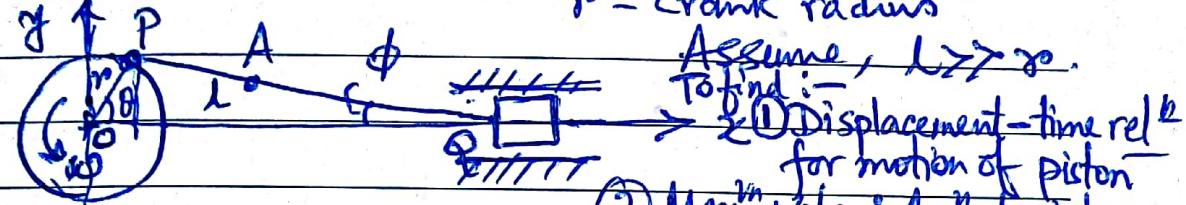


Slider Crank mechanism: ω = Angular velo. of crank
 r^o = Crank radius



Assume, $l \gg r^o$.

To find:
① Displacement-time rel^k for motion of piston
② Max^m Velo. & Accⁿ of piston
③ Motion of point A.

Here, $\theta = \omega t$

$$x = r \cos \theta + l \cos \phi \quad \dots \dots (1)$$

from ΔOPA : $\frac{\sin \phi}{r^o} = \frac{\sin \theta}{l}$ or, $\sin \phi = \frac{r^o}{l} \sin \theta$.

or,

So, $x = r \cos \theta + l \left[1 - \frac{r^o}{l} \sin^2 \theta \right]^{\frac{1}{2}}$ or, $\cos \phi = \sqrt{1 - \frac{r^o}{l} \sin^2 \theta}$

As $l \gg r^o$, $\frac{r^o}{l} \approx 0$.

$$x = r \cos \theta + l = r \cos \omega t + l.$$

$$v = \dot{x} = -r \omega \sin \omega t$$

$$a = \ddot{x} = \ddot{v} = -r \omega^2 \cos \omega t$$

Velocity will be max^m when $\theta = \omega t = 90^\circ$

Acceleration will be max^m when $\theta = \omega t = 0^\circ$

Let point A is p distance away from point P

$$\begin{aligned} x_A &= r \cos \theta + p \cos \phi \\ &= r \cos \omega t + p \left(1 - \frac{r^o}{l} \sin^2 \omega t \right)^{\frac{1}{2}} \end{aligned}$$

$$\begin{aligned} y_A &= (l - p) \sin \phi = (l - p) \cdot \frac{x}{l} \cdot \sin \theta \\ &= \left(\frac{l-p}{l} \right) r \sin \omega t \end{aligned}$$

At $l \approx \infty$,

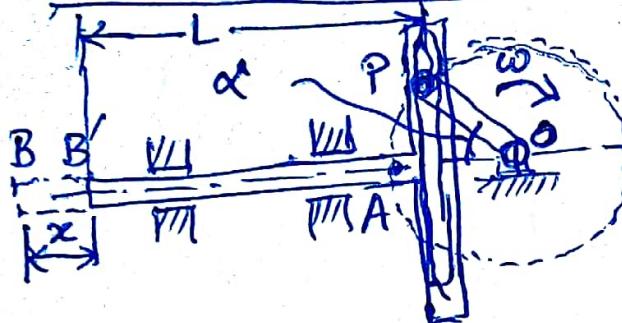
$$x_A = (r + p) \cos \omega t$$

$$y_A = (r - p) \sin \omega t$$

$$\frac{x_A^2}{(r+p)^2} + \frac{y_A^2}{(r-p)^2} = \cos^2 \omega t + \sin^2 \omega t = 1$$

This is the eqn of ellipse.

Scotch Yoke Mechanism: ω = Angular velocity of crank
 Find displacement, velocity (OP) and Accⁿ of Sliding Bar AB.



At horizontal position, point P and point A become same ie, at $\alpha = 0^\circ$

At other position, some displacement $BB' = x$

$$\text{Now, } x = OB - OB' = OB - OP \cos \alpha - AB'$$

$$x = OB - OP \cos \alpha - L \quad \dots (1)$$

$$\text{At } \alpha = 0^\circ, \quad OB = L + OA = L + OP$$

$$\text{from Eqn (1)} \quad x = OB - OP \cos \alpha - L$$

$$\begin{aligned} &= (L + OP) - OP \cos \alpha - L \\ &= OP(1 - \cos \alpha) \end{aligned}$$

Here, length of crank OP can be taken as radius R of rotating circle.

$$x = R(1 - \cos \alpha) = R(1 - \cos \omega t)$$

$$v = \dot{x} = R \omega \sin \omega t$$

$$a = \ddot{x} = R \omega^2 \cos \omega t$$

Prob A stuntman likes to cross the ditch as shown in figure. Find the maxth velocity reqd at point P.
 4m Determine the direction & magnitude of velocity of the stuntman just at the instant of clearing ditch.
 Let, V = Velo. at point P
 ie, $Vt = 7$ and $\frac{1}{2}gt^2 = 4$

At clearing the ditch,

$$V_y = \sqrt{2gx} = \sqrt{2 \times 9.81 \times 4} = 8.858 \text{ m/s}$$

At point P ..

$$V = \sqrt{V_x^2 + V_y^2} = \sqrt{7.75^2 + 8.858^2} = 11.77 \text{ m/s} = 48.8^\circ \text{ to } V_x$$

Prob Show that curvature of plane curve at any point can be expressed as .

$$\frac{1}{r} = \frac{\ddot{x}\dot{y} - \ddot{y}\dot{x}}{(x^2 + y^2)^{3/2}}$$

We know,

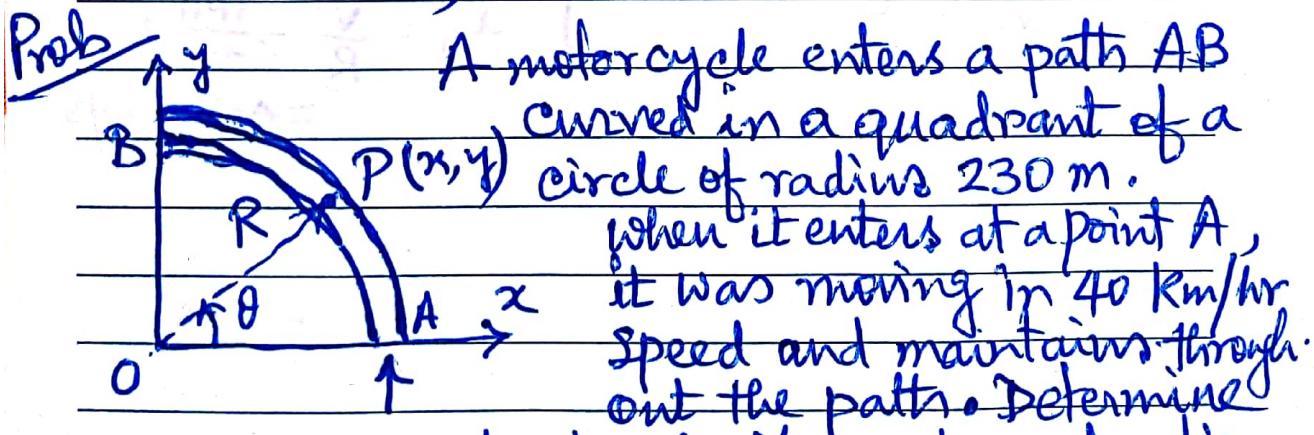
$$\frac{1}{r} = \frac{\dot{x}^2}{[1 + (\frac{dy}{dx})^2]^{3/2}} \quad (1)$$

