Design of an Electromagnetic Calorimeter for use at a future e^+e^- Linear Collider

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Abstract

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Keywords: Particle flow calorimetry, ECAL, Linear Collider, CLIC

1. Introduction

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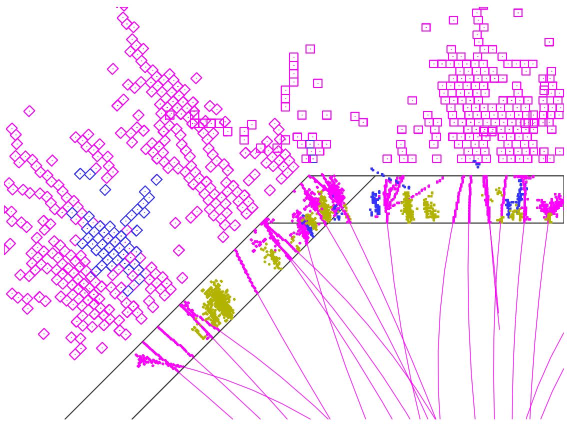

Figure 1

2. Implementation

2.1. Simulation

The simulation of detector model response to various physics events was performed using MOKKA.

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MOKKA uses Geant4 and the geometry information for a given detector model to produce detailed simulations of detector response for various ILC detector concepts. The flexibility in MOKKA allows various detector parameters to be modified and so optimisation studies can be performed. In this study the optimisation was performed with respect to the ILD detector model.

2.2. *Reconstruction*

The reconstruction of physics events was performed using the MARLIN reconstruction framework, which allows for modular implementation of various c++ programs each tasked with one aspect of the reconstruction. While several programs are used in the full reconstruction it is important to highlight firstly to the pattern recognition implementation of particle flow calorimetry, which is done using PandoraPFA, and secondly the digitisation of calorimeter hits, implemented in the ILDCaloDigi processor.

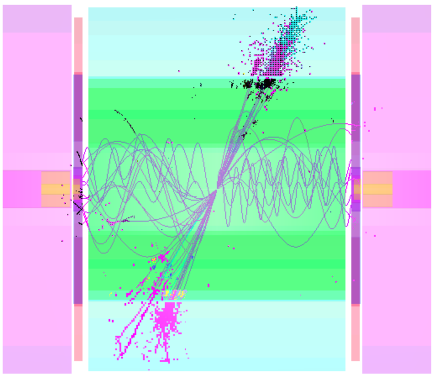

Figure 2: Typical topologies 250GeV jets simulated in the ILD detector model.

2.2.1. *PandoraPFA*

The PandoraPFA software package implements the pattern recognition side of particle flow calorimetry. It is essential that calorimeter hits are correctly assigned to charged particle tracks to avoid double counting of energy in the particle flow paradigm. This can be extremely challenging given the complex topologies associated to particle showers at high energies such as those shown in figure 2.

Due to the varied nature of detector models being simulated in these studies it is necessary to have reusable and flexible software that is isolated from the detector model. This is achieved using the Pandora Software Development Kit (PandoraSDK), which is an independent framework for applying sophisticated topological algorithms designed to perform the association of calorimeter hits to charged particle tracks. The algorithm logic applied in PandoraSDK is independent of detector model.

2.2.2. *ILDCaloDigi*

The ILDCaloDigi processor is designed to perform the digitisation of calorimeter hits for the ILD detector model. Digitisation of calorimeter hits is the process by which the energy deposits in the absorber (non-active) region of a calorimeter cell is estimated using the measured value of the energy deposited in the active region of the cell. Accurate energy estimators for calorimeter cells is crucial for estimating detector performance and so calibration at this stage is essential for all simulations.

A number of realistic effects can be simulated using the ILDCaloDigi processor, the full details of which can be found here. For these studies electrical noise and a limited electronics read out range was simulated and the effect of timing cuts was analysed in detail.

