

TAU FUNCTION AND MODULI OF MEROMORPHIC FORMS ON ALGEBRAIC CURVES

D. KOROTKIN AND P. ZOGRAF

ABSTRACT. We study the moduli space of meromorphic 1-forms on complex algebraic curves having at most simple poles with fixed nonzero residues. We interpret the Bergman tau function on this moduli space as a section of a line bundle and study its asymptotic behavior near the boundary and the locus of forms with non-simple zeros. As an application, we decompose the projection of this locus to the moduli space of curves into a linear combination of standard generators of the rational Picard group with explicit coefficients that depend on the residues.

The present note is a continuation of a series of papers by the authors devoted to applications of tau functions to geometry of moduli spaces and can be considered as an extension of [13], where the case of holomorphic 1-forms (Abelian differentials of the first kind) was treated. The case of meromorphic 1-forms with simple poles (Abelian differentials of the third kind) basically follows the same lines, but involves some new technicalities like the dependence on relations between residues at poles and the appearance of tautological ψ -classes similar to [3].

We begin with introducing the necessary notation. Denote by $\mathcal{H}_{g,n}$ the space of equivalence classes of pairs (C, v) , where C is a smooth complex curve of genus g with n labeled marked points p_1, \dots, p_n , and v is a meromorphic 1-form on C with at most simple poles at p_i and no other poles; by Riemann-Roch, $\dim \mathcal{H}_{g,n} = 4g - 4 + 2n$. The forgetful map $\mathcal{H}_{g,n} \rightarrow \mathcal{M}_{g,n}$ is a vector bundle (in the orbifold sense) with fiber $H^0(C, K_C(p_1 + \dots + p_n))$ over (C, v) of dimension $g - 1 + n$, where K_C is the canonical class of C . The bundle $\mathcal{H}_{g,n}$ naturally extends to the Deligne-Mumford compactification $\overline{\mathcal{M}}_{g,n}$ of $\mathcal{M}_{g,n}$. More precisely, $\overline{\mathcal{H}}_{g,n} = \pi_*(\omega_{g,n}(p_1 + \dots + p_n))$, where $\omega_{g,n}$ is the relative dualizing sheaf of the universal curve $\pi : \overline{\mathcal{C}}_{g,n} \rightarrow \overline{\mathcal{M}}_{g,n}$, and p_1, \dots, p_n are the structure sections of π .

A classical result that can be traced back to Riemann claims that for any $(C; p_1, \dots, p_n) \in \mathcal{M}_{g,n}$ and any $(r_1, \dots, r_n) \in \mathbb{C}^n$ there exists a meromorphic 1-form v on C having at most simple poles at the points $p_1, \dots, p_n \in C$ with residues $r_1, \dots, r_n \in \mathbb{C}$ and smooth elsewhere if and only if $r_1 + \dots + r_n = 0$. Such a 1-form is defined uniquely up to an arbitrary holomorphic 1-form on C . This result naturally extends to the case of stable curves (where the 1-forms are allowed to have simple poles at the nodes, see below) and has two important consequences:

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(1) The *residue map* $\text{res} : \overline{\mathcal{H}}_{g,n} \longrightarrow \mathbb{C}^n$, $(C, v) \mapsto (r_1, \dots, r_n)$, where $r_i = \text{res}_{|p_i} v$, $i = 1, \dots, n$, fits into the exact sequence

$$\overline{\mathcal{H}}_{g,n} \xrightarrow{\text{res}} \mathbb{C}^n \xrightarrow{\sum r_i} \mathbb{C}.$$

(2) Any fiber $\overline{\mathcal{H}}_{g,n}[\mathbf{r}] = \text{res}^{-1}(r_1, \dots, r_n)$ of the residue map, where $\mathbf{r} = (r_1, \dots, r_n)$, is an affine bundle on $\overline{\mathcal{M}}_{g,n}$ of dimension g . For smooth curves, the fibers of the projection $\overline{\mathcal{H}}_{g,n}[\mathbf{r}] \longrightarrow \overline{\mathcal{M}}_{g,n}$ are affine spaces over the vector spaces of holomorphic 1-forms on C , and a minor adjustment is needed for nodal curves.

By a result of Grothendieck [7] (cf. also [5], p. 69), the projection

$$\text{pr} : \overline{\mathcal{H}}_{g,n}[\mathbf{r}] \longrightarrow \overline{\mathcal{M}}_{g,n}$$

induces an isomorphism between the rational Picard groups $\text{Pic}(\overline{\mathcal{H}}_{g,n}[\mathbf{r}]) \otimes \mathbb{Q}$ and $\text{Pic}(\overline{\mathcal{M}}_{g,n}) \otimes \mathbb{Q}$.

Let us recall the standard set of generators of $\text{Pic}(\overline{\mathcal{M}}_{g,n}) \otimes \mathbb{Q}$ (cf. [1]). It consists of:

- the Hodge class $\lambda = c_1(\pi_* \omega_{g,n})$,
- the classes $\psi_i = c_1(\mathcal{L}_i)$ of the tautological line bundles $\mathcal{L}_i \rightarrow \overline{\mathcal{M}}_{g,n}$, $i = 1, \dots, n$,
- the class δ_{irr} of the boundary divisor parameterizing irreducible nodal curves, and
- the classes $\delta_{j,I}$ of the boundary divisors parameterizing reducible curves with components of genera j and $g-j$ containing k and $n-k$ marked points p_i respectively (here $j = 0, \dots, [g/2]$, and the multi-index $I = \{i_1, \dots, i_k\}$ runs over all subsets of $[n] = \{1, \dots, n\}$; the unstable cases $j = 0$, $k = 0, 1$ are excluded).¹

The moduli space $\mathcal{H}_{g,n}$ admits a natural stratification according to the multiplicities of zeros and poles. Each stratum has a linear structure given by the *period coordinates*, where the transition maps are linear with integer coefficients. In the case of the principal stratum (i.e. the locus of 1-forms with only simple zeros) these coordinates are the integrals of v over a set of generators of the relative homology group $H_1(C \setminus \{z_i\}_{i=1}^{2g-2+n} \cup \{p_j\}_{j=1}^n)$ that includes the periods of v over a basis of canonical cycles on C , the integrals of v between its zeros z_i , and the residues of v at the poles p_j .

The principal stratum is open and dense in $\overline{\mathcal{H}}_{g,n}$, and its complement is a divisor whose class we denote by δ_{deg} (the subscript *deg* stands for *degenerate* as opposed to generic). The moduli space $\overline{\mathcal{H}}_{g,n}$ fibers into affine spaces parameterizing 1-forms with fixed residues $\mathbf{r} = (r_1, \dots, r_n)$, where $r_1 + \dots + r_n = 0$. By linearity, the intersection $\delta_{deg} \cdot \overline{\mathcal{H}}_{g,n}[\mathbf{r}]$ has codimension 1 in $\overline{\mathcal{H}}_{g,n}[\mathbf{r}]$ and is locally described by vanishing of the integral of v along a path connecting two coalescing zeros of v , cf. [12], Section 5.2, for a relevant discussion.

