

# **Something something something physics**

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# Abstract

This thesis describes the optimisation of the calorimeter design for collider experiments at the future Compact Linear Collider (CLIC) and the International Linear Collider (ILC). The detector design of these experiments is built around high-granularity Particle Flow Calorimetry that, in contrast to traditional calorimetry, uses the energy measurements for charged particles from the tracking detectors. This can only be realised if calorimetric energy deposits from charged particles can be separated from those of neutral particles. This is made possible with fine granularity calorimeters and sophisticated pattern recognition software, which is provided by the PandoraPFA algorithm. This thesis presents results on Particle Flow calorimetry performance for a number of detector configurations. To obtain these results a new calibration procedure was developed and applied to the detector simulation and reconstruction to ensure optimal performance was achieved for each detector configuration considered.

This thesis also describes the development of a software compensation technique that vastly improves the intrinsic energy resolution of a Particle Flow Calorimetry detector. This technique is implemented within the PandoraPFA framework and demonstrates the gains that can be made by fully exploiting the information provided by the fine granularity calorimeters envisaged at a future linear collider.

A study of the sensitivity of the CLIC experiment to anomalous gauge couplings that effect vector boson scattering processes is presented. These anomalous couplings provide insight into possible beyond standard model physics. This study, which utilises the excellent jet energy resolution from Particle Flow Calorimetry, was performed at centre-of-mass energies of 1.4 TeV and 3 TeV with integrated lumi-

nosities of  $1.5\text{ab}^{-1}$  and  $2\text{ab}^{-1}$  respectively. The precision achievable at CLIC is shown to be approximately one to two orders of magnitude better than that currently offered by the LHC.

Finally, a study into various technology options for the CLIC vertex detector is described.

## Declaration

This dissertation is the result of my own work, except where explicit reference is made to the work of others, and has not been submitted for another qualification to this or any other university. This dissertation does not exceed the word limit for the respective Degree Committee.

Andy Buckley



## Acknowledgements

Of the many people who deserve thanks, some are particularly prominent, such as my supervisor...



## **Preface**

This thesis describes my research on various aspects of the LHCb particle physics program, centred around the LHCb detector and LHC accelerator at CERN in Geneva.

For this example, I'll just mention Chapter ?? and Chapter ??.



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*“Writing in English is the most ingenious torture  
ever devised for sins committed in previous lives.”*

— James Joyce



# **Chapter 1.**

## **Future Linear Collider Experiments**

### **1.1. The International Linear Collider**

### **1.2. The Compact Linear Collider**

#### **1.2.1. Experimental Conditions at CLIC**

The CLIC experiment will operate in a unique environment in comparison to previous generations of lepton colliders and this must be properly accounted for to get an accurate measure of the physics potential that CLIC has to offer. The following aspects of the CLIC experiment present the largest challenges to the physics potential for the CLIC experiment:

- The high bunch charge density. The small beam size at the impact point produces very large electromagnetic fields. These fields can interact with the opposite beam particles causing them to radiate photons in an effect known as beamstrahlung. Beamstrahlung acts to reduce the collision energy of the  $e^+e^-$  pairs.
- Beam related backgrounds. Beamstrahlung photons can subsequently interact to produce background events that must be accounted for. Dominant backgrounds of this form that cannot be easily vetoed in the reconstruction include incoherent pair production of  $e^+e^-$  and  $\gamma\gamma \rightarrow \text{Hadron}$ .
- Fast readout technology is crucial. The CLIC bunch train consists of 312 bunches with a repetition rate of 50 Hz. Each bunch is separated by 0.5ns, therefore, it

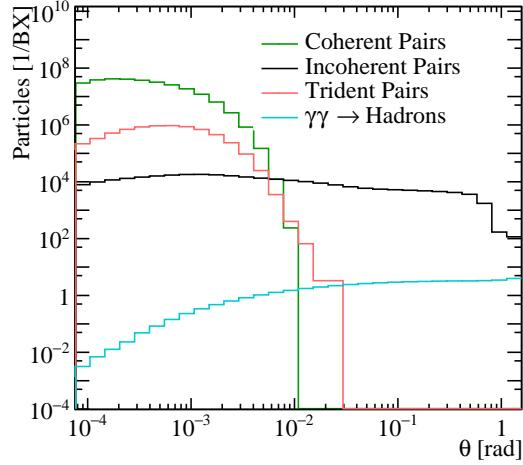
will be necessary to integrate over multiple bunch crossing when reading out the detectors. This places tight constraints on all detector electrical readout speeds and time resolutions.

## Beam-Related Backgrounds at CLIC

The primary sources of background for the CLIC experiment are as follows:

- $e^+e^-$  pair creation from the interaction of a beamstrahlung photons with the opposing beam. The different mechanisms for pair creation are as follows:
  - **Coherent pair production.** This mechanism involves the interaction of a real beamstrahlung photon with the electromagnetic field from the opposing beam.
  - **Trident pair production.** This mechanism involves the interaction of a virtual beamstrahlung photon with the electromagnetic field from the opposing beam.
  - **Incoherent pair production.** This mechanism involves the interaction of a real or virtual beamstrahlung photon with the individual particles in the opposing beam.
- $\gamma\gamma \rightarrow \text{Hadron}$  from the interaction of real or virtual beamstrahlung photons with each other. Example Feynman diagrams for such processes is shown in figure ??.
- Beam halo muons that arise from interactions of the beam particles during collimation. The dominant mechanisms producing beam halo muons are photon conversions into muon pairs ( $\gamma e^- \rightarrow \mu^+ \mu^- e^-$ ) and annihilation of positrons with atomic  $e^-$  into muon pairs ( $e^+ e^- \rightarrow \mu^+ \mu^-$ ) [12].

Each of these has to be properly addressed to get a true measure of the physics potential at CLIC. Coherent and trident pair production is not a dominant source of background as they are produced at low transverse momenta, as figure 1.1 shows, and a simple cut would veto these backgrounds. This is not the case for incoherent pair production of  $e^+e^-$ , which are dominant in the forward regions of the detector, and  $\gamma\gamma \rightarrow \text{Hadron}$ , which are dominant in the tracker and the calorimeters (with the exception of low radii in the calorimeter endcaps) [8, 13]. Beam halo muons are not a major source of background either as they can be easily removed during the



**Figure 1.1:** Angular distribution of number of particles for beam induced backgrounds for CLIC at  $\sqrt{s} = 3$  TeV. Taken from CLIC CDR.

reconstruction due to the clear signal they create in the detector. An algorithm was developed within the PandoraPFA framework for this purpose and it was found to be highly effective at removing the beam halo muons background [8].

$\gamma\gamma \rightarrow$  Hadron events are the most dominant source of background to consider at CLIC as these events deposit more energy throughout the detector than incoherent pair production of  $e^+e^-$  events [8]. The effect of the  $\gamma\gamma \rightarrow$  Hadron background is incorporated into this analysis by overlaying  $\gamma\gamma \rightarrow$  Hadron events onto the event samples used in this analysis. The overlaid backgrounds are added prior to reconstruction so that their effect on the reconstruction is fully accounted for. For a given event the exact number of background events overlaid is drawn from a Poisson distribution with a mean of 3.2 (1.3) events per bunch crossing at 3 (1.4) TeV. While incoherent pairs are still a source of background they will produce a second order effect in comparison to the  $\gamma\gamma \rightarrow$  Hadron events.

The PFO choices described in section 2.9 are applied to veto the effect of PFOs that arise from the overlaid  $\gamma\gamma \rightarrow$  Hadron events.



# **Chapter 2.**

## **The Sensitivity of CLIC to Anomalous Gauge Couplings through Vector Boson Scattering**

*“Kids, you tried your best, and you failed miserably. The lesson is, never try.”*

— Homer Simpson

### **2.1. Motivation**

A potential source of beyond the standard model physics that could be studied at the CLIC experiment is that of the anomalous gauge couplings  $\alpha_4$  and  $\alpha_5$ . The theoretical basis for these couplings is described in section THEORY REF. CLIC would show sensitivity to these couplings through vector boson scattering processes that are summarised in figures ??, ??, ?? and ?. An analysis of the sensitivity of CLIC to this processes is presented in the following chapter.

## 2.2. Event Generation, Simulation and Reconstruction

The event generation software used by the CLIC experiment is Whizard [7, 11]. Whizard version 1.97 was used for generating the new samples, while version 1.95 is used for the official CLIC samples. It was recommended by the Whizard authors to use version 1.97 as it contains a unitarisation scheme that ensures the probabilities remain physical up to high energies when considering the effect of anomalous gauge couplings.

The hadronic channels are the dominant decay modes of the W and the Z boson, with branching fractions of the order of 70% for both (REFERENCE PDG), and as the vector boson scattering is the desired signal channel, the focus of this analysis will be upon the hadronic decays of the W and Z. The vector boson scattering dominated signal final states containing hadronic decay products for the bosons are  $\nu\nu\text{qqqq}$ ,  $l\nu\text{qqqq}$  and  $ll\text{qqqq}$ .

For all samples considered in this analysis, the CLID\_ILD detector [1] was used. The detector was simulated using MOKKA [10], a GEANT4 [2] wrapper providing detailed geometric descriptions of detector concepts for the linear collider. Events were reconstructed using MARLIN [5], a c++ framework designed for reconstruction at the linear collider. PandoraPFA [9, 14] is used to apply particle flow calorimetry in this reconstruction.

The CLIC\_ILD is a variant of the ILD detector described in section REFERENCE. The only significant difference between the models is that CLIC\_ILD has a 60 layer scintillator-tungsten HCal in comparison to the 48 layers found in the default ILD detector. The thicknesses of the layers in the HCal models are identical, so the extra layers correspond to an increase in the total thickness of the HCal. This is needed to compensate for the effects of leakage at the higher energies seen by the CLIC experiment in comparison to the ILC.

Final State	Cross Section 1.4 TeV [fb]
$e^+e^- \rightarrow \nu\nuqqqq$	24.7
$e^+e^- \rightarrow l\nuqqqq$	110.4
$e^+e^- \rightarrow llqqqq$	62.1
$e^+e^- \rightarrow qqqq$	1245.1
$e^+e^- \rightarrow \nu\nu qq$	787.7
$e^+e^- \rightarrow l\nu qq$	4309.7
$e^+e^- \rightarrow llqq$	2725.8
$e^+e^- \rightarrow qq$	4009.5
$\gamma_{\text{EPA}} e^- \rightarrow qqqqe^-$	287.1
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	1160.7
$e^+ \gamma_{\text{EPA}} \rightarrow qqqqe^+$	286.9
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	1156.3
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	32.6
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	136.9
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	32.6
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	136.4
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow qqqq$	753.0
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	4034.8
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	4018.7
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow qqqq$	21406.2

**Table 2.1.:** Cross sections of signal and background processes at 1.4 TeV. In the above table  $q \in u, \bar{u}, d, \bar{d}, s, \bar{s}, c, \bar{c}, b$  or  $\bar{b}$  while  $l \in e^\pm, \mu^\pm$  or  $\tau^\pm$  and  $\nu \in \nu_e, \nu_\mu$  and  $\nu_\tau$ . The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.

## 2.3. Modelling of Anomalous Gauge Couplings

### 2.3.1. Cross Section Sensitivity

To determine which final states are sensitive to  $\alpha_4$  and  $\alpha_5$  a comparison was made between the cross section using the standard model values of  $\alpha_4$  and  $\alpha_5$ , i.e. 0, and the same calculation using non-zero values of these couplings. This comparison was performed on all final states that would be relevant either as signal or background

processes, for an analysis involving the purely hadronic decay channels of a vector boson scattering process. In full the states that were tested are:

- Vector boson scattering signal final states that are expected to show sensitivity to the anomalous couplings:  $e^+e^- \rightarrow \nu\nu qqqq$ ,  $e^+e^- \rightarrow l\nu qqqq$  and  $e^+e^- \rightarrow llqqqq$
- Four jet final states arising from  $e^+e^-$  interactions:  $e^+e^- \rightarrow qqqq$ .
- Two jet final states arising from  $e^+e^-$  interactions:  $e^+e^- \rightarrow \nu\nu qq$ ,  $e^+e^- \rightarrow l\nu qq$ ,  $e^+e^- \rightarrow llqq$  and  $e^+e^- \rightarrow qq$ .
- Four jet final states arising from the interactions of either  $e^+$  or  $e^-$  with a beamstrahlung photon:  $\gamma_{BS}e^- \rightarrow qqqqe^-$ ,  $e^+\gamma_{BS} \rightarrow qqqqe^+$ ,  $\gamma_{BS}e^- \rightarrow qqqq\nu$  and  $e^+\gamma_{BS} \rightarrow qqqq\nu$ .
- Four jet final states arising from the interactions of either  $e^+$  or  $e^-$  with the electromagnetic field of the opposing beam particle. These cross sections are calculated using equivalent photon approximation (EPA), which represents the electromagnetic field of the opposing beam particle as a series of photon and so the final states appear as interactions of  $e^+$  or  $e^-$  with photons:  $\gamma_{EPA}e^- \rightarrow qqqqe^-$ ,  $e^+\gamma_{EPA} \rightarrow qqqqe^+$ ,  $\gamma_{EPA}e^- \rightarrow qqqq\nu$  and  $e^+\gamma_{EPA} \rightarrow qqqq\nu$ .
- Four jet final states arising from the interaction of the electromagnetic fields of opposing beam particles using the EPA approximation:  $\gamma_{EPA}\gamma_{EPA} \rightarrow qqqq$ .
- Four jet final states arising from the interaction of the electromagnetic field of either  $e^+$  or  $e^-$  using the EPA approximation with a beamstrahlung photon:  $\gamma_{EPA}\gamma_{BS} \rightarrow qqqq$  or  $\gamma_{BS}\gamma_{EPA} \rightarrow qqqq$ .
- Four jet final states arising from the interaction of two beamstrahlung photons:  $\gamma_{BS}\gamma_{BS} \rightarrow qqqq$ .

Note: In the above list  $q = u, d, s, c$  and  $b, l = e, \mu, \tau$  and  $\nu = \nu_e, \nu_\mu$  and  $\nu_\tau$ .

The cross section was found to differ when using non-zero values for the anomalous couplings in comparison to the standard model prediction for the vector boson scattering signal final states  $\nu\nu qqqq$ ,  $l\nu qqqq$  and  $llqqqq$ . The cross section comparisons for these final states can be found in table 2.2 for 1.4 TeV samples. In reality, non-zero anomalous couplings would change the cross sections of all processes considered, however, the sensitivity would only arise from high order terms in the Lagrangian. Such terms would not be dominant in determining the cross section and so are omitted

from the generator making certain final states appear invariant to changes in the anomalous couplings.