2.3. *Calibration*

To ensure reliability in the conclusions drawn from these optimisation studies, it was necessary to calibrate the response of each detector model considered. This occurred in three stages, the full details of which can be found here.

Initially, the detector response to minimum ionising particles (MIPs) was determined by looking at the detector response to muons. This set the MIP scale in both the digitiser and inside PandoraPFA, which is needed as a reference energy unit for applying thresholds and cuts.

Secondly, the digitisation stage of the reconstruction was tuned for events of photons and long lived neutral kaons for events that were contained within the ECAL and HCAL respectively. This tuning was an iterative procedure where constants in the digitisation processor, ILDCaloDigi, were varied and simulations repeated until the sum of calorimeter hit energies matched the Monte-Carlo energy of the photons and long lived neutral kaons being simulated.

Finally, the electromagnetic and hadronic energy scales within PandoraPFA must be correctly set. This is done by independently scaling the energy of particle flow objects (PFOs), the output reconstructed particles from PandoraPFA, for PFOs originating from electromagnetic and hadronic showers separately. This is again done in an iterative procedure involving changing the inputs to PandoraPFA and repeating simulations until the PFO energy matches the Monte-Carlo energy for the photons (electromagnetic showers) and long lived neutral kaons (hadronic showers) being simulated.

2.4. *Parameterising Detector Performance*

The primary figure of merit used in these optimisation studies is the jet energy resolution, as extensively described in (Pandora paper chapter 5). The jets used in these studies are from the decay of off-shell mass Z bosons decaying at rest into a pair of light quarks (u,d,s). Typically, such decays form mono-energetic jets back to back as can be seen in figure 2.

It is possible to decompose the jet energy resolution into various components by cheating various parts of the reconstruction as shown in figures 3 and 4. This provides a wealth of information that can be used to distinguish performance changes related to pattern recognition from changes related to intrinsic energy resolution.

3. Default Detector Performance

In these studies the reference detector model will be the ILD detector model. However, it is necessary to clarify both the timing cut placed on the simulation and the hadronic energy truncation applied in PandoraPFA to digitised calorimeter cells before accurate comparisons between detector models can be made. Both of these reconstruction variables have a significant impact on detector performance, which is elaborated upon below.

3.1. *Timing Cuts*

The detector at a future linear collider will have timing cuts applied of the order of tens of nano seconds. Timing cuts are necessary to prevent saturation of the detector model when reading out signals from the detector when integrating over multiple collisions. However, a balance has to be struck between having large enough timing cuts to record all of the signal for a collision and not saturating the detector. Currently, technology options for the read out electronics also apply a lower limit of approximately ten nano seconds to the timing cut window.

As it is clear that timing cuts will have to be applied the impact of the choice of timing cuts has been analysed and results shown in figure 5.

3.2. *Hadronic Energy Truncation*

4. ECAL Parameter Scan

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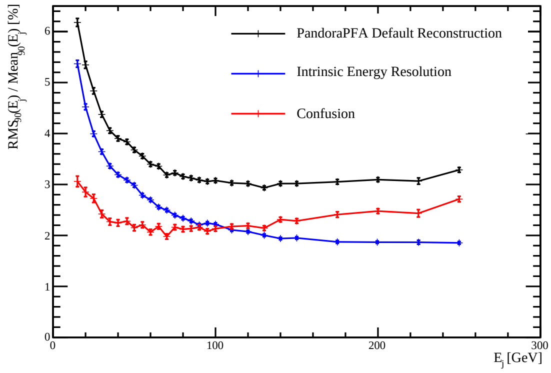

Figure 3: Jet energy resolution as a function of jet energy for the ILD detector model. The jet energy resolution has been decomposed into the intrinsic energy resolution and the confusion terms.