The main objective of this note is to decompose the pushforward class

$$\delta_{deg}[\mathbf{r}] = \text{pr}_* (\delta_{deg} \cdot \overline{\mathcal{H}}_{g,n}[\mathbf{r}])$$

¹Slightly abusing notation, we denote a divisor and its class in $\text{Pic}(\overline{\mathcal{M}}_{g,n}) \otimes \mathbb{Q}$ by the same symbol.

into a linear combination of the standard generators of $\text{Pic}(\overline{\mathcal{M}}_{g,n}) \otimes \mathbb{Q}$, see Theorem 1. Rather unexpectedly, the result depends on partial sums of the residues being zero or not.

Now we remind the basic facts about the tau function, cf. [11]. Choose a canonical basis $\{a_1, b_1, \dots, a_g, b_g\}$ in $H_1(C, \mathbb{Z})$. It defines a normalized basis $\{v_1, \dots, v_g\}$ in the space of holomorphic 1-forms on C by the condition $\int_{a_i} v_j = \delta_{ij}$, $i, j = 1, \dots, g$. Denote by Ω the period matrix $\Omega_{ij} = \int_{b_i} v_j$ of C .

Consider the following multi-valued $g(1-g)/2$ -differential on C :

$$c(x) = \frac{1}{W(x)} \left(\sum_{i=1}^g v_i(x) \frac{\partial}{\partial w_i} \right)^g \Theta(w, \Omega) \Big|_{w=K^x} \quad (1)$$

(see [4]). Here Θ is the theta function on C , K^x is the vector of Riemann constants with respect to the base point $x \in C$ defined by

$$K_j^x = \frac{1 + \Omega_{jj}}{2} - \sum_{i=1, i \neq j}^g \int_{y=a_i} \left(v_i(y) \int_x^y v_j \right)$$

(see [4], p. 6), and $W(x)$ is the Wronskian of holomorphic 1-forms defined as the usual Wronskian

$$W(x) = W(v_1, \dots, v_g) = W(f_1(\xi), \dots, f_g(\xi))(d\xi)^{g(g+1)/2}$$

of g functions $f_1(\xi), \dots, f_g(\xi)$, where ξ is a local coordinate near $x \in C$ and $v_i(\xi) = f_i(\xi)d\xi$, $i = 1, \dots, g$.

The differential v gives rise to a natural set of the so-called *distinguished* local parameters on C (see [11, 9]) defined as follows:

- near any regular point y (that is, $y \notin (v)$, where (v) is the divisor of v on C), the distinguished local parameter is $\zeta(x) = \int_x^y v$.
- near a simple zero z_i of v the distinguished local parameter is

$$\xi_i(x) = \left(\int_{z_i}^x v \right)^{1/2}, \quad i = 1, \dots, 2g - 2 + n \quad (2)$$

(here we assume that all zeros of v are simple and numbered), and

- near a simple pole p_j of v the distinguished local parameter is

$$\zeta_j(x) = \exp \left(\frac{2\pi\sqrt{-1}}{r_j} \int_{z_1}^x v \right), \quad j = 1, \dots, n, \quad (3)$$

where the contour of integration in (3) connects the first zero z_1 with a neighborhood of the pole p_j .

The Schottky–Klein *prime form* $E(x, y)$ is a (normalized) holomorphic section of the line bundle with the divisor $\{x = y\}$ on $C \times C$ (see [4]). It is explicitly given by the formula

$$E(x, y) = \frac{\Theta[*](\mathcal{A}(x) - \mathcal{A}(y))}{\sqrt{h(x)}\sqrt{h(y)}}, \quad (4)$$

where $\Theta[*]$ the theta function with an arbitrary half-integer characteristic $[*]$, $\mathcal{A}(x)$ is the Abel map, and $h(x)$ is a holomorphic 1-form on C with only

double zeros defined by

$$h(x) = \sum_{j=1}^g \frac{\partial}{\partial w_j} \Theta[*](w) \Big|_{w=0} v_j(x).$$

The values of the prime forms $E(x, z_i)$, $E(x, p_j)$, $E(z_i, z_k)$, $E(p_j, p_l)$ and $E(z_i, p_j)$, where one of or both arguments coincide with the points of divisor (v) , can be computed in terms of the distinguished local coordinates (2), (3). Namely,

$$\begin{aligned} E(x, z_i) &= E(x, y) \sqrt{d\xi_i(y)}|_{y=z_i}, \\ E(x, p_j) &= E(x, y) \sqrt{d\zeta_j(y)}|_{y=p_j}, \\ E(z_i, z_j) &= E(x, y) \sqrt{d\xi_i(x)d\xi_j(y)}|_{x=z_i, y=z_j}, \\ E(p_i, p_j) &= E(x, y) \sqrt{d\zeta_i(x)d\zeta_j(y)}|_{x=p_i, y=p_j}, \\ E(p_i, z_j) &= E(x, y) \sqrt{d\zeta_i(x)d\xi_j(y)}|_{x=p_i, y=z_j}. \end{aligned} \quad (5)$$

Since all zeros of v are assumed to be simple, we can choose for this computation a fundamental domain of C such that $\mathcal{A}_x((v)) + 2K^x = 0$, where \mathcal{A} is the Abel map (see [11], Lemma 6).

Now we are in a position to introduce the following

Definition 1. *The tau function is given by the formula*

$$\tau(C, v) = c^{16}(x) \left(\frac{v(x) \prod_{i=1}^n E(x, p_i)}{\prod_{i=1}^{2g-2+n} E(x, z_i)} \right)^{8(g-1)} \left(\frac{\prod_{i<j} E(z_i, z_j) \prod_{i<j} E(p_i, p_j)}{\prod_{i=1}^n \prod_{j=1}^{2g-2+n} E(p_i, z_j)} \right)^4. \quad (6)$$

Here and below we adopt the convention that the product over $i < j$ means that it is taken over all pairs of indices i and j in their respective range (i.e. from 1 to $2g - 2 + n$ in the case of $E(z_i, z_j)$ and from 1 to n in the case of $E(p_i, p_j)$).