Final State	Cross Section [fb] ( $\alpha_4 = \alpha_5 = 0.00$ )	Cross Section [fb] ( $\alpha_4 = \alpha_5 = 0.05$ )	Percentage Change[%]	CLIC Cross Section [fb]
$e^+ e^- \rightarrow \nu\nu qqqq$	20.8	34.6	+66.3	24.7
$e^+ e^- \rightarrow l\nu qqqq$	112	113	+0.9	115.3
$e^+ e^- \rightarrow llqqqq$	59.7	68.6	+14.9	62.1

**Table 2.2.:** Cross section for selected processes for given value of  $\alpha_4$  and  $\alpha_5$  at 1.4 TeV.

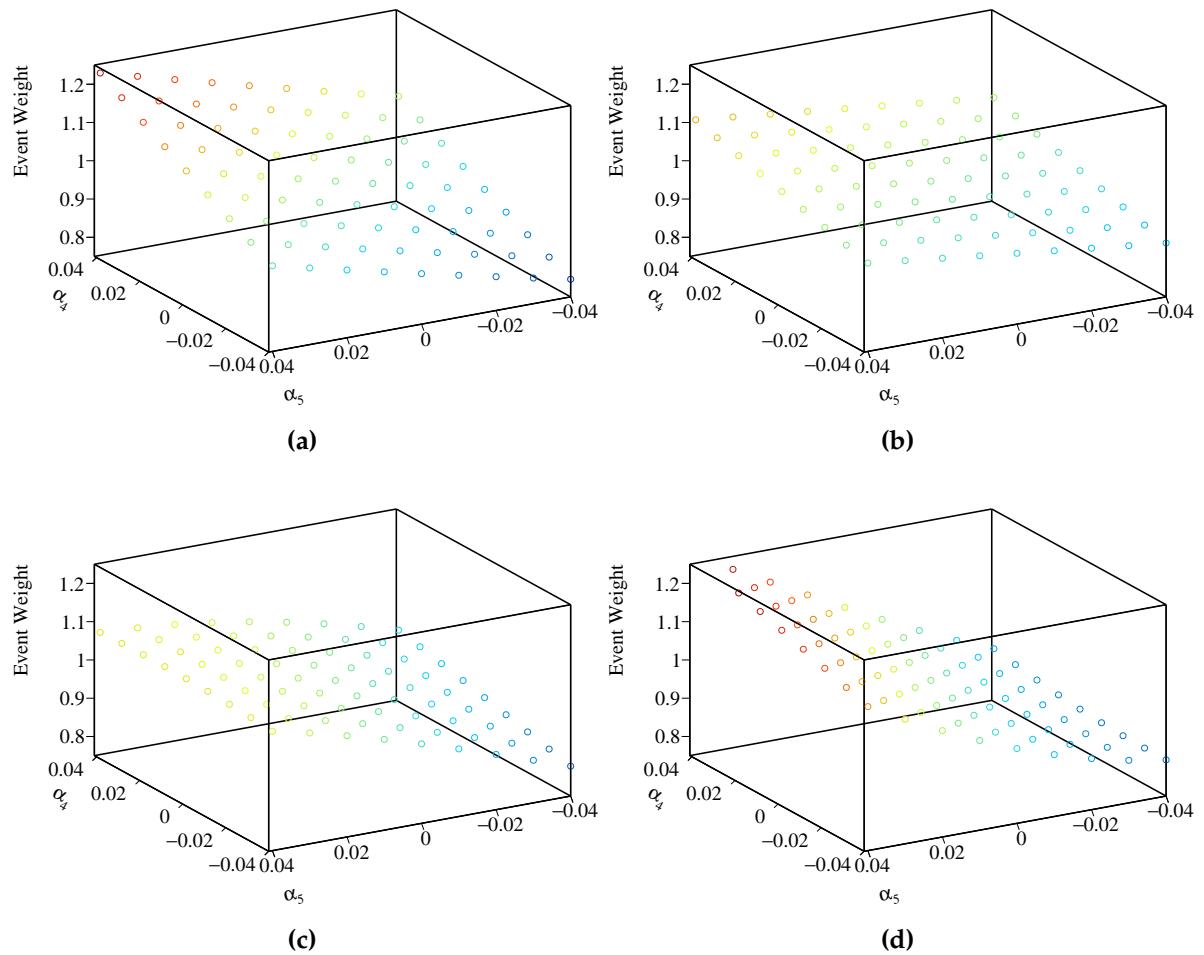
The cross section calculations show that the most sensitive final state to the anomalous gauge couplings is  $\nu\nu qqqq$ , therefore, this analysis will focus solely upon this final state. Furthermore, as the  $l\nu qqqq$  final state has a much reduced sensitivity in comparison to the  $\nu\nu qqqq$  state and as the  $llqqqq$  can be easily vetoed from the analysis, as will be shown in subsequent chapters, it is only necessary to consider the sensitivity of the  $\nu\nu qqqq$  final state. For the aforementioned reasons the  $l\nu qqqq$  and  $llqqqq$  final states will be treated as backgrounds that are invariant to changes in the anomalous couplings  $\alpha_4$  and  $\alpha_5$ .

The sensitivity of an individual event to the anomalous gauge couplings is determined through an event weight. This weight is given by the ratio of the squares of the matrix element used in the cross section calculation, one matrix element using non-zero values of  $\alpha_4$  and  $\alpha_5$  and the other matrix element using the standard model values of  $\alpha_4$  and  $\alpha_5$ , i.e. 0. The weight varies as a function of  $\alpha_4$  and  $\alpha_5$  as well as varying on an event by event basis as the kinematics of the final state changes. Examples of the event weights as a function of  $\alpha_4$  and  $\alpha_5$  for selected events is shown in figure 2.1 for 1.4 TeV events.

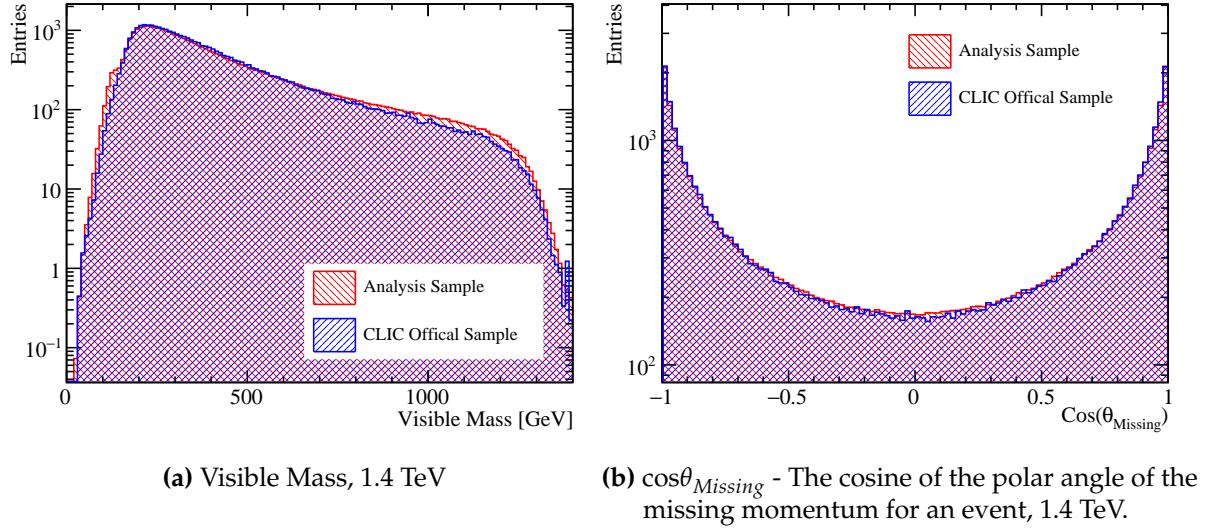
This reweighting procedure has many advantages over the alternative of generating new samples with fixed  $\alpha_4$  and  $\alpha_5$ , notably the absence of systematic errors arising from new event generation, simulation and reconstruction.

### 2.3.2. Validation of Samples

The CLIC experiment has a repository of simulated and reconstructed samples that can be used for physics analyses, however, it is not possible to calculate the event weights



**Figure 2.1.:** A selection of plots showing how the event weight changes when varying the anomalous couplings  $\alpha_4$  and  $\alpha_5$  for 1.4TeV  $\nu\nu\text{qqqq}$  final state events.



**Figure 2.2.:** Comparison of various distributions between samples used in this analysis and the official CLIC samples for the  $\nu\nu\text{qqqq}$  final state at 1.4 TeV. Both samples have been normalised to the correct luminosity for CLIC running at 1.4 TeV.

for these samples as the raw Whizard format event files are missing. Therefore, a new  $e^+e^- \rightarrow \nu\nu\text{qqqq}$  sample was generated with the relevant files to make reweighting as a function of  $\alpha_4$  and  $\alpha_5$  possible. An identical setup to that used for the production of official CLIC samples was used for the event generation, detector simulation and reconstruction. Mimicking this production chain made it possible to use the official CLIC samples for the background final states used in this analysis.

Several reconstructed level distributions were compared to the official CLIC samples and were found to be comparable to each other. A selection of these distributions is shown in figure 2.10.

In order to determine the event weights it was necessary to use the anomalous gauge coupling model in Whizard, which in turn enforces a unit CKM matrix. In the context of vector boson scattering this will restrict the hadronic decays of the  $W^-$  boson to  $d\bar{u}$  and  $s\bar{c}$ , the  $W^+$  boson to  $u\bar{d}$  and  $c\bar{s}$  and the  $Z$  boson to  $u\bar{u}$ ,  $d\bar{d}$ ,  $s\bar{s}$ ,  $c\bar{c}$  and  $b\bar{b}$ . In contrast the official CLIC samples use a non-unit CKM matrix, which gives rise to alternative hadronic decay modes for the  $W$  and  $Z$  bosons. When comparing the unit CKM matrix and the non-unit CKM  $\nu\nu\text{qqqq}$  final state samples it was found that there were negligible differences in a variety of reconstructed level distributions, such as those found in figure 2.10. Furthermore, as flavour tagging information is not used in this analysis this difference was deemed insignificant.

## 2.4. Data Analysis

The focus of this section is to describe the post reconstruction procedure applied to the signal and background events, described in 2.3.1, to extract the relevant information needed for this sensitivity study.

### 2.4.1. Jet Finding

The MarlinFastJet processor, a wrapper for the FastJet [4] processor, is used to cluster the signal and background events into four jets. In the case of the signal final state,  $\nu\nu\text{qqqq}$ , it is assumed these four jets have arisen from the hadronic decays of the bosons involved in the vector scattering process. These jets are paired up on the assumption that the correct pairing is achieved when the invariant masses of the two pairs are closest together.

The longitudinally invariant  $k_t$  jet algorithm in exclusive mode was used for the jet clustering. In contrast to the inclusive mode, the exclusive mode allows the user to request a fixed number of jets in the output from MarlinFastJet. The longitudinally invariant  $k_t$  algorithm proceeds as follows:

- For each pair of particles, i and j, the  $k_t$  distance,  $d_{ij}$ , and beam distance,  $d_{iB} = p_t^2$ , are calculated.

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \Delta R_{ij}^2 / R^2 \quad (2.1)$$

where  $\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$ ,  $p_t$  is the transverse momentum of the particle with respect to the beam axis,  $y_i$  is the rapidity of particle i and  $\phi_i$  is the azimuthal angle of particle i.  $R$  is a configurable parameter that typically is of the order of 1.

- The minimum distance,  $d_{\min}$ , of all the  $k_t$  and beam distances is found. If the minimum occurs for a  $k_t$  distance, merge particles i and j, summing their 4-momenta in the energy combination scheme. If the beam distance is the minimum, declare particle i to be apart of the "beam" jet. Remove this particle from the list of particles and do not included in the final jet output.
- Repeat until the requested number of jets have been created, or inclusive mode no particles are left in the event.

Two other clustering algorithms were considered, but, as figure 2.11 shows, were found to be inappropriate for the experimental conditions at CLIC. These alternative algorithm choices are applied in the same manor as the longitudinally invariant  $k_t$  algorithm, however, they differ in their for the  $k_t$  distance,  $d_{ij}$ , and beam distance,  $d_{iB}$ .

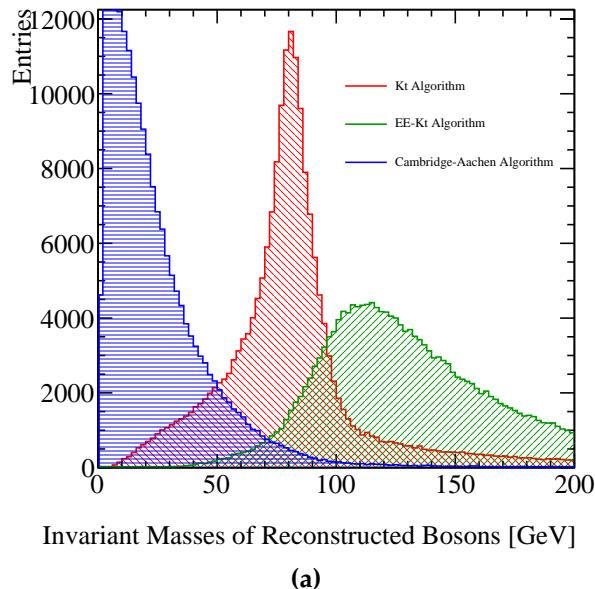
The first alternative jet algorithm considered was the  $k_t$  algorithm for  $e^+e^-$  colliders (or Durham algorithm) where  $d_{ij} = 2\min(E_i^2, E_j^2)(1 - \cos\theta_{ij})$  and  $d_{iB}$  is not used.  $\theta_{ij}$  is the opening angle of particles i and j meaning that in the collinear limit  $d_{ij}$  corresponds to the relative transverse momenta of the particles. The major failure of this algorithm when applied to CLIC is the absence of  $d_{iB}$ , which leads to many beam related background particles being associated to jets. As figure 2.11 shows, the invariant mass of the paired jets, which is expected to be centred around the W and Z boson masses, is much larger than expected, due to the presence of the beam related backgrounds in the jets. Also this algorithm is not invariant to boosts along the beam direction making it inappropriate to use at CLIC as the beam induced backgrounds modify the nominal collision kinematics.