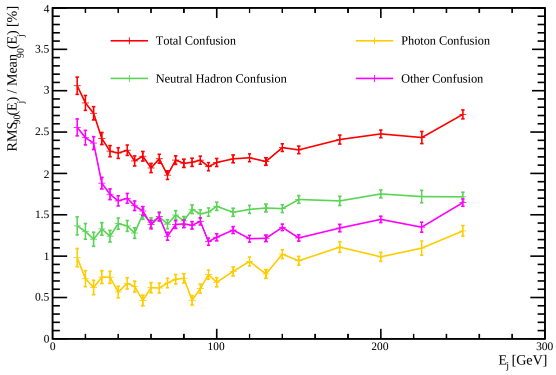

Figure 4: Confusion decomposed into various terms as a function of jet energy. Simulation was performed using the ILD detector model.

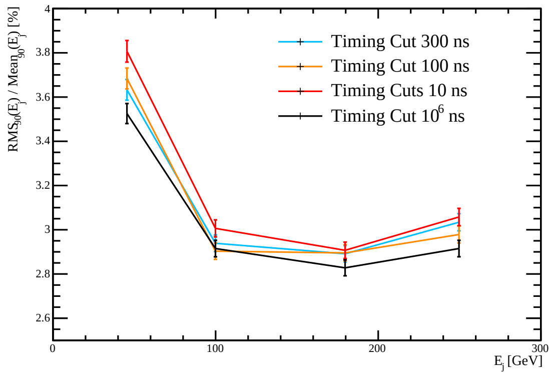

Figure 5: Jet energy resolution as a function of jet energy for the ILD detector model with various timings cuts applied to the reconstruction. The hadronic energy truncation applied to these simulations is 1 GeV.