Lemma 1. *The tau function defined by (6) is independent of $x \in C$.*

Proof. For convenience, we write the divisor of v as

$$(v) = \sum_{i=1}^{2g-2+n} z_i - \sum_{i=1}^n p_i = \sum_{j=1}^N k_j x_j, \quad (7)$$

where $N = 2g - 2 + 2n$ and $k_i = \pm 1$ ($+1$ for z_i 's and -1 for p_j 's). Then the expression (6) can be rewritten as follows:

$$\tau = c^{16}(x) \left(\frac{v(x)}{\prod_{i=1}^N E^{k_i}(x, x_i)} \right)^{8(g-1)} \prod_{i<j} E^{4k_i k_j}(x_i, x_j) \quad (8)$$

(recall that it holds when the fundamental polygon is chosen such that $\mathcal{A}_x((v)) + 2K^x = 0$). Note that (8) formally coincides with (3.24) and (3.25) of [11] except that some of k_j 's can now be negative. Observe that the right-hand side of (8) is non-singular and non-vanishing on C : all zeros

and poles of v are cancelled by the corresponding poles or zeros of $E(x, x_j)$. Moreover, the tensor weight of τ with respect to x is 0 since

- $c^{16}(x)$ is a differential of weight $8g(1 - g)$,
- $v(x)^{8(g-1)}$ is a differential of weight $8(g - 1)$,
- the product of prime forms in the denominator of (8) is a differential of weight $(-1/2)(8(g - 1)) \sum_{j=1}^N k_j = -8(g - 1)^2$.

It remains to verify that τ , as a function of x , is single-valued on C . When x goes around any of a -cycles a_j , both $c(x)$ and the prime forms $E(x, x_j)$ remain invariant, possibly up to a sign, thus τ does not change. To compute the change of τ around a cycle b_j we recall that the prime form $E(x, y)$ transforms as follows [4]:

$$E(b_j(x), y) = \pm e^{-\pi\sqrt{-1}\Omega_{jj} - 2\pi\sqrt{-1}(\mathcal{A}_j(x) - \mathcal{A}_j(y))} E(x, y),$$

so that

$$E(b_j(x), x_k) = \pm e^{-\pi\sqrt{-1}\Omega_{jj} + 2\pi\sqrt{-1}(\mathcal{A}_x(x_k))_j} E(x, y),$$

while for the factor $c(x)$ we have

$$c(b_j(x)) = e^{-\pi\sqrt{-1}(g-1)^2\Omega_{jj} - 2\pi\sqrt{-1}(g-1)K_j^x} c(x)$$

(here $b_j(x)$ means that the argument x goes along b_j once). Thus, we have

$$\begin{aligned} \log \frac{\tau(b_j(x))}{\tau(x)} &= 16(-\pi\sqrt{-1}(g-1)^2\Omega_{jj} - 2\pi\sqrt{-1}(g-1)K_j^x) \\ &\quad - 8(g-1)(-\pi\sqrt{-1}\Omega_{jj} \sum_{i=1}^N k_i + 2\pi\sqrt{-1} \sum_{i=1}^N k_i(\mathcal{A}_x(x_i))_j) = 0 \end{aligned}$$

since $\sum_{i=1}^N k_i = 2g - 2$ and $\sum_{i=1}^N \mathcal{A}_x(x_i) = -2K_\alpha^x$. Therefore, $\tau(x)$ is a holomorphic function on C and as such is x -independent. \square

Let us make a few more comments about definition (6)

- The tau function on the space of holomorphic 1-forms [11, 13] was originally defined as a covariant constant section of the trivial line bundle on the corresponding moduli space with respect to a certain flat connection. The equations for the tau function were then explicitly integrated that led to expression (6) with $n = 0$. In the case of the moduli space of meromorphic 1-forms, a derivation of defining equations for the tau function was just outlined in [9]. In order to keep exposition as self-contained as possible, we decided not to use them in the present paper and to derive all the required properties of the tau function starting from the explicit formula (6). This would also provide alternative proofs to the results of [11, 13].
- Expression (6) is a scalar that seemingly depend on the choice of distinguished local coordinates near z_i and p_j . Notice that the parameters ξ_i in (3) are defined up to a sign, but since the powers of all prime forms in (6) are multiples of 4, the formula for τ is invariant with respect to the choice of these signs.

- In turn, the parameters ζ_j given by (3) depend on the choice of the first zero z_1 and the choice of integration paths connecting z_1 with the poles p_j . This dependence propagates to the tau function (6), and an easy computation shows that the total power of the differential $d\zeta_i$ in (6) is equal to -2 . Namely, it enters $4(g-1)$ times to the power of $E(x, p_i)$ and $2(n-1)$ times to the power of $E(p_i, p_j)$ in the numerator of $\tau(C, v)$, and $-2(2g-2+n)$ times to the power of $E(p_i, z_j)$ in the denominator (taken with the minus sign), summing up to -2 . Thus, the product

$$\tau(C, v) \left(\prod_{i=1}^n d\zeta_i(p_i) \right)^2 \quad (9)$$

does not depend on the choice of local coordinate ζ_i near p_i , $i = 1, \dots, n$, given by (3).

Yet another expression for the tau function uses the bidifferential

$$\sigma(x, y) = \left(\frac{c(x)}{c(y)} \right)^{1/(1-g)}, \quad (10)$$

where $\sigma(x, y)$ is a multivalued differential of weight $-g/2$ in x and weight $g/2$ in y , cf. [4]. We have

Proposition 1. *Up to a multiplicative constant, the tau function (8) is given by the formula*

$$\tau(C, v) = c^{16}(x) \left(\prod_{j=1}^N \sigma^{4k_j}(x, x_j) \right) \left(\frac{v(x)}{\prod_{j=1}^N E^{k_j}(x, x_j)} \right)^{4(g-1)}, \quad (11)$$

where

$$\sigma(x, x_l) = \sigma(x, y) (d\xi_k(y))^{-g/2} \Big|_{y=x_l},$$

and ξ_l is a distinguished local coordinate near x_l .

Proof. To begin with, the expression (11) was obtained in [10], formula (1.8), when all $k_i = 1$, and in a preliminary version of [11], arXiv:math/0405042v1, formulas (4.4) and (4.5), when all $k_i > 0$. Here we give a proof for arbitrary k_i . Using the independence of (8) on x , we can express the ratio $c(x)/c(y)$ via the prime forms like follows:

$$\sigma(x, y) = \left(\frac{v(x)}{v(y)} \prod_{i=1}^N \frac{E^{k_i}(y, x_i)}{E^{k_i}(x, x_i)} \right)^{1/2}.$$

From here we get

$$\begin{aligned} \sigma(x, x_l) &= \left(\frac{v(x)}{\prod_{i=1}^N E^{k_i}(x, x_i)} \right)^{1/2} \prod_{i \neq l} E^{1/2}(x_l, x_i) \\ &\times \left(\frac{E^{k_l}(y, x_l) (d\xi_l(y))^{1+k_l/2}}{v(y)} \Big|_{y=x_l} \right)^{1/2} \end{aligned} \quad (12)$$

