

# Classical Mechanics

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## 1 Preliminaries

**Definition 1.1** (Mechanical System). A **Mechanical System** is a tuple  $(M, L)$  where

1.  $M$  is the configuration space - a manifold in a vector space of generalized coordinates
2.  $L$  is the Lagrangian - a function  $L : TM \times \mathbb{R}_t \rightarrow \mathbb{R}$  where  $TM$  is the tangent bundle to  $M$  and  $\mathbb{R}_t \subseteq \mathbb{R}$

In many applications, the Lagrangian is equal to the kinetic energy minus the potential energy of the system.

**Definition 1.2** (Principle of Stationary Action). Given time instances  $t_1$  and  $t_2$  ( $t_1 < t_2$ ), the time evolution of the mechanical system  $(M, L)$  is given by a path  $q_0 : [t_1, t_2] \rightarrow M$  that minimizes

$$A(q) = \int_{t_1}^{t_2} L(q, \dot{q}, t) dt \quad (1)$$

The path  $q_0$  is a **stationary point**.

## 2 Deriving Lagrange's Equations

Consider adding a variation  $\eta$  to a stationary point  $q_0$ . Then

$$\begin{aligned} 0 &= \frac{d}{d\epsilon} A(q_0 + \epsilon\eta)|_{\epsilon=0} = \frac{d}{d\epsilon} \int_{t_1}^{t_2} L(q_0 + \epsilon\eta, \dot{q}_0 + \epsilon\dot{\eta}, t) dt|_{\epsilon=0} \\ &= \int_{t_1}^{t_2} [L_q(q, \dot{q}, t) \eta + L_{\dot{q}}(q, \dot{q}, t) \dot{\eta}] dt \\ &= \int_{t_1}^{t_2} L_q(q, \dot{q}, t) \eta dt + L_{\dot{q}}(q, \dot{q}, t) \eta|_{t_1}^{t_2} - \int_{t_1}^{t_2} \frac{\partial L_{\dot{q}}}{\partial t}(q, \dot{q}, t) \eta dt \\ &= \int_{t_1}^{t_2} \left[ L_q(q, \dot{q}, t) - \frac{\partial L_{\dot{q}}}{\partial t}(q, \dot{q}, t) \right] \eta dt \end{aligned} \quad (2)$$

By the fundamental lemma of the calculus of variations, it follows that

$$L_q(q, \dot{q}, t) = \frac{\partial L_{\dot{q}}}{\partial t}(q, \dot{q}, t) \quad (3)$$

### 3 Deriving Hamilton's Equations

Consider a physical system comprised of trajectories  $q$  in configuration space, velocities  $\dot{q}$  and Lagrangian  $L(q, \dot{q})$ . The conjugate momenta to the trajectories is

$$p = \frac{\partial L}{\partial \dot{q}}(q, \dot{q}) \quad (4)$$

Under certain regularity conditions (the conditions of the Inverse Function Theorem), we may invert this map to recover  $\dot{q}$  from  $p$  and  $q$ :

$$\dot{q} = f(q, p) \quad (5)$$

The Hamiltonian for this system is

$$H(q, p) = \sum_i p_i f_i(q, p) - L(q, f(q, p)) \quad (6)$$

By adding a variation to individual components of the phase space, we find

$$\begin{aligned} \frac{\partial H}{\partial p_j} \eta &= \frac{\partial}{\partial \epsilon} H(q, p_{-j}, p_j + \epsilon \eta) |_{\epsilon=0} \\ &= \sum_i p_i \frac{\partial f_i}{\partial p_j} \eta + f_j \eta - \sum_i \frac{\partial L}{\partial \dot{q}_i} \frac{\partial f_i}{\partial p_j} \eta \\ &= f_j \eta \end{aligned} \quad (7)$$

$$\begin{aligned} \frac{\partial H}{\partial q_j} \eta &= \frac{\partial}{\partial \epsilon} H(q_{-j}, q_j + \epsilon \eta, p) |_{\epsilon=0} \\ &= \sum_i p_i \frac{\partial f_i}{\partial q_j} \eta - \sum_i \frac{\partial L}{\partial \dot{q}_i} \frac{\partial f_i}{\partial q_j} \eta - \frac{\partial L}{\partial q_j} \eta \\ &= -\frac{\partial L}{\partial q_j} \eta \end{aligned} \quad (8)$$