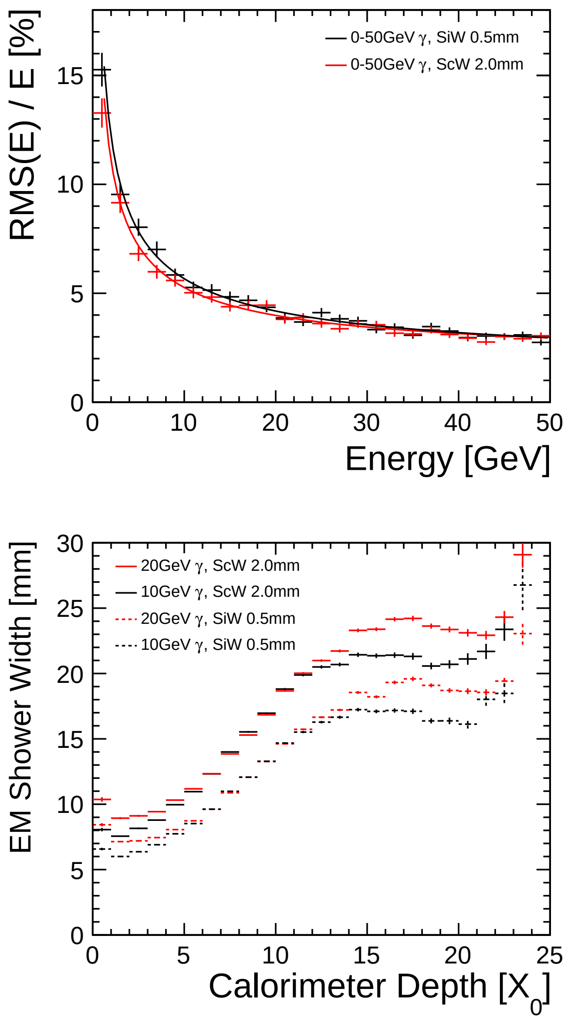

Figure 6

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4.1. Transverse Granularity

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4.2. Number of Layers

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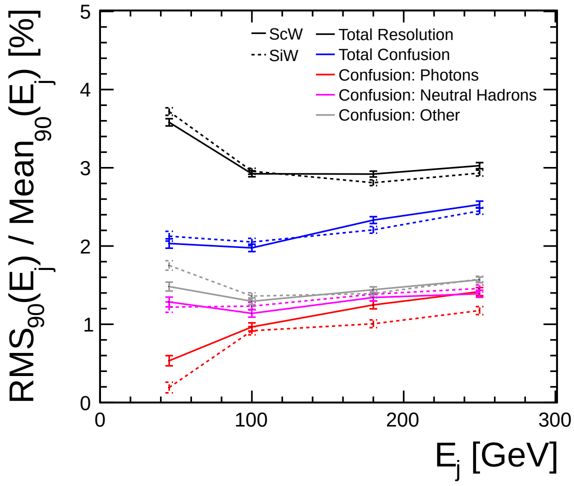
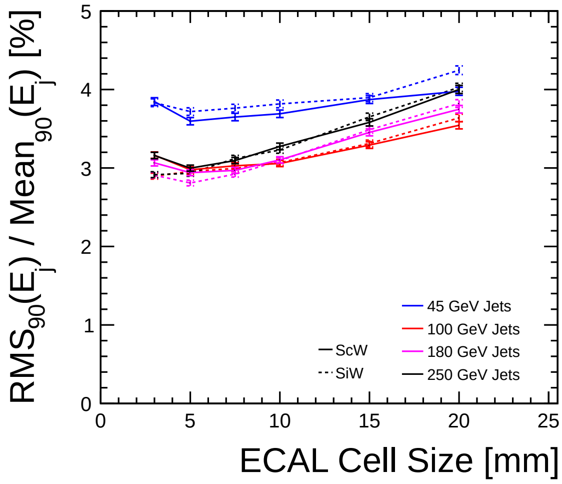


Figure 7



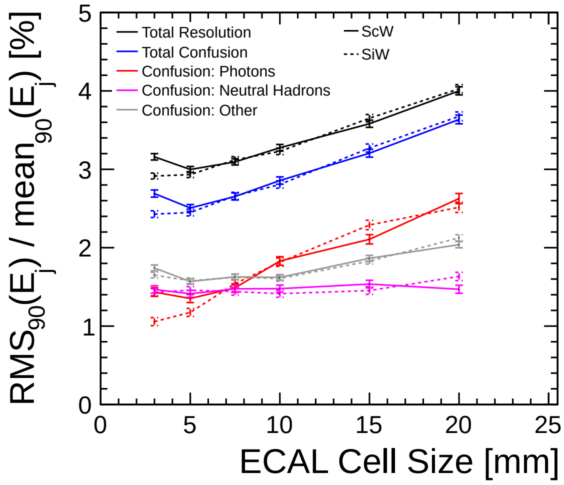


Figure 8

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5. HCAL Parameter Scan

ILD detector model used. Timing cuts used are 10ns in HCal and 20ns in ECal. No MaxHCalHitHadronicEnergy. Full calibration procedure applied for all detector models. QGSP_BERT physics list used in all cases except Fe and W comparison where both QGSP_BERT and QGSP_BERT_HP used.

5.1. Fe/W HCAL Comparison

The feasible options for HCal absorber material are steel (Iron) and tungsten (WMod). The thicknesses of the scintillator and absorber layers of the HCal were scaled to maintain the total number of nuclear interaction lengths in the HCal. The ratio of scintillator and absorber thicknesses in the HCal was held constant to maintain the sampling fraction. Otherwise default ILD DBD detector parameters were used.

Both the QGSP_BERT and QGSP_BERT_HP physics lists were used in this analysis. The high precision neutron package, included in QGSP_BERT_HP, offers more realistic modelling of the transportation of neutrons below 20 MeV down to thermal energies. The compact nature of the hadronic showers showering in tungsten was the primary reason for the inclusion of the high precision neutron package in this analysis.

HCals using steel as an absorber material were found to outperform those using tungsten across the entire range of jet energies considered. This trend was found to be more prominent for high energy jets. No strong dependance on the choice of physics list was observed.

5.2. Transverse Granularity

The transverse granularity of the HCal, the HCal cell size, was varied. Otherwise, default ILD DBD detector parameters were used in this analysis.