The last term in the product (12) is constant since ξ_l is a distinguished local coordinate near x_l . Consider first the case $k_l = 1$, i.e. $x_l = z_i$ is a simple zero of v . Then $v(y) = 2\xi_i(y)d\xi_i(y)$ and $E(y, x_i) \sim \xi_i(y)(d\xi_i(y))^{-1/2}$ so that

$$\frac{E(y, z_i)(d\xi_i(y))^{3/2}}{v(y)} \Big|_{y=z_i} = \frac{1}{2}.$$

Another case relevant for this paper is $k_l = -1$, i.e. $x_l = p_j$ is a simple pole of v . Then $v(y) = \frac{r_j}{2\pi\sqrt{-1}} \frac{d\zeta_j(y)}{\zeta_j(y)}$, $E(z_j, y) \sim \zeta_j(y)(d\zeta_j(y))^{-1/2}$ and

$$\frac{(d\zeta_j(y))^{1/2}}{v(y)E(y, \zeta_j)} \Big|_{y=\zeta_j} = \frac{2\pi\sqrt{-1}}{r_j}.$$

Therefore,

$$\prod_{j=1}^N \sigma^{4k_j}(x, x_j) = \text{const} \left(\frac{v(x)}{\prod_{j=1}^N E^{k_j}(x, x_j)} \right)^{4g-4} \prod_{i < j} E^{4k_i k_j}(x_i, x_j),$$

which, being substituted into (11), gives (8). \square

Now we are in a position to formulate

Proposition 2. *Fix $r_1, \dots, r_n \in \mathbb{C}$ such that $\sum_{i=1}^n r_i = 0$ and $r_i \neq 0$, $i = 1, \dots, n$. Then the tau function $\tau(C, v)$ is a holomorphic section of the line bundle $(\det \pi_* \omega_{g,n})^{\otimes 24} \otimes \bigotimes_{i=1}^n \mathcal{L}_i^{\otimes 2}$, nowhere vanishing over the principal stratum $\overline{\mathcal{H}}_{g,n}[\mathbf{r}] \cap (\mathcal{H}_{g,n} \setminus \delta_{deg})$ in the moduli space $\overline{\mathcal{H}}_{g,n}[\mathbf{r}]$.*

A proof of this Proposition based on transformation properties of the prime form and the differential $c(x)$ [4] is given in Appendix A.

Now let us compute the divisor of the tau function $\tau(C, v)$ on $\overline{\mathcal{H}}_{g,n}[\mathbf{r}]$, or, equivalently, the asymptotics of $\tau(C, v)$ near $\delta_{deg} \cdot \overline{\mathcal{H}}_{g,n}[\mathbf{r}]$, $\text{pr}^* \delta_{irr}$, and $\text{pr}^* \delta_{j,I}$.

Asymptotics of τ near $\delta_{deg} \cdot \overline{\mathcal{H}}_{g,n}[\mathbf{r}]$. We begin with computing the asymptotics of the tau function $\tau(C, v)$ near the divisor δ_{deg} in the total moduli space $\overline{\mathcal{H}}_{g,n}$. Let z_1, z_2 be any two simple zeros of v on C that coalesce to a double zero.

Lemma 2. *The tau function has the following asymptotics near δ_{deg} :*

$$\tau(C, v) = \text{const} \cdot t + o(t) \quad \text{as } t \rightarrow 0 \quad (13)$$

with some $\text{const} \neq 0$, where t is a transversal local coordinate on $\overline{\mathcal{H}}_{g,n}$ near δ_{deg} .

Proof. Consider a domain D on C with local coordinate s containing both z_1 and z_2 ; we assume that

$$v = s(s - \varepsilon)ds.$$

Then δ_{deg} is locally defined by equation $\varepsilon = 0$, and the local transversal coordinate t near δ_{deg} can be chosen to be $t = \varepsilon^2$ (since z_1 and z_2 are interchangeable; notice that the period coordinate $\int_{z_1}^{z_2} v = -\varepsilon^3/6$). The distinguished local coordinates near z_1 and z_2 are given by $\xi_1(s) = (\int_0^s v)^{1/2}$ and $\xi_2(s) = (\int_\varepsilon^s v)^{1/2}$ respectively, so that

$$\frac{d\xi_1(s)}{ds} \Big|_{s=0} = \left(-\frac{\varepsilon}{2}\right)^{1/2}, \quad \frac{d\xi_2(s)}{ds} \Big|_{s=\varepsilon} = \left(\frac{\varepsilon}{2}\right)^{1/2}. \quad (14)$$

The following terms in (6) do not have finite nonzero limit as $z_1 \rightarrow z_2$.

First, this is the theta function in the numerator of the prime form $E^4(z_1, z_2)$, which contributes the term ε^4 to the asymptotics of τ as $\varepsilon \rightarrow 0$ (since the theta-function with non-singular odd characteristics has simple zero at the origin).

Second, these are derivatives $\frac{d\xi_1}{ds}(z_1)$ and $\frac{d\xi_2}{ds}(z_2)$, each of which enter with power -2 (that appears in the same way as the power of $d\zeta_i$ computed before (9)). Therefore, these derivatives contribute the power ε^{-2} to the asymptotics of τ as $\varepsilon \rightarrow 0$, and in total we get $\tau \sim \varepsilon^2$ as $\varepsilon \rightarrow 0$ which implies (13) since $t = \varepsilon^2$. \square

Asymptotics of τ near $\text{pr}^*\delta_{irr}$. This asymptotics is also computed using only the explicit formula for the tau function. As the curve C degenerates to an irreducible curve with a node, it acquires two new simple poles at the preimages of the node with opposite residues, while its genus drops by 1.

Lemma 3. *The tau function (6) has the following asymptotics near $\text{pr}^*\delta_{irr}$:*

$$\tau(C, v) = \text{const} \cdot t^2 + o(t^2) \quad \text{as } t \rightarrow 0 \quad (15)$$

with some $\text{const} \neq 0$, where t is a transversal local coordinate on $\overline{\mathcal{H}}_{g,n}$ near $\text{pr}^*\delta_{irr}$.

Proof. The expression (6) for the tau function is written under such choice of the fundamental polygon that $\mathcal{A}_x((v)) + 2K^x = 0$; however, this is not the choice suitable for analyzing the asymptotics near $\text{pr}^*\delta_{irr}$. Instead, we are going to choose the fundamental domain such that

$$\mathcal{A}_x((v)) + 2K^x - \Omega e_g = 0 \quad (16)$$

where $e_g = {}^t(0, \dots, 0, 1)$. Under such a choice of the fundamental domain the formula (6) needs to be multiplied by an extra exponential factor (see (3.24), (3.25) of [11] with $\mathbf{p} = -e_g$ and $\mathbf{q} = 0$):

$$\begin{aligned} \tau(C, v) &= \exp(-4\pi\sqrt{-1}\Omega_{gg}) \\ &\times c^{16}(x) \left(\frac{v(x) \prod_{i=1}^n E(x, p_i)}{\prod_{i=1}^{2g-2+n} E(x, z_i)} \right)^{8(g-1)} \left(\frac{\prod_{i < j} E(z_i, z_j) \prod_{i < j} E(p_i, p_j)}{\prod_{i=1}^n \prod_{j=1}^{2g-2+n} E(p_i, z_j)} \right)^4. \end{aligned} \quad (17)$$

