n\pi x}{L}\right)$
0	$-\frac{9}{n^2\pi^2}\sin\left(\frac{n\pi x}{L}\right)$

<u>Note:</u> $L = 3$.

$$= \frac{1}{3}\Big[-\frac{3x}{n\pi}\cos\left(\frac{n\pi x}{3}\right)\Big]_{-3}^{3} + \Big[\left(\frac{3}{n\pi}\right)^2\sin\left(\frac{n\pi x}{3}\right)\Big]_{-3}^{3} \tag{75}$$

$$= \frac{1}{3}\Bigg[-\frac{9}{n\pi}\cos(n\pi) + \frac{9}{n^2\pi^2}\sin(n\pi) - \Big(+\frac{9}{n\pi}\cos(-n\pi) + \frac{9}{n^2\pi^2}\sin(-n\pi)\Big)\Bigg] \tag{76}$$

$$= -\frac{6}{n\pi}\cos(n\pi) \tag{77}$$

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The Fourier Series is not valid at $x = \pm 3$ since it is not continuous at ± 3.

Let's say $f(x)$ is odd, then $f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)$

Let's say $f(x)$ is even, then $f(x) = \sum_{n=1}^{\infty} b_n \cos\left(\frac{n\pi x}{L}\right)$

If we are only interested in the behavior of $f(x)$ on $[0, L$, then we can either use a Fourier Sine Series $f(x) = \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi x}{L}\right)$ or a Fourier series $f(x) = \sum_{n=0}^{\infty} b_n \cos\left(\frac{n\pi x}{L}\right)$.

<u>Solving the Heat Equation</u>

$$u_t = \alpha^2 u_{xx} \tag{78}$$

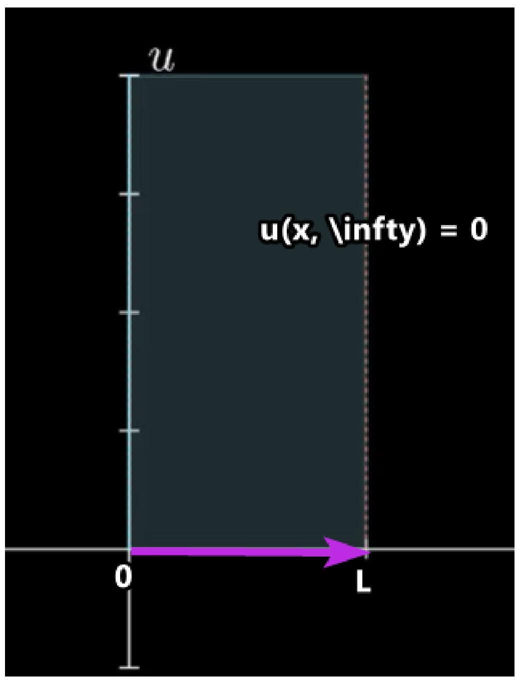
0 o============o L $u(x,t)$: temp

Initial condition:

$$u(x,0) = f(x) \tag{79}$$

Boundary conditions

$$u(0,t) = u(L,t) = 0 \tag{80}$$



So, whatever we get, we better have $\lim_{t\to\infty} u(x,t) = 0$. The method of reflection? relies on two things:

1. Fourier Series
2. Linearity

<u>Method</u>

1. Try a solution of the form

$$u(x,t) = X(x)T(t) \leftarrow \text{ Assume the solution is separable} \tag{81}$$

   Boundary Conditions: Here, we conclude $X(0)$ is 0 because we want $T(t)$ to change as t changes.

$$u(0,t) = 0 \Rightarrow X(0)T(t) = 0 \Rightarrow X(0) = 0 \tag{82}$$
$$u(L,t) = 0 \Rightarrow X(L)T(t) = 0 \Rightarrow X(L) = 0 \tag{83}$$
$$U_t = \alpha^2 u_{xx} \Rightarrow X(x)T'(t) = \alpha^2 X''(x)T(t) \tag{84}$$

   Here, we divide by X, T, α^2.

$$\Rightarrow \frac{T'(t)}{\alpha^2 T(t)} = \frac{X''}{X(x)} = -\lambda \tag{85}$$

   Here, λ is a constant.

2.

$$\frac{X''}{X(x)} = -\lambda \tag{86}$$
$$X''(x) = -\lambda X(x) \tag{87}$$

   Here, we know $(x) = X(L) = 0$. We call every $(\lambda, X(x))$ pair that satisfies this equation an eigenvalue/eigenfunction pair for the differential equation.

Assume $\lambda > 0$

$$X'' = -\lambda x \tag{88}$$
$$x(0) = x(L) = 0 \tag{89}$$
$$\Rightarrow X(x) = A\cos(\sqrt{\lambda}x) + B\sin(\sqrt{\lambda}x) \tag{90}$$
$$\Rightarrow A\cos 0 + B\sin 0 = 0 \tag{91}$$
$$\Rightarrow A = 0 \tag{92}$$
$$\Rightarrow X(x) = B\sin(\sqrt{\lambda}x) \tag{93}$$
$$X(L) = 0 \Rightarrow B\sin(\sqrt{\lambda}L) = 0 \tag{94}$$
$$\Rightarrow \sin(\sqrt{\lambda}L) = 0 \tag{95}$$
$$\Rightarrow \sqrt{\lambda}L = n\pi, n \in \mathbb{Z}^+ \tag{96}$$
$$\Rightarrow \lambda_n = \left(\frac{n\pi}{L}\right)^2 \tag{97}$$
$$\Rightarrow_n (x) = \sin\left(\frac{n\pi x}{L}\right) \tag{98}$$

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$$u(x,t) = X(x)T(t) \tag{99}$$
$$\frac{X''}{x} = \frac{T'}{\alpha^2 T} = -\lambda \tag{100}$$
$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \tag{101}$$
$$X_n(x) = \sin\left(\frac{n\pi x}{L}\right) \tag{102}$$

Here, we have a bsis function for Fourier Sine Series.

3. Solve for T

$$\frac{T'}{\alpha^2 T} = -\lambda \tag{103}$$
$$T' = -\alpha^2 \lambda T \tag{104}$$
$$T_n' = -\alpha^2 \lambda_n T_n \tag{105}$$
$$= -\alpha^2 \left(\frac{n\pi}{L}\right)^2 T \tag{106}$$

If we have something like $y' = ky$, we know that this derives from $y = e^{kx}$.

$$T_n(t) = e^{-\alpha^2 \left(\frac{n\pi}{L}\right)^2 T} \tag{107}$$

4. Combine for u_n

$$u_n(x,t) = X_n(x)T_n(t) \tag{108}$$
$$= \sin\left(\frac{n\pi x}{L}\right) e^{-\alpha^2 \left(\frac{n\pi}{L}\right)^2 T} \tag{109}$$

Each one of the $n's$ will yield a different u. We also know that $n \in \mathbb{N}$. We can take as many u's and add them all together. We find our u_n's and use it to find u.

By linearity,

$$u(x,t) = \sum_{n=1}^{\infty} A_n \tag{110}$$
$$= \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) e^{-\alpha^2 \left(\frac{n\pi}{L}\right)^2 T} \tag{111}$$

5. Satisfy the initial condition

$$u(x,0) = f(x) \tag{112}$$
$$u(x,0) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) = f(x) \tag{113}$$

Line 113) is considered the Fourier Sine Series.

The A_n's are the coefficients of the Fourier Sine Series of $f(x)$.

$$A_n = \frac{1}{L}\int_{-L}^{L} f(x)\sin\left(\frac{n\pi x}{L}\right)dx \tag{114}$$
$$= \frac{2}{L}\int_{0}^{L} f(x)\sin\left(\frac{n\pi x}{L}\right)dx \tag{115}$$

<u>Ex:</u> Solve with the following conditions:

1. $u_t = 4u_{xx}$
2. $u(0,t) = u(3,t) = 0$

3. $u(x,0) = 5\sin\left(\frac{2\pi x}{3}\right) - 7\sin(4\pi x)$

Now, let us perform the five steps to solve our equation:

1. Assume $u(x,t) = X(x)T(t)$
   Boundary Conditions :

   - $u(0,t) = 0$, then $X(0)T(t) = 0$. Here either $X(0)$ or $T(t)$ is 0, and we want $X(0) = 0$ here.
   - $(3,t) = 0$, then $X(3)T(t) = 0$, following the same logic, we have $X(3) = 0$.

$$u_t = 4u_{xx} \tag{116}$$
$$XT' = 4X''T \tag{117}$$
$$\frac{T'}{4T} = \frac{X''}{X} = -\lambda \tag{118}$$

2. Now, since we know more information regarding X, let us solve for X.

$$\frac{X''}{X} = -\lambda \tag{119}$$
$$X'' = -\lambda X, \quad X(0) = X(3) = 0 \tag{120}$$

   Let us assume $\lambda > 0$. Here, we want an X'' where deriving twice gives us $-X$. Assume $\lambda > 0$

$$X = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x) \tag{121}$$

   Set $X(0) = 0$

$$X = A \tag{122}$$

   Now, let us find $X(3) = 0$:

$$0 = A\sin(\sqrt{\lambda}3) \tag{123}$$
$$\sqrt{\lambda}3 = n\pi \tag{124}$$
$$\lambda_n = \left(\frac{n\pi}{3}\right)^2 \tag{125}$$
$$X_n(x) = \sin\left(\frac{n\pi x}{3}\right) \tag{126}$$

3. Now, let us find T.

$$\frac{T'}{4T} = -\lambda \tag{127}$$
$$T_n' = -4\left(\frac{n\pi}{3}\right)^2 T_n \tag{128}$$
$$T_n(t) = e^{-4\left(\frac{n^2\pi^2}{9}\right)t} \tag{129}$$

4. Combine to find u_n and u

$$u_n(x,t) = X_n(x)T_n(t) \tag{130}$$
$$= \sin\left(\frac{n\pi x}{3}\right) e^{-4\left(\frac{n^2\pi^2}{9}\right)t} \tag{131}$$

   By linearity,

$$u(x,t) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{3}\right) e^{-4\left(\frac{n^2\pi^2}{9}\right)t} \tag{132}$$

5. Use the initial conditions to find A_n's

$$u(x,0) = 5\sin(\frac{2\pi x}{3}) - \sin(4\pi x) \tag{133}$$
$$u(x,0) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{3}\right) \tag{134}$$
$$A_n = \frac{2}{3}\int_0^3 5\left[\sin\left(\frac{2\pi x}{3}\right) - 7\sin(4\pi x)\right]\sin(\frac{n\pi x}{3})\ \mathrm{dx} \tag{135}$$

   Lets look at our initial condition on line 133). The first one is $n = 2$, so $A_2 = 5$. In addition, the second term is at $A_12 = -7$. Therefore, we have $A_n = 0 \forall n$ except $n = 2, 12$.