The second alternative jet algortihm considered was the Cambridge/Aachen jet algorithm where  $d_{ij} = \Delta R_{ij}^2 / R^2$  and  $d_{iB} = 1$ . This algorithm gave poor performance as does not account for the transverse momentum or energy of the particles being clustered. In essence, this is a cone clustering algorithm with a cone radius defined through  $\Delta R_{ij} = R$ , which even for large R was found to discard too much energy in the event to be useful for this analysis. This can be seen in figure 2.11 as the invariant mass of the paired jets is much lower than expected. This algorithm is useful for events with highly boosted jets, but at CLIC the jets are too disperse for this algorithm to be successfully applied.

## 2.4.2. Lepton Finding

Brief Unsure of details needed in this section.

An isolated lepton finder is included in the analysis chain in an attempt to reject background events containing leptons.



**Figure 2.3.:** Reconstructed masses for different choices of jet algorithm for 1.4 TeV  $\nu\nu\text{qqqq}$  events. These masses arise by forcing the reconstructed events into 4 jets and then pairing up the jets into pairs such that the reconstructed invariant masses of the pairs are closest to each other. These samples should be dominated by vector boson scattering involving pairs of W bosons and so it is expected that a peak at the W boson true mass should be observed. As this does not occur for the Cambridge-Aachen algorithm or the ee\_kt algorithm they were deemed unsuitable for this analysis at both 1.4. In the case of the kt algorithm and the ee\_kt algorithm an R parameter of 0.7 was used.

### 2.4.3. Discriminant Variables

**Best way to show this?**

Finally, an analysis processor is run, which calculates a number of variables used downstream in the analysis. Included in these are:

- Number of PFOs in the jets and the paired up bosons.
- Number of charged PFOs in the jets and paired up bosons.
- Highest energy PFO: energy, momentum, transverse momentum,  $\cos\theta$ .
- Highest energy electron PFO: energy, momentum, transverse momentum,  $\cos\theta$ .
- Highest energy muon PFO: energy, momentum, transverse momentum,  $\cos\theta$ .
- Highest energy photon PFO: energy, momentum, transverse momentum,  $\cos\theta$ .
- (If in existence) Highest and second highest energy isolated lepton: energy, momentum, transverse momentum,  $\cos\theta$ .
- Bosons: energy, momentum, transverse momentum,  $\cos\theta$ .
- Invariant mass of the boson pair.
- Jets: energy, momentum, transverse momentum,  $\cos\theta$ .
- $\cos\theta$  Of the missing 3-momentum vector.
- Recoil mass.
- Invariant mass of the visible system.
- $y_i, y_{i+1}$ . Jet clustering parameters ranging from  $i = 0$  to 6.
- $\cos\theta_{jet}^*$ . This is the opening angle of a pair of jets, assumed to be from a signle boson, in the rest frame of the boson.
- $\cos\theta_{Boson}^*$ . This is the opening angle of a pair of bosons, assumed to be from vector boson scattering, in the rest frame of the di-boson pair.
- Transverse momentum and energy of the event.
- Acolinearity of the jet pairs forming the bosons and the acoilinearity of the boson pair.

- Principle thrust  $T$  and the thrust axes  $\bar{\mathbf{n}}$ . Note  $\bar{\mathbf{n}}$  is a unit vector. These are defined by the following equation

$$T = \max_{\bar{\mathbf{n}}} \left( \frac{\sum_i \mathbf{p}_i \cdot \bar{\mathbf{n}}}{\sum_i |\mathbf{p}_i|^2} \right) \quad (2.2)$$

- The major and minor thrust values. These are defined with respect to the thrust axes  $\bar{\mathbf{n}}$  in the following way:

$$T = \max_{\bar{\mathbf{n}}_{major}} \left( \frac{\sum_i \mathbf{p}_i \cdot \bar{\mathbf{n}}_{major}}{\sum_i |\mathbf{p}_i|^2} \right) \quad (2.3)$$

where  $\bar{\mathbf{n}}_{major} \cdot \bar{\mathbf{n}} = \mathbf{0}$ . Similarly the minor thrust value is defined as

$$T = \frac{\sum_i \mathbf{p}_i \cdot \bar{\mathbf{n}}_{minor}}{\sum_i |\mathbf{p}_i|^2} \quad (2.4)$$

where  $\bar{\mathbf{n}}_{minor} \cdot \bar{\mathbf{n}} = \bar{\mathbf{n}}_{minor} \cdot \bar{\mathbf{n}}_{major} = \mathbf{0}$

- Sphericity. This is defined using the sphericity tensor  $S^{ab}$  defined as:

$$S^{ab} = \frac{\sum_i p_i^\alpha p_j^\alpha}{\sum_{i,\alpha=1,2,3} |p_i^\alpha|^2} \quad (2.5)$$

Where  $p_i$  are the components of the momenta of particle i in the frame of the detector and the sum runs over all particles in the event. Sphericity is defined as  $S = \frac{3}{2}(\lambda_2 + \lambda_3)$ , where  $\lambda_i$  are the eigenvalues of the sphericity tensor defined such  $\lambda_1 \geq \lambda_2 \geq \lambda_3$ . This provides a measure of how spherical the reconstructed event topology is with isotropic events having  $S \approx 1$ , while two jet events have  $S \approx 0$ . (Also  $\lambda_1 + \lambda_2 + \lambda_3 = 1$ .)

- Aplanarity. Aplanarity is defined as  $\frac{3}{2}\lambda_3$  where  $\lambda_3$  is an eigenvalue of the sphericity tensor. This provides a measure of whether an event is linear or planar.
- B and C tag values for the jets, the min and max B and C tag values for the bosons.

Alongside these variables, for the  $\nu\nu\text{qqqq}$  final state a number of Monte-Carlo variables are calculated for informative purposes and are not used in the analysis. These include:

- The quark and neutrino 4 momenta.

- Invariant mass of boson pair using MC pairing and MC energy.
- Invariant mass of boson pair using MC pairing and reconstructed jet energy.

## 2.5. Event Selection

As discussed earlier the signal events for this analysis contain the  $\nu\nu\text{qqqq}$  final state. The processes to be considered in this analysis alongside the signal are events that would topologically look similar to signal in the detector. This includes events that could be confused with 4 jet events with missing energy, while excluding those events with large numbers of high energy leptons that could be vetoed easily during the analysis stage. In full the list includes:

Equivalent Photon Approximation (EPA) processes model the electromagnetic field of a charged particle as virtual photons. BS (beamstrahlung) processes involve photons that have been radiated from incoming charged particles due to interactions with the electromagnetic field of the opposite beam. The energy spectrum of the incoming particles for CLIC at the relevant operating energy is used to model the energy of these incoming photons. Included in this study are photon-photon interactions from photons appearing from the EPA and beamstrahlung processes.

### 2.5.1. Pre Selection

The primary selection of the  $\nu\nu\text{qqqq}$  signal will be done using a multivariate analysis, however, in an attempt to veto trivial backgrounds a simple cut based preselection is applied. Cuts are applied to the transverse momentum, invariant mass of the visible system and the number of isolated leptons. The raw distributions of these variables is shown in figure 2.4. Based on these distributions the following cuts were applied:

- Transverse momentum  $> 100 \text{ GeV}$ . This cut is effective due to the presence of missing energy in the form of neutrinos in the signal final state.
- Visible mass of the system  $> 200 \text{ GeV}$ . This cut is effective for accounting for the missing energy of the neutrinos in the final state along the longitudinal direction of the detector instead.

Final State	Cross Section 1.4 TeV [fb]
$e^+e^- \rightarrow \nu\nu qqqq$	24.7
$e^+e^- \rightarrow l\nu qqqq$	110.4
$e^+e^- \rightarrow llqqqq$	62.1
$e^+e^- \rightarrow qqqq$	1245.1
$e^+e^- \rightarrow \nu\nu qq$	787.7
$e^+e^- \rightarrow l\nu qq$	4309.7
$e^+e^- \rightarrow llqq$	2725.8
$e^+e^- \rightarrow qq$	4009.5
$\gamma_{\text{EPA}} e^- \rightarrow qqqqe^-$	287.1
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	1160.7
$e^+ \gamma_{\text{EPA}} \rightarrow qqqqe^+$	286.9
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	1156.3
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	32.6
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	136.9
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	32.6
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	136.4
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow qqqq$	753.0
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	4034.8
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	4018.7
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow qqqq$	21406.2

**Table 2.3.:** Cross sections of signal and background processes at 1.4 TeV. In the above table  $q \in u, \bar{u}, d, \bar{d}, s, \bar{s}, c, \bar{c}, b$  or  $\bar{b}$  while  $l \in e^\pm, \mu^\pm$  or  $\tau^\pm$  and  $\nu \in \nu_e, \nu_\mu$  and  $\nu_\tau$ . The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.

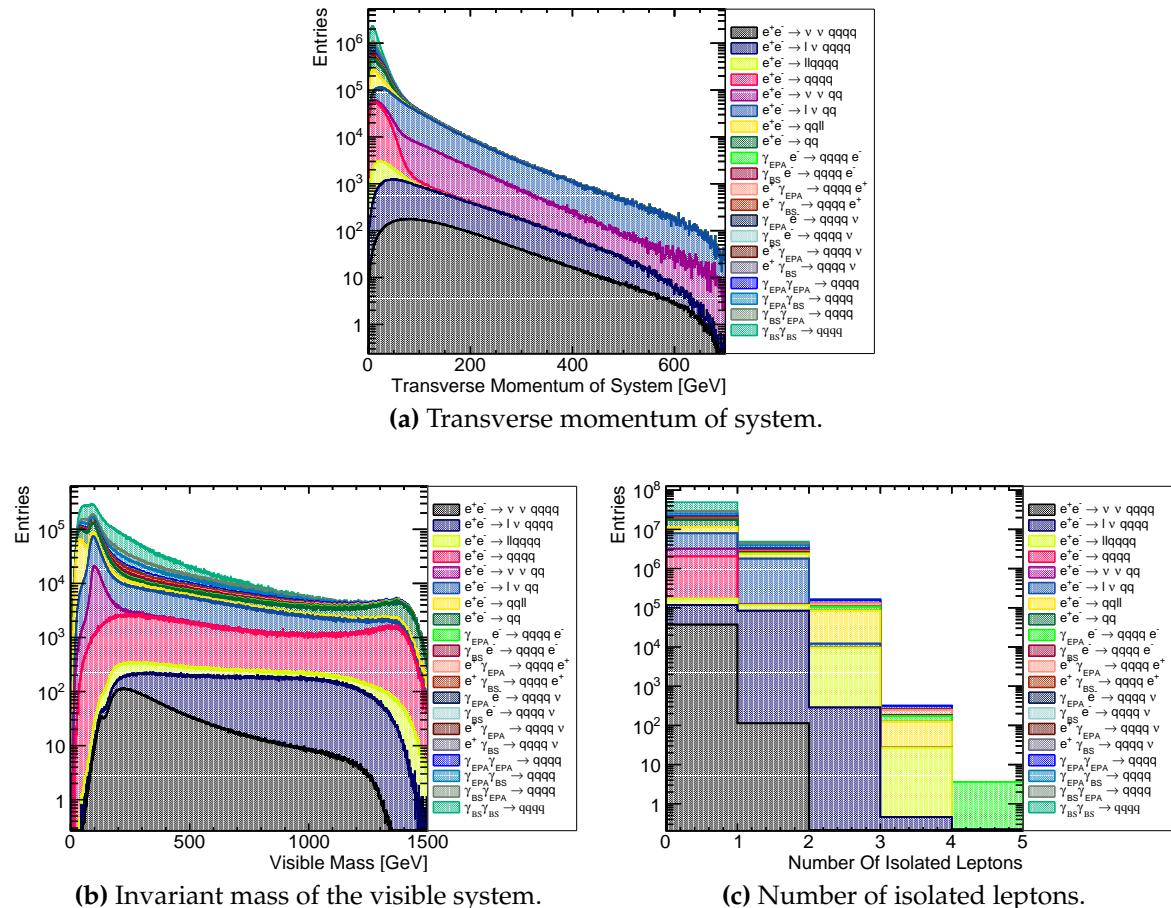
- Number of isolated leptons = 0. This cut vetoes a large number of events with leptons in the final state. The effect of these preselection cuts can be found in table 1.3. While a large fraction of the signal events are lost through these cuts, particularly the transverse momentum cuts, a much large fraction of background events are removed justifying the cut.

The event numbers for the signal and background are shown in table 2.4 as these cuts are cumulatively applied. These numbers are normalised to the correct luminosity for CLIC running at 1.4 TeV. As is expected the large transverse momentum cut removes practically all backgrounds containing no missing energy. The invariant mass

cut removes significant fractions of two quark and missing energy events. Finally, the isolated lepton finder cut removes backgrounds containing visible leptonic final states.

Final State	Raw Event Numbers	$p_T > 100 \text{ GeV}$	$p_T > 100 \text{ GeV} \& M_{\text{Vis}} > 200 \text{ GeV}$	$p_T > 100 \text{ GeV} \& M_{\text{Vis}} > 200 \text{ GeV} \& N_{\text{Isolated Leptons}} = 0$
$e^+ e^- \rightarrow \nu\nu qqqq$	37,050	23,800	21,080	21,020
$e^+ e^- \rightarrow l\nu qqqq$	165,600	81,620	80,840	42,410
$e^+ e^- \rightarrow llqqqq$	93,150	1,151	1,140	700
$e^+ e^- \rightarrow qqqq$	1,868,000	6,487	6,467	6,445
$e^+ e^- \rightarrow \nu\nu qq$	1,181,000	514,100	50,260	50,150
$e^+ e^- \rightarrow l\nu qq$	6,464,000	2,003,000	1,259,000	567,600
$e^+ e^- \rightarrow llqq$	4,088,000	7,754	7,351	5,643
$e^+ e^- \rightarrow qq$	6,011,000	34,610	34,130	34,070
$\gamma_{\text{EPA}} e^- \rightarrow qqqqe^-$	430,600	2,463	2,446	865
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	1,306,000	1,382	1,340	1,002
$e^+ \gamma_{\text{EPA}} \rightarrow qqqqe^+$	430,300	2,846	2,823	1,121
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	1,301,000	654	643	469
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	48,890	17,450	13,490	8,852
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	154,000	56,380	36,350	35,900
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	48,890	17,520	13,550	8,928
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	153,400	56,280	36,340	35,900
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow qqqq$	1,129,000	3,160	3,079	1,563
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	4,539,000	5,325	5,270	3,987
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	4,521,000	3,810	3,730	2,318
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow qqqq$	20,550,000	2,445	2,445	1,673

**Table 2.4.:** Number of events passing the various cuts applied in the preselection at 1.4 TeV. Event numbers are normalised to the correct luminosity for CLIC at 1.4 TeV.  $p_T$  is the transverse momentum of the event,  $M_{\text{Vis}}$  is the visible mass and  $N_{\text{Isolated Leptons}}$  is the number of isolated leptons in the event. In the above table  $q \in u, \bar{u}, d, \bar{d}, s, \bar{s}, c, \bar{c}, b$  or  $\bar{b}$  while  $l \in e^\pm, \mu^\pm$  or  $\tau^\pm$  and  $\nu \in \nu_e, \nu_\mu$  and  $\nu_\tau$ . The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.