For the following facts we refer to p.13 of [4]. In the limit $t \rightarrow 0$ all entries of the period matrix remain finite except the entry Ω_{gg} which behaves as

$$\Omega_{gg} = \frac{1}{2\pi\sqrt{-1}} \log t + O(t) \quad (18)$$

Assuming that in the limit $t \rightarrow 0$ the point x stays far from the nodes a and b (i.e. outside of the “plumbing zone”), all the components of the vector of Riemann constants K^x remain finite except K_g^x that has the asymptotics

$$K_g^x = \frac{1}{4\pi\sqrt{-1}} \log t + O(1). \quad (19)$$

At the same time, $c(x)$ and the prime form $E(x, y)$ have finite limits assuming again that x and y , as well as the path connecting them remain outside of the plumbing zone. The vector $K^x - \Omega e_g$ in (16) also remains finite as

$t \rightarrow 0$. Therefore, all the paths connecting the point x with the points of divisor (v) can be drawn outside of the plumbing zone. Thus, all the factors in the expression (17) remain finite as $t \rightarrow 0$, except the exponential factor which behaves like t^2 due to (18). \square

Asymptotics of τ near $\text{pr}^*\delta_{j,I}$. Two different cases are covered by the following two lemmas.

Lemma 4. *Let C degenerate to a reducible curve with components C_1 of genus j and C_2 of genus $g-j$ such that the poles p_{i_1}, \dots, p_{i_k} belong to C_1 and the poles $p_{i_{k+1}}, \dots, p_{i_n}$ belong to C_2 . Let also assume that*

$$\sum_{l=1}^k r_{i_l} = \sum_{l=k+1}^n r_{i_l} = 0.$$

Then in terms of any transverse local coordinate t near $\text{pr}^\delta_{j,I}$ in $\overline{\mathcal{H}}_{g,n}[\mathbf{r}]$ the asymptotics of τ is given by*

$$\tau(C, v) = \text{const} \cdot t^3 + o(t^3) \quad \text{as } t \rightarrow 0 \quad (20)$$

with some $\text{const} \neq 0$.

Proof. In this case two zeros of v , say, z_1 and z_2 , coalesce and disappear at the node that becomes a regular point for the both of components C_1 and C_2 . Then $t = \left(\int_{z_1}^{z_2} v\right)^2$ can be taken as a transversal local coordinate near $\text{pr}^*\delta_{j,I}$.

More precisely, the degeneration picture is qualitatively described by the standard fixture $\{zw = t \mid |z|, |w|, |t| < 1\}$ fibered over the disc $\{|t| < 1\}$. The fiber over $t \neq 0$ is the annulus $\{|t| < |z| < 1\}$ (identified with $\{|t| < |w| < 1\}$ by the map $w = t/z$), while the fiber over $t = 0$ is the union of two discs $\{z = 0\}$ and $\{w = 0\}$ intersecting at the origin. Up to rescaling, the 1-form v locally looks like $v \sim (1 - t/z^2)dz = (1 - t/w^2)dw$. The two coalescing zeros are, approximately, $z_{1,2} = \pm t^{1/2}$, and $v(z) \rightarrow 1$ as $t \rightarrow 0$. This local behavior of v is demonstrated by the model case of $\overline{\mathcal{H}}_{0,4}$, where the differential v is given by the formula (26). Near the boundary point, say, $p = 0$, the plumbing parameter t can be identified with p , and zeros of v as functions of t can be explicitly found from (26)).

Now we will use the alternative expression for the tau function given by (11). Assuming that x on C_1 and y on C_2 stay outside of the plumbing zone as $t \rightarrow 0$, we can analyze the asymptotics of all terms in (11). Namely, we have

$$c(x) \sim t^{-(g-j)(g-j-1)/2}, \quad E(x, y) \sim t^{-1/2}, \quad \sigma(x, y) \sim t^{(g-2j)/2} \quad (21)$$

(see [4], p. 14). If both x and y stay on C_1 or C_2 outside of the plumbing zone, then $E(x, y)$ and $\sigma(x, y)$ remain finite as $t \rightarrow 0$.

The asymptotics of $E(x, z_1)$, $E(x, z_2)$, $\sigma(x, z_1)$ and $\sigma(x, z_2)$ as $t \rightarrow 0$ are less trivial since both z_1 and z_2 tend to the nodal point. These asymptotics were computed in [10], Section 2.3. Namely,

$$E(x, z_i) \sim t^{-1/8}, \quad i = 1, 2,$$

see (2.35) in [10], and

$$\sigma(x, z_i) \sim t^{(3g-4j)/8}, \quad i = 1, 2,$$

where j is the genus of C_1 , see (2.40) in [10].

Now assume that the poles p_i, \dots, p_k stay on C_1 and the poles p_{k+1}, \dots, p_n stay on C_2 as $t \rightarrow 0$; the zeros $z_{1,2}$ tend to the node while the zeros z_3, \dots, z_{2j+k} stay on C_2 and the remaining $2(g-j) - 2 + (n-k)$ zeros stay on C_2 . We use the above asymptotic formulas to compute the total power of t in the asymptotics of $\tau(C, v)$ as $t \rightarrow 0$:

- the term $c^{16}(x)$ gives the power $-8(g-j)(g-j-1)$ of t ;
- the product of $\sigma^{4k_i}(x, x_i)$ when x_i remains on C_2 gives the power $4(g-2j)(g-j-1)$ of t ;
- the product $(\sigma(x, z_1)\sigma(x, z_2))^4$ gives the power $3(g-j) - j$ of t ;
- the product of prime forms $E^{4(1-g)k_j}(x_i, x)$ for x_i 's staying on C_2 gives the power $2(2(g-j) - 2)(g+2j-1)$ of t ;
- the product $(E(x, z_1)E(x, z_2))^{4(1-g)}$ gives the power $g-1$ of t .

Summing up all of these powers we get 3, as claimed in (20). \square

Lemma 5. *Let C degenerate to a reducible curve with components C_1 of genus j and C_2 of genus $g-j$ such that the poles p_{i_1}, \dots, p_{i_k} belong to C_1 and the poles $p_{i_{k+1}}, \dots, p_{i_n}$ belong to C_2 . Let also assume that*

$$\sum_{l=1}^k r_{i_l} = - \sum_{l=k+1}^n r_{i_l} \neq 0.$$

Then in terms of any transverse local coordinate t near $\text{pr}^ \delta_{j,I}$ in $\overline{\mathcal{H}}_{g,n}[\mathbf{r}]$ the asymptotics of τ is given by*

$$\tau(C, v) = \text{const} \cdot t^2 + o(t^2) \quad \text{as } t \rightarrow 0 \quad (22)$$

with some $\text{const} \neq 0$.