   Now, let us look at our linearity equation.

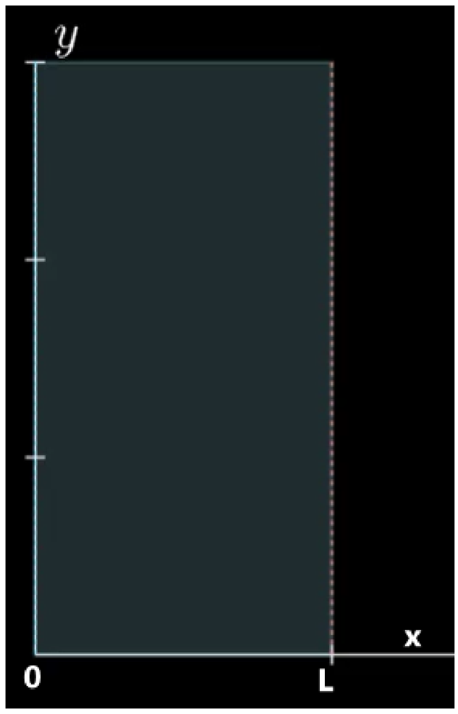
$$u(x,t) = 5\sin\left(\frac{2\pi x}{3}\right) e^{-\frac{16\pi^2}{9}t} - 7\sin(4\pi x)\, e^{-64\pi^2 t} \tag{136}$$

   Here, this is our final solution.

Laplace's Equation 1-D : $u_{xx} = 0 \Rightarrow u = ax + b$

If $u(0) = u(L) = 0$, then that would force our function to be $u = 0$. This is the steady state solution. If our function is in the form of $ax + b$, then $u = 0$ is the only solution for the function to hit 0 twice in this fashion.

2-D: $\Delta u = 0 \Rightarrow u_{xx} + u_{yy} = 0$



We have two types of boundary conditions:

(a) Specify u on the perimeter

- $u(0, y) = 0$
- $u(L, y) = 0$
- $u(x, 0) = 0$
- $u(x, M) = f(x)$

(b) Nuemann conditions : Specify the direction derivative in the normal direction on the boundary.

- $u_x(0, y) = 0$
- $u_x(L, y) = 0$
- $u_y(x, 0) = 0$
- $u_y(x, M) = \widetilde{f}(x)$
- $u(0, 0) \ = T$

This is if we know the heat flux $\vec{q} \cdot \vec{n}$ on the boundary.

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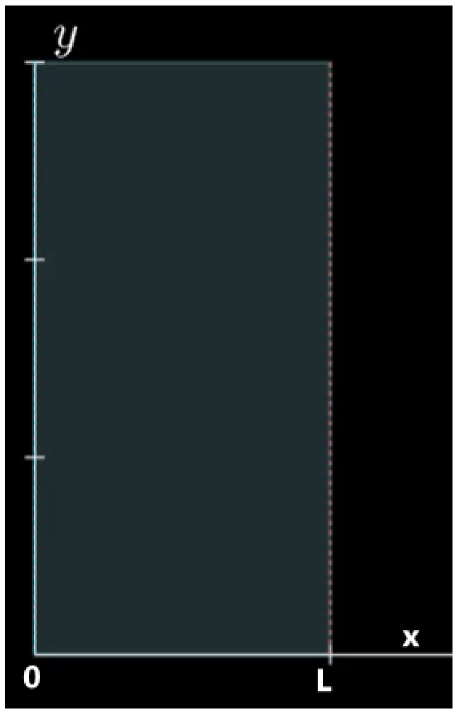
Solving Laplace's Equation



$$u_{xx} + u_{yy} = 0$$

- $u_x(0, y) = 0$
- $u_x(L, y) = 0$
- $u(x, 0) = 0$
- $u(x, M) = f(x)$

1. Assume $u(x, y) = X(x)Y(y)$

   Boundary Conditions

$$u(x,y) = X(x)Y(y) \tag{137}$$
$$\Rightarrow X'(x)Y(y) \tag{138}$$

Now, let us write our boundary condition:

$$U_x(0,y) = 0 \tag{139}$$
$$\Rightarrow X'(0)Y(y) = 0 \tag{140}$$
$$\Rightarrow X'(0) = 0 \tag{141}$$

Now, let us find the next item,

$$u_x(L,y) = 0 \tag{142}$$
$$\Rightarrow X'(L)Y(y) = 0 \tag{143}$$
$$\Rightarrow X'(L) = 0 \tag{144}$$

Now, the next two items do not have a derivative:

$$u(x,0) = 0 \tag{145}$$
$$\Rightarrow X(x)Y(0) = 0 \tag{146}$$
$$\Rightarrow Y(0) = 0 \tag{147}$$

Now, let us write:

$$u_{xx} + u_{yy} = 0 \tag{148}$$
$$\Rightarrow X''Y + XY'' = 0 \tag{149}$$
$$\Rightarrow X''Y = -XY'' \tag{150}$$
$$\Rightarrow \frac{X''}{X} = -\frac{Y''}{Y} = -\lambda \tag{151}$$

2. Solve for X (Note: We solve for X first here, since we have more information about X).

$$\frac{X''}{X} = -\lambda \tag{152}$$
$$\Rightarrow X'' = -\lambda X, \quad X'(0) = X'(L) = 0 \tag{153}$$
$$\lambda > 0 \Rightarrow x(x) = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x) \tag{154}$$
$$\Rightarrow X'(x) = A\sqrt{\lambda}\cos(\sqrt{\lambda}x) - B\sqrt{\lambda}\sin(\sqrt{\lambda}x) \tag{155}$$
$$X'(0) = 0 \Rightarrow A\sqrt{\lambda} = 0 \tag{156}$$

Now, if we rewrite out equation, we have:

$$X(x) = B\cos(\sqrt{\lambda}x) \tag{157}$$

Next, we want to find $X'(L) = 0$:

$$0 = -B\sqrt{\lambda}\sin(\sqrt{\lambda}L) \tag{158}$$
$$\sqrt{\lambda}L = n\pi \tag{159}$$
$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \tag{160}$$
$$\Rightarrow X_n(x) = \cos\left(\frac{n\pi x}{L}\right) \tag{161}$$

If $\lambda = 0$

$$\frac{X_0''}{X_0} \Rightarrow X_0'' = 0 \tag{162}$$
$$\Rightarrow X_0(x) = Ax + B \tag{163}$$
$$\Rightarrow X_0'(x) = A \tag{164}$$
$$\Rightarrow X_0'(0) = 0 \tag{165}$$
$$\Rightarrow A = 0 \tag{166}$$
$$\Rightarrow X_0'(L) = 0 \tag{167}$$
$$\Rightarrow A = 0 \tag{168}$$

Neither conditions tell us more information about B,

$$\Rightarrow X_0(x) = B_0 \tag{169}$$

3. Now, we want to solve for Y: $-\frac{Y''}{Y} = -\lambda$

$$Y'' = \lambda y \tag{170}$$
$$Y'' = \left(\frac{n\pi}{L}\right)^2 Y_n, \quad Y_n(0) = 0 \tag{171}$$
$$Y_n(y) = Ce^{\frac{n\pi}{L}y} + De^{-\frac{n\pi}{L}}y \tag{172}$$
$$Y_n(0) = 0 \Rightarrow C + D = 0 \tag{173}$$

Here, we do not have an additional condition that could help use solve this equality. Let us consider the hyperbolic sin and cos:

$$\sinh(x) = \frac{e^x - e^{-x}}{2} \tag{174}$$
$$\cosh(x) = \frac{e^x + e^{-x}}{2} \tag{175}$$

Instead of writing Y in the same fashion we solved for X, we use the hyperbolic sinh and *cosh*

$$Y_n(y) = C\sinh\left(\frac{n\pi y}{L}\right) + D\cosh\left(\frac{n\pi y}{L}\right) \tag{176}$$
$$Y_n(0) = 0 \Rightarrow D = 0 \tag{177}$$
$$Y_n(y) = \sinh\left(\frac{n\pi y}{L}\right) \tag{178}$$

Now, let us write:

$$\frac{Y_0''}{Y_0} = \lambda_0 \tag{179}$$
$$\Rightarrow Y_0'' = 0 \tag{180}$$
$$\Rightarrow Y_0 = Cy + D \tag{181}$$
$$\Rightarrow Y_0(0) = 0 \tag{182}$$
$$\Rightarrow D = 0 \tag{183}$$
$$\Rightarrow Y_0(y) = C_0 y \tag{184}$$

4. Combine to find u_n and u:

$$u_n(x,y) = X_n(x)Y_n(y) = \begin{cases} \cos\left(\frac{n\pi x}{L}\right)\sinh\left(\frac{n\pi y}{L}\right) & n \geq 1 \\ B_0 C_0 y & n = 0 \end{cases} \tag{185}$$

By linearity,

$$u(x,y) = \widetilde{B}_0 y + \sum_{n=1}^{\infty} B_n \cos\left(\frac{n\pi x}{L}\right)\sinh\left(\frac{n\pi y}{L}\right) \tag{186}$$

5. Here, use the final boundary condition to find the coefficients.

$$u(x,M) = f(x) \tag{187}$$
$$u(x,M) = \widetilde{B}_0 M + \sum_{n=1}^{\infty} B_n \cos\left(\frac{n\pi x}{L}\right)\sinh\left(\frac{n\pi M}{L}\right) \tag{188}$$

This is our Fourier Cosine Series for $f(x)$. Here, we can say a few things about this equation,

- $b_0 = \widetilde{B}_0 M$
- $b_n = B_n \sinh\left(\frac{n\pi M}{L}\right)$

$$\widetilde{B}_0 M = \frac{2}{2L}\int_0^L f(x)\ \mathrm{dx} \tag{189}$$

$$\widetilde{B}_0 = \frac{1}{ML}\int_0^L f(x)\ \mathrm{dx} \tag{190}$$

Next, let us find:

$$B_n \sinh\left(\frac{n\pi M}{L}\right) = \frac{2}{L}\int_0^L f(x)\cos\left(\frac{n\pi x}{L}\right)\ \mathrm{dx} \tag{191}$$

$$= \frac{2}{L\sinh\left(\frac{n\pi M}{L}\right)}\int_0^L f(x)\cos\left(\frac{n\pi x}{L}\right)\ \mathrm{dx} \tag{192}$$