$$\frac{dy}{dx} = \frac{dy}{dt} \cdot \frac{dt}{dx} = \frac{(\frac{dy}{dt})}{(\frac{dx}{dt})} = \frac{\dot{y}}{\dot{x}}$$

$$\frac{d^2y}{dx^2} = \frac{d}{dx} \left(\frac{dy}{dx} \right) = \frac{d}{dt} \left(\frac{dy}{dx} \right) \cdot \frac{dt}{dx} = \frac{d}{dt} \left(\frac{\dot{y}}{\dot{x}} \right) \cdot \frac{1}{\dot{x}}$$

Substituting in eq(1)

$$\begin{aligned} \frac{1}{r} &= \frac{\ddot{x}\dot{y} - \ddot{y}\dot{x}}{\dot{x}^3} \times \frac{1}{\left(1 + \frac{\dot{y}^2}{\dot{x}^2}\right)^{3/2}} \\ &= \frac{\ddot{x}\dot{y} - \ddot{y}\dot{x}}{\dot{x}^3} \\ &= \frac{\ddot{x}\dot{y} - \ddot{y}\dot{x}}{(x^2 + y^2)^{3/2}} \end{aligned}$$



Express the velocity and acceleration in terms of n-t coordinate system.

Linear velocity of motorcycle $V = 40 \text{ Km/hr.}$

$$= \frac{40 \times 1000}{3600}$$

$$= 11.11 \text{ m/s.}$$

Angular velocity

$$\omega = \frac{V}{R} = \frac{11.11}{230} = 0.0483 \text{ rad/s.}$$

Let's assume after t sec the motorcycle reaches at point 'P' subtending angle $\theta = \omega t$

$$x = R \cos \theta = R \cos \omega t$$

$$y = R \sin \theta = R \sin \omega t$$

$$\dot{x} = -R \omega \sin \omega t$$

$$\dot{y} = R \omega \cos \omega t$$

$$\ddot{x} = -R \omega^2 \cos \omega t$$

$$\ddot{y} = -R \omega^2 \sin \omega t$$

Putting $R = 230$, $\omega = 0.0483 \text{ rad/s}$, $t = 17 \text{ sec}$ we get.

$$\dot{x} = -0.1592 \text{ m/s}$$

$$\dot{y} = 11.1078 \text{ m/s.}$$

$$\ddot{x} = -0.5365 \text{ m/s}^2$$

$$\ddot{y} = -7.689 \times 10^{-3} \text{ m/s}^2$$

In n-t co-ordinate system

$$\left. \begin{array}{l} V_t = v = 11.11 \text{ m/s} \\ V_n = 0 \end{array} \right\} \quad \begin{array}{l} a_t = \frac{dv}{dt} = 0 \\ a_n = \frac{v^2}{R} = \frac{11.11^2}{230} \\ \qquad \qquad \qquad = 0.536 \text{ m/s}^2 \end{array}$$

pounce (P)

florence (C)

jounce (S)

jerk (J)

Acc² (a)

velo. (v)

displ. (x)

absement (A)

absity (B)

abseleration (C)

abserte (D)

$$\left(\frac{d}{dt} \right)$$

$$\int dt$$

Power (P)

Workdone (W)

Actergy (X)

Energy (E)

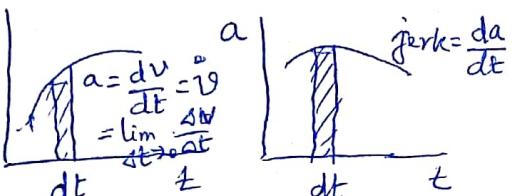
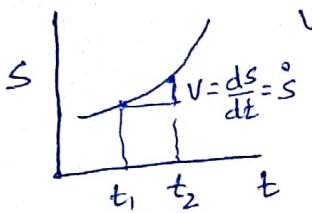
Power (P)

Workdone (W)

Actergy (X)

Energy (E)

Rectilinear Motion:



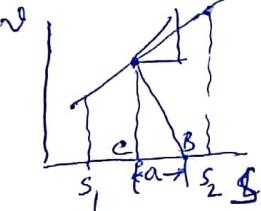
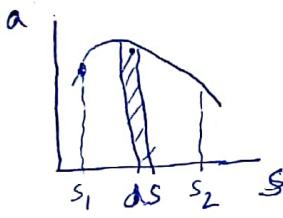
$$v = \frac{ds}{dt} = \ddot{s}$$

$$a = \frac{dv}{dt} = \ddot{v} = \ddot{\dot{s}}$$

$$= \frac{d^2s}{dt^2} = \frac{dv}{dt} = \frac{dv}{ds} \cdot \frac{ds}{dt} = v \frac{dv}{ds}$$

$$vdv = ads$$

$$\ddot{s} ds = \ddot{\dot{s}} ds$$



$$\frac{ds}{dt} = v$$

$$ds = v dt$$

$$\int_{s_1}^{s_2} ds = \int_{t_1}^{t_2} v dt$$

$(s_2 - s_1)$ = Area Under v-t curve

$$\frac{dv}{dt} = a$$

$$v_2 dt = a dt$$

$$\int_{v_1}^{v_2} dv = \int_{t_1}^{t_2} a dt$$

$(v_2 - v_1)$ = Area Under a-t curve

$$v_2 \boxed{vdv = ads}$$

$$\int_{v_1}^{v_2} v dv = \int_{s_1}^{s_2} ads$$

$\frac{1}{2}(v_2^2 - v_1^2)$ = Area Under a-s curve

Case I: Constant Accn.