**Figure 2.4.: Distribution of variables cut on in the preselection at 1.4 TeV.**

### 2.5.2. MVA

A multivariate analysis was applied to the data set to refine the selection using the TMVA toolkit [6]. The following variables were used for training the TMVA selection.

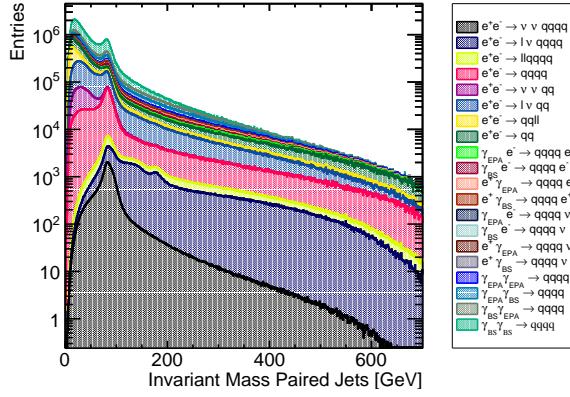
- Number of PFOs in the event.
- Highest energy PFO type.
- Transverse momentum of the event.
- $\cos\theta_{Missing}$ . The cosine of the polar angle of the missing momentum.
- $\cos\theta_{\text{Highest Energy Track}}$ . The cosine of the polar angle of the track with the largest momentum.
- $y_i, y_{i+1}$ . Jet clustering parameters ranging from  $i = 0$  to 6.
- Principle thrust, sphericity and aplanarity as defined in section BLAH.
- Energy of the highest energy electron in the event.
- Energy of the highest energy PFO in the event.
- Energy of the reconstructed bosons.
- Acolinearity of the reconstructed boson pair.
- Invariant mass of the reconstructed bosons.
- Acolinearity of the jets forming the reconstructed bosons.

It was found that the best MVA algorithm for both performance and speed was the booted decision tree (BDT) when comparing different methods using the default settings. Add plot here.

The BDT was further optimised by varying the number of trees used, the depth of the trees and the number of cuts applied. The results shown in the rest of this section use the optimal configuration. For the optimal BDT configuration a significance of  $S/\sqrt{(S+B)} = 53.6$  was obtained.

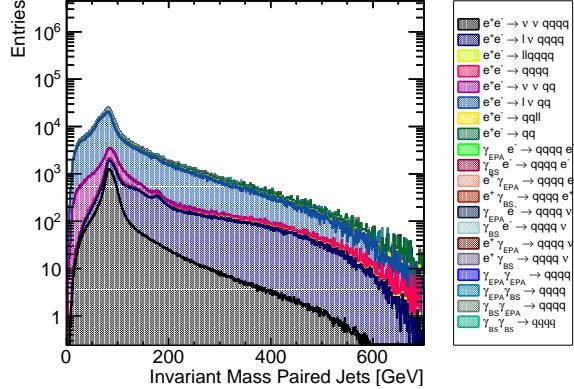
The event numbers passing the BDT cut can be found in table 2.5. The performance of the BDT is shown in figure 2.5, which shows the change in the distribution of the invariant mass of the reconstructed bosons as the MVA is applied. As expected

(a)

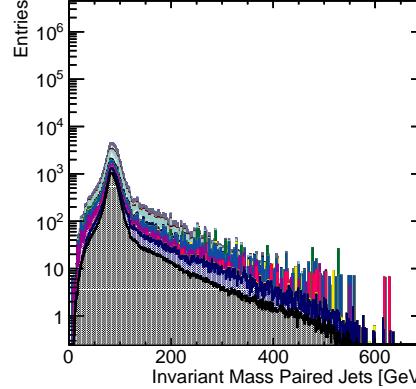


No cuts applied.

(b)



Preselection cuts applied.



Preselection cuts and MVA applied.

**Figure 2.5.:** Impact of preselection and MVA on the reconstructed invariant mass of the bosons arising from jet pairing at 1.4 TeV.

the dominant background processes after the MVA is applied are those that will look identical to the visible signal process i.e.  $qqqq$  and missing energy. Two smaller sources of background that pass the MVA exists, those where two jets and missing energy are confused as four jets and missing energy and those where a lepton is not properly reconstructed and the events look like four jets and missing energy.

The summary of the selection procedure is given in table 2.6.

Final State	Raw Event Numbers	Post MVA Selection Numbers
$e^+e^- \rightarrow \nu\nu\text{qqqq}$	37,050	14,770
$e^+e^- \rightarrow l\nu\text{qqqq}$	165,600	6,159
$e^+e^- \rightarrow ll\text{qqqq}$	93,150	80
$e^+e^- \rightarrow \text{qqqq}$	1,868,000	1,264
$e^+e^- \rightarrow \nu\nu\text{qq}$	1,181,000	3,286
$e^+e^- \rightarrow l\nu\text{qq}$	6,464,000	6,262
$e^+e^- \rightarrow ll\text{qq}$	4,088,000	234
$e^+e^- \rightarrow \text{qq}$	6,011,000	1,016
$\gamma_{\text{EPA}} e^- \rightarrow \text{qqqq}e^-$	430,300	20
$\gamma_{\text{BS}} e^- \rightarrow \text{qqqq}e^-$	1,306,000	42
$e^+ \gamma_{\text{EPA}} \rightarrow \text{qqqq}e^+$	430,300	19
$e^+ \gamma_{\text{BS}} \rightarrow \text{qqqq}e^+$	1,301,000	44
$\gamma_{\text{EPA}} e^- \rightarrow \text{qqqq}\nu$	48,890	3,552
$\gamma_{\text{BS}} e^- \rightarrow \text{qqqq}\nu$	154,000	18,540
$e^+ \gamma_{\text{EPA}} \rightarrow \text{qqqq}\nu$	48,890	3,652
$e^+ \gamma_{\text{BS}} \rightarrow \text{qqqq}\nu$	153,400	18,770
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow \text{qqqq}$	1,129,000	68
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow \text{qqqq}$	4,539,000	55
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow \text{qqqq}$	4,521,000	0
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow \text{qqqq}$	20,550,000	0

**Table 2.5.:** Number of events passing the MVA selection at 1.4TeV. Event numbers are normalised to the correct luminosity for CLIC at 1.4 TeV. The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.

## 2.6. Effect of Anomalous Coupling/Fitting Methodology

The metric used to determine the optimal jet algorithm for this analysis is the sensitivity of the pure signal final state,  $\nu\nu\text{qqqq}$ , to the anomalous gauge couplings. Pure signal was used in the optimisation to avoid the need to process the large number of background files for each iteration of the jet algorithm considered, while still basing the optimisation on the physics of interest. The focus of this section in the description of the fitting technique that will be used for both jet algorithm optimisation and in the final sensitivity study.

Final State	$\epsilon_{\text{presel}}$	$\epsilon_{\text{BDT}}$	$N_{\text{BDT}}$
$e^+e^- \rightarrow \nu\nu qqqq$	56.7%	39.9%	14,770
$e^+e^- \rightarrow l\nu qqqq$	25.7%	3.7%	6,159
$e^+e^- \rightarrow \nu\nu qq$	4.3%	0.3%	3,286
$e^+e^- \rightarrow l\nu qq$	8.8%	0.1%	6,262
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	18.0%	7.3%	3,552
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	23.2%	12.0%	18,540
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	18.2%	7.5%	3,652
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	23.4%	12.2%	18,770

**Table 2.6.:** Selection summary at 1.4TeV. The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung. Channels omitted from this table have less than 1,500 events in the post MVA selection.

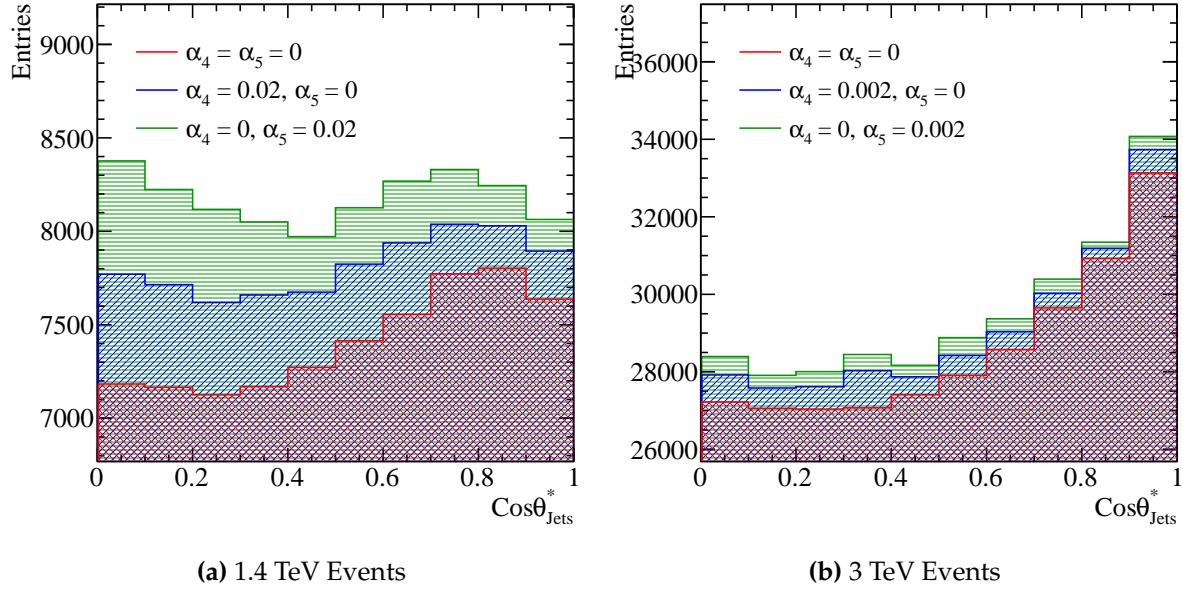
### 2.6.1. Choice of Fitting Distribution

The sensitivity of CLIC to the anomalous gauge couplings is determined through the use of a  $\chi^2$  fit to the distribution of  $\cos\theta_{\text{jets}}^*$ . For a given event, the jet clustering and pairing proceeds as described in section ?? and leads to the event being clustered into four jets, which are then paired up to give two candidate bosons.  $\theta_{\text{jets}}^*$  is defined as the opening angle of the jets in the rest frame of these candidate bosons. The distribution of  $\cos\theta_{\text{jets}}^*$  proved to be highly sensitive to the anomalous gauge couplings as shown in figure 2.6.

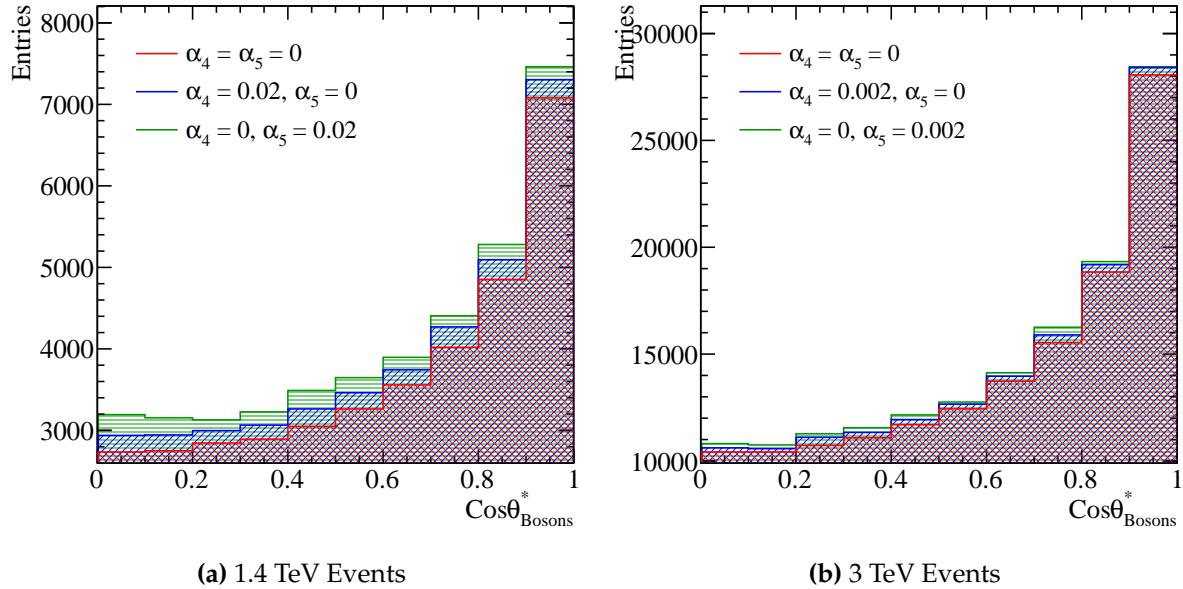
The distribution of  $\cos\theta_{\text{Bosons}}^*$  was also considered for this sensitivity study, however, it proved to be less sensitive than  $\cos\theta_{\text{jets}}^*$ . This can be seen when comparing figures 2.6 and 2.7.  $\theta_{\text{Bosons}}^*$  is defined as the opening angle between the two candidate bosons two bosons in the rest frame of the candidate boson pair. Furthermore, it was found that the  $\chi^2$  distribution formed from the two dimensional distribution of  $\cos\theta_{\text{jets}}^*$  against  $\cos\theta_{\text{Bosons}}^*$  did not significantly benefit the sensitivity in comparison using the one dimensional distribution of  $\cos\theta_{\text{jets}}^*$  and therefore was not considered for this analysis.

### 2.6.2. Application Of Anomalous Gauge Coupling Event Weights

As described in section 2.3, event weights are used to determine the sensitivity of CLIC to the anomalous gauge couplings. These event weights are extracted on an event



**Figure 2.6.:** Sensitivity of  $\text{cos}\theta_{\text{jets}}^*$  to anomalous couplings at 1.4 and 3 TeV. The jet algorithm used for this example was the longitudinally invariant kt algorithm with an R parameter of 0.7. This sample corresponds to pure signal of hadronic decays in vector boson scattering i.e.  $\nu\nu\text{qqqq}$ .