<u>Ex:</u> Solve $\Delta u = 0$

- $(x, y) = 0$
- $(2, y) = 0$
- $(x, 0) = 0$
- $(x, 3) = 4\sin(5x)$

1. Assume $u(x, y) = X(x)Y(y)$

   Here, let us look at our boundary conditions:

$$u(0, y) = 0 \tag{193}$$

$$X(0)Y(y) = 0 \tag{194}$$

$$X(0) = 0 \tag{195}$$

   Here, let us look at our next boundary conditions:

$$u(2, y) = 0 \tag{196}$$

$$X(2)Y(y) = 0 \tag{197}$$

$$X(2) = 0 \tag{198}$$

   Here, let us look at our next boundary conditions:

$$u(x, 0) = 0 \tag{199}$$

$$X(x)Y(0) = 0 \tag{200}$$

$$Y(x) = 0 \tag{201}$$

   Now, we can write:

$$u_{xx} + u_{yy} = 0 \tag{202}$$

$$X''Y + XY'' = 0 \tag{203}$$

$$\frac{X''}{X} = -\frac{Y''}{Y} = -\lambda \tag{204}$$

2. Now, let us solve for x:

$$\frac{X''}{X} = -\lambda \tag{205}$$

$$X'' = -\lambda X, \quad X(0) = X(2) = 0 \tag{206}$$

$$\lambda > 0 \Rightarrow X(x) = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x) \tag{207}$$

$$X(0) = B = 0 \tag{208}$$

$$X(2) = A\sin(\sqrt{\lambda}2) = 0 \tag{209}$$

$$= \lambda 2 = n\pi \tag{210}$$

$$= \lambda_n = \left(\frac{n\pi}{2}\right)^2 \tag{211}$$

$$= X_n(x) = \sin(\frac{n\pi x}{2}) \tag{212}$$

3. Let us solve for y:

$$\frac{Y_n''}{Y_n} = \lambda_n \tag{213}$$

$$Y_n'' = \left(\frac{n\pi}{2}\right)^2 Y_n, \quad Y_n(0) = 0 \tag{214}$$

$$Y_n(y) = C\sinh\left(\frac{n\pi y}{2}\right) + D\cosh\left(\frac{n\pi y}{2}\right) \tag{215}$$

$$= Y_n(0) = 0 \Rightarrow D = 0 \tag{216}$$

$$Y_n(y) = \sinh\left(\frac{n\pi y}{2}\right) \tag{217}$$

   We are picking a constant for this last term later, so we can drop C.

Here, let us write out our equation for the following function,

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- $\Delta u = 0$
- $u(0, y) = 0$
- $u(2, y) = 0$
- $u(x, 0) = 0$
- $u(x, 3) = 4$
- $\lambda_n = \left(\frac{n\pi}{2}\right)^2$
- $X_n(x) = \sin\left(\frac{n\pi x}{2}\right)$
- $Y_n(x) = \sinh\left(\frac{n\pi y}{2}\right)$

4. Combine to find u_n and u

$$u_n(x, y) = \sin\left(\frac{n\pi x}{2}\right)\sinh\left(\frac{n\pi y}{2}\right) \tag{218}$$

   By linearity,

$$u(x, y) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{2}\right)\sinh\left(\frac{n\pi y}{2}\right) \tag{219}$$

5. Find coefficients using last boundary conditions

$$u(x, y) = 4\sin(5\pi x) \tag{220}$$
$$u(x, y) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{2}\right)\sinh\left(\frac{n\pi 3}{2}\right) \tag{221}$$
$$= 4\sin(5\pi x) \tag{222}$$

   Recall, the coefficient of the Fourier Sine Series is anything but sin. Notice our last line with discrete number, our coeficient must equal 4 and our n is 10, therefore, we write:

$$A_{10}\sinh\left(\frac{10\pi 3}{2}\right) = 4 \tag{223}$$
$$\Rightarrow A_1 0 \frac{4}{\sinh(15\pi)} \tag{224}$$

The Wave Equation

$$u_{tt} = c^2 u_{xx} \tag{225}$$

Boundary conditions:

$$u(0, t) = u(L, t) = 0 \tag{226}$$

Initial Conditions:

$$u(x, 0) = f(x) \tag{227}$$
$$u_t(x, 0) = g(x) \tag{228}$$

Here, $g \in C^2$, $g(0) = g(L)$.

1. Assume separable:

$$u(x, t) \tag{229}$$

   Boundary conditions:

$$u(0, t) = 0 \Rightarrow X(0)T(t) = 0 \Rightarrow X(0) = 0 \tag{230}$$
$$u(L, t) = 0 \Rightarrow X(L)T(t) = 0 \Rightarrow X(L) = 0 \tag{231}$$

   Now, let us rewrite our variables:

$$u_{tt} = c^2 u_{xx} \tag{232}$$
$$XT'' = c^2 X'' T \tag{233}$$
$$\frac{T''}{c^2 T} = \frac{X''}{X} = -\lambda \tag{234}$$

2. Solve for X:

$$\frac{X''}{X} = -\lambda \tag{235}$$
$$X'' = -\lambda x \tag{236}$$
$$X(0) = X(L) = 0 \tag{237}$$

   Here, let us write our general equation:

$$X(x) = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x) \tag{238}$$
$$X(0) = 0 \Rightarrow B = 0 \tag{239}$$
$$X(L) = A\sin(\sqrt{\lambda}L) = 0 \tag{240}$$
$$= \sqrt{\lambda}L = n\pi \tag{241}$$

3. Solve for T:

$$\frac{T''}{c^2T_n} = -\lambda_n \tag{242}$$
$$T_n'' = -c^2\left(\frac{n\pi}{L}\right)^2 T_n \tag{243}$$
$$T_n(t) = C_n \sin\left(\frac{cn\pi t}{L}\right) + D_n \cos\left(\frac{cn\pi t}{L}\right) \tag{244}$$

4. Combine to find u_n and u

$$u_n(x,t) = \sin\left(\frac{n\pi x}{L}\right)\left[C_n \sin\left(\frac{cn\pi z}{L}\right) + D_n \cos\left(\frac{cn\pi z}{L}\right)\right] \tag{245}$$
$$u(x,t) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi x}{L}\right)\sin\left(\frac{cn\pi z}{L}\right) + D_n \sin\left(\frac{n\pi x}{L}\right)\cos\left(\frac{cn\pi z}{L}\right) \tag{246}$$

5. Find coefficients using Inidial Conditions

$$u(x,0) = f(x) \tag{247}$$
$$= \sum_{n=1}^{\infty} D_n \sin\left(\frac{n\pi x}{L}\right) = f(x) \tag{248}$$
$$D_n = \frac{2}{L}\int_0^L f(x)\sin\left(\frac{n\pi x}{L}\right)\,\mathrm{dx} \tag{249}$$
$$u_t(x,0) = g(x) \tag{250}$$
$$u_t(x,t) \tag{251}$$

Here, we took the t partial from line 246.

$$u_t(x,t) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{L}\right)\frac{cn\pi}{L}\cos\left(\frac{cn\pi t}{L}\right) - D_n \sin\left(\frac{n\pi x}{L}\right)\frac{cn\pi}{L}\sin\left(\frac{cn\pi t}{L}\right) \tag{252}$$
$$u_t(x,0) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{L}\right)\frac{cn\pi}{L} = g(x) \tag{253}$$

Here, the non-sin terms are the coefficients of the Fourier Sine Series.

$$C_n\frac{cn\pi}{L} = \frac{2}{L}\int_0^L g(x)\sin\left(\frac{n\pi x}{L}\right)\,\mathrm{dx} \tag{254}$$
$$C_n = \frac{2}{cn\pi}\int_0^L g(x)\sin\left(\frac{n\pi x}{L}\right)\,\mathrm{dx} \tag{255}$$

<u>Theorem:</u> You can differentiate a Fourier Series if it represents a C^2 function that is the same as the endpoints.

<u>Ex:</u> Solve $u_{tt} = u_{xx}$

- $u(0,t) = 0$
- $u(4,t) = 0$
- $u(x,0) = 2\sin(3\pi x) - \frac{1}{5}\sin\left(\frac{7\pi x}{2}\right)$
- $u_t(x,0) = 0$

1. Assume separable

$$u(x,t) = X(x)T(t) \tag{256}$$

Here, we have our boundary conditions,

$$u(0,t) = 0 \Rightarrow X(0)T(t) = 0 \Rightarrow X(0) = 0 \tag{257}$$
$$u(4,t) = 0 \Rightarrow X(4)T(t) = 0 \Rightarrow X(4) = 0 \tag{258}$$

Now, let us separate:

$$u_{tt} = u_{xx} \tag{259}$$
$$XT'' = X''T \tag{260}$$
$$\frac{T''}{T} = \frac{X''}{X} = -\lambda \tag{261}$$

2. Solve for x:

$$\frac{X''}{X} = -\lambda \tag{262}$$
$$X'' = -\lambda x \tag{263}$$
$$X(0) = X(4) = 0 \tag{264}$$
$$X(x) = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x) \tag{265}$$
$$X(0) = B = 0 \tag{266}$$
$$X(4) = A\sin(\sqrt{\lambda}4) = 0 \tag{267}$$
$$= \sqrt{\lambda}4 = n\pi \tag{268}$$
$$= \lambda_n = \left(\frac{n\pi}{4}\right)^2 \tag{269}$$
$$= X_n(x) = \sin\left(\frac{n\pi x}{4}\right) \tag{270}$$

3. Solve for T:

$$\frac{T_n''}{T_n} = -\lambda_n \tag{271}$$

$$T_n'' = -\left(\frac{n\pi}{4}\right)^2 T_n \tag{272}$$

Here, we have the negative sign, therefore we use sine and cosine:

$$T_n(t) = C_n \sin\left(\frac{n\pi t}{4}\right) \tag{273}$$

4. Combine to get u_n and u

$$u_n(x,t) = \sin\left(\frac{n\pi x}{4}\right) + D_n \cos\left(\frac{n\pi t}{4}\right) \tag{274}$$

By linearity,

$$u(x,t) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{4}\right) + D_n \sin\left(\frac{n\pi t}{4}\right) \cos\left(\frac{n\pi t}{4}\right) \tag{275}$$

5. Use the initial conditions to find the coefficients

$$u(x,0) = 2\sin(3\pi x) - \frac{1}{5}\sin\left(\frac{7\pi x}{2}\right) \tag{276}$$

$$u(x,0) = \sum_{n=1}^{\infty} D_n \sin\left(\frac{n\pi x}{4}\right) \tag{277}$$

$$D_{12} = 2, D_{14} = -\frac{1}{5}, D_n = 0\ \forall n, n \neq 12, 14 \tag{278}$$

$$u_t(x,t) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{4}\right)\frac{n\pi}{4}\cos\left(\frac{n\pi t}{4}\right) - D_n \sin\left(\frac{n\pi x}{4}\right)\frac{n\pi}{4}\sin\left(\frac{n\pi t}{4}\right) \tag{279}$$

$$u_t(x,0) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{4}\right)\frac{n\pi}{4} = g(x) = 0 \tag{280}$$

$$C_n = 0\ \forall n \tag{281}$$

$$u(x,t) = 2\sin(3\pi x)\cos(3\pi t) - \frac{1}{5}\sin\left(\frac{7\pi x}{2}\right)\cos\left(\frac{7\pi t}{2}\right) \tag{282}$$

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More general Boundary Conditions

Before, we used to deal with boundary conditions where u starts and end at 0. Now, let us consider boundary conditions where u can start at any number.