$$\frac{dv}{dt} = a$$

$$vdv = a dt$$

$$\int_{v_0}^v dv = a \int_{t_0}^t dt$$

$$v - v_0 = at$$

$$v = v_0 + at$$

$$vdv = ads$$

$$\int v dv = \int ads$$

$$\frac{1}{2}(v^2 - v_0^2) = a(s - s_0)$$

$$v^2 = v_0^2 + 2a(s - s_0)$$

$$\frac{ds}{dt} = v_0 + at$$

$$ds = (v_0 + at) dt$$

$$\int_{s_0}^s ds = \int_0^t (v_0 + at) dt$$

$$(s - s_0) = v_0 t + \frac{1}{2} a t^2$$

$$s = s_0 + v_0 t + \frac{1}{2} a t^2$$

Case II: $a = f(t)$

$$f(t) = \frac{dv}{dt}$$

$$dv = f(t) \cdot dt$$

$$\int_{v_0}^v dv = \int_0^t f(t) dt$$

$$v - v_0 = \int_0^t f(t) dt$$

$$v = v_0 + \int_0^t f(t) dt$$

Case III: $a = f(v)$

$$f(v) = \frac{dv}{dt}$$

$$dt = \frac{dv}{f(v)}$$

$$t = \int_{v_0}^v \frac{dv}{f(v)} = t$$

$$vdv = ads$$

$$vdv = f(v) \cdot ds$$

$$\frac{vdv}{f(v)} = ds$$

$$\int_{v_0}^v \frac{vdv}{f(v)} = \int_{s_0}^s ds$$

$$v - v_0 = \int_{v_0}^v \frac{vdv}{f(v)}$$

$$s = s_0 + \int_{v_0}^v \frac{vdv}{f(v)}$$

Case IV: $a = f(s)$

$$vdv = ads$$

$$vdv = f(s) ds$$

$$\int_{v_0}^v v dv = \int_{s_0}^s f(s) ds$$

$$\frac{1}{2}(v^2 - v_0^2) = \int_{s_0}^s f(s) ds$$

$$v^2 = v_0^2 + 2 \int_{s_0}^s f(s) ds$$

If: $v = g(s)$

$$\frac{ds}{dt} = v = g(s)$$

$$\frac{ds}{g(s)} = dt$$

$$t = \int_{s_0}^s \frac{ds}{g(s)}$$

$$s = s_0 + \int_{s_0}^s \frac{ds}{g(s)}$$

$$t = \int_{s_0}^s \frac{ds}{g(s)}$$

$$s = s_0 + \int_{s_0}^s \frac{ds}{g(s)}$$

$$t = \int_{s_0}^s \frac{ds}{g(s)}$$

$$s = s_0 + \int_{s_0}^s \frac{ds}{g(s)}$$

$$t = \int_{s_0}^s \frac{ds}{g(s)}$$

$$s = s_0 + \int_{s_0}^s \frac{ds}{g(s)}$$

$$s = ut + \frac{1}{2} \frac{f t^2}{g}$$

$$v = u + gt$$

$$v^2 = u^2 + 2gs$$

$$[f = g]$$

when gravitational accn is considered)

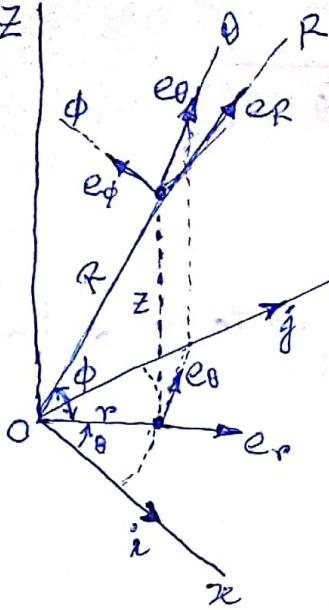
Example - SHM

$$x = a \cos \omega t$$

$$\dot{x} = -a \omega \sin \omega t = v$$

$$\ddot{x} = -a \omega^2 \cos \omega t = f = -a \omega^2 x$$

Space Curvilinear Motion: (a) Rectangular Co-ordinate System (x-y-z)

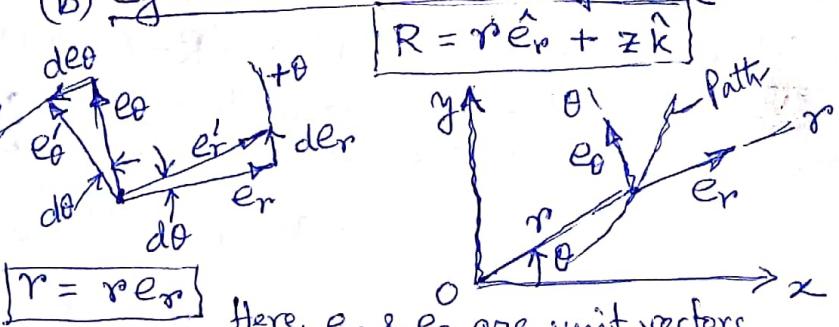


$$R = x \hat{i} + y \hat{j} + z \hat{k} \quad [R = \text{Position vector}]$$

$$v = \dot{R} = \dot{x} \hat{i} + \dot{y} \hat{j} + \dot{z} \hat{k}$$

$$a = \ddot{v} = \ddot{R} = \ddot{x} \hat{i} + \ddot{y} \hat{j} + \ddot{z} \hat{k}$$

(b) Cylindrical Co-ordinate System (r-θ-z)



$$r = r e_r$$

Here, e_r & e_θ are unit vectors

$$de_r = e_\theta \cdot d\theta$$

$$\frac{de_r}{d\theta} = e_\theta$$

$$\therefore \frac{de_r}{dt} = \left(\frac{d\theta}{dt} \right) e_\theta$$

$$\therefore \dot{e}_r = \dot{\theta} e_\theta$$

$$de_\theta = -e_r d\theta$$

$$\frac{de_\theta}{d\theta} = -e_r$$

$$\frac{de_\theta}{dt} = -\left(\frac{d\theta}{dt} \right) e_r$$

$$\dot{e}_\theta = -\dot{\theta} e_r$$

$$v = \dot{r} e_r + r \dot{\theta} e_\theta$$

$$[v = \dot{r} e_r + r \dot{\theta} e_\theta] \quad [\because \dot{e}_r = \dot{\theta} e_\theta] \quad v = V_r \hat{e}_r + V_\theta \hat{e}_\theta$$

Where, $V_r = \dot{r}$ } $V = \sqrt{V_r^2 + V_\theta^2}$
 $V_\theta = r \dot{\theta}$

$$a = \ddot{v} = (\ddot{r} e_r + \dot{r} \dot{e}_r) + (\dot{r} \dot{\theta} e_\theta + r \ddot{\theta} e_\theta + r \dot{\theta} \dot{e}_\theta)$$

Putting $\dot{e}_r = \dot{\theta} e_\theta$ and $\dot{e}_\theta = -\dot{\theta} e_r$, we get

$$a = (\ddot{r} - r \dot{\theta}^2) e_r + (r \ddot{\theta} + 2 \dot{r} \dot{\theta}) e_\theta \quad \text{ie, } a = a_r \hat{e}_r + a_\theta \hat{e}_\theta$$

Where, $a_r = \ddot{r} - r \dot{\theta}^2$ } $a = \sqrt{a_r^2 + a_\theta^2}$
 $a_\theta = r \ddot{\theta} + 2 \dot{r} \dot{\theta}$

So, $R = r e_r + z \hat{k}$

$$v = \dot{R} = \dot{r} e_r + r \dot{\theta} e_\theta + \dot{z} \hat{k}$$

Where, $V_r = \dot{r}$
 $V_\theta = r \dot{\theta}$
 $V_z = \dot{z}$

$$a = \ddot{R} = \dot{V} = (\ddot{r} - r\dot{\theta}^2) \hat{e}_r + (r\ddot{\theta} + 2r\dot{\theta}) \hat{e}_{\theta} + \ddot{z} \hat{k}$$

$$\left. \begin{aligned} a_r &= \ddot{r} - r\dot{\theta}^2 \\ a_{\theta} &= r\ddot{\theta} + 2r\dot{\theta} \\ a_z &= \ddot{z} \end{aligned} \right] \quad a = \sqrt{a_r^2 + a_{\theta}^2 + a_z^2}$$

(c) Spherical Co-ordinate System ($R - \theta - \phi$)

When a radial distance and two angles are used to specify the position of a particle as in radar measurement.