**Figure 2.7.:** Sensitivity of  $\text{cos}\theta_{\text{Bosons}}^*$  to anomalous couplings at 1.4 and 3 TeV. The jet algorithm used for this example was the longitudinally invariant kt algorithm with an R parameter of 0.7. This sample corresponds to pure signal of hadronic decays in vector boson scattering i.e.  $\nu\nu\text{qqqq}$ .

by event basis for the signal final state  $\nu\nu\text{qqqq}$  from the generator software Whizard. To achieve a smooth  $\chi^2$  distribution a fine sampling of the  $\cos\theta_{jets}^*$  distribution in the  $\alpha_4$  and  $\alpha_5$  space is needed. However, as extracting the event weights is highly CPU intensive, it is unfeasible to produce a finely sampled grid of event weights on an event by event basis by calling the generator. To resolve this issue, an interpolation scheme was applied to determine the event weights within a sampled region of the  $\alpha_4$  and  $\alpha_5$  space. This allows for an infinite sampling of the  $\cos\theta_{jets}^*$  distribution in the space of  $\alpha_4$  and  $\alpha_5$  within the sampled region, without having to call the generator an infinite number of times.

A bicubic interpolation scheme, cubic interpolation along two dimensions, was applied to the event weights that were extracted from the generator. This procedure is best illustrated by showing the interpolated surface superimposed with the raw event weights from the generator, which is shown for several  $\nu\nu\text{qqqq}$  events at 1.4 TeV in figure 2.8 (ADD 3 TEV). This interpolation scheme produces a smooth and continuous surface that is sufficiently accurate for the fitting procedure applied in this analysis.

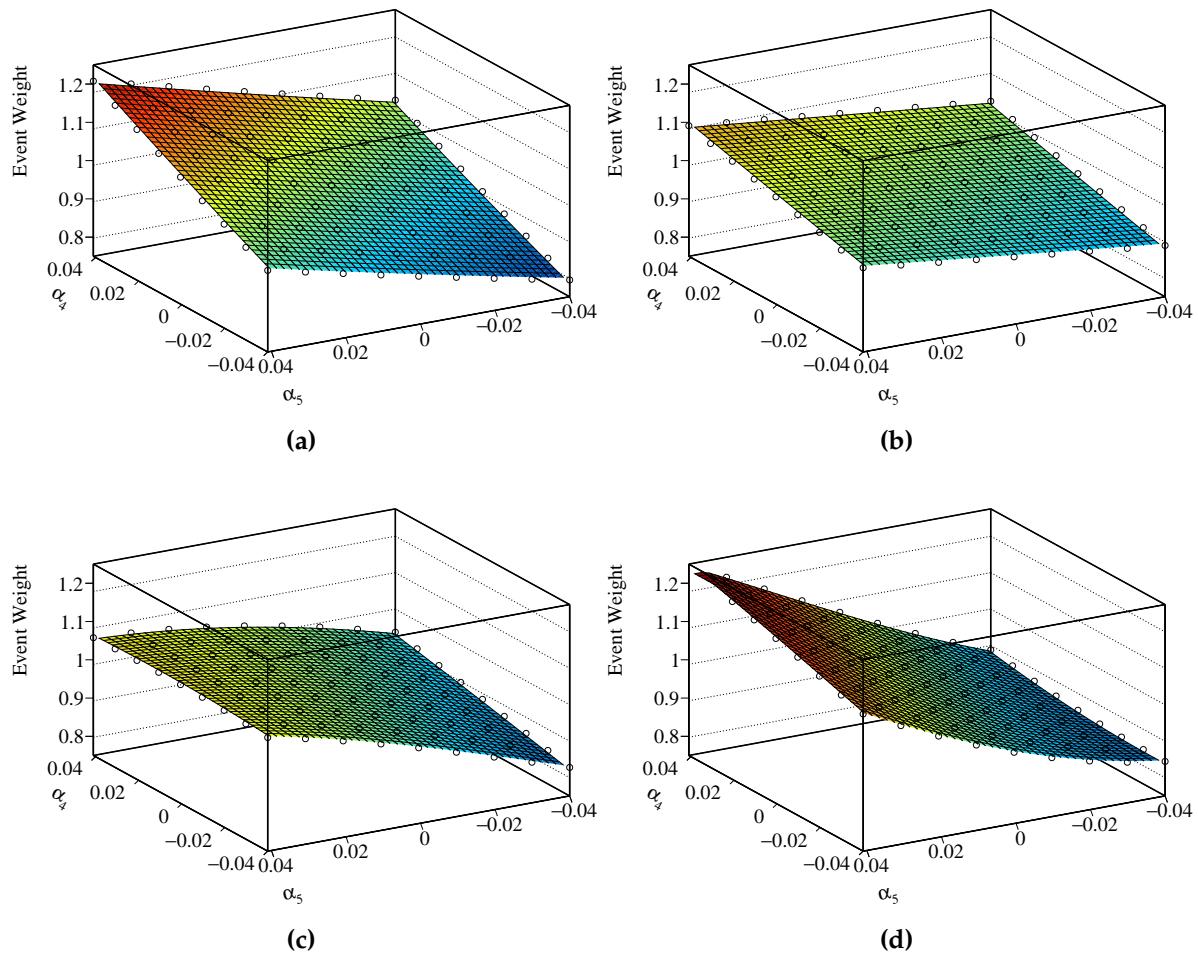
For reference at 1.4 TeV event weights were produced from the generator, Whizard, by stepping along  $\alpha_4$  and  $\alpha_5$  in steps of 0.01 ranging from -0.07 to 0.07, as shown in figure 2.1, while at 3 TeV event weights were samples in steps of 0.00025 from -0.0045 to 0.0045. These ranges proved to be sufficient for the contours of interest for the CLIC sensitivity analysis at these energies.

### 2.6.3. Determination of Sensitivity

The sensitivity of CLIC to the anomalous gauge couplings  $\alpha_4$  and  $\alpha_5$  was determined using a  $\chi^2$  of the following form:

$$\chi^2 = \sum_i \frac{(O_i - E_i)^2}{E_i} \quad (2.6)$$

Here, where  $O_i$  is the observed, or data, bin content for bin  $i$  in the distribution of  $\cos\theta_{jets}^*$  produced with event weights corresponding to zero  $\alpha_4$  and  $\alpha_5$  and  $E_i$  is the expected, or Monte-Carlo, bin content for bin  $i$  produced with event weights corresponding to non-zero  $\alpha_4$  and  $\alpha_5$ . The distribution of  $\cos\theta_{jets}^*$  was binned in a



**Figure 2.8:** A selection of plots showing how the event weight changes when varying the anomalous couplings  $\alpha_4$  and  $\alpha_5$  for 1.4TeV  $\nu\nu\text{qqqq}$  final state events. The hollow circles show the event weight produced from the generator while the surface shown is found using bicubic interpolation between those points.

histograms containing 10 bins ranging from 0 to 1, as shown in figure 2.6. This binning was selected to maximise the sensitivity of the distribution, while minimising the effect of large bin by bin fluctuations arising from individual events with large event weights.

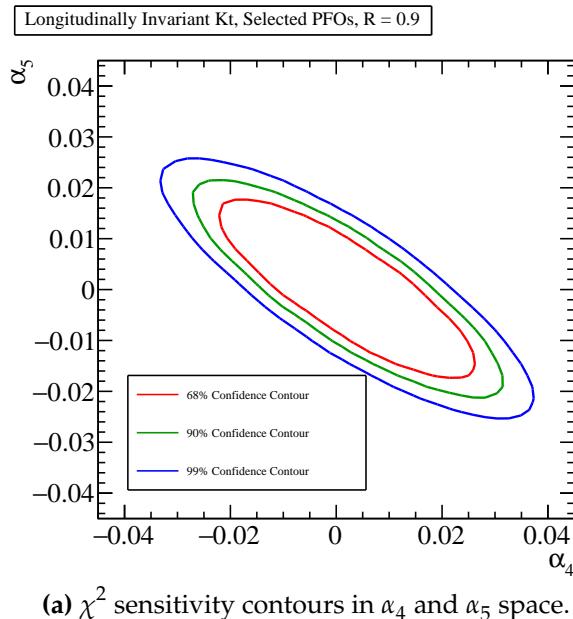
Confidence limits that describe the sensitivity of CLIC to the anomalous gauge couplings, were found by examining the  $\chi^2$  surface in  $\alpha_4$  and  $\alpha_5$  space. Deviations about the minima of this surface, which by construction occurs at  $\alpha_4 = \alpha_5 = 0$ , yield confidence limits that indicate the probability of observing a particular value of  $\alpha_4$  and  $\alpha_5$  based on the  $\cos\theta_{jets}^*$  distribution. The confidence limits used in subsequent sections, 68%, 90% and 99%, are defined using fixed deviations from the minima of  $\chi^2$  contours of 2.28, 4.61 and 9.21 respectively. These numbers arise from the integral of the two dimensional  $\chi^2$  function.

It proved useful to consider the sensitivities to the individual parameters  $\alpha_4$  and  $\alpha_5$  independently. This was done by projecting out the  $\alpha_4 = 0$  or  $\alpha_5 = 0$  one dimensional  $\chi^2$  distribution from the two dimensional  $\chi^2$  previously discussed. It was then possible to extract the sensitivity to an individual parameters using confidence limits arising from the integral of the one dimensional  $\chi^2$  function i.e. 68% confidence limit occurs for  $\chi^2 = 0.989$ . In subsequent chapters these are the sensitivities quoted for individual anomalous gauge coupling parameters.

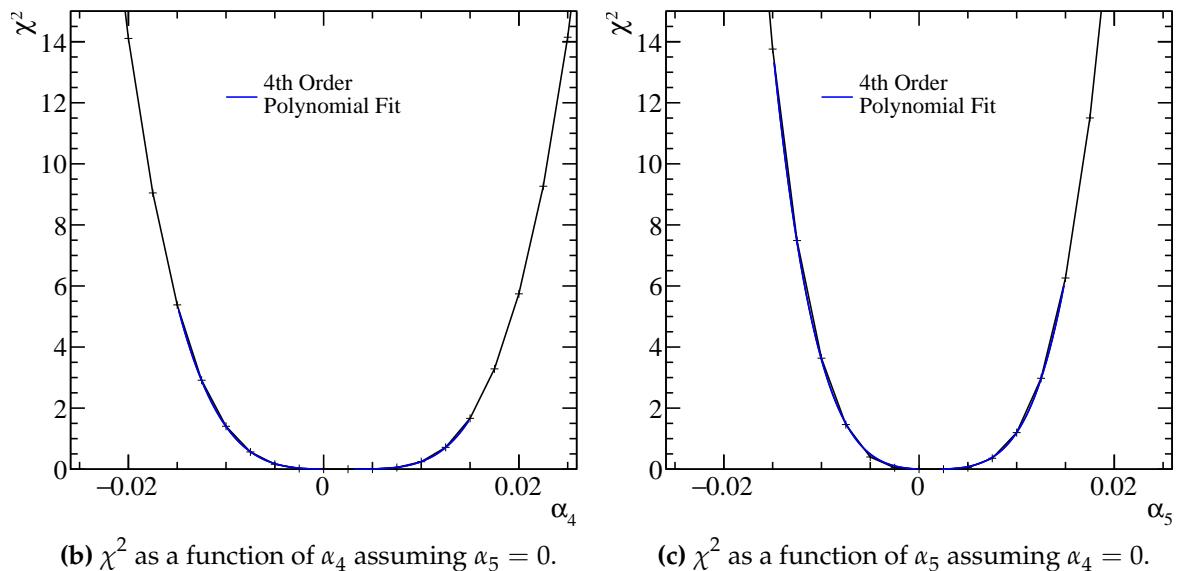
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## 2.7. Results

The sensitivity of the CLIC experiment to the anomalous gauge couplings  $\alpha_4$  and  $\alpha_5$  at 1.4 TeV is shown in figure 2.9a. This result shows the sensitivity after the application of preselection and MVA described in sections 2.5.1 and 2.5.2 purposed to remove the included background channels, described in section 2.2. These contours yield the one  $\sigma$  confidence limit on the measurement of  $\alpha_4$  to the range -0.00831, 0.0130 and similarly for the measurement of  $\alpha_5$  the range is -0.00606, 0.00904.



(a)  $\chi^2$  sensitivity contours in  $\alpha_4$  and  $\alpha_5$  space.



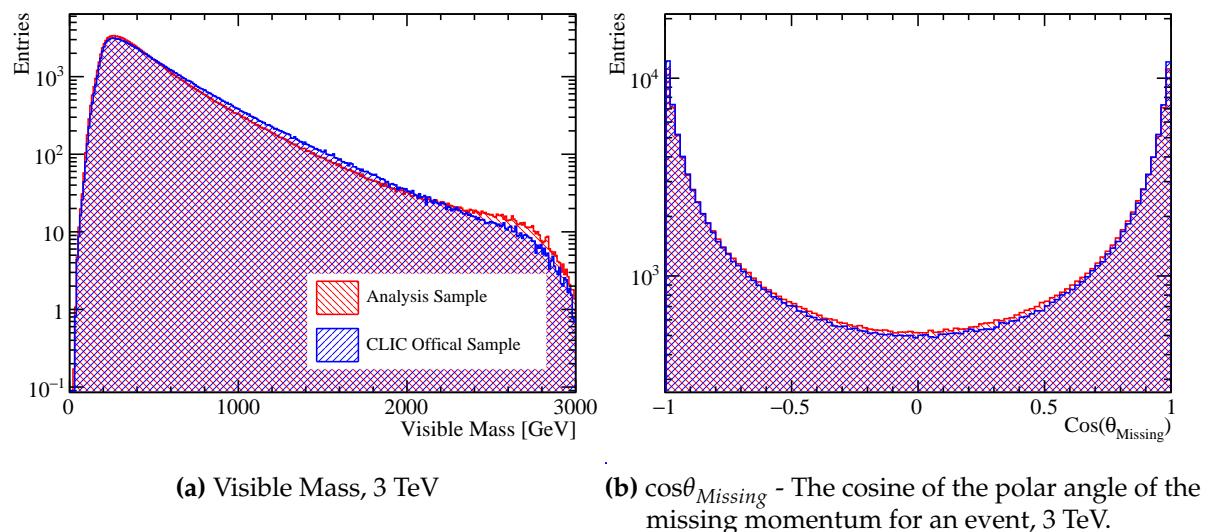
**Figure 2.9.:**  $\chi^2$  sensitivity distributions at 1.4 TeV arising from a fit to  $\cos\theta_{\text{jets}}^*$ . Results include the effect of backgrounds after the application of preselection and MVA.