The steady state solution is the following:

$$\frac{T_2 - T_1}{L}x + T_1 \tag{283}$$

Ideas, we want to try the following: $u(x,t) = w(x,t) + u(x,\infty)$. Note that ∞ is the steady state. Here, as time goes to infinity, $w(x,t)$ cancels out. We want to solve for w:

We must specify w.

$$u_t = \alpha^2 u_{xx} \tag{284}$$
$$w_t = \alpha^2 w_{xx} \tag{285}$$

We want to find out more about the boundary conditions. We also need the initial conditions to solve this.

Boundary Conditions

- $u(0,t) = T_1$
- $u(L,t) = T_2$

Let us consider the first boundary condition.

$$w(0,t) + u(0,\infty) = T_1 \tag{286}$$

Here, we know that for the steady state, x is T_1 at $x = 0$. Therefore,

$$w(0,t) = 0 \tag{287}$$

We repeat with our second bullet.

$$u(L,t) = T_2 \Rightarrow \tag{288}$$
$$w(L,t) + u(L,\infty) = T_2 \tag{289}$$
$$w(L,t) = 0 \tag{290}$$

Initial Conditions

$$u(x,0) = f(x) \Rightarrow \tag{291}$$
$$w(x,0) + u(x,\infty) = f(x) \tag{292}$$
$$w(x,0) = f(x) - u(x,\infty) \tag{293}$$
$$w(x,0) = f(x) - \left[\frac{T_2 - T_1}{L}x + T_1\right] \tag{294}$$

Ex:

Solve $u_t = u_{xx}$, $u(0,t) = 2$, $u(4,t) = 3$, $u(x,0) = -6\sin(\pi x) + \frac{x}{4} + 2$.

First, find the steady-state solution:

$$u(x,\infty) = \frac{3-2}{4}x + 2 \tag{295}$$
$$= \frac{x}{4} + 2 \tag{296}$$

Now, we assume $u(x,t) = w(x,t) + u(x,\infty)$. We can make the following assumption:

$$u_t = u_{xx} \Rightarrow w_t = w_{xx} \tag{297}$$

Boundary Conditions

$$u(0,t) = 2 \Rightarrow w(0,t) = u(0,t) - u(0,\infty) = 2 - 2 = 0 \tag{298}$$
$$u(4,t) = 3 \Rightarrow w(4,t) = u(4,t) - u(4,\infty) = 3 - 3 = 0 \tag{299}$$

Here, we plug in our x into our steady-state solution and get 2, 3.

Initial Conditions

$$w(x,0) = u(x,0) - u(x,\infty) \tag{300}$$
$$= -\sin(\pi x) + \frac{x}{4} + 2 - \left(\frac{x}{4} + 2\right) \tag{301}$$
$$= -\sin(\pi x) \tag{302}$$

Now, solve for w:

1. Assume $w(x,t) = X(x)T(t)$

   Boundary Conditions

$$w(0,t) = 0 \Rightarrow X(0)T(t) = 0 \Rightarrow X(0) = 0 \tag{303}$$
$$w(4,t) = 0 \Rightarrow X(4)T(t) = 0 \Rightarrow X(4) = 0 \tag{304}$$
$$w_t = w_{xx} \Rightarrow XT' = X''T \Rightarrow \frac{T'}{T} = \frac{X''}{X} = -\lambda \tag{305}$$

2. Solve for X :

$$\frac{X''}{X} = -\lambda \Rightarrow \tag{306}$$
$$X'' = -\lambda X \tag{307}$$
$$X(0) = X(4) = 0 \tag{308}$$

   Here, let us write our general equation:

$$X(x) = A\sin(\sqrt{\lambda}x) + B\cos(\sqrt{\lambda}x) \tag{309}$$
$$X(0) = 0 \Rightarrow B = 0 \tag{310}$$
$$X(4) = 0 \Rightarrow A\sin(\sqrt{\lambda}4) = 0 \tag{311}$$
$$\Rightarrow \sqrt{\lambda}4 = n\pi \tag{312}$$
$$\Rightarrow \lambda_n = \left(\frac{n\pi}{4}\right)^2 \tag{313}$$
$$\Rightarrow X_n(x) = \sin\left(\frac{n\pi x}{4}\right) \tag{314}$$

3. Solve for T:

$$\frac{T_n'}{T_n} = -\lambda_n \tag{315}$$
$$T_n' = -\left(\frac{n\pi}{4}\right)^2 T_n \tag{316}$$
$$T_n(t) = e^{-\left(\frac{n\pi}{4}\right)^2 t} \tag{317}$$

4. Combine to find w_n and w:

$$w_n(x,t) = \sin\left(\frac{n\pi x}{4}\right) e^{-\left(\frac{n\pi}{4}\right)^2 t} \tag{318}$$

   By linearity,

$$w(x,t) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{4}\right) e^{-\left(\frac{n\pi}{4}\right)^2 t} \tag{319}$$

5. Find coefficients using Initial Condition

$$w(x,0) = -\sin(\pi x) \tag{320}$$
$$w(x,0) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{4}\right) \tag{321}$$
$$= -6\sin(\pi x) \tag{322}$$
$$w(x,t) = -6\sin(\pi x)e^{-\pi^2 t} \tag{323}$$

   Here, $a_4 = -6$.

$$u(x,t) = -6\sin(\pi x)e^{-\pi^2 t} + \frac{x}{4} + 2 \tag{324}$$

Laplace's Equation General Dirichlet Boundary Conditions

- $u(x,0) = f_1(x)$
- $u(x,M) = f_2(x)$
- $u(0,y) = f_3(y)$
- $u(L,y) = f_4(y)$

Write our solution as the following:

$$u(x,y) = u_1(x,y) + u_2(x,y) + u_3(x,y) + u_4(x,y) \tag{325}$$

- $\Delta u_1 = 0$
  - $u_1(x,0) = f_1(x)$
  - $u_1(x,M) = 0$
  - $u_1(0,y) = 0$
  - $u_1(L,y) = 0$
- $\Delta u_1 = 0$
  - $u_2(x,0) = 0$
  - $u_2(x,M) = f_2(x)$

  - $u_2(0, y) = 0$
  - $u_2(L, y) = 0$
- $\Delta u_1 = 0$
  - $u_3(x, 0) = 0$
  - $u_3(x, M) = 0$
  - $u_3(0, y) = f_3(y)$
  - $u_3(L, y) = 0$
- $\Delta u_1 = 0$
  - $u_4(x, 0) = 0$
  - $u_4(x, M) = 0$
  - $u_4(0, y) = 0$
  - $u_4(L, y) = f_4(y)$

This method works because Laplace's equation is linear.

We have already seen that for u_2 :

$$u_2(x, y) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) \sinh\left(\frac{n\pi y}{L}\right) \tag{326}$$

$$u_2(x, M) = f_2(x) \tag{327}$$

$$\Rightarrow \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right) \sinh\left(\frac{n\pi M}{L}\right) \tag{328}$$

$$= f(x) \tag{329}$$

Here, $B_n \sinh\left(\frac{n\pi M}{L}\right)$ is the coefficient for Laplace.

$$B_n \sinh\left(\frac{n\pi M}{L}\right) = \frac{2}{L} \int_0^L f_2(x) \sin\left(\frac{n\pi x}{L}\right) \mathrm{dx} \tag{330}$$

$$B_n = \frac{2}{L \sinh\left(\frac{n\pi M}{L}\right)} \int_0^L f_2(x) \sin\left(\frac{n\pi x}{L}\right) \mathrm{dx} \tag{331}$$

Similarly, for u_4,

$$u_4(x, y) = \sum_{n=1}^{\infty} D_n \sin\left(\frac{n\pi y}{M}\right) \sinh\left(\frac{n\pi x}{L}\right) \tag{332}$$

$$u_4(L, y) = f_4(y) \tag{333}$$

$$\Rightarrow \sum_{n=1}^{\infty} D_n \sin\left(\frac{n\pi y}{M}\right) \sinh\left(\frac{n\pi L}{M}\right) \tag{334}$$

$$= f_4(y) \tag{335}$$

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Recall we are consider $u = u_1 + u + 2 + u_3 + u_4$. Let us write:

$$u_4(x, y) = \sum_{n=1}^{\infty} D_n \sin\left(\frac{n\pi y}{M}\right) \sinh\left(\frac{n\pi x}{M}\right) \tag{336}$$

$$u_4(L, y) = f_4(y) \tag{337}$$

$$= \sum_{n=1}^{\infty} D_n \sin\left(\frac{n\pi y}{M}\right) \sinh\left(\frac{n\pi L}{M}\right) = f_4(y) \tag{338}$$

Recall, our coefficient is D_n and the sinh function.

$$D_n \sinh\left(\frac{n\pi L}{M}\right) = \frac{2}{M} \int_0^M f_4(y) \sin\left(\frac{n\pi y}{M}\right) dy \tag{339}$$

$$D_n = \frac{2}{M \sinh\left(\frac{n\pi L}{M}\right)} \int_0^M f_4(y) \sin\left(\frac{n\pi y}{M}\right) dy \tag{340}$$

Let's look at u_1:

- $\Delta u_1 = 0$
- $u_1(x, 0) = f_1(x)$
- $u_1(x, M) = 0$
- $u_1(0, y) = 0$
- $u_1(L, y) = 0$