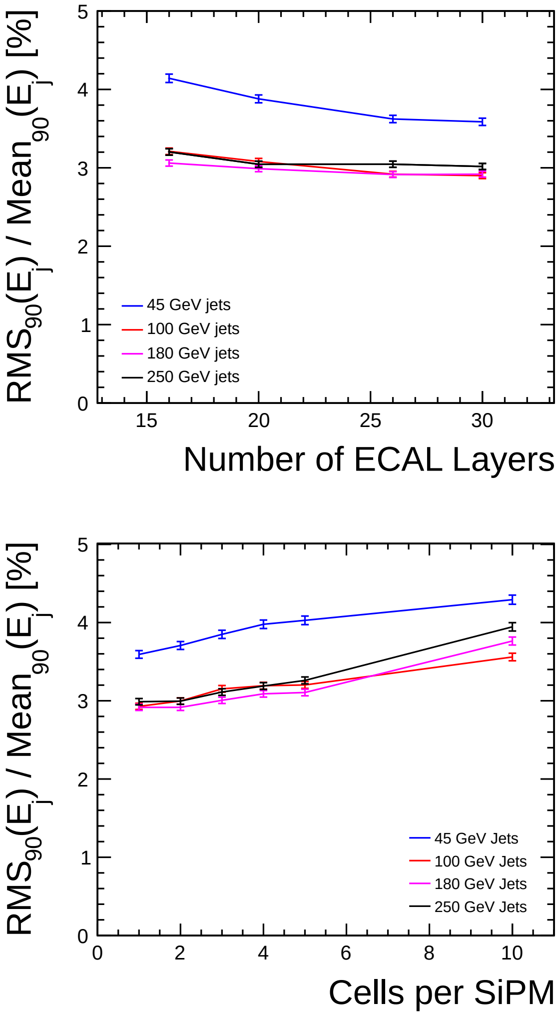

Figure 9

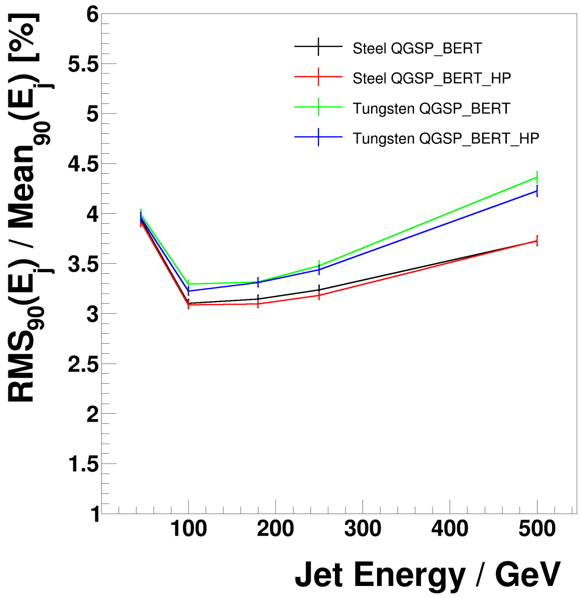

Figure 10: The jet energy resolution as a function of jet energy. Results are shown for detectors using both steel and tungsten HCal absorber materials and using both the QGSP_BERT and QGSP_BERT_HP physics lists.

Figure 11 shows the jet energy resolution as a function of HCal cell size for various energy jets. It was found that smaller HCal cell sizes benefited the jet energy resolution and that this trend became more prominent for higher energy jets. The jet energy resolution was decomposed into the confusion and intrinsic energy resolution terms. The results of this decomposition for the 250 GeV jets are shown in figure 11.

The intrinsic energy resolution was found to be invariant under changes to HCal cell size. This indicates the performance trend observed for HCal cell size are dictated by the confusion term. This is expected as changes to HCal cell size will aid the association of calorimeter cell clusters to charged particle tracks, but will not affect intrinsic energy resolution.