Final State	Cross Section 1.4 TeV [fb]	Cross Section 3 TeV [fb]
$e^+ e^- \rightarrow \nu\nu qqqq$	24.7	71.5
$e^+ e^- \rightarrow l\nu qqqq$	110.4	106.6
$e^+ e^- \rightarrow llqqqq$	62.1	169.3
$e^+ e^- \rightarrow qqqq$	1245.1	546.5
$e^+ e^- \rightarrow \nu\nu qq$	787.7	1317.5
$e^+ e^- \rightarrow l\nu qq$	4309.7	5560.9
$e^+ e^- \rightarrow llqq$	2725.8	3319.6
$e^+ e^- \rightarrow qq$	4009.5	2948.9
$\gamma_{\text{EPA}} e^- \rightarrow qqqqe^-$	287.1	287.8
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	1160.7	1268.6
$e^+ \gamma_{\text{EPA}} \rightarrow qqqqe^+$	286.9	287.8
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	1156.3	1267.3
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	32.6	54.2
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	136.9	262.5
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	32.6	54.2
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	136.4	262.3
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow qqqq$	753.0	402.7
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	4034.8	2423.1
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	4018.7	2420.6
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow qqqq$	21406.2	13050.3

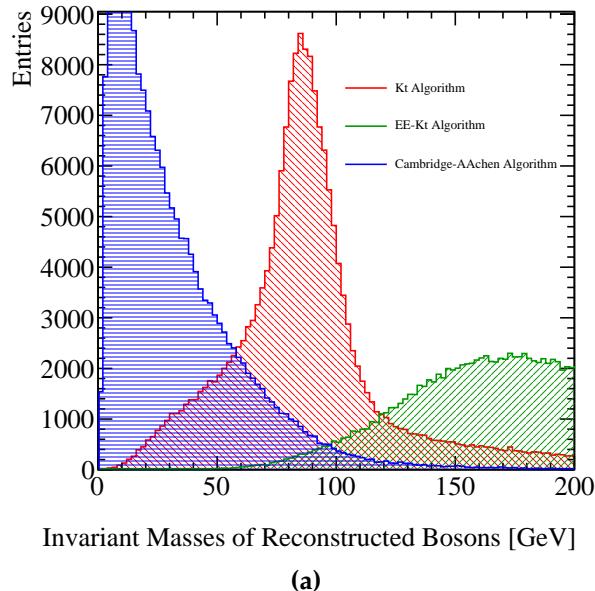
**Table 2.7.:** Cross sections of signal and background processes at 1.4 and 3 TeV. In the above table  $q \in u, \bar{u}, d, \bar{d}, s, \bar{s}, c, \bar{c}, b$  or  $\bar{b}$  while  $l \in e^\pm, \mu^\pm$  or  $\tau^\pm$  and  $\nu \in \nu_e, \nu_\mu$  and  $\nu_\tau$ . The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.

Final State	Cross Section [fb] ( $\alpha_4 = \alpha_5 = 0.000$ )	Cross Section [fb] ( $\alpha_4 = \alpha_5 = 0.005$ )	Percentage Change[%]	CLIC Cross Section [fb]
$e^+ e^- \rightarrow \nu\nu qqqq$	51.2	77.7	+51.8	71.5
$e^+ e^- \rightarrow l\nu qqqq$	111.9	115.9	+3.6	106.6
$e^+ e^- \rightarrow llqqqq$	169.7	161.7	-4.9	169.3

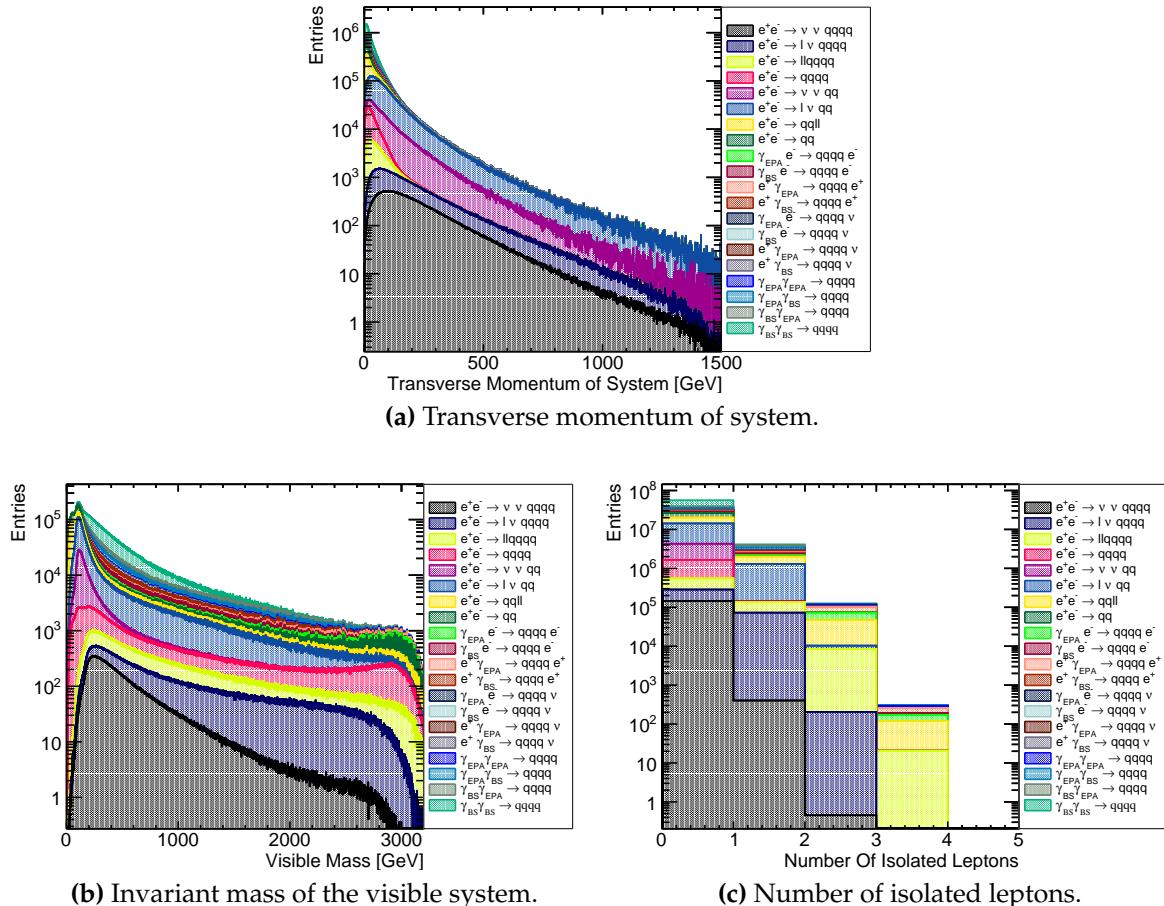
**Table 2.8.:** Cross section for selected processes for given value of  $\alpha_4$  and  $\alpha_5$  at 3 TeV.



**Figure 2.10.:** Comparison of various distributions between samples used in this analysis and the official CLIC samples for the  $\nu\nu qqqq$  final state at 3 TeV. Both samples have been normalised to the correct luminosity for CLIC running at 3 TeV.



**Figure 2.11.**: Reconstructed masses for different choices of jet algorithm for 3 TeV  $\nu\nu\text{qqqq}$  events. These masses arise by forcing the reconstructed events into 4 jets and then pairing up the jets into pairs such that the reconstructed invariant masses of the pairs are closest to each other. These samples should be dominated by vector boson scattering involving pairs of W bosons and so it is expected that a peak at the W boson true mass should be observed. As this does not occur for the Cambridge-Aachen algorithm or the ee\_kt algorithm they were deemed unsuitable for this analysis at both 3 TeV. In the case of the kt algorithm and the ee\_kt algorithm an R parameter of 0.7 was used.



**Figure 2.12.: Distribution of variables cut on in the preselection at 3 TeV.**

Final State	Raw Event Numbers	$p_T > 100 \text{ GeV}$	$p_T > 100 \text{ GeV} \&$ $M_{\text{Vis}} > 200 \text{ GeV}$	$p_T > 100 \text{ GeV} \&$ $M_{\text{Vis}} > 200 \text{ GeV} \&$ $N_{\text{Isolated Leptons}} = 0$
$e^+ e^- \rightarrow \nu\nu qqqq$	143,000	106,600	99,390	99,130
$e^+ e^- \rightarrow l\nu qqqq$	213,200	129,800	127,300	82,880
$e^+ e^- \rightarrow llqqqq$	338,600	32,750	31,010	23,550
$e^+ e^- \rightarrow qqqq$	1,093,000	40,180	37,360	37,300
$e^+ e^- \rightarrow \nu\nu qq$	2,634,000	1,333,000	380,100	379,500
$e^+ e^- \rightarrow l\nu qq$	11,120,000	4,240,000	2,479,000	1,836,000
$e^+ e^- \rightarrow llqq$	6,639,000	131,400	84,980	54,780
$e^+ e^- \rightarrow qq$	5,897,000	79,440	66,790	66,730
$\gamma_{\text{EPA}} e^- \rightarrow qqqqe^-$	575,600	57,920	54,640	34,480
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	2,004,000	99,930	90,750	83,440
$e^+ \gamma_{\text{EPA}} \rightarrow qqqqe^+$	575,600	57,990	54,290	34,190
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	2,002,000	100,300	90,830	83,960
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	108,400	63,780	60,660	46,380
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	414,700	233,800	215,600	213,600
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	108,400	64,230	61,130	46,720
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	414,400	236,400	219,000	217,000
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow qqqq$	805,400	54,010	48,720	37,730
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	3,828,000	150,800	131,600	114,500
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	3,825,000	150,600	133,600	116,900
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow qqqq$	18,010,000	123,500	115,400	105,000

**Table 2.9.:** Number of events passing the various cuts applied in the preselection at 3 TeV. Event numbers are normalised to the correct luminosity for CLIC at 3 TeV.  $p_T$  is the transverse momentum of the event,  $M_{\text{Vis}}$  is the visible mass and  $N_{\text{Isolated Leptons}}$  is the number of isolated leptons in the event. In the above table  $q \in u, \bar{u}, d, \bar{d}, s, \bar{s}, c, \bar{c}, b$  or  $\bar{b}$  while  $l \in e^\pm, \mu^\pm$  or  $\tau^\pm$  and  $\nu \in \nu_e, \nu_\mu$  and  $\nu_\tau$ . The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.

## 2.8. Sensitivity at 3 TeV

### 2.8.1. Pre Selection - 3 TeV

### 2.8.2. MVA - 3 TeV

Final State	Raw Event Numbers	Post MVA Selection Numbers
$e^+e^- \rightarrow \nu\nuqqqq$	143,000	64,750
$e^+e^- \rightarrow l\nuqqqq$	213,200	23,310
$e^+e^- \rightarrow llqqqq$	338,600	2,409
$e^+e^- \rightarrow qqqq$	1,093,000	3,069
$e^+e^- \rightarrow \nu\nu qq$	2,634,000	19,040
$e^+e^- \rightarrow l\nu qq$	11,120,000	27,910
$e^+e^- \rightarrow llqq$	6,639,000	786
$e^+e^- \rightarrow qq$	5,897,000	1,335
$\gamma_{\text{EPA}} e^- \rightarrow qqqqe^-$	575,600	2,860
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	2,004,000	8,352
$e^+ \gamma_{\text{EPA}} \rightarrow qqqqe^+$	575,600	3,063
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	2,002,000	8,090
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	108,400	17,950
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	414,700	108,000
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	108,400	17,980
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	414,400	109,700
$\gamma_{\text{EPA}} \gamma_{\text{EPA}} \rightarrow qqqq$	805,400	3,058
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	3,828,000	9,812
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	3,825,000	8,880
$\gamma_{\text{BS}} \gamma_{\text{BS}} \rightarrow qqqq$	18,010,000	2,213

**Table 2.10.:** Number of events passing the MVA selection at 3 TeV. Event numbers are normalised to the correct luminosity for CLIC at 3 TeV. The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung.

In attempt to remove the effects of the overlaid  $\gamma\gamma \rightarrow \text{Hadron}$  background event a number of cuts [9] are applied to the transverse momenta and the timing information of the PFOs produced by PandoraPFA to reduce the PFOs into a subset that are believed to originate from the desired interaction. Different values of these cuts give rise to the tight, default and loose selected PFOs that were considered in this optimisation.

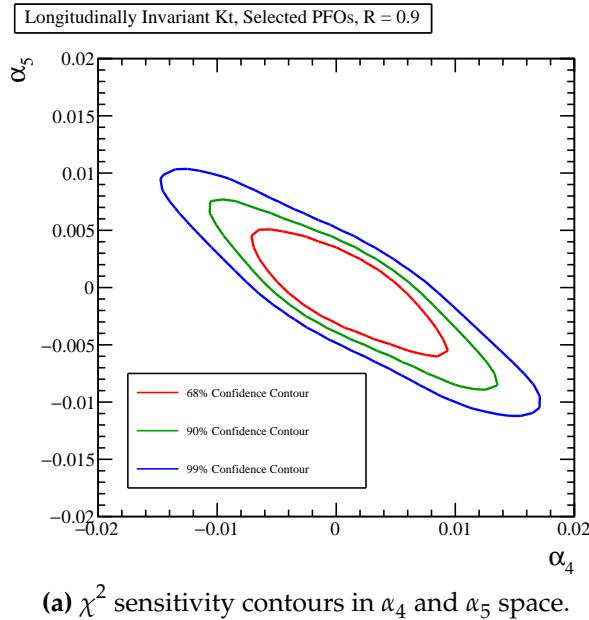
Final State	$\epsilon_{\text{presel}}$	$\epsilon_{\text{BDT}}$	$N_{\text{BDT}}$
$e^+e^- \rightarrow \nu\nu qqqq$	69.4%	45.3%	64,750
$e^+e^- \rightarrow l\nu qqqq$	38.9%	10.9%	23,310
$e^+e^- \rightarrow \nu\nu qq$	14.4%	0.7%	19,040
$e^+e^- \rightarrow l\nu qq$	16.5%	0.3%	27,910
$\gamma_{\text{BS}} e^- \rightarrow qqqqe^-$	4.1%	0.4%	8,352
$e^+ \gamma_{\text{BS}} \rightarrow qqqqe^+$	4.2%	0.4%	8,090
$\gamma_{\text{EPA}} e^- \rightarrow qqqq\nu$	42.8%	16.6%	17,950
$\gamma_{\text{BS}} e^- \rightarrow qqqq\nu$	51.6%	26.0%	108,000
$e^+ \gamma_{\text{EPA}} \rightarrow qqqq\nu$	43.1%	16.6%	17,980
$e^+ \gamma_{\text{BS}} \rightarrow qqqq\nu$	52.3%	26.5%	109,700
$\gamma_{\text{EPA}} \gamma_{\text{BS}} \rightarrow qqqq$	3.0%	0.3%	9,812
$\gamma_{\text{BS}} \gamma_{\text{EPA}} \rightarrow qqqq$	3.1%	0.2%	8,880

**Table 2.11.:** Selection summary at 3 TeV. The subscript EPA or BS for the incoming photons indicate whether the photon is generated from the equivalent photon approximation or beamstrahlung. Channels omitted from this table have less than 6,000 events in the post MVA selection.