1. Here, let us consider $\Delta u_1 = 0$:

$$\frac{X''}{X} = -\frac{Y''}{y} = -\lambda \tag{341}$$

$$u_1(x, y) = X(x)Y(y) \tag{342}$$

Boundary Conditions

$$u_(x, M) = 0 \Rightarrow X(x)Y(M) = 0 \Rightarrow Y(M) = 0 \tag{343}$$
$$u_(0, M) = 0 \Rightarrow X(0)Y(M) = 0 \Rightarrow X(0) = 0 \tag{344}$$
$$u_(L, M) = 0 \Rightarrow X(L)Y(M) = 0 \Rightarrow X(L) = 0 \tag{345}$$
$$\tag{346}$$

2. $\lambda_n = \left(\frac{n\pi}{L}\right)^2$, $X_n(x) = \sin\left(\frac{n\pi x}{L}\right)$
3. Solve for y:

$$\frac{Y''}{Y_n} = \lambda_n \tag{347}$$
$$Y_n'' = \left(\frac{n\pi}{L}\right)^2 Y_n \tag{348}$$
$$Y_n(y) = C\sinh\left(\frac{n\pi y}{L}\right) + D\cosh\left(\frac{n\pi y}{L}\right) \tag{349}$$

Let us see what we have with $Y(m) = 0$:

$$C\sinh\left(\frac{n\pi M}{L}\right) + D\cosh\left(\frac{n\pi M}{L}\right) = 0 \tag{350}$$

Here, this does not work for us. Let us go back and change our $Y_n(y)$:

$$Y_n(y) = C\sinh\left(\frac{n\pi(M-y)}{L}\right) + D\cosh\left(\frac{n\pi(M-y)}{L}\right) \tag{351}$$

Now, let us use our Y:

$$Y_n(M) = C\sinh\left(\frac{n\pi(M-M)}{L}\right) + D\cosh\left(\frac{n\pi(M-M)}{L}\right) \tag{352}$$
$$= D = 0 \tag{353}$$
$$Y_n(y) = \sinh\left(\frac{n\pi(M-y)}{L}\right) \tag{354}$$

4. Let us combine:

$$u_m(x, y) = \sin\left(\frac{n\pi x}{L}\right)\sinh\left(\frac{n\pi(M-y)}{L}\right) \tag{355}$$

By linearity

$$u_1(x, y) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right)\sinh\left(\frac{n\pi(M-y)}{L}\right) \tag{356}$$

5. Find coefficients:

$$u_1(x, 0) = f_1(x) \tag{357}$$
$$= \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right)\sinh\left(\frac{n\pi M}{L}\right) = f_1(x) \tag{358}$$

Once more, we have our coefficient with A_n and sinh.

$$A_n \sinh\left(\frac{n\pi M}{L}\right) = \frac{2}{L}\int_0^L f_1(x)\sin\left(\frac{n\pi x}{L}\right)\,\mathrm{dx} \tag{359}$$
$$A_n = \frac{2}{L\sinh\left(\frac{n\pi M}{L}\right)}\int_0^L f_1(x)\sin\left(\frac{n\pi x}{L}\right)\,\mathrm{dx} \tag{360}$$

Wave Equation

$$\begin{array}{r|c|l} \text{t} & \# & \\ & \# & \\ u = H_1 & \# & u = H_2 \\ & \# & \\ \hline & 0 \quad \text{L} & \end{array}$$

$u(x, 0) = f(x)$
$u_t(x, 0) = g(x)$

Steady-state:

$$u_t = 0 \Rightarrow u_{tt} = 0 \Rightarrow u_{xx} = 0 \Rightarrow u = \frac{H_2 - H_1}{L}x + H_1 \tag{361}$$

Try a solution of the form : $u(x, t) = w(x, t) + u(x, \infty)$. Therefore, $w(x, t) = u(x, t) - u(x, \infty)$.

$$u_{tt} = c^2 u_{xx} \Rightarrow w_{tt} = c^2 w_{xx} \tag{362}$$

Boundary conditions:

$$w(0,t) = u(0,t) - u(0,\infty) = H_1 - H_1 = 0 \tag{363}$$
$$w(L,t) = u(L,t) - u(L,\infty) = H_2 - H_2 = 0 \tag{364}$$

Initial Conditions

$$w(x,0) = u(x,0) - u(x,\infty) = f(X) - \frac{H_2 - H_1}{L}x + H_1 \tag{365}$$
$$w_t(x,t) = u_t(x,t) \Rightarrow w_t(x,0) = u_t(x,0) = g(x) \tag{366}$$

Laplace's Equation in Polar Coordinates

Let's say we want to solve $\Delta u = 0$ with Dirichlet Boundary Conditions on a disk or annulus.

Problem: $\Delta u = u_{xx} + u_{yy} \leftarrow$ in terms of x and y.

We must find it in terms of r and θ.

$$u(x,y) \to u(r,\theta) \tag{367}$$
$$x = r\cos\theta \tag{368}$$
$$y = r\sin\theta \tag{369}$$
$$r = \sqrt{x^2+y^2} \tag{370}$$
$$\theta = \arctan\frac{y}{x} \tag{371}$$
$$\tan\theta = \frac{y}{x} \tag{372}$$

We are going to find : u_x, u_{xx}, u_{yy}

1. u_x :

$$u(x,y) = u(x(r,\theta), y(r,\theta)) = u(r,\theta) = u(r(x,y),\theta(x,y)) \tag{373}$$

Here, we break our chain rule as the following:

$$\begin{array}{cccc} & & u & \\ & | & & | \\ & r & & \theta \\ | & | & | & | \\ x & y & x & y \end{array}$$

According to our tree, we have two routes.

$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial r}\frac{\partial r}{\partial x} + \frac{\partial u}{\partial \theta}\frac{\partial \theta}{\partial x} \tag{374}$$
$$= u_r\frac{x}{\sqrt{x^2+y^2}} + u_\theta\frac{-\frac{y}{x^2}}{1+\left(\frac{y}{x}\right)^2} \tag{375}$$

Note that we know r from line 370. We can rewrite r as $(x^2+y^2)^{\frac{1}{2}}$.

Now, let us multiply our equation by $\frac{x^2}{x^2}$.

$$= u_r\frac{x}{\sqrt{x^2+y^2}} + u_\theta\frac{-\frac{y}{x^2}}{1+\left(\frac{y}{x}\right)^2}\cdot\frac{x^2}{x^2} \tag{376}$$
$$= u_r\frac{r\cos\theta}{r} - u_\theta\frac{r\sin\theta}{r^2} \tag{377}$$
$$= u_r\cos\theta - u_\theta\frac{\sin\theta}{r} \tag{378}$$

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Idea: What if we are on a disk?

$$\Delta u = 0 \Rightarrow u_{xx} + u_{yy} = 0 \tag{379}$$
$$u_x = u_r\cos\theta u_\theta\frac{\sin\theta}{r} \tag{380}$$
$$u_{xx} = \tag{381}$$
$$= \frac{\partial}{\partial x}[u_x] \tag{382}$$
$$= \frac{\partial u_x}{\partial r}\frac{\partial r}{\partial x} + \frac{\partial u_x}{\partial \theta}\frac{\partial \theta}{\partial x} \tag{383}$$
$$= \frac{\partial}{\partial r}\left[u_r\cos\theta - u_\theta\frac{\sin\theta}{r}\right] + \frac{\partial}{\partial\theta}\left[u_r\cos\theta - u_\theta\frac{\sin\theta}{r}\right]\left[-\frac{\sin\theta}{r}\right] \tag{384}$$
$$= \left[u_{rr}\cos\theta - u_{\theta r}\frac{\sin\theta}{r} + u_\theta\frac{\sin\theta}{r^2}\right]\cos\theta + \left[u_{r\theta}\cos\theta - u_r\sin\theta - u_{\theta\theta}\frac{\sin\theta}{r} - u_\theta\frac{\cos\theta}{r}\right]\frac{\sin\theta}{r} \tag{385}$$
$$= u_{rr}\cos^2\theta - 2u_{\theta r}\frac{\sin\theta\cos\theta}{r} + 2u_\theta\frac{\sin\theta\cos\theta}{r^2} + u_r\frac{\sin^2\theta}{r} + u_{\theta\theta}\frac{\sin^2\theta}{r^2} \tag{386}$$
$$u_{yy} = u_{rr}\sin^2\theta + 2u_{\theta r}\frac{\sin\theta\cos\theta}{r} - 2u_\theta\frac{\sin\theta\cos\theta}{r^2} + u_r\frac{\cos^2\theta}{r} + u_{\theta\theta}\frac{\cos^2\theta}{r^2} \tag{387}$$
$$\Delta u = u_{xx} + u_{yy} \tag{388}$$
$$= u_{rr} + \frac{u_r}{r} + \frac{u_{\theta\theta}}{r^2} = 0 \tag{389}$$

Solving Laplace's Equation in Polar Coordinates

If we have an open disk, akin to a washer, we have two boundaries: The innter boundary and outer boundary.

1. Assume $u(r,\theta) = R(r)\Theta(\theta)$

$$u_{rr} + \frac{u_r}{r} + \frac{u_{\theta\theta}}{r^2} = 0 \tag{390}$$
$$R''\Theta + \frac{R'\Theta}{r} + \frac{R\Theta''}{r^2} = 0 \tag{391}$$
$$R''\Theta + \frac{R'\Theta}{r} = -\frac{R\Theta''}{r^2} \tag{392}$$
$$\frac{r^2}{R''}R + \frac{rR'}{R} = -\frac{\Theta''}{\Theta} = \lambda \tag{393}$$

2. Solve for $\Theta : -\frac{\Theta''}{\Theta} = \lambda \Rightarrow \Theta'' = -\lambda\Theta$ If $\lambda > 0$, then we have:

$$\Theta(\theta) = A\sin(\sqrt{\lambda}\theta) + B\cos(\sqrt{\lambda}\theta) \tag{394}$$

Since the solution is periodic in terms of θ : $\begin{cases} \Theta(0) & = \Theta(2\pi) \\ \Theta'(0) & = \Theta'(2\pi) \end{cases}$

$$\Theta(0) = \Theta(2\pi) \Rightarrow B = A\sin(\sqrt{\lambda}2\pi) + B\cos(\sqrt{\lambda}2\pi) \tag{395}$$
$$\Theta'(\theta) = A\sqrt{\lambda}\cos(\sqrt{\lambda}\theta) - B\sqrt{\lambda}\sin(\sqrt{\lambda}\theta) \tag{396}$$
$$\Theta'(0) = \Theta'(2\pi) \Rightarrow A\sqrt{\lambda} = A\sqrt{\lambda}\cos(\sqrt{\lambda}2\pi) - B\sqrt{\lambda}\sin(\sqrt{\lambda}2\pi) \tag{397}$$

These equations are satisfied when $\sqrt{\lambda}2\pi = 2\pi n$, $n \in \mathbb{Z}^+$. Then, $\lambda_n = n^2$. If we consider n, then we can also write $\Theta_n(\theta) = A_n \sin(n\theta) + B_n \cos(n\theta)$.