5.3. *Number of Layers*

The number of layers in the HCal was varied. The scintillator and absorber thicknesses were scaled to maintain the total number of nuclear interaction lengths in the HCal. The ratio of scintillator and absorber thicknesses was held constant to maintain the sampling fraction. Otherwise, default ILD DBD detector parameters were used in this analysis.

Figure 12 shows the jet energy resolution as a function of the number of layers in the HCal for various energy jets. It was found that reducing the number of layers in the HCal degrades the jet energy resolution and that this trend applies to all jet energies considered.

Insufficient sampling of particle showers in the HCal is likely to be the primary cause of the degradation at low numbers of layers in the HCal.

5.4. *Depth*

The total number of nuclear interaction lengths in the HCal was varied by changing both the absorber and scintillator thicknesses in the HCal. The ratio of absorber to scintillator thicknesses in the HCal was held constant to maintain the sampling fraction. Otherwise,

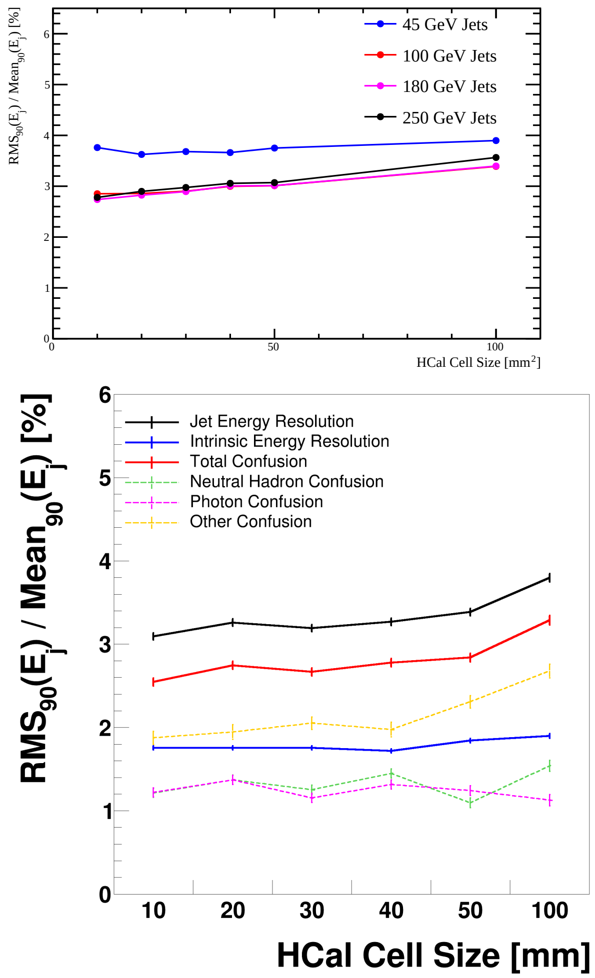

Figure 11: (top) The jet energy resolution as a function of HCal cell size. Results are shown for various jet energies ranging from 45 GeV to 500 GeV. (bottom) The jet energy resolution decomposition as a function of HCal cell size. The results shown are for a jet energy of 250 GeV.

default ILD DBD detector parameters were used in this analysis.

Figure shows the jet energy resolution as a function of the total number of nuclear interaction lengths in the HCal. It was found that increasing the number of nuclear interaction lengths in the HCal improved the jet energy resolution for high energy jets.

The number of nuclear interaction lengths in the HCal will determine the impact of leakage of energy out of the back of the calorimeters. Energy leaked from the back of the calorimeters will be measured in the muon chamber. The energy resolution in the muon chamber is significantly worse than in the calorimeters, therefore, leakage degrades the jet energy resolution. As jet energy increases so does the fraction of the total energy leaking out of the back of the detector. Therefore, reducing leakage is more beneficial to higher energy jets, which is what is observed.