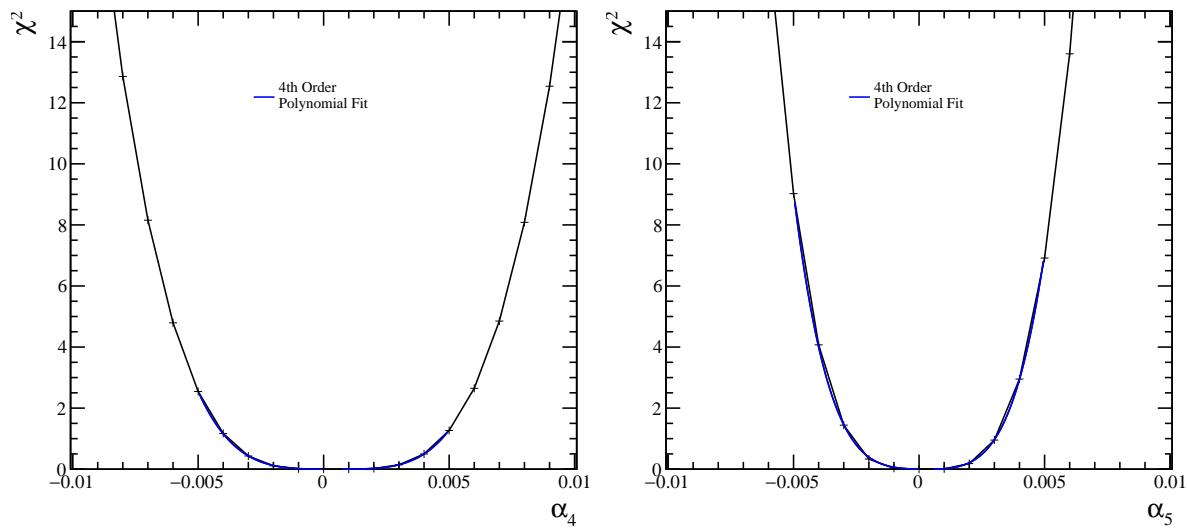
### 2.9.1. 1.4 TeV Optimal Jet Reconstruction

At 1.4 TeV the optimal sensitivity is achieved for either loose selected PFOs with an R parameter of 0.7 or default selected PFOs with an R parameter of 0.9 as can be seen from tables 2.12 and 2.13. As a tie breaker between these options the separation power, the fraction of events misidentified as either arising from a WW pair or a ZZ pair, was considered. Again performance was similar, but there was a slight preference towards the use of selected PFOs and an R parameter of 0.9. While not used in the primary analysis the separation of samples into WW and ZZ events is important for an extension analysis found in section BLAH.

The optimal contours can be found in figure 2.13a and the optimal 1D plot used to produce the errors references in the tables 2.12 and 2.13 can be found in figures 2.13b and 2.13c respectively. All other contours and plots for this optimisation can be found in the appendices. There are minimal performance differences between the various jet algorithm configurations at 1.4 TeV.



(a)  $\chi^2$  sensitivity contours in  $\alpha_4$  and  $\alpha_5$  space.



(b)  $\chi^2$  as a function of  $\alpha_4$  assuming  $\alpha_5 = 0$ .

(c)  $\chi^2$  as a function of  $\alpha_5$  assuming  $\alpha_4 = 0$ .

**Figure 2.13.:**  $\chi^2$  sensitivity distributions for the  $qqqq\nu\nu$  final state arising from a fit to  $\cos\theta_{\text{jets}}^*$  at 1.4 TeV for the optimal jet reconstruction parameters.

PFO Selection	Tight Selected PFOs	Selected PFOs	Loose Selected PFOs
R Parameter			
0.7	$-0.00391 + 0.00497$	$-0.00385 + 0.00500$	$-0.00368 + 0.00465$
0.9	$-0.00405 + 0.00508$	$-0.00375 + 0.00464$	$-0.00379 + 0.00475$
1.1	$-0.00406 + 0.00509$	$-0.00392 + 0.00502$	$-0.00403 + 0.00496$

**Table 2.12.:** Precision on measurement of  $\alpha_4$  at 1.4 TeV for different jet reconstruction parameters considering pure signal and applying a  $\chi^2$  fit to  $\cos\theta_{jets}^*$ .

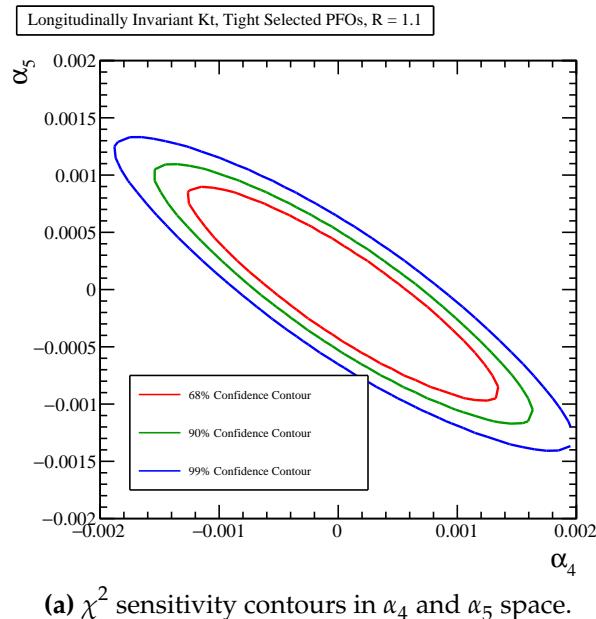
PFO Selection	Tight Selected PFOs	Selected PFOs	Loose Selected PFOs
R Parameter			
0.7	$-0.00267 + 0.00313$	$-0.00266 + 0.00318$	$-0.00255 + 0.00302$
0.9	$-0.00280 + 0.00320$	$-0.00259 + 0.00302$	$-0.00258 + 0.00303$
1.1	$-0.00281 + 0.00321$	$-0.00272 + 0.00319$	$-0.00282 + 0.00315$

**Table 2.13.:** Precision on measurement of  $\alpha_5$  at 1.4 TeV for different jet reconstruction parameters considering pure signal and applying a  $\chi^2$  fit to  $\cos\theta_{jets}^*$ .

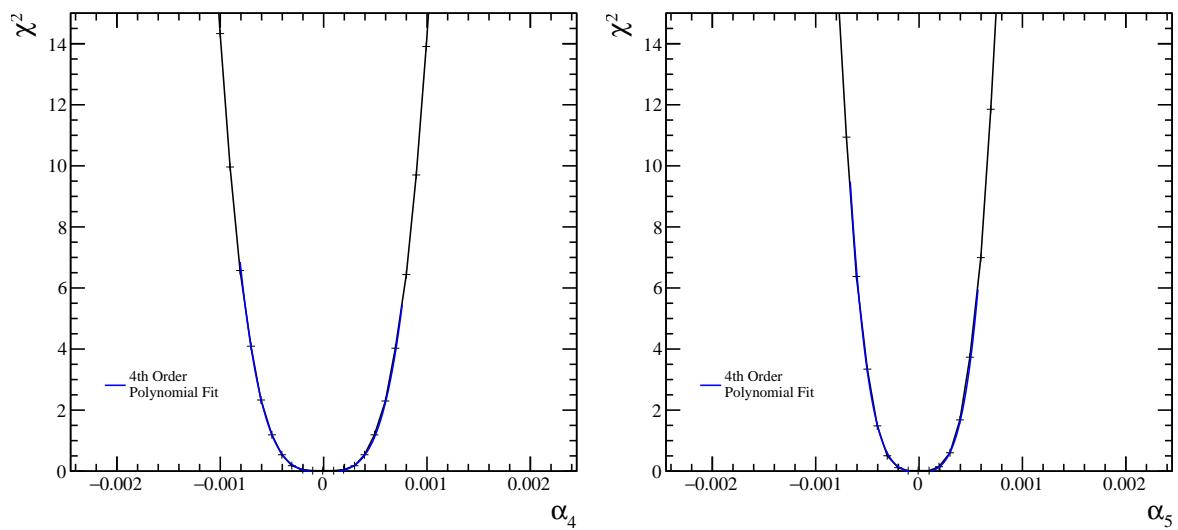
### 2.9.2. 3 TeV Optimal Jet Reconstruction

At 3 TeV the optimal sensitivity for the reconstructions considered is achieved for tight selected PFOs with an R parameter of 1.1 as can be seen from tables 2.14 and 2.15. The optimal contours can be found in figure 2.14a and the optimal 1D plot used to produce the errors references in the tables 2.14 and 2.15 can be found in figures 2.14b and 2.14c respectively. All other contours and plots for this optimisation can be found in the appendices.

The gains in optimising the jet algorithm at 3 TeV are larger than those found at 1.4 TeV. The preference for the tight selected PFOs is to be expected as this configuration minimises the effect of beam induced backgrounds, which are more prominent at higher energies.



(a)  $\chi^2$  sensitivity contours in  $\alpha_4$  and  $\alpha_5$  space.



(b)  $\chi^2$  as a function of  $\alpha_4$  assuming  $\alpha_5 = 0$ .

(c)  $\chi^2$  as a function of  $\alpha_5$  assuming  $\alpha_4 = 0$ .

**Figure 2.14.:**  $\chi^2$  sensitivity distributions for the  $qqqq\nu\nu$  final state arising from a fit to  $\cos\theta_{\text{jets}}^*$  at 3 TeV for the optimal jet reconstruction parameters.

PFO Selection	Tight Selected PFOs	Selected PFOs	Loose Selected PFOs
R Parameter			
0.7	$-0.000530 + 0.000525$	$-0.000502 + 0.000507$	$-0.000547 + 0.000555$
0.9	$-0.000566 + 0.000555$	$-0.000539 + 0.000520$	$-0.000568 + 0.000553$
1.1	$-0.000472 + 0.000472$	$-0.000508 + 0.000492$	$-0.000504 + 0.000490$

**Table 2.14.:** Precision on measurement of  $\alpha_4$  at 3 TeV for different jet reconstruction parameters considering pure signal and applying a  $\chi^2$  fit to  $\cos\theta_{Jets}^*$ .

PFO Selection	Tight Selected PFOs	Selected PFOs	Loose Selected PFOs
R Parameter			
0.7	$-0.000393 + 0.000370$	$-0.000356 + 0.000348$	$-0.000357 + 0.000348$
0.9	$-0.000394 + 0.000365$	$-0.000392 + 0.000361$	$-0.000396 + 0.000368$
1.1	$-0.000351 + 0.000337$	$-0.000374 + 0.000354$	$-0.000353 + 0.000336$

**Table 2.15.:** Precision on measurement of  $\alpha_5$  at 3 TeV for different jet reconstruction parameters considering pure signal and applying a  $\chi^2$  fit to  $\cos\theta_{Jets}^*$ .

# Appendix A.

## Pointless extras

*“Le savant n’étudie pas la nature parce que cela est utile;  
il l’étudie parce qu’il y prend plaisir,  
et il y prend plaisir parce qu’elle est belle.”*

— Henri Poincaré, 1854–1912

Appendices (or should that be “appendices”?) make you look really clever, ‘cos it’s like you had more clever stuff to say than could be fitted into the main bit of your thesis. Yeah. So everyone should have at least three of them...

### A.1. Anomalous Gauge Coupling Quartic Vertices Of Relevance in Vector Boson Scattering

The anomalous gauge couplings involving  $\alpha_4$  and  $\alpha_5$  arise in EFT through the addition of the following terms to the Lagrangian.

$$\text{Tr}(V^\mu V_\nu) \text{Tr}(V^\nu V_\mu) \text{ and } [\text{Tr}(V^\mu V_\mu)]^2 \quad (\text{A.1})$$

Where  $V_\mu$  is defined in the following way.

$$V_\mu = \Sigma (D_\mu \Sigma)^\dagger \quad (\text{A.2})$$

and  $\Sigma$ , the Higgs field matrix, is defined as.

$$\Sigma = \exp\left(-\frac{i}{v}\mathbf{w}\right) \quad (\text{A.3})$$

Where  $\mathbf{w} = w^a \sigma^a$ .  $w^a$  are the ... and  $\sigma^a$  are the Pauli spin matrices. The covariant derivative of the Higgs field matrix is

$$D_\mu \Sigma = \left( \partial_\mu + \frac{ig}{2} W_\mu - \frac{ig'}{2} B_\mu \sigma^3 \right) \Sigma \quad (\text{A.4})$$

For clarity consider the unitarity gauge where  $\mathbf{w} = 0$ , which implies  $\Sigma = 1$ . In this gauge  $V_\mu$  takes the following form.

$$\begin{aligned} V_\mu &= \frac{i}{2} (g W_\mu^i \sigma^i - g' B_\mu \sigma^3) = \frac{i}{2} \begin{pmatrix} g W_\mu^3 - g' B_\mu & g(W_\mu^1 - iW_\mu^2) \\ g(W_\mu^1 + iW_\mu^2) & -g W_\mu^3 + g' B_\mu \end{pmatrix} \\ &= \frac{i}{2} \begin{pmatrix} \sqrt{g^2 + g'^2} Z_\mu & g\sqrt{2} W_\mu^+ \\ g\sqrt{2} W_\mu^- & \sqrt{g^2 + g'^2} Z_\mu \end{pmatrix} \end{aligned}$$

Where the relationship between the mass and gauge symmetry basis are as follows.