If $\lambda = 0 : \Theta_0'' = A\theta + B$.

$$\Theta_0(0) = \Theta_0(2\pi) \Rightarrow B = A2\pi + \theta \tag{398}$$
$$\Rightarrow A = 0 \tag{399}$$
$$\Theta_0'(\theta) = 0 \Rightarrow \Theta'(0) = \Theta'(2\pi) \tag{400}$$
$$\Rightarrow \Theta_0(\theta) = B_0 \tag{401}$$

Note: $\lambda < 0$ yields the trivial solution.

3. Solve for R

$$r^2\frac{R_n''}{R_n} + r\frac{R_n'}{R_n} = \lambda_n \Rightarrow r^2R + rR_n' = n^2R_n \tag{402}$$
$$\Rightarrow r^2R_n'' + rR_n' - n^2R_n = 0 \tag{403}$$

Here, let our guess be $R_n(r) = r^m$. Let us plug our guess in:

$$r^2m(m-1)r^{m-2} + rmr^{m-1} - n^2r^m = 0 \tag{404}$$
$$r^2\left(m(m-1) + m - n^2\right) = 0 \tag{405}$$
$$r^2\left(m^2 - n^2\right) = 0 \Rightarrow m = \pm n \tag{406}$$

If $n > 0$,

$$R_n(r) = C_nr^n + D_nr^{-n} \tag{407}$$

If $n = 0$,

$$r^2R_0'' + rR_0' = 0 \tag{408}$$
$$R_0(r) = C_0r^0 + D_0\ln r \tag{409}$$
$$= C_0 + D_0\ln r \tag{410}$$

4. Combine to find u_n and u:

$$u_n(r,\theta) = \begin{cases} B_0(C_0 + D_0\ln r) & n = 0 \\ C_nr^n + D_nr^{-n}\left(A_n\cos(n\theta) + B_n\cos(n\theta)\right) & n \in Z^+ \end{cases} \tag{411}$$

By linearity,

$$u(n\theta) = B_0(C_0 + D_0\ln r) + \sum_{n=1}^{\infty}\left(C_nr^n + D_nr^{-n}\right)\left(A_n\sin(n\theta) + B_n\cos(n\theta)\right) \tag{412}$$
$$= c_0 + d_0\ln r + \sum_{n=1}^{\infty}(a_nr^n + b_nr^{-n})\sin(n\theta) + (c_nr^n + d_nr^{-n})\cos(n\theta) \tag{413}$$

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5. Let us find the oefficients using the boundary conditions.

$$u(R_1,\theta) = g_1(\theta) \tag{414}$$
$$\Rightarrow g_1(\theta) = \underline{C_0 + d_0 \ln R_1} + \sum_{n=1}^{\infty} \left[(\underline{a_n R_1^n + b_n R_1^{-n}}) \sin(n\theta) + (\underline{C_n R_1^n + d_n R_1^{-n}}) \cos(n\theta) \right] \tag{415}$$
$$u(R_2,\theta) = G_2(\theta) \tag{416}$$
$$\Rightarrow g_2(\theta) = \underline{C_0 + d_0 \ln R_2} + \sum_{n=1}^{\infty} \left[(\underline{a_n R_2^n + b_n R_2^{-n}}) \sin(n\theta) + (\underline{C_n R_2^n + d_n R_2^{-n}}) \cos(n\theta) \right] \tag{417}$$

Underlines book-scan: $B_0, A_n, B_n, \widetilde{B}_0, \widetilde{A}_n, \widetilde{B}_n$:

$$\begin{cases} B_0 & = c_0 + d_0 \ln R_1 \\ \widetilde{B}_0 & = c_0 + d_0 \ln R_1 2 \end{cases} \tag{418}$$
$$\begin{cases} A_n & = a_n R_1^n + b_n R_1^{-n} \\ \widetilde{A}_n & = a_n R_2^n + b_n R_2^{-n} \end{cases} \tag{419}$$
$$\begin{cases} B_n & = c_n R_1^n + d_n R_1^{-n} \\ \widetilde{B}_n & = c_n R_2^n + d_n R_2^{-n} \end{cases} \tag{420}$$

<u>Ex:</u> Solve $\Delta u = 0$, where

- $u(1,\theta) = 3\sin(2\theta)$
- $u(2,\theta) = 7\cos(5\theta)$

1. Assume $u(r,\theta) = R(r)\Theta(\theta)$

$$\Delta u = u_{rr} + \frac{u_r}{r} + \frac{u_{\theta\theta}}{r^2} = R''\Theta + \frac{R'\Theta}{r} + \frac{R\Theta''}{r^2} = 0 \tag{421}$$
$$\Rightarrow R''\Theta + \frac{R'\Theta}{r} = -\frac{R\Theta''}{r^2} \tag{422}$$
$$\Rightarrow r^2\frac{R''}{R} + r\frac{R'}{R} = -\frac{\Theta''}{\Theta} = \lambda \tag{423}$$

2. Solve for $\Theta : -\frac{\Theta''}{\Theta} = \lambda \Rightarrow \Theta'' = -\lambda\Theta$.

   If $\lambda > 0$, then

$$\Theta(\theta) = A\sin(\sqrt{\lambda}\theta) + B\cos(\sqrt{\lambda}\theta) \tag{424}$$
$$\Theta'(\theta) = A\sqrt{\lambda}\cos(\sqrt{\lambda}\theta) - B\sqrt{\lambda}\sin(\sqrt{\lambda}\theta) \tag{425}$$
$$\sqrt{\lambda}2\pi = 2n\pi \Rightarrow \lambda_n = n^2, n \in \mathbb{Z}^+ \begin{cases} \Theta(0) = \Theta(2\pi) & \Rightarrow B = A\sin(\sqrt{\lambda}2\pi) + B\cos(\sqrt{\lambda}2\pi) \\ \Theta' = \Theta'(2\pi) & \Rightarrow A\sqrt{\lambda} = A\sqrt{\lambda}\cos(\sqrt{\lambda}2\pi) - B\sqrt{\lambda}\sin(\sqrt{\lambda}2\pi) \end{cases} \tag{426}$$
$$= n^2 \Rightarrow \Theta(n)(\theta) = A_n\sin(n\theta) + B_n\cos(n\theta) \tag{427}$$

   If $\lambda = 0$, then the second derivative is 0.

$$\Theta_0'' \Rightarrow \Theta_0(\theta) = A_0\Theta + B_0 \tag{428}$$
$$\Rightarrow \Theta_0'(\theta) = A_0 \tag{429}$$
$$\Rightarrow \Theta_0(0) = \Theta_0(2\pi) \Rightarrow B_0 = 2\pi A_0 + B_0 \Rightarrow A_0 = 0 \tag{430}$$
$$\Rightarrow \Theta_0'(0) = \Theta_0'(2\pi) = 0 \tag{431}$$

3. Solve for $R : r^2\frac{R_n''}{R_n} + \frac{rR_n'}{R_n} = \lambda_n$

$$r^2R_n'' + rR_n' - n^2R_n = 0 \tag{432}$$
$$\tag{433}$$

   Try $R_n(r) = R^m$, then

$$r^2m(m-1)r^{m-2} + rmr^{m-1} - n^2r^m = 0 \tag{434}$$
$$r^m\left[m(m-1) + m - n^2\right] = 0 \tag{435}$$
$$m^n - n^2 = 0 \tag{436}$$
$$m = \pm n \tag{437}$$

   Next, let us write:

$$\Rightarrow \begin{cases} R_n(r) & = C_nr^n + D_nr^{-n}, n \in \mathbb{Z}^+ \\ R_n(r) & = C_0 + D_0\ln r \end{cases} \tag{438}$$

4. Combine to obtain u_n and u,

$$u_n(r,\theta) = \begin{cases} B_0(C_0 + D_0\ln r) & n = 0 \\ (C_nr^n + D_nr^{-n})(A_n\sin(n\theta) + B_n\cos(n\theta)) & n \in Z^+ \end{cases} \tag{439}$$

By linearity,

$$u(r,\theta) = c_0 + d_0 \ln r + \sum_{n=1}^{\infty} \big((C_n r^n + D_n r^{-n})(A_n \sin(n\theta) + B_n \cos(n\theta))\big) \tag{440}$$
$$= c_0 + d_0 \ln r + \sum_{n=1}^{\infty} \big((a_n r^n + b_n r^{-n}) \sin(n\theta) + (c_n r^n + d_n r^{-n}) \cos(n\theta)\big) \tag{441}$$

5. Find coefficients using BCs :

$$u(1,\theta) = 3\sin(2\theta) \tag{442}$$
$$u(2,\theta) = 7\cos(5\theta) \tag{443}$$
$$u(1,\theta) = c_0 + d_0 \ln(1) + \sum_{n=1}^{\infty} [(a_n + b_n)\sin(n\theta) + (c_n + d_n)\cos(n\theta)] \tag{444}$$
$$\begin{cases} c_0 & = 0 \\ c_n + d_n & = 0 \ \ \forall n \\ a_2 + b_2 & = 3 \\ a_n + b_n & = 0 \ \ \forall n, n \neq 2 \end{cases} \tag{445}$$

Now, let us write:

$$u(2,\theta) = 7\cos(5\theta) \tag{446}$$
$$u(2,\theta) = c_0 + d_0 \ln 2 + \sum_{n=1}^{\infty} \big[(a_n 2^n + b_n 2^{-n})\sin(n\theta) + (c_n 2^n + d_n 2^{-n})\cos(n\theta)\big] = 7\cos(5\theta) \tag{447}$$
$$\begin{cases} c_0 + d_0 \ln 2 & = 0 \Rightarrow d_0 = 0 \\ a_n 2^n + b_n 2^{-n} & = 0 \ \forall n \\ c_5 2^5 + d_5 2^{-5} & = 7 \\ c_n 2^n + d_n 2^{-n} & = 0 \ \forall n, n \neq 5 \end{cases} \tag{448}$$

If $n \neq 5$:

$$\begin{cases} c_n + d_n & = 0 \\ c_n 2^n + d_n 2^{-n} & = 0 \end{cases} \Rightarrow c_n = d_n = 0 \tag{449}$$