Here, $\boxed{V = v_R \hat{e}_R + v_{\theta} \hat{e}_{\theta} + v_{\phi} \hat{e}_{\phi}}$

Where, $\hat{e}_R, \hat{e}_{\theta}$ and \hat{e}_{ϕ} are unit vectors

$$v_R = \dot{R}$$

$$v_{\theta} = R\dot{\theta} \cos\phi$$

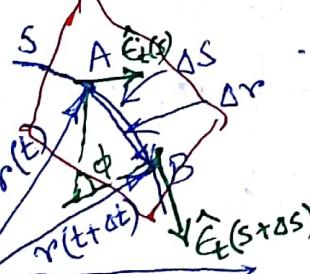
$$v_{\phi} = R\dot{\phi}$$

Similarly,

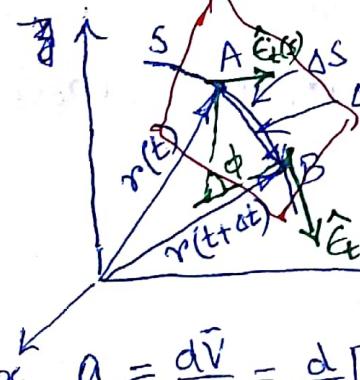
$$\boxed{a = a_R \hat{e}_R + a_{\theta} \hat{e}_{\theta} + a_{\phi} \hat{e}_{\phi}}$$

where, $a_R = \ddot{R} - R\dot{\phi}^2 - R\dot{\theta}^2 \cos\phi$

$$a_{\theta} = \frac{\cos\phi}{R} \frac{d}{dt}(R^2 \dot{\theta}) - 2R\dot{\theta}\dot{\phi} \sin\phi$$

Osculating Plane 

$$a_{\phi} = \frac{1}{R} \frac{d}{dt}(R^2 \dot{\phi}) + R\dot{\theta}^2 \sin\phi \cos\phi$$



$V = \frac{d\bar{r}}{dt} = \lim_{\Delta t \rightarrow 0} \frac{r(t+\Delta t) - r(t)}{\Delta t} = \lim_{\Delta t \rightarrow 0} \frac{\Delta r}{\Delta t}$

$= \lim_{\Delta t \rightarrow 0} \left[\left(\frac{\Delta r}{\Delta s} \right) \left(\frac{\Delta s}{\Delta t} \right) \right] = \left(\frac{ds}{dt} \right) \hat{E}_t$

where \hat{E}_t = Unit vector with dir^n along tangent \Rightarrow Speed

$$a = \frac{d\bar{V}}{dt} = \frac{d}{dt} \left[\left(\frac{ds}{dt} \right) \hat{E}_t \right] = \frac{d^2 s}{dt^2} \hat{E}_t + \left(\frac{ds}{dt} \right) \left(\frac{d\hat{E}_t}{dt} \right) = \frac{d^2 s}{dt^2} \hat{E}_t + \left(\frac{ds}{dt} \right) \left(\frac{d\hat{E}_t}{ds} \right) \left(\frac{ds}{dt} \right)$$

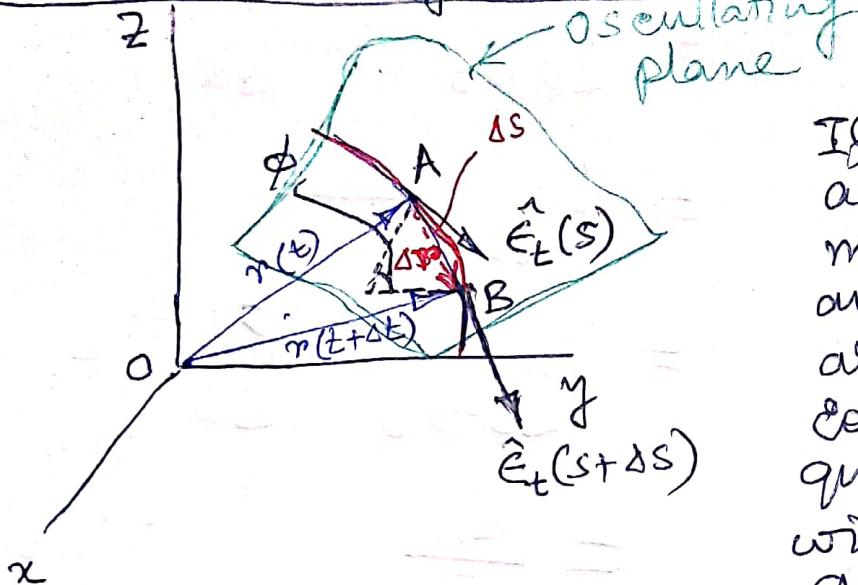
$$= \frac{d^2 s}{dt^2} \hat{E}_t + \left(\frac{ds}{dt} \right)^2 \left(\frac{d\hat{E}_t}{ds} \right)$$

$$a = \left(\frac{d^2 s}{dt^2} \right) \hat{E}_t + \left(\frac{ds}{dt} \right)^2 \frac{1}{R} \hat{n}$$

$R = \sqrt{1 + \dot{y}_t^2}^{1/2}$
 $= \text{Radius of curvature}$

$$\frac{d\hat{E}_t}{ds} = \lim_{\Delta s \rightarrow 0} \frac{\hat{E}_t(s+\Delta s) - \hat{E}_t(s)}{\Delta s} = \frac{f\Delta E_t}{\Delta s} \quad | \text{Now } \Delta E_t \approx E_t | \Delta \phi = \Delta \phi = \frac{\Delta s}{R} \quad | \text{Hence, } E_t = \frac{\Delta s}{R} \hat{E}_n$$

Normal and Tangential Reference System:



If the velocity & acc² in curvilinear motion along normal and tangential dirⁿ are needed to express, consider two subsequent points A & B with radius of curvature as 'p'.

Let, the unit vectors along the tangential & normal dirⁿ are \hat{E}_t and \hat{E}_n , respectively.