Proof. Contracting a separating loop on C we get a nodal curve and a meromorphic 1-form on it with simple poles at the preimages of the node $a \in C_1$ and $b \in C_2$ with residues r_0 and $-r_0$, respectively, where $r_0 = -\sum_{l=1}^k r_{i_l} \neq 0$. As before, we assume that the poles p_{i_1}, \dots, p_{i_k} belong to C_1 together with $2j-1+k$ zeros z_i . The remaining poles $p_{i_{k+1}}, \dots, p_{i_n}$ belong to C_2 together with $2(g-j)-1+n-k$ zeros z_i ; all zeros and poles (except for a, b) stay away from the node.

Like in the previous lemma, the local degeneration picture is described by the standard fixture $\{zw = t \mid |z|, |w|, |t| < 1\}$ fibered over the disc $\{|t| < 1\}$. We remind that the fiber over $t \neq 0$ is the annulus $\{|t| < |z| < 1\}$ (identified with $\{|t| < |w| < 1\}$ by the map $w = t/z$), while the fiber over $t = 0$ is the union of two discs $\{z = 0\}$ and $\{w = 0\}$ intersecting at the origin. Then the 1-form v asymptotically behaves like $v \sim r_0 \frac{dz}{z} = -r_0 \frac{dw}{w}$ as $t \rightarrow 0$ and develops simple poles with opposite residues at the points $a \in C_1$ and $b \in C_2$. Again, this can be seen from the formula (26).

Following [4], p. 14, Case II, we compute the asymptotics of all terms in (11) as $t \rightarrow 0$. Assuming that $x \in C_1$ and using the asymptotics (21) we have

$$c(x)^{16} \sim t^{-8((g-j)^2 + g-j)}$$

In the product of $\sigma^{k_i}(x, x_i)$ only the terms with $x_i \in C_2$ contribute to non-trivial powers of t . Thus

$$\prod_{i=1}^N \sigma^{4k_i}(x, x_i) = \frac{\prod_{i=1}^{2(g-j)-1+n-k} \sigma^4(x, z_i)}{\prod_{i=1}^{n-k} \sigma^4(x, p_i)} \sim t^{2(g-2j)(2(g-j)-1)}$$

and, finally,

$$\prod_{i=1}^N E^{k_i}(x, x_i) = \frac{\prod_{i=1}^{2(g-j)-1+n-k} E(x, z_i)}{\prod_{i=1}^{n-k} E(x, p_i)} \sim t^{1/2(2(g-j)-1)}.$$

Substituting these asymptotics into (11) we get (22). \square

Combining Proposition 2 with asymptotic formulas (13), (15), (20) and (22) we get

Theorem 1. Fix $r_1, \dots, r_n \in \mathbb{C}$ such that $\sum_{i=1}^n r_i = 0$ and $r_i \neq 0$, $i = 1, \dots, n$. Then the following relation holds in the Picard group $\text{Pic}(\overline{\mathcal{M}}_{g,n}) \otimes \mathbb{Q}$

$$\delta_{deg}[\mathbf{r}] = 24\lambda + 2 \sum_{i=1}^n \psi_i - 2\delta_{irr} - \frac{1}{2} \sum_{j=0}^g \sum_{I \subset [n]} \alpha_I \delta_{j,I}, \quad (23)$$

where $I = \{i_1, \dots, i_k\}$ runs over all subsets of $[n] = \{1, \dots, n\}$, and

$$\alpha_I = \begin{cases} 2, & \text{if } \sum_{j=1}^k r_{i_j} \neq 0, \\ 3, & \text{if } \sum_{j=1}^k r_{i_j} = 0. \end{cases} \quad (24)$$

Remark 1. The condition $r_i \neq 0$, $i = 1, \dots, n$, is crucial for our approach. When $r_i \rightarrow 0$, the tau function (6) acquires an essential singularity that makes its asymptotic analysis complicated, if at all doable.

Let us illustrate Theorem 1 on two simple examples.

Example 1. Let $g = 0$ and $n = 4$. Then $\overline{\mathcal{M}}_{0,4} = \mathbb{C}P^1$, its boundary divisor is $\delta = \{0, 1, \infty\}$, and in this case $\lambda = 0$, $\delta_{irr} = 0$. Moreover, combining Mumford's formula

$$\kappa_1 = 12\lambda + \sum_{i=1}^n \psi_i - \delta \quad (25)$$

(cf., e.g., [2], Chapter XIII, Eq. (7.7)) with the formula

$$\kappa_1 = \frac{1}{2(n-1)} \sum_{k=2}^{n-2} \sum_{|I|=k} (k-1)(n-k-1) \delta_{0,I}$$

from [16], Section 5, we see that in our case $\sum_{i=1}^4 \psi_i = \frac{4}{3} \delta$; here κ_1 is the first Mumford's class (also known as the Weil-Petersson class).

A general meromorphic 1-form is then given by the formula

$$v(z)dz = \left(\frac{r_1}{z-p} + \frac{r_2}{z} + \frac{r_3}{z-1} \right) dz, \quad z \in \mathbb{C}, \quad (26)$$

where $p, 0, 1$ and ∞ are the simple poles of v with residues r_1, r_2, r_3 and $r_4 = -r_1 - r_2 - r_3$ respectively. There are three distinct cases to consider.

Case 1. This is generic situation, when $r_1 + r_2 \neq 0$, $r_1 + r_3 \neq 0$, $r_2 + r_3 \neq 0$. The divisor δ_{deg} intersects each space $\overline{\mathcal{H}}_{0,4}[\mathbf{r}]$ at precisely two points with coordinates

$$p_{1,2} = -\frac{r_1 r_3 + r_2 r_4 \pm 2\sqrt{r_1 r_2 r_3 r_4}}{(r_2 + r_3)^2}$$

corresponding to the 1-forms with a double zero, so that in this case $\delta_{deg}[\mathbf{r}] = \frac{2}{3}\delta$. Since here $\sum_{I \subset [4]} \alpha_I \delta_{0,I} = 6 \cdot \frac{1}{3}\delta = 2\delta$, this agrees with (23).

Case 2. Here $r_1 + r_2 \neq 0$, $r_1 + r_3 \neq 0$, $r_2 + r_3 = 0$. The divisor δ_{deg} intersects each space $\overline{\mathcal{H}}_{0,4}[\mathbf{r}]$ at precisely one point with coordinate

$$p = \frac{(r_1 + r_2)^2}{4r_1 r_2},$$

corresponding to the 1-form with a double zero, so that in this case $\delta_{deg}[\mathbf{r}] = \frac{1}{3}\delta$. Since here $\sum_{I \subset [4]} \alpha_I \delta_{0,I} = (4 + 3) \cdot \frac{1}{3}\delta = \frac{7}{3}\delta$, this also agrees with (23).