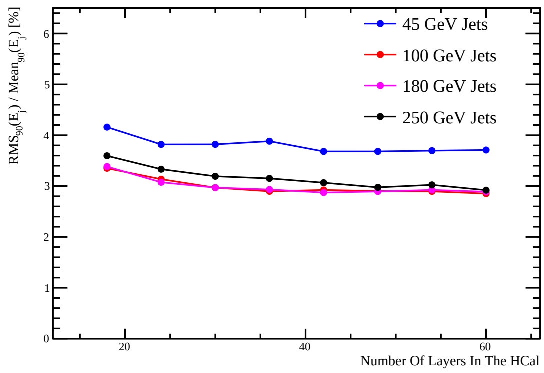

Figure 12: The jet energy resolution as a function of number of layers in the HCal. Results are shown for various jet energies ranging from 45 GeV to 500 GeV.

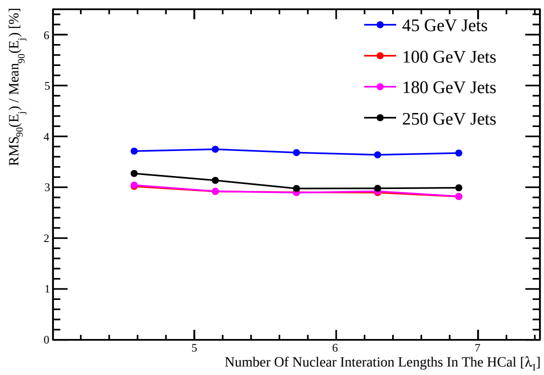

Figure 13: The jet energy resolution as a function of the total number of nuclear interaction lengths in the HCal. Results are shown for various jet energies ranging from 45 GeV to 500 GeV.

5.5. Sampling Fraction in HCal

Sampling fraction varied. Total number of nuclear interaction lengths in HCal held constant.

No picture needed in this section, all jet energy resolutions are flat wrt sampling fraction except at 0.05 (range considered 0.05 to 0.25 steps 0.05) where they start to deteriorate across all jet energies considered.

6. Global Parameter Scan

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6.1. Inner Radius

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6.2. B-Field Strength

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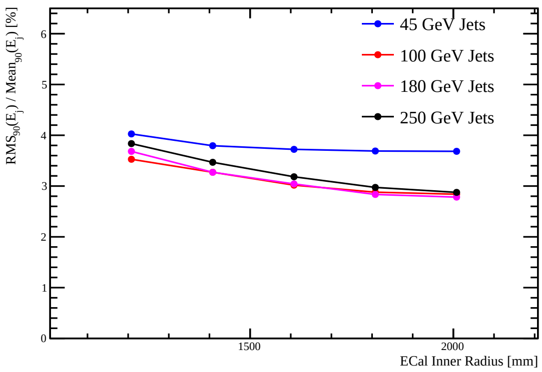

Figure 14: The jet energy resolution as a function of the outer TPC radius. The outer TPC radius can be interpreted as the inner ECal radius. Results are shown for various jet energies ranging from 45 GeV to 500 GeV.

6.3. Scintillator Thickness

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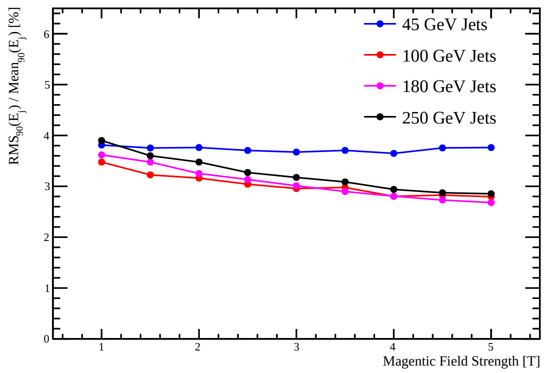

Figure 15: The jet energy resolution as a function of the magnetic field strength within the detector. Results are shown for various jet energies ranging from 45 GeV to 500 GeV.