Now, the instantaneous velocity can be expressed as,

$$V = \lim_{\Delta t \rightarrow 0} \frac{\Delta r}{\Delta t} = \frac{dr}{dt} = \frac{dr}{ds} \left(\frac{ds}{dt} \right)$$

$$\lim_{\Delta t \rightarrow 0} \frac{r(t+\Delta t) - r(t)}{\Delta t} = \left(\frac{ds}{dt} \right) \cdot \hat{E}_t$$

In this case, the normal component of velocity is zero i.e., $\hat{E}_n = 0$

Now, the instantaneous acceleration

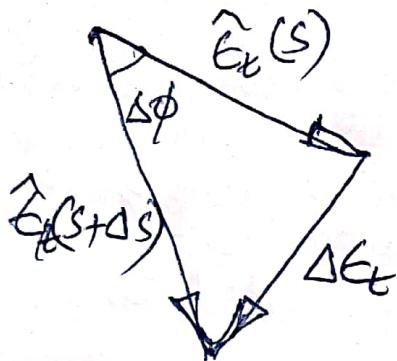
$$a = \frac{dV}{dt} = \frac{d}{dt} \left[\left(\frac{ds}{dt} \right) \hat{E}_t \right]$$

$$= \frac{d^2 s}{dt^2} \cdot \hat{E}_t + \left(\frac{ds}{dt} \right) \left(\frac{d \hat{E}_t}{dt} \right)$$

$$= \frac{d^2 s}{dt^2} \cdot \hat{E}_t + \left(\frac{ds}{dt} \right) \left(\frac{ds}{dt} \right) \left(\frac{d \hat{E}_t}{ds} \right)$$

$$= \left(\frac{d^2 s}{dt^2} \right) \hat{E}_t + \left(\frac{ds}{dt} \right)^2 \cdot \left(\frac{d \hat{E}_t}{ds} \right)$$

Now,



$$\frac{dE_t}{ds} = \lim_{\Delta s \rightarrow 0} \frac{E_t(s+\Delta s) - E_t(s)}{\Delta s}$$
$$= \lim_{\Delta s \rightarrow 0} \frac{\Delta E_t}{\Delta s}$$
$$= \lim_{\Delta s \rightarrow 0} \frac{(\Delta s/R) E_n}{\Delta s}$$

Now,

$$\Delta E_t \approx |E_t| \Delta \phi$$

$$\approx \Delta \phi$$

$$= \frac{\Delta s}{R}$$

Hence, $\hat{E}_t = \frac{\Delta s}{R} \cdot \hat{E}_n$.

$$= \frac{E_n}{R}$$

[Putting the value of eqz(A)]

Thus,

$$a = \left(\frac{d^2 s}{dt^2} \right) \cdot \hat{E}_t + \left(\frac{ds}{dt} \right)^2 \cdot \left(\frac{dE_t}{ds} \right)$$

$$= \left(\frac{d^2 s}{dt^2} \right) \hat{E}_t + \left(\frac{ds}{dt} \right)^2 \frac{1}{R} \cdot \hat{E}_n$$

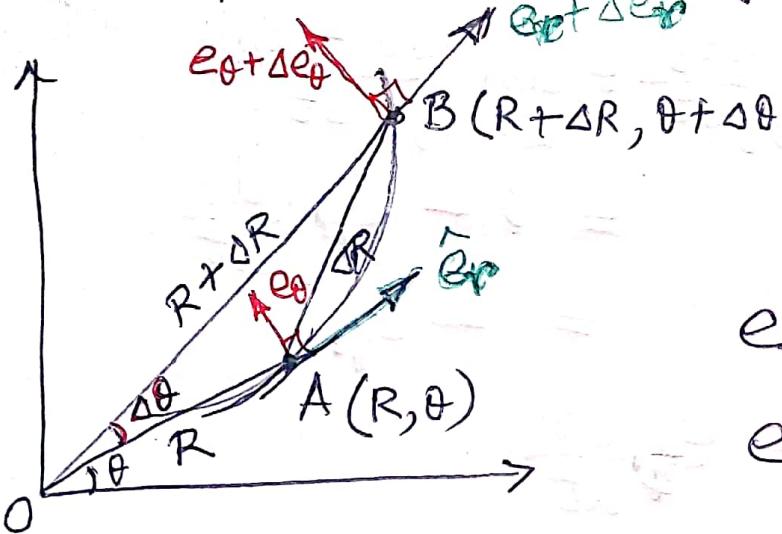
Where, $\overset{at}{P}$ = Radius of curvature

$$= \frac{(1 + y_1'^2)^{3/2}}{y_2}$$

a_n and a_t are the normal acceleration and tangential acceleration components, respectively.

Velocity and Acceleration in Polar reference System :

In polar coordinate system, initial position of a particle at



of a particle at $A(R, \theta)$ and after time(Δt) it reaches to point $B(R+\Delta R, \theta+\Delta\theta)$

e_r → Unit vector at radial dir^o

e_θ → Unit vector at θ -dir^o.

from geometry, we can write

$$e_{\theta} = \Delta\theta e_\theta \quad \text{and} \quad \Delta e_\theta = \Delta\theta (-e_r)$$

The position vector of points A and B are.

At point A
 $R = (r e_r)$

At point B
 $R + \Delta R = (r + \Delta r) (e_r + \Delta e_r)$

or, $R + \Delta R = r e_r + r \Delta e_r + \Delta r e_r + \Delta r \Delta e_r$ (as very small)

or, $R + \Delta R = (r e_r) + r \Delta e_r + e_r \Delta r$

or, $R + \Delta R = R + r \Delta e_r + e_r \Delta r$

or, $\Delta R = r \Delta e_r + e_r \Delta r$

The instantaneous velocity

$$V = \lim_{\Delta t \rightarrow 0} \frac{\Delta R}{\Delta t} = \lim_{\Delta t \rightarrow 0} \frac{r \Delta e_r + e_r \Delta r}{\Delta t} = r \frac{de_r}{dt} + e_r \frac{dr}{dt}$$

$$\text{Now, } \dot{\vec{e}_\theta} = \frac{d\vec{e}_\theta}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \vec{e}_\theta}{\Delta t} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \theta \vec{e}_\theta}{\Delta t} = \vec{e}_\theta \frac{d\theta}{dt} \\ = \vec{e}_\theta \cdot \dot{\theta}$$

$$\dot{\vec{e}_r} = \frac{d\vec{e}_r}{dt} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \vec{e}_r}{\Delta t} = \lim_{\Delta t \rightarrow 0} \frac{\Delta \theta (-\vec{e}_\theta)}{\Delta t} \\ \leq \dot{\theta} - \vec{e}_\theta \cdot \frac{d\theta}{dt} \\ = -\vec{e}_\theta \cdot \dot{\theta}$$

Finally,
velocity

$$V = r \frac{d\vec{e}_r}{dt} + \vec{e}_\theta \frac{dr}{dt} \\ = \underbrace{r \dot{\theta} \vec{e}_\theta}_{V_\theta} + \underbrace{r \dot{r} \vec{e}_r}_{V_r}$$