Case 3. Here $r_1 + r_2 = 0$, $r_1 + r_3 \neq 0$, $r_2 + r_3 = 0$. The intersection of $\overline{\mathcal{H}}_{0,4}[\mathbf{r}]$ with δ_{deg} is trivial, so that in this case $\delta_{deg}[\mathbf{r}] = 0$. Since here $\sum_{I \subset [4]} \alpha_I \delta_{0,I} = (2 + 6) \cdot \frac{1}{3}\delta = \frac{8}{3}\delta$, this again agrees with (23).

Example 2. Take $g = 1$, $n = 2$, $r_1 = 1$, $r_2 = -1$, and compute the pushforward class $\delta_{deg}[\pm 1]$ in the Picard group $\text{Pic}(\overline{\mathcal{M}}_{1,2}) \otimes \mathbb{Q}$ in two different ways.

On the one hand, formula (23) adapted for this case reads

$$\delta_{deg}[\pm 1] = 24\lambda + 2(\psi_1 + \psi_2) - 2\delta_0 - 3\delta_1, \quad (27)$$

where $\delta_0 = \frac{1}{2}[\overline{\mathcal{M}}_{0,4}]$ and $\delta_1 = [\overline{\mathcal{M}}_{1,1}]$ are the classes of the boundary divisors of $\overline{\mathcal{M}}_{1,2}$.

On the other hand, we can modify to our case (we briefly sketch it below) an approach of D. Zvonkine [17] originally developed for the case of holomorphic 1-forms (some details can be found in [15], Section 1.6, An example) which leads to the same result.

Namely, denote by $\overline{\mathcal{H}}_{1,2}^1[\pm 1]$ the moduli space of meromorphic 1-forms on elliptic curves with two simple poles of residues ± 1 and one marked point, say, $x \in C$, $x \neq p_1, p_2$. Let \mathcal{L} be the tautological line bundle on $\overline{\mathcal{H}}_{1,2}^1[\pm 1]$, where $\mathcal{L}|_{(C,v,x)} = T_x^* C$. It comes with a canonical section that assigns to the 1-form v the cotangent vector $v(x) \in \mathcal{L}$. This section vanishes precisely when v has a zero at x , and we denote the closure of its zero set by Δ_1 . Clearly, $[\Delta_1] = \psi$, where $\psi = c_1(\mathcal{L})$.

Let us turn our attention to the locus Δ_1 . The restriction of $\mathcal{L}^{\otimes 2}$ to Δ_1 has a section that vanishes when a 1-form has a second order zero at the marked point x . Denote the corresponding sub-locus by Δ_2 ; clearly, $[\Delta_2] = \psi \cdot (2\psi)$. The locus Δ_2 consists of components of three types:

- (1) The closure $\Delta_2^{(1)}$ of the locus where the marked point is a zero of order exactly 2.
- (2) The locus $\Delta_2^{(2)}$, where the marked point belongs to an elliptic tail attached to the rest of the curve by a single node.
- (3) The locus $\Delta_2^{(3)}$, where the marked point belongs to a rational curve connecting two otherwise disconnected components of the curve.

Push the equality $[\Delta_2^{(1)}] + [\Delta_2^{(2)}] + [\Delta_2^{(3)}] = 2\psi^2$ forward using the natural forgetful map $\text{pr} : \overline{\mathcal{H}}_{1,2}^1[\pm 1] \rightarrow \overline{\mathcal{M}}_{1,2}$. We get

$$\text{pr}_*([\Delta_2^{(1)}]) = \delta_{deg}[\pm 1], \quad \text{pr}_*([\Delta_2^{(2)}]) = 0, \quad \text{pr}_*([\Delta_2^{(3)}]) = \delta_1, \quad 2\text{pr}_*(\psi^2) = 2\kappa_1$$

that yields

$$\delta_{deg}[\pm 1] = 2\kappa_1 - \delta_1. \quad (28)$$

Substituting Mumford's formula (25) into the last equality we get formula (27).

Moreover, Mumford's computation of the Chow ring of $\overline{\mathcal{M}}_2$ (cf. [14], Part III) implies, in particular, that

$$\lambda = \frac{1}{12} \delta_0, \quad \psi_{1,2} = \frac{1}{12} \delta_0 + \delta_1, \quad \kappa_1 = \frac{1}{6} \delta_0 + \delta_1,$$

and, therefore, $\delta_{deg}[\pm 1] = \frac{1}{3} \delta_0 + \delta_1$ in the Picard group $\text{Pic}(\overline{\mathcal{M}}_{1,2}) \otimes \mathbb{Q}$.

Remark 2. After this paper was completed, A. Sauvaget informed us that the methods of his paper [15] would, in principle, allow one to express the divisor class δ_{deg} in terms of the standard generators of the rational Picard group of the moduli space $\overline{\mathcal{H}}_{g,n}$.

Remark 3. The tau function on moduli spaces of holomorphic differentials (when $n = 0$) gives the holomorphic factorization formula for the determinant of Laplacian on a Riemann surface with flat conical metric $|v|^2$ (see [11]). For the meromorphic case $n > 0$ the volume of a Riemann surface in flat metric $|v|^2$ is infinite; this makes the spectral theory more complicated. Nevertheless, the determinant of Laplacian can be defined on the so-called Mandelstam diagrams (see [6]) that correspond to the space of pairs (C, v) , where all periods of the differential v are imaginary. On this real slice of the moduli space the determinant of Laplacian was defined and computed in terms of the tau function (6) in [8], Corollary 1.

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APPENDIX A. TRANSFORMATION OF TAU FUNCTION UNDER A CHANGE OF CANONICAL BASIS OF CYCLES

Here we prove that the expression (9) is a section of 24th power of the Hodge line bundle, i. e. it transforms properly under symplectic changes of canonical bases of cycles on C . Let a canonical basis of cycles $\{a_j, b_j\}_{j=1}^g$ transform to a new basis $\{\hat{a}_j, \hat{b}_j\}_{j=1}^g$ by a matrix $G \in Sp(2g, \mathbb{Z})$:

$$\begin{pmatrix} \hat{b} \\ \hat{a} \end{pmatrix} = G \begin{pmatrix} b \\ a \end{pmatrix}, \quad (29)$$

where

$$G = \begin{pmatrix} A & B \\ C & D \end{pmatrix}.$$

Then the column vector of v of (normalized) Abelian differentials v_1, \dots, v_g , the period matrix Ω , and the prime form $E(x, y)$ transform as follows ([4], pp.11-12):

$$\hat{v} = {}^t(C\Omega + D)^{-1}v, \quad (30)$$

$$\begin{aligned} \hat{\Omega} &= (A\Omega + B)(C\Omega + D)^{-1}, \\ \hat{E}(x, y) &= E(x, y)e^{\pi\sqrt{-1}\langle Q\mathcal{A}|_x^y, \mathcal{A}|_x^y \rangle}, \end{aligned} \quad (31)$$