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6.4. Parameterisation of Results

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7. Novel ECAL Models

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8. Performance for Higher Energy Jets

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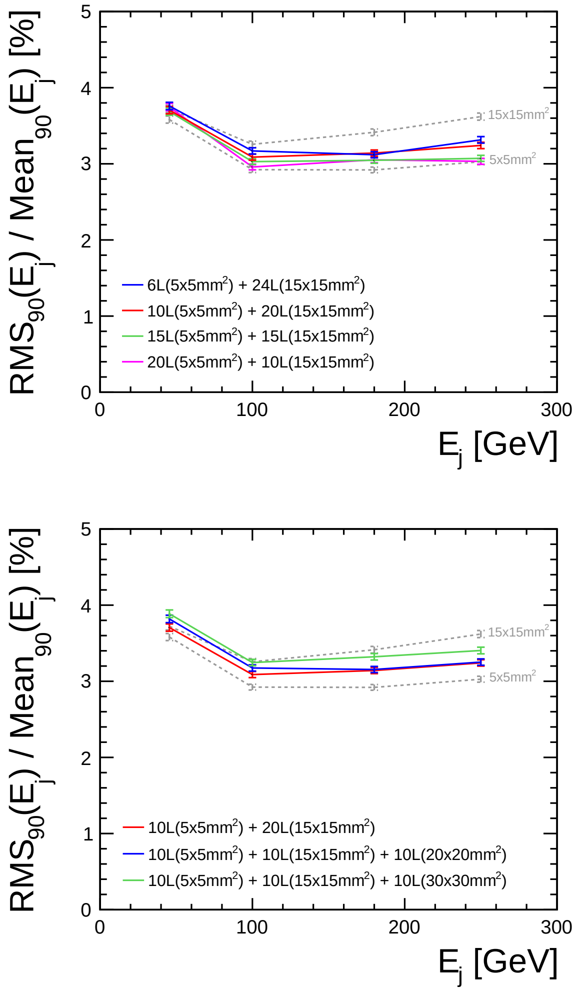

Figure 16

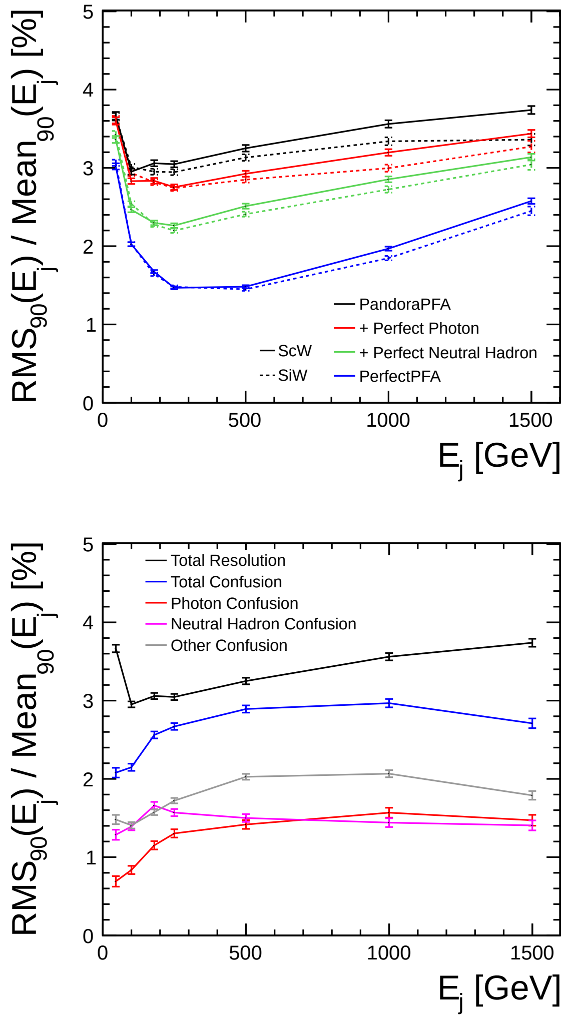

Figure 17

Eodem modo typi, qui nunc nobis videntur parum clari, fiant sollemnes in futurum.

9. Summary

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Acknowledgements

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References