Acceleration

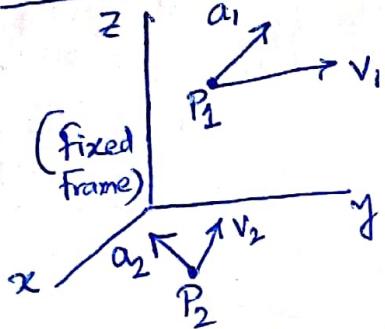
$$a = \frac{dV}{dt} = \frac{d}{dt} (r \dot{\theta} \vec{e}_\theta + r \dot{r} \vec{e}_r) \\ = \dot{r} \vec{e}_\theta + r(\dot{\theta}) \vec{e}_\theta + r \dot{r} \vec{e}_r + r \ddot{r} \vec{e}_r \\ = \underbrace{\dot{r} \vec{e}_\theta}_{a_r} + \underbrace{r \dot{\theta} (\dot{\theta} \vec{e}_\theta)}_{a_\theta} + r \ddot{\theta} \vec{e}_\theta \\ + \underbrace{\dot{r} \vec{e}_r}_{a_r} + \underbrace{r(\dot{\theta}) \vec{e}_r}_{a_\theta} \\ = \left(\ddot{r} - r \dot{\theta}^2 \right) \vec{e}_r + \left(r \ddot{\theta} + 2r \dot{\theta} \right) \vec{e}_\theta$$

Case 1: When rotation is in constant angular speed,
Here, $\dot{\theta} = \text{constant}$ then, $\ddot{\theta} = 0$

Case 2: When rotation is a plane circular motion, Here,
 $r = \text{constant}$, i.e., $\dot{r} = \ddot{r} = 0$ then $V = V_\theta = (r \dot{\theta}) \vec{e}_\theta$.
and $a = a_r + a_\theta$ where, $a_r = -r \dot{\theta}^2 \vec{e}_r$ and
 $a_\theta = r \ddot{\theta} \vec{e}_\theta$.

Motion referred to moving frames of reference:

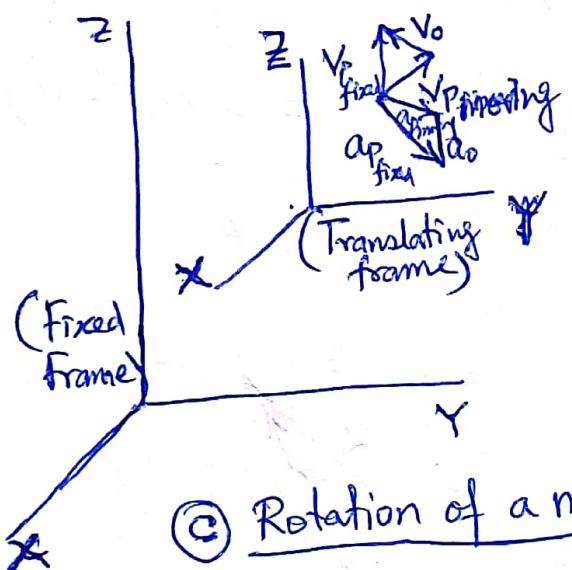
(a) Relative motion of two points:



$$v_{12} = v_1 - v_2 \quad \text{and} \quad v_{21} = v_2 - v_1 = -v_{12}$$

$$a_{12} = a_1 - a_2 \quad \text{and} \quad a_{21} = a_2 - a_1 = -a_{12}$$

(b) Translation of moving frame:



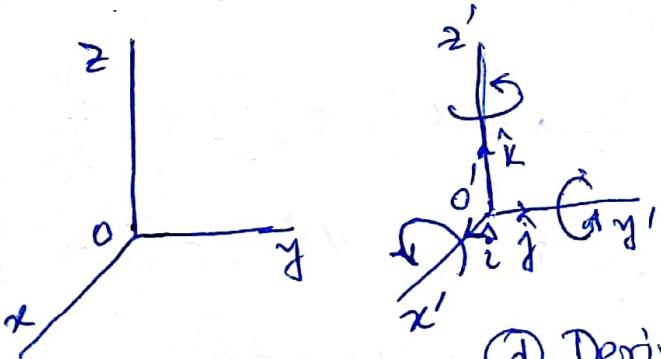
$$v_0 = \text{Velocity of moving frame}$$

$$a_0 = \text{Accel of "}$$

$$v_{P\text{fixed}} = v_{P\text{moving}} + v_0$$

$$a_{P\text{fixed}} = a_{P\text{moving}} + a_0$$

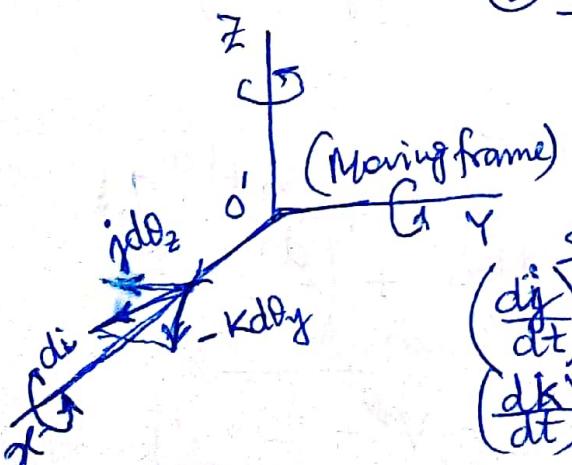
(c) Rotation of a moving frame:



$$\omega = \omega_x \hat{i} + \omega_y \hat{j} + \omega_z \hat{k}$$

$$= \frac{d\theta_x}{dt} \hat{i} + \frac{d\theta_y}{dt} \hat{j} + \frac{d\theta_z}{dt} \hat{k}$$

(d) Derivatives of moving unit vector:



$$\left(\frac{di}{dt} \right)_f = j \frac{d\theta_z}{dt} - k \frac{d\theta_y}{dt}$$

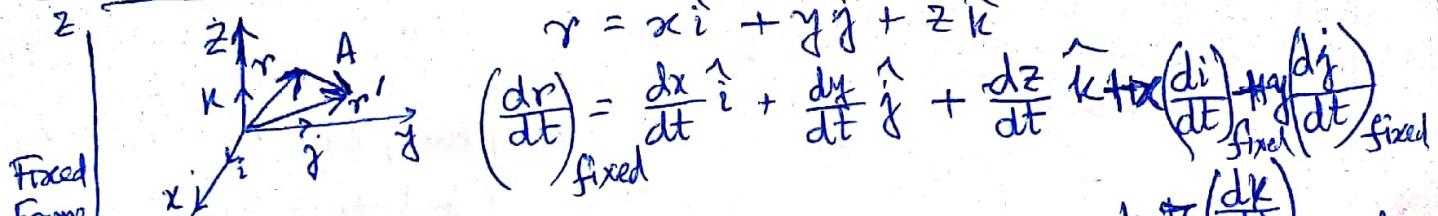
$$= j \omega_z - k \omega_y$$

$$\left(\frac{dj}{dt} \right)_f = \omega \times j$$

$$\left(\frac{dk}{dt} \right)_f = \omega \times k$$

Similarly,

② Derivative of Constant vector in moving frame:



$$r = xi\hat{i} + yj\hat{j} + zk\hat{k}$$

$$\left(\frac{dr}{dt}\right)_{\text{fixed}} = \frac{dx}{dt}\hat{i} + \frac{dy}{dt}\hat{j} + \frac{dz}{dt}\hat{k} \cancel{+ \omega \times \left(\frac{dr}{dt}\right)_{\text{fixed}}}$$

$$\left(\frac{dr}{dt}\right)_{\text{moving}} = \frac{dx}{dt}\hat{i} + \frac{dy}{dt}\hat{j} + \frac{dz}{dt}\hat{k} = 0$$

$$\begin{aligned} \left(\frac{dr}{dt}\right)_{\text{fixed}} &= x(\omega \times i) + y(\omega \times j) + z(\omega \times k) \\ &= \omega \times (xi + yj + zk) = \omega \times r \end{aligned}$$