where

$$Q = (C\Omega + D)^{-1}C.$$

Transformation formulas for the vector of Riemann constants K^x and the multi-differential $c(x)$ also involve two vectors $\alpha_0, \beta_0 \in \frac{1}{2}\mathbb{Z}^g$ defined by the requirement that

$$\alpha_0 - \frac{1}{2}(C^t D)_0 \in \mathbb{Z}^g, \quad \beta_0 - \frac{1}{2}(A^t B)_0 \in \mathbb{Z}^g,$$

see Lemma 1.5 of [4]; here the subscript 0 denotes the vector consisting of diagonal entries of a square matrix. The above conditions fix vectors α_0 and β_0 up to a period of the Jacobian; this period is determined by the transformation law of the vector of the Riemann constants:

$$\hat{K}^x = M^{-1}K^x + \hat{\Omega}\alpha_0 + \beta_0 \quad (32)$$

with

$$M = {}^t(C\Omega + D).$$

Finally, the transformation of the differential $c(x)$ looks as follows (see (1.23) of [4]):

$$\begin{aligned} \hat{c}(x) &= \epsilon (\det M)^{3/2} c(x) \\ &\times \exp(\pi\sqrt{-1}\langle K^x, QK^x \rangle - \pi\sqrt{-1}\langle \alpha_0, \hat{\Omega}\alpha_0 \rangle - 2\pi\sqrt{-1}\langle \alpha_0, M^{-1}K^x \rangle), \end{aligned} \quad (33)$$

where $\epsilon^8 = 1$.

Now we will need the following

Lemma 6. *Under symplectic transformation (29) of the canonical basis of cycles the tau function (6) transforms as follows:*

$$\hat{\tau} = (\det M)^{24}\tau \quad (34)$$

Proof. The proof is obtained by substituting the transformation formulas (31), (32) and (33) into the formula for the tau function (8); recall that it holds when the fundamental polygon is chosen such that $\mathcal{A}_x((v)) + 2K^x = 0$.

Under an arbitrary choice of the fundamental polygon this relation holds up to an integer combination of the Jacobian lattice vectors, i.e. there exist two vectors $\mathbf{r}, \mathbf{q} \in \mathbb{Z}^g$ such that

$$\mathcal{A}_x((v)) + 2K^x + \Omega\mathbf{p} + \mathbf{q} = 0 \quad (35)$$

If the divisor (v) satisfies (35), then the expression (8) gets an extra exponential factor given by

$$\exp\{-4\pi\sqrt{-1}\langle \mathbf{r}, \Omega\mathbf{p} \rangle - 16\pi\sqrt{-1}\langle \mathbf{p}, K^x \rangle\} \quad (36)$$

(see (3.24) and (3.25) of [11]).

Let us now apply the symplectic transformation (29) to the right hand side of (8) and denote the result by $\hat{\tau}$. To write down the formula for $\hat{\tau}$ in terms of \hat{E} and \hat{c} we observe that the Abel map $\hat{\mathcal{A}}_x$ of divisor (v) can be computed using the relation $\mathcal{A}_x((v)) + 2K^x = 0$ combined with (30) and (32):

$$\hat{\mathcal{A}}_x((v)) + 2\hat{K}^x - 2\hat{\Omega}\alpha_0 - 2\beta_0 = 0, \quad (37)$$

i.e. in (35) we get $\mathbf{p} = -2\alpha_0$, $\mathbf{q} = -2\beta_0$.

Therefore, according to (36),

$$\begin{aligned} \hat{\tau} = & \hat{c}^{16}(x) \left(\frac{v(x)}{\prod_{i=1}^N \hat{E}^{k_i}(x, x_i)} \right)^{8(g-1)} \prod_{i < j} \hat{E}^{4k_i k_j}(x_i, x_j) \\ & \times \exp(-16\pi\sqrt{-1}\langle\alpha_0, \hat{\Omega}\alpha_0\rangle + 32\pi\sqrt{-1}\langle\alpha_0, \hat{K}^x\rangle). \end{aligned} \quad (38)$$

Due to (32) the exponential factor in (38) can be rewritten as

$$\exp(16\pi\sqrt{-1}\langle\alpha_0, \hat{\Omega}\alpha_0\rangle + 32\pi\sqrt{-1}\langle\alpha_0, M^{-1}K^x\rangle), \quad (39)$$

where we used the equality $\exp(32\pi\sqrt{-1}\langle\alpha_0, \beta_0\rangle) = 1$ that holds because α_0 and β_0 are vectors with half-integer components.

Compute now $\log \frac{\hat{\tau}}{\tau}$ using (31), (33). Recall that $\sum_{i=1}^N k_i = 2g - 2$, and put $u_i = \mathcal{A}_x(x_i)$. Then the contribution of the product of exponents in the denominator of $\log \frac{\hat{\tau}}{\tau}$ equals to

$$8\pi\sqrt{-1}(1-g) \sum_{i=1}^N k_i \langle Qu_i, u_i \rangle, \quad (40)$$

while the contribution of the prime forms in the numerator of $\hat{\tau}$ and τ is

$$4\pi\sqrt{-1} \sum_{i < j} k_i k_j \langle Q(u_i - u_j), u_i - u_j \rangle. \quad (41)$$

The sum of two sums (40) and (41) equals to

$$-4\pi\sqrt{-1} \langle Q \sum_{i=1}^N k_i u_i, \sum_{j=1}^N k_j u_j \rangle = -16\pi\sqrt{-1} \langle QK^x, K^x \rangle \quad (42)$$

(we use that $\sum_{i=1}^N k_i = 2g - 2$ here).

Moreover, due to (33),

$$\begin{aligned} 16 \log \frac{\hat{c}(x)}{c(x)} = & 24 \log \det(C\Omega + D) + 16\pi\sqrt{-1} \langle QK^x, K^x \rangle \\ & - 16\pi\sqrt{-1} \langle \alpha_0, \hat{\Omega}\alpha_0 \rangle - 32\pi\sqrt{-1} \langle \alpha_0, M^{-1}K^x \rangle. \end{aligned} \quad (43)$$

Notice that the exponential factor (39) in the expression for $\log(\hat{\tau}/\tau)$ is canceled by the last two terms in (43), while the second term in the right hand side of (43) cancels against (42). Finally, we arrive at (34). \square

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DEPARTMENT OF MATHEMATICS AND STATISTICS, CONCORDIA UNIVERSITY, 1455 DE MAISONNEUVE WEST, MONTREAL, H3G 1M8 QUEBEC, CANADA
Email address: dmitry.korotkin@concordia.ca

ST. PETERSBURG DEPARTMENT OF STEKLOV MATHEMATICAL INSTITUTE, FONTANKA 27, ST. PETERSBURG 191023 RUSSIA, AND
 CHEBYSHEV LABORATORY, ST. PETERSBURG STATE UNIVERSITY, 14TH LINE V.O. 29, SAINT PETERSBURG 199178 RUSSIA
Email address: zograp@pdmi.ras.ru