It follows that

$$\frac{\partial H}{\partial p_j} = f_j = \dot{q}_j \quad (9)$$

In addition, if the trajectories through phase space minimize the action of the system, then

$$\frac{\partial L}{\partial q_i}(q, \dot{q}) = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i}(q, \dot{q}) \quad (10)$$

so that

$$\frac{\partial H}{\partial q_i} = -\dot{p}_i \quad (11)$$

We have derived Hamilton's equations:

$$\frac{\partial H}{\partial p_j} = \dot{q}_j \quad \text{and} \quad \frac{\partial H}{\partial q_i} = -\dot{p}_i \quad (12)$$

## 4 Poisson Brackets

Two equivalent definitions of the Poisson Bracket:

$$\{F, G\} = \sum_i \frac{\partial F}{\partial q_i} \frac{\partial G}{\partial p_i} - \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} \quad (13)$$

Axiomatic definition of Poisson Bracket:

1. Anticommutativity:  $\forall f, g : \{f, g\} = -\{g, f\}$
2. Linearity:  $\forall a, b \in \mathbb{R}, \forall f, g, h : \{af + bg, h\} = a\{f, h\} + b\{g, h\}$
3. Leibniz's Rule:  $\forall f, g, h : \{fg, h\} = f\{g, h\} + g\{f, h\}$
4. Canonical Coordinates:  $\forall i, j : \{q_i, q_j\} = 0, \{p_i, p_j\} = 0, \{q_i, p_j\} = \delta_{ij}$

## 5 Curl and Divergence

Suppose  $\vec{v} = (f(x, y, z), g(x, y, z), h(x, y, z))$  is a vector field such that  $\vec{\nabla} \cdot \vec{v} = f_x + g_y + h_z = 0$ . Then there is another vector field  $\vec{w}$ , such that  $\vec{v} = \vec{\nabla} \times \vec{w}$ . Let

$$\vec{w}_1 = \left( 0, \int_0^x h(t, y, z) dt, -\int_0^x g(t, y, z) dt \right)$$

Now

$$\begin{aligned} \vec{\nabla} \times \vec{w}_1 &= \left( \frac{-\partial \int_0^x g(t, y, z) dt}{\partial y} - \frac{\partial \int_0^x h(t, y, z) dt}{\partial z}, \frac{\partial \int_0^x g(t, y, z) dt}{\partial x}, \frac{\partial \int_0^x h(t, y, z) dt}{\partial x} \right) \\ &= \left( -\int_0^x \left[ \frac{\partial g}{\partial y}(t, y, z) + \frac{\partial h}{\partial x}(t, y, z) \right] dt, g(x, y, z), h(x, y, z) \right) \\ &= (A(x, y, z), g(x, y, z), h(x, y, z)) \end{aligned} \quad (14)$$

where

$$A(x, y, z) = - \int_0^x \left[ \frac{\partial g}{\partial y}(t, y, z) + \frac{\partial h}{\partial x}(t, y, z) \right] dt \quad (15)$$

Let

$$\vec{v}_1 = \vec{v} - \vec{\nabla} \times \vec{w}_1 = (f - A, 0, 0) \quad (16)$$

Then

$$\frac{\partial}{\partial x}(f - A) = \vec{\nabla} \cdot \vec{v}_1 = \vec{\nabla} \cdot \vec{v} - \vec{\nabla} \cdot \vec{\nabla} \times \vec{w}_1 = 0 \quad (17)$$

so that  $f - A$  does NOT depend on  $x$ . Let  $k(y, z) = f - A$  and

$$\vec{w}_2 = \left( 0, 0, \int_0^y k(t, z) dt \right) \quad (18)$$

Then

$$\vec{\nabla} \times \vec{w}_2 = (k(y, z), 0, 0) = (f - A, 0, 0) \quad (19)$$

Finally,

$$\begin{aligned} \vec{\nabla} \times (\vec{w}_1 + \vec{w}_2) &= \vec{\nabla} \times \vec{w}_1 + \vec{\nabla} \times \vec{w}_2 \\ &= (A, g, h) + (f - A, 0, 0) \\ &= (f, g, h) \end{aligned} \quad (20)$$

so that  $\vec{v} = \vec{\nabla} \times (\vec{w}_1 + \vec{w}_2)$