Similarly $\left(\frac{dr'}{dt}\right)_{\text{fixed}} = \omega \times r'$ $\therefore \left\{ \frac{d}{dt}(r' - r) \right\}_{\text{fixed}} = \omega \times (r' - r)$

$$\left(\frac{dA}{dt}\right)_{\text{fixed}} = \omega \times A.$$

③ Derivative of position vector for Different reference:

$$r = xi\hat{i} + yj\hat{j} + zk\hat{k}$$

$$R = R_0 + r$$

$$\left(\frac{dr}{dt}\right)_{\text{moving}} = \frac{dx}{dt}\hat{i} + \frac{dy}{dt}\hat{j} + \frac{dz}{dt}\hat{k}$$

$$\begin{aligned} \left(\frac{dr}{dt}\right)_{\text{fixed}} &= \frac{dx}{dt}\hat{i} + \frac{dy}{dt}\hat{j} + \frac{dz}{dt}\hat{k} + x \cdot \frac{di}{dt} + y \cdot \frac{dj}{dt} + z \cdot \frac{dk}{dt} \\ &= \frac{dx}{dt}\hat{i} + \frac{dy}{dt}\hat{j} + \frac{dz}{dt}\hat{k} + x(\omega_x i) + y(\omega_y j) + z(\omega_z k) \end{aligned}$$

$$= \left(\frac{dr}{dt}\right)_{\text{moving}} + \omega \times (xi + yj + zk)$$

$$= \left(\frac{dr}{dt}\right)_{\text{moving}} + \omega \times r$$

$$v_p = v_o + \left[\left(\frac{dr}{dt}\right)_{\text{moving}} + \omega \times r \right]$$

④ Velocity of a point:

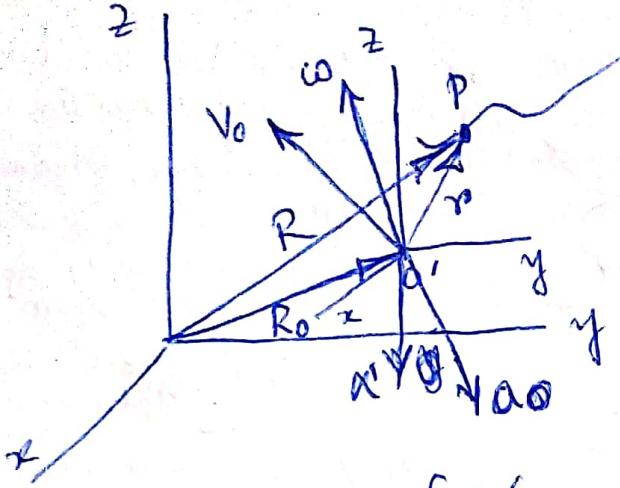
$$v_{P_{\text{fixed}}} = \left(\frac{dR}{dt}\right)_{\text{fixed}}$$

$$v_{P_{\text{moving}}} = \left(\frac{dr}{dt}\right)_{\text{moving}}$$

$$\left(\frac{dR}{dt}\right)_{\text{fixed}} = \left(\frac{dR_0}{dt}\right)_{\text{fixed}} + \left(\frac{dr}{dt}\right)_{\text{fixed}}$$

$$v_{P_{\text{fixed}}} = v_{P_{\text{moving}}} + v_o + \omega \times r$$

⑥ Acceleration of a point



$$a_{P \text{fixed}} = \left[\frac{d}{dt} (V_{P \text{fixed}}) \right]_{\text{fixed}}$$

$$a_{P \text{moving}} = \left[\frac{d}{dt} (V_{P \text{moving}}) \right]_{\text{moving}}$$

$$= \left[\frac{d^2 r}{dt^2} \right]_{\text{moving}}$$

$$a_{P \text{fixed}} = \left[\frac{d}{dt} (V_{P \text{fixed}}) \right]_{\text{fixed}}$$

$$= a_0 + \left[\frac{d}{dt} (V_{P \text{moving}}) \right]_{\text{fixed}} + \left[\frac{d}{dt} (\omega \times r) \right]_{\text{fixed}}$$

$$= \left[\frac{d}{dt} (V_{P \text{moving}}) \right]_{\text{moving}} + \omega \times V_{P \text{moving}}$$

$$= a_{P \text{moving}} + \omega \times V_{P \text{moving}}$$

$$= \omega \times \left(\frac{dr}{dt} \right)_{\text{fixed}} + \left(\frac{d\omega}{dt} \right) \times r_{\text{fixed}}$$

$$= \omega \times \left[\left(\frac{dr}{dt} \right)_{\text{moving}} + \omega \times r \right] + \alpha \times r$$

$$= \omega \times [V_{P \text{moving}} + \omega \times r] + \alpha \times r$$

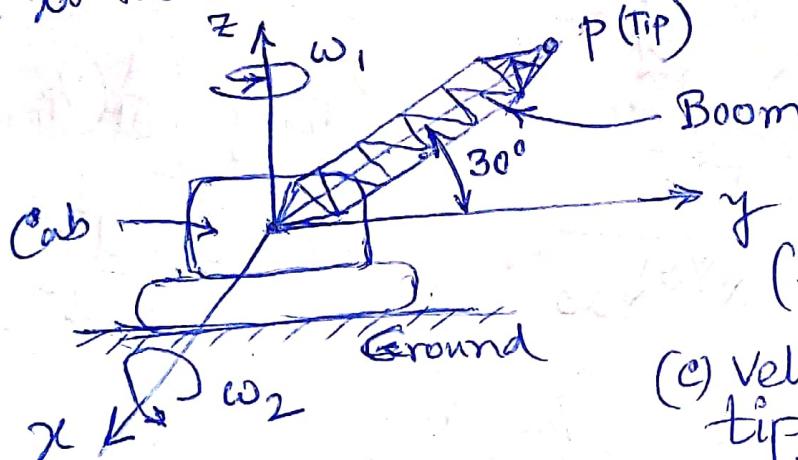
$$= \omega \times V_{P \text{moving}} + \omega \times (\omega \times r) + \alpha \times r$$

Acceleration of a point P
w.r.t fixed frame of reference

$$a_{P \text{fixed}} = a_0 + a_{P \text{moving}} + \alpha(\alpha \times r) + \omega \times (\omega \times r) + 2\omega \times V_{P \text{moving}}$$