$$W_\mu^+ = \frac{1}{\sqrt{2}} (W_\mu^1 - iW_\mu^2) \quad (\text{A.5})$$

$$W_\mu^- = \frac{1}{\sqrt{2}} (W_\mu^1 + iW_\mu^2) \quad (\text{A.6})$$

$$Z_\mu = c_w W_\mu^3 - s_w B_\mu \quad (\text{A.7})$$

$$A_\mu = s_w W_\mu^3 + c_w B_\mu \quad (\text{A.8})$$

With  $c_w = \frac{g}{\sqrt{g^2 + g'^2}}$  and  $s_w = \frac{g'}{\sqrt{g^2 + g'^2}}$ . Consider the expansion of the terms to be included in the Lagrangian.

$$V^\mu V_\nu = \frac{-1}{4} \begin{pmatrix} \sqrt{g^2 + g'^2} Z^\mu & g\sqrt{2}W^{+\mu} \\ g\sqrt{2}W^{-\mu} & \sqrt{g^2 + g'^2} Z^\mu \end{pmatrix} \begin{pmatrix} \sqrt{g^2 + g'^2} Z_\nu & g\sqrt{2}W_\nu^+ \\ g\sqrt{2}W_\nu^- & \sqrt{g^2 + g'^2} Z_\nu \end{pmatrix} \quad (\text{A.9})$$

$$\text{Tr}[V^\mu V_\nu] = \frac{-1}{2}((g^2 + g'^2)Z^\mu Z_\nu + g^2 W^{+\mu} W_\nu^- + g^2 W^{-\mu} W_\nu^+) \quad (\text{A.10})$$

$$\text{Tr}[V^\mu V_\nu] \text{Tr}[V_\mu V^\nu] = \frac{(g^2 + g'^2)^2}{4}(Z^\mu Z_\mu)^2 + g^2(g^2 + g'^2)(Z^\mu Z^\nu W_\mu^- W_\nu^+) \quad (\text{A.11})$$

$$+ \frac{g^4}{2}(W^{-\mu} W_\mu^+)^2 + \frac{g^4}{2}(W^{-\mu} W^{+\nu} W_\mu^- W_\nu^+) \quad (\text{A.12})$$

$$\text{Tr}[V^\mu V_\mu]^2 = \frac{(g^2 + g'^2)^2}{4}(Z^\mu Z_\mu)^2 + g^2(g^2 + g'^2)(Z^\mu Z^\nu W_\mu^- W_\nu^+) \quad (\text{A.13})$$

$$+ g^4(W^{-\mu} W_\mu^+)^2 \quad (\text{A.14})$$

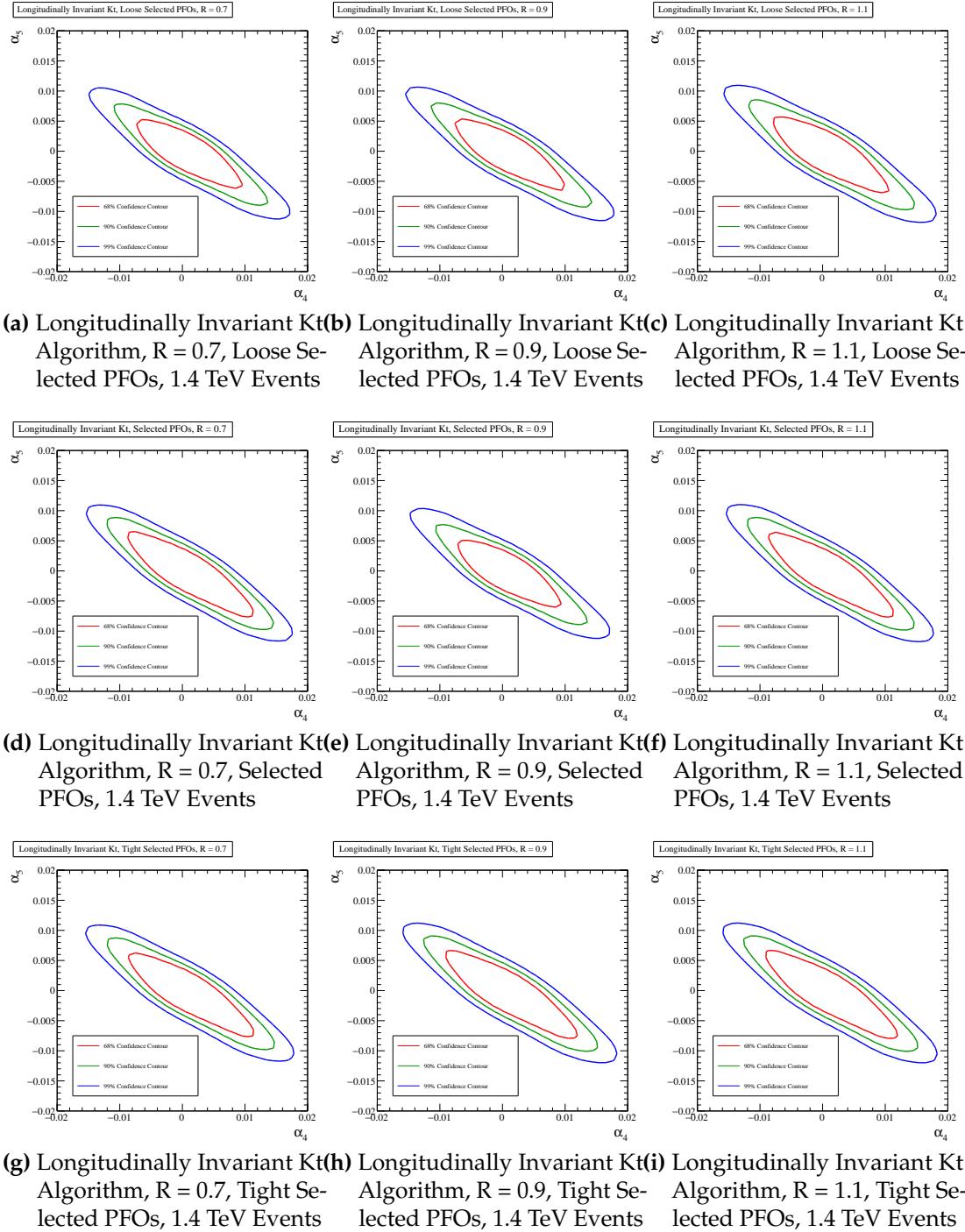
These two terms change the cross section for the vector boson scattering processes at CLIC that involve  $ZZ \rightarrow ZZ$ ,  $W^+W^- \rightarrow ZZ$ ,  $ZZ \rightarrow W^+W^-$  and  $W^+W^- \rightarrow W^+W^-$ .

$$\text{Z} \quad \text{Z} \quad \text{Z} \quad \text{Z} \subset (\alpha_4 + \alpha_5) \frac{(g^2 + g'^2)^2}{4} \quad (\text{A.15})$$

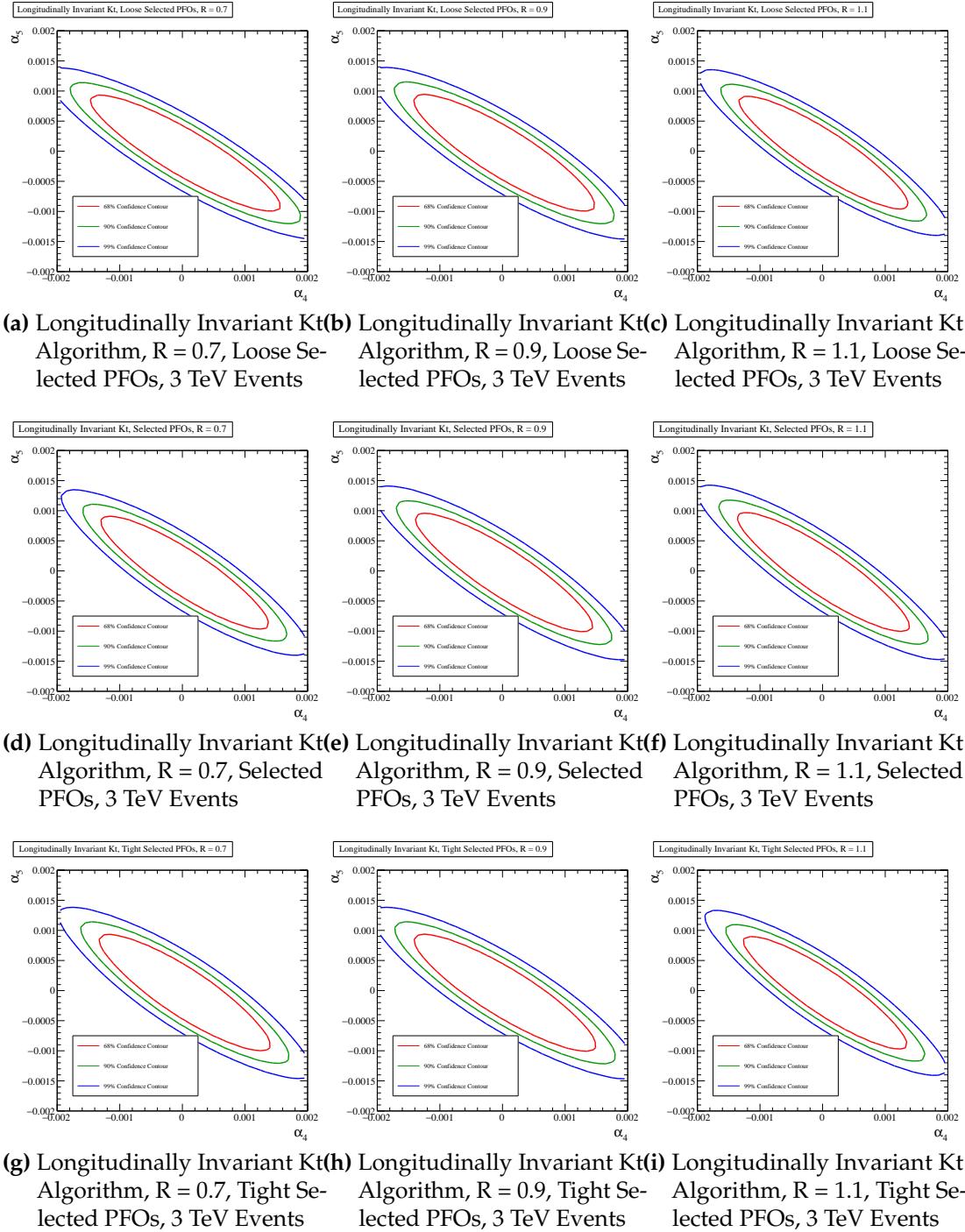
$W$                    $W$   
 $Z$                    $Z$      $\subset (\alpha_4 + \alpha_5)g^2(g^2 + g'^2)$                   (A.16)

$W$                    $W$   
 $W$                    $W$      $\subset (\alpha_4 + 2\alpha_5)\frac{g^4}{2}$  and  $\frac{g^4}{2}\alpha_4$                   (A.17)

## A.2. $\chi^2$ Contour Plots for Jet Algorithm Optimisation



**Figure A.1:**  $\chi^2$  Sensitivity contours for the  $qqqq\nu\nu$  final state arising from a fit to  $\cos\theta_{\text{jets}}^*$  at 1.4 TeV for different values of jet reconstruction parameters.



**Figure A.2:**  $\chi^2$  Sensitivity contours for the  $qqqq\nu\nu$  final state arising from a fit to  $\cos\theta_{\text{jets}}^*$  at 1.4 TeV for different values of jet reconstruction parameters.

# Colophon

This thesis was made in L<sup>A</sup>T<sub>E</sub>X 2 <sub>$\epsilon$</sub>  using the “hepthesis” class [3].



# Bibliography

- [1] Toshinori Abe et al. The International Large Detector: Letter of Intent. 2010, 1006.3396.
- [2] S. Agostinelli et al. GEANT4: A Simulation toolkit. *Nucl. Instrum. Meth.*, A506:250–303, 2003.
- [3] Andy Buckley. The heptesis L<sup>A</sup>T<sub>E</sub>X class.
- [4] Matteo Cacciari, Gavin P. Salam, and Gregory Soyez. FastJet User Manual. *Eur. Phys. J.*, C72:1896, 2012, 1111.6097.
- [5] F. Gaede. Marlin and LCCD: Software tools for the ILC. *Nucl. Instrum. Meth.*, A559:177–180, 2006.
- [6] Andreas Hoecker, Peter Speckmayer, Joerg Stelzer, Jan Therhaag, Eckhard von Toerne, and Helge Voss. TMVA: Toolkit for Multivariate Data Analysis. *PoS*, ACAT:040, 2007, physics/0703039.
- [7] Wolfgang Kilian, Thorsten Ohl, and Jurgen Reuter. WHIZARD: Simulating Multi-Particle Processes at LHC and ILC. *Eur. Phys. J.*, C71:1742, 2011, 0708.4233.
- [8] Lucie Linssen, Akiya Miyamoto, Marcel Stanitzki, and Harry Weerts. Physics and Detectors at CLIC: CLIC Conceptual Design Report. 2012, 1202.5940.
- [9] J. S. Marshall, A. M<sup>Ã</sup>ajnnich, and M. A. Thomson. Performance of Particle Flow Calorimetry at CLIC. *Nucl. Instrum. Meth.*, A700:153–162, 2013, 1209.4039.
- [10] P. Mora de Freitas and H. Videau. Detector simulation with MOKKA / GEANT4: Present and future. In *Linear colliders. Proceedings, International Workshop on physics and experiments with future electron-positron linear colliders, LCWS 2002, Seogwipo, Jeju Island, Korea, August 26-30, 2002*, pages 623–627, 2002.
- [11] Mauro Moretti, Thorsten Ohl, and Jurgen Reuter. O'Mega: An Optimizing matrix

- element generator. 2001, hep-ph/0102195.
- [12] B. Pilicer, Helmut Burkhardt, Laurent Gatignon, Andrea Latina, Ercan Pilicer, Daniel Schulte, İlhan Tapan, and Rogelio Tomás. Study of Muon Backgrounds in the CLIC Beam Delivery System. In *Proceedings, 6th International Particle Accelerator Conference (IPAC 2015): Richmond, Virginia, USA, May 3-8, 2015*, page TUPTY032, 2015.
  - [13] András Sailer. *Radiation and Background Levels in a CLIC Detector due to Beam-Beam Effects*. PhD thesis, CERN, 2012-05-09.
  - [14] M. A. Thomson. Particle Flow Calorimetry and the PandoraPFA Algorithm. *Nucl. Instrum. Meth.*, A611:25–40, 2009, 0907.3577.

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