If $n \neq 2$:

$$\begin{cases} a_n + b_n & = 0 \\ a_n 2^n + b_n 2^{-n} & = 0 \end{cases} \Rightarrow a_n = b_n = 0 \tag{450}$$

If $n = 5$,

$$\begin{cases} c_5 2^5 + d_5 2^{-5} & = 7 \\ c_5 + d_5 & = 0 \Rightarrow d_5 = -c_5 \end{cases} \tag{451}$$
$$c_5 2^5 - c_5 2^{-5} = 7 \tag{452}$$
$$c_5 (32 - \frac{1}{32}) = 7 \tag{453}$$
$$c_5 = \frac{7}{32 - \frac{1}{32}} \tag{454}$$

If $n = 2$,

$$4(a_2 + b_2 = 3) \tag{455}$$
$$-a_2 2^2 + b_2 2^{-3} = 0 \tag{456}$$

$$\frac{15}{4} b_2 = 12 \Rightarrow b_2 = \frac{48}{15} = \frac{16}{5} \tag{457}$$
$$\Rightarrow a_2 = 3 - b_2 = 3 - \frac{16}{5} = -\frac{1}{5} \tag{458}$$

6. $u_y(x,0) = 0 \Rightarrow X(x)Y'(0) = 0 \Rightarrow Y'(0) = 0$

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Heat and Wave Equations in Polar Coordinates

Heat Equation:

$$u_t = \alpha^2 \Delta u \tag{459}$$

Let $\alpha = 1$, then let us write:

$$u_t = u_{rr} + \frac{u_r}{r} + \frac{u_{\theta\theta}}{r^2} \tag{460}$$

Wave Equation:

$$u_{tt} = c^2 \Delta u \tag{461}$$

Here, let $c = 1$:

$$u_{tt} = u_r r + \frac{u_r}{r} + \frac{u_{\theta\theta}}{r^2} \tag{462}$$

In the last two equations, we worked with three variables: t, r, θ.

$$u_t = u_{rr} + \frac{u_r}{r} + \frac{u_{\theta\theta}}{r^2} \tag{463}$$

Here, assume $u(r, \theta, t) = R(r)\Theta(\theta)T(t)$

$$R\Theta T' = R''\Theta T + \frac{R'\Theta T}{r} + \frac{R\Theta'' T}{r^2} \tag{464}$$

$$\frac{T'}{T} = \frac{R''}{R} + \frac{R'}{rR} + \frac{\Theta''}{r^2\Theta} = -\lambda \tag{465}$$

$$\frac{T'}{T} = -\lambda \tag{466}$$

Here, let us put a pin on T and solve for the second part of the equation:

$$\frac{R''}{R} + \frac{R'}{rR} + \frac{\Theta''}{r^2\Theta} = -\lambda \tag{467}$$

$$\frac{r^2R''}{R} + \frac{rR'}{R} + \frac{\Theta''}{\Theta} = -\lambda r^2 \tag{468}$$

$$\frac{r^2R''}{R} + \frac{rR'}{R} + \lambda r^2 = -\frac{\Theta''}{\Theta} = \mu \tag{469}$$

We now have separate ODEs for each of the functions T, R, Θ. The solution for Θ looks like Laplace in polar. Recall, we set λ as n^2:

$$\frac{r^2R''}{R} + \frac{rR'}{R} + \lambda r^2 n^2 \tag{470}$$

$$r^2R'' + rR' + (\lambda r^2 - n^2)R = 0R'' + \frac{R'}{r} + (\lambda - \frac{n^2}{r^2})R \qquad = 0 \tag{471}$$

Here, this is Bessel's Equation.

We use the power series to solve Bessel's Equation and obtain:

$$R_n(r) = \sum_{i=1}^{\infty} \frac{(-\lambda)^i r^{2i+n}}{2^{n+2}(i+n)!i!} \tag{472}$$

If $\lambda = 1$, we get the Bessel function:

$$J_n(x) = \sum_{i=1}^{\infty} \frac{(-1)^i x^{2i+n}}{2^{n+2}(i+n)!i!} \tag{473}$$

Laplace in Spherical Coordinates

$$\Delta u = u_{xx} + u_{yy} + u_{zz} \tag{474}$$

Using the chain rule, we obtain:

$$\Delta u = u_{\rho\rho} + \frac{2}{\rho}u_\rho + \frac{1}{p^2}u_{\phi\phi} + \frac{\cot\theta}{\rho^2}u_\phi + \frac{1}{\rho^2\sin^2\phi}u_{\theta\theta} = 0 \tag{475}$$

Here, assume $u(\rho, \theta, \Phi) = P(\rho)\Theta(\theta)\Phi(\phi)$:

Equation for Θ

$$(1 - x^2)\Theta''(x) - 2\Theta'(x) + n(n+1)\Theta(x) = 0 \tag{476}$$

Where $x = \cos(\theta)$. This equation is Legendre's Equation.

The solutions are Legendre Polynomials $P_n(x)$:

n	$P_n(x)$
0	1
1	x
2	$\frac{3x^2-1}{2}$
3	$\frac{5x^3-3x}{2}$
4	$\frac{35x^4-30x^2+3}{8}$

Legendre Polynomials ? make an orthogonal set on $[-1, 1]$

$$\int_{-1}^{1} P_m(x)P_n(x)\ \mathrm{dx} = 0, m \neq n \tag{477}$$

Laplace's Equation: $\Delta u = 0$, $u \in \Omega \subset \mathbb{R}$, $x \in \tilde{\Omega} \subseteq \mathbb{R}^n$

Definition: A function that satisfies Laplace's Equation is called a harmonic function.

Definition: Let the ball of radius r centered on point x_0 be:

$$B_r(x_0) = \{x : ||x - x_0||_2 \leq r\} \tag{478}$$

Here, let us revisit what the norm notation indicates:

$$x = \begin{bmatrix} x_1 \\ x_2 \\ \dots \\ x_n \end{bmatrix} \tag{479}$$

$$||x||_{\ell_2} = \sqrt[2]{x_1^2 + x_2^2 + \ldots + x_n^2} \text{ Standard/Euclidean Norm} \tag{480}$$

$$||x||_{\ell_1} = |x_1| + |x_2| + \ldots + |x_n| \text{ Take } |x_n| \text{ for odd powers} \tag{481}$$

$$||x||_{\ell_4} = \sqrt[4]{x_1^4 + x_2^4 + \ldots + x_n^4} \tag{482}$$

$$||x||_{\ell_\infty} = \max\{|x_1|, |x_2|, \ldots, |x_n|\} \tag{483}$$

Here, let us write:

$$||f||_{L_1} = \int_\Omega |f(x)|\ \mathrm{dx} \tag{484}$$

$$||f||_{L_2} = \sqrt{\int_\Omega f(x)^2\ \mathrm{dx}} \tag{485}$$

$$||f||_{L_n} = \sqrt[n]{\int_\Omega f(x)^n\ \mathrm{dx}} \tag{486}$$

$$||f||_{L_\infty} = \text{essup } |f(x)| \tag{487}$$

Essential Supremum, $x \in \Omega$.

ℓ_2 and L_2 are the only two that correspond to a Hilbert space.

Theorem: If u is harmonic and $B_r(x_0) \subset \Omega$, then then average value of u in the ball equals $u(x_0)$.

$$u(x_0) = \frac{\int_{B_r(x_0)} u(x)\ \mathrm{dx}}{\int_{B_r(x_0)} 1\ \mathrm{dr}} \tag{488}$$

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Note: These theorems are t? for any ball of any radians in dimensions n. The theorems don't care about the shape of Ω or Boundary Conditions.

Ex: $n = 1$, $\Delta u = u_{xx} = 0$

Here, the function is a linear function

$$u = Ax + B \tag{489}$$

When $n = 1$, we are working with an interval $[x_0 - r, x_0 + r]$.

Theorem 1:

$$u(x_0) = \frac{\int_{x_0-r}^{x_0+r} u(x)\ \mathrm{dx}}{\int_{x_0-r}^{x_0+r} 1\ \mathrm{dx}} \tag{490}$$

$$= \frac{\int_{x_0-r}^{x_0+r} u(x)\ \mathrm{dx}}{2r} \tag{491}$$

Here, we have the definition for average.

Theorem:

$$u(x_0) = \frac{u(x_0 - r) + u(x_0 + r)}{2} \tag{492}$$

Gauss' MVT in Complex Analysis

If $f(z)$ is analytic, then

$$f(z_0) = \frac{\oint_c f(z)\ \mathrm{dz}}{2\pi r} \tag{493}$$

Fourier Series are for $f(x)$ where x is defined over a finite interval. Fourier Transforms are for $f(x)$ where x is defined on $(-\infty, \infty)$.

There is a different form of the Fourier Series.

Here, let us consider Euler's Formula

$$e^{ix} = \cos x + i \sin x \tag{494}$$

$$e^{ix} = \cos x - i \sin x \tag{495}$$

When we combine both formulas, we get

$$\begin{cases} \cos x & = \frac{e^{ix}+e^{-ix}}{2} \\ \sin x & = \frac{e^{ix}-e^{-ix}}{2i} = \frac{\sin ix}{2i} \end{cases} \tag{496}$$

Now, let us rewrite Fourier Series:

$$f(X) = \sum_{n=0}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) + B_n \cos\left(\frac{n\pi x}{L}\right) \tag{497}$$

Here, let us replace our Fourier Series with terms we found,

$$f(x) = \sum_{n=0}^{\infty} A_n \frac{e^{i\frac{n\pi x}{L}} - e^{-i\frac{n\pi x}{L}}}{2i} + B_n \frac{e^{i\frac{n\pi x}{L}} + e^{-i\frac{n\pi x}{L}}}{2} \tag{498}$$

$$= \sum_{n=0}^{\infty} \left[\frac{A_n}{2i} + \frac{B_n}{2}\right] e^{\frac{in\pi x}{L}} + \left[-\frac{A_n}{2i} + \frac{B_n}{2}\right] e^{-\frac{in\pi x}{L}} \tag{499}$$

$$= \sum_{n=-\infty}^{\infty} \alpha_n e^{\frac{in\pi x}{L}} \tag{500}$$

Here, we found an alternative Fourier Series where

$$\alpha_n = \frac{A_n}{2i} + \frac{B_n}{2} n = 0, 1, 2, \ldots \tag{501}$$

$$\alpha_n = -\frac{A_n}{2i} + \frac{B_n}{2} n = 0, -1, -2, \ldots \tag{502}$$

In the alternative Fourier Series, there is a basis function aside α_n

The basis functions are almost orthogonal

$$\int_{-L}^{L} e^{\frac{im\pi x}{L}} e^{\frac{in\pi x}{L}}\ \mathrm{dx} = \int_{-L}^{L} e^{\frac{i\pi x(m+n)}{L}}\ \mathrm{dx} \tag{503}$$

$$= \frac{L}{i\pi(m+n)} e^{\frac{i\pi x(m+n)}{L}} \Big|_{-L}^{L} \tag{504}$$

$$= \frac{L}{i\pi(n+m)} \left[e^{i\pi(m+n)} - e^{-i\pi(m+n)}\right] \tag{505}$$

$$= \frac{L}{i\pi(m+n)} 2i \sin(\pi(m+n)) = 0 \tag{506}$$

If $m = -n$, then we get:

$$\int_{-L}^{L} e^{\frac{im\pi x}{L}} e^{\frac{-im\pi x}{L}} \text{ dx} = \int_{-L}^{L} 1 \text{ dx} = 2L \tag{507}$$

To find α_n:

$$f(x) = \sum_{n=-\infty}^{\infty} \alpha_n e^{\frac{in\pi x}{L}} \tag{508}$$

Here, let us multiply by $e^{-\frac{in\pi x}{L}}$ and integrate.

$$\int_{-L}^{L} f(x) e^{-\frac{ik\pi x}{L}} \text{ dx} = \sum_{n=-\infty}^{\infty} \alpha_n \int_{-L}^{L} e^{\frac{in\pi x}{L}} e^{\frac{-ik\pi x}{L}} \text{ dx} \tag{509}$$

Here, the integral is 0 except when $k = n$.

$$\int_{-L}^{L} f(x) e^{-\frac{n\pi x}{L}} \text{ dx} = \alpha_n 2L \tag{510}$$

$$\alpha_n = \frac{1}{2L} \int_{-L}^{L} f(x) e^{-\frac{in\pi x}{L}} \text{ dx} \tag{511}$$

$$f(x) = \sum_{n=-\infty}^{\infty} \alpha_n e^{\frac{in\pi x}{L}} \tag{512}$$

$$= \sum_{n=-\infty}^{\infty} \frac{1}{2L} \int_{-L}^{L} f(x) e^{-\frac{in\pi x}{L}} \text{ dx} e^{\frac{in\pi x}{L}} \tag{513}$$

Fourier Transform

Define $\xi_n = \frac{n\pi}{L}$, $\Delta\xi = \frac{\pi}{L}$.

$$f(x) = \sum_{n=-\infty}^{\infty} \frac{\Delta\xi}{2\pi} \int_{-L}^{L} f(x) \text{ dx} e^{i\xi nx} \tag{514}$$

This is a Riemann Sum.

Now let $L \to \infty, \Delta\xi \to d\xi$, replace ξ_n with ξ and $\sum \to \int$.

$$f(x) = \int_{-\infty}^{\infty} \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} \text{ dx } e^{i\xi x} d\xi \tag{515}$$

Define the Fourier Transform

$$F[f] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\xi x} \text{ dx} = \hat{f}(\xi) \tag{516}$$

$$f^{-1}[\hat{f}] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\xi) e^{i\xi x} \text{ d}\xi = f(x) \tag{517}$$

The first line is the Fourier Transform, the second line is the Inverse Fourier Transform.

Note: Laplace Transform

$$F(s) = \frac{1}{\sqrt{2\pi}} \int_{0}^{\infty} f(t) e^{-st} \text{ dt} \tag{518}$$

$$f(t) = \frac{1}{\sqrt{2\pi} i} \int_{c-i\infty}^{c+i\infty} F(s) e^{st} \text{ ds} \tag{519}$$

The first line is the Laplace Transform, whereas the second line is the Inverse Laplace Transform.

We use Laplace Transforms on $[0, \infty)$.

Note: We use Fourier Transforms for functions $f(x)$ where

$$\int_{-\infty}^{\infty} |f(x)| \text{ dx } < \infty \tag{520}$$

Note: $c < \infty$ indicates finite.

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Fourier Transform

$$F[f(x)] = \hat{f}(\xi) \tag{521}$$

Here, $F[f]$ represents the frequencies in f.

Panseval's Equality

1. if $x \in [-L, L]$

$$\frac{1}{2L}\int_{-L}^{L} [f(x)]^2 \text{ dx} = \sum_{n=-\infty}^{\infty} |\alpha_n|^2 \tag{522}$$

   On the left integral, we have the inner product of f with itself. On the right side, we have the coefficients of Fourier Series

2. If $x \in (-\infty, \infty)$

$$\int_{-\infty}^{\infty} [f(x)]^2 \text{ dx} = \int_{-\infty}^{\infty} [\hat{f}(\xi)]^2 \text{ d}\xi \tag{523}$$

Key Property of the Fourier Transform

$$\hat{f}(\xi) = \frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty} f(x)e^{-ix\xi} \text{ dx} \tag{524}$$

$$F\left[\frac{df}{dx}\right] = \frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty} \frac{df}{dx}e^{-ix\xi} \text{ dx} \tag{525}$$

Here, let us integrate our derivative to get $f(x)$.

$$\begin{array}{c|c} u = e^{-ix\xi} & f(x) \\ \hline du = -i\xi e^{-ix\xi} & \frac{df}{dx} \text{ dx} \end{array}$$

$$F\left[\frac{df}{dx}\right] = \frac{1}{\sqrt{2\pi}}\left[f(x)e^{-ix\xi}\Big|_{-\infty}^{\infty} + i\xi\int_{-\infty}^{\infty} f(x)e^{-ix\xi} \text{ dx}\right] \tag{526}$$

Recall, the L_1 norm of f is finite, allowing us to remove the term on the left in our brackets. Here, we are left with the integral:

$$F\left[\frac{df}{dx}\right] = \frac{i\xi}{\sqrt{2\pi}}\int_{-\infty}^{\infty} f(x)e^{-ix\xi} \text{ dx} \tag{527}$$

$$= i\xi\hat{f}(\xi) \tag{528}$$

So a derivative in real space corresponds to multiplication in Fourier Space.

$$F\left[\frac{d^n f}{dx^n}\right] = (i\xi)^n\hat{f}(\xi) \tag{529}$$

We can use the Fourier Transform to help solve any linear PDE where the domain of a spatial variable is $(-\infty, \infty)$.

Linear Equations with Infinite Domains

1. The Transport Equation

$$u_t = cu_x \tag{530}$$

   (a) First order equation
   (b) $x \in (-\infty, \infty)$
   (c) $t \in [0, \infty)$
   (d) In essence, $u(x, 0) = f(x)$

Here, let us guess $u(x, t) = v(x + ct)$. Solutions of this form are called travelling wave equations.

Here, let us establish $\eta = x + ct$

When finding our partials, we run through the following tree:

$$\begin{array}{c} u \\ | \\ v \\ \eta \\ | \quad | \\ x \quad t \end{array}$$

Let's show that this satisfies $u_t = cu_x$.

$$\frac{\partial u}{\partial t} = \frac{\partial v}{\partial t} = \frac{dv}{d\eta}\frac{\partial \eta}{\partial t} = \frac{dv}{d\eta}\cdot c \tag{531}$$

$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial t} = \frac{dv}{d\eta}\frac{\partial \eta}{\partial x} = \frac{dv}{d\eta}\cdot 1 \tag{532}$$

$$u_t = cu_x \tag{533}$$

Any function of the form $u = v(x+ct)$ is a solution to $u_t = cu_x$.

Let's look at the initial condition:

$$u(x,0) = f(x) \tag{534}$$

$$u(x,0) = v(x) \tag{535}$$

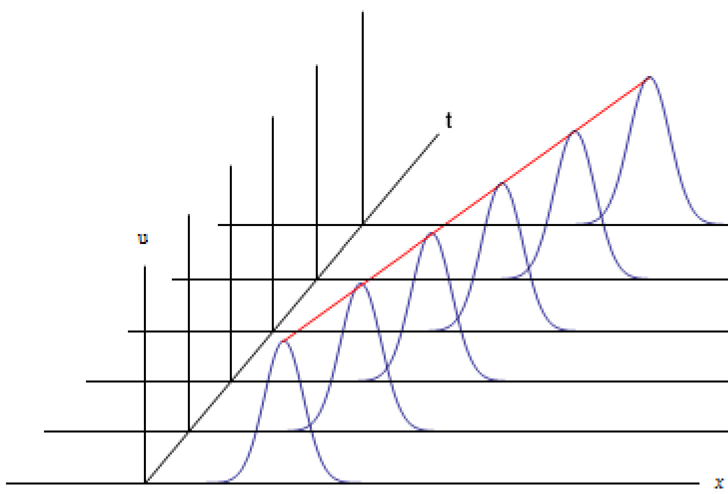
Here, $v(x) = f(X)$. In addition, $u = v(x+ct) = f(x+ct)$.

Ex:

$$u_t = -3u_X \tag{536}$$

$$u(x,0) = e^{-x^2} \tag{537}$$

$$u(x,t) = e^{-(x-3t)^2} \tag{538}$$



The solution translates as time increases, which is why it's called a travelling wave solution.

In this particular case, if $x - 3t =$ constant, then u is fixed.

Remarks

1. The parallel lines are called characteristic curves
2. The slope of the characteristic line is $-\frac{1}{c}$ for $u_t = cu_x$. c is the speed of the solution. It tells us how fast the waves are translated in the $x-$direction.

Lemma

$$\int_{-\infty}^{\infty} e^{-x^2}\ \mathrm{dx} = \sqrt{\pi} \tag{539}$$

Proof. Let us look at our equation squared:

$$I^2 = \int_{-\infty}^{\infty} e^{-x^2}\ \mathrm{dx} \cdot \int_{-\infty}^{\infty} e^{-y^2}\ \mathrm{dy} \tag{540}$$

$$= \int_{-\infty}^{\infty}\int_{-\infty}^{\infty} e^{-x^2}\cdot e^{-y^2}\ \mathrm{dx\ dy} \tag{541}$$

$$= \int_{-\infty}^{\infty}\int_{-\infty}^{\infty} e^{-x^2-y^2}\ \mathrm{dx\ dy} \tag{542}$$

$$= \int_{0}^{2\pi}\int_{0}^{\infty} e^{-r^2} r\ \mathrm{dr\ d}\theta \tag{543}$$

Here, use polar to find our integral. □