↓
 Translational accⁿ Acc of pt P w.r.t moving frame of reference acc due to moving frame accelerating with alpha normal component of accⁿ (centripetal accⁿ) Coriolis component of accⁿ

Problem A crane rotates about a vertical axis with a constant angular velocity of 0.4 rad/s . While the boom is being raised with a constant angular velocity of 0.5 rad/s relative to the cab as shown in figure. If the length of boom is 10m , determine



- (a) Angular velocity of the boom
- (b) Angular accⁿ of the boom
- (c) Velocity of the tip of the boom

Hints: $\gamma = 10(\cos 30\hat{i} + \sin 30\hat{k})$ (d) Accⁿ of the tip of the boom

$$\omega_{\text{boom}} = \omega_1 + \omega_2$$

$$= 0.4\hat{k} + 0.5\hat{i} \text{ rad/s}$$

$w_1 \rightarrow$ of Cab
 $w_2 \rightarrow$ of boom w.r.t Cab

$$\alpha_{\text{boom}} = \left(\frac{d\omega_{\text{boom}}}{dt} \right)_f = \left(\frac{d\omega_1}{dt} \right)_f + \left(\frac{d\omega_2}{dt} \right)_f$$

Now, $\left(\frac{d\omega_1}{dt} \right)_f = 0$ as cab rotates at constant angular veloc.

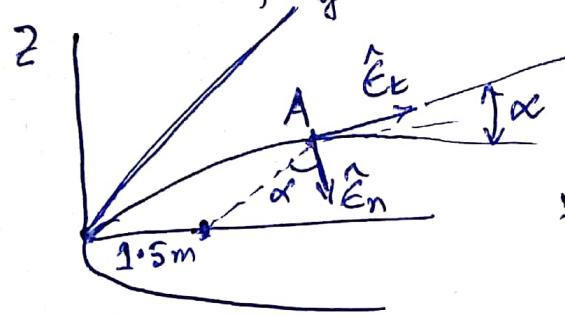
$$\alpha_{\text{boom}} = \left(\frac{d\omega_2}{dt} \right)_f = \omega_1 \times \omega_2 = 0.4\hat{k} \times 0.5\hat{i} \\ = 0.2\hat{j} \text{ rad/s}^2$$

$$\text{Velocity of Tip}(P) (V_P) = \omega \times \gamma \\ = (0.4\hat{k} + 0.5\hat{i}) \times (0.865\hat{j} + 0.5\hat{k}) \\ = -3.465\hat{i} - 2.5\hat{j} + 4.33\hat{k} \text{ m/s}$$

Acceleration of Tip(P)

$$\alpha_p = \alpha_{\text{boom}} \times \gamma + \omega_{\text{boom}} \times (\omega_{\text{boom}} \times \gamma) \\ = 0.2\hat{j} \times (0.865\hat{j} + 0.5\hat{k}) + (0.4\hat{k} + 0.5\hat{i}) \times [(0.4\hat{k} + 0.5\hat{i}) \times (0.865\hat{j} + 0.5\hat{k})] \\ = 2\hat{i} - 3.55\hat{j} - 1.25\hat{k} \text{ m/s}^2$$

Problem A particle is moving in the xy plane along a parabolic path on $y = 1.22\sqrt{x}$. At position A , ($x=1.5 \text{ m}$) the particle has a speed of 3 m/sec along the path. What is the acceleration vector of the particle at this position?



Now,

$$\vec{e}_t = \cos \alpha \hat{i} + \sin \alpha \hat{j}$$

$$\text{Here, } \tan \alpha = \frac{dy}{dx} = \frac{d}{dx}(1.22\sqrt{x}) \\ = \frac{0.610}{\sqrt{x}}$$

$$\text{At } x = 1.5 \text{ m,}$$

$$\tan \alpha = \frac{0.610}{\sqrt{1.5}} = \frac{1}{2}$$

$$\therefore \alpha = 26.5^\circ$$

Hence,

$$\vec{e}_t = 0.895 \hat{i} + 0.446 \hat{j} \quad \dots \dots \dots (1)$$

$$\text{And from diagram, } \vec{e}_n = \sin \alpha \hat{i} - \cos \alpha \hat{j}$$

$$\therefore \vec{e}_n = 0.446 \hat{i} - 0.895 \hat{j} \quad \dots \dots \dots (2)$$

Now,

$$y = 1.22\sqrt{x}$$

$$y_1 = \frac{dy}{dx} = 0.61x^{-\frac{1}{2}}$$

$$y_2 = \frac{d^2y}{dx^2} = -0.305x^{-\frac{3}{2}}$$

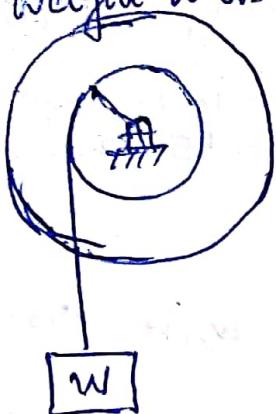
$$\therefore R = \frac{(1+y_1^2)^{\frac{3}{2}}}{y_2} = 8.4 \text{ m}$$

$$\text{We know, } \vec{a} = \left(\frac{d^2s}{dt^2} \right) \vec{e}_t + \left(\frac{ds}{dt} \right)^2 \cdot \frac{\vec{e}_n}{R}$$

$$= 3(0.895 \hat{i} + 0.446 \hat{j}) + \frac{3^2}{8.4} (0.446 \hat{i} - 0.895 \hat{j})$$

$$\vec{a} = 3.16 \hat{i} + 0.379 \hat{j} \text{ m/s}^2$$

Problem A string is wound around a pulley of 0.28 m dia. One of its ends is fixed to the pulley while the other end is fixed to a weight W as shown in figure. This weight moves a distance of 7.7 m starting from rest in 4.3 sec.



Determine the angular speed of the pulley. Also determine the total length traversed by the weight to make the pulley rotate by 425 rpm.

$$s = ut + \frac{1}{2}at^2$$

$$7.7 = 0 \times t + \frac{1}{2}a \times 4.3^2$$

$$a = 0.833 \text{ m/s}^2$$

$$\text{Angular accrl. of pulley, } \alpha = \frac{a}{R} = \frac{0.833}{\left(\frac{0.28}{2}\right)} = 5.95 \text{ rad/s}^2$$

$$\begin{aligned} \omega &= \omega_0 + \alpha t \\ &= 0 + 5.95 \times 4.3 \\ &= 25.585 \text{ rad/s} \end{aligned}$$

$$\text{Final speed of pulley, } \omega_{\text{final}} = \frac{2\pi \times 425}{60} =$$

$$\text{Time reqd to reach the final speed, } t = \frac{\omega_{\text{final}} - \omega_0}{\alpha} = \frac{\frac{2\pi \times 425}{60} - 0}{5.95} = 7.48 \text{ sec.}$$

Total length traversed by the weight to make pulley rotate by 425 rpm

$$\begin{aligned} s &= ut + \frac{1}{2}at^2 \\ &= 0 \times t + \frac{1}{2} \times 0.833 \times 7.48^2 = 23.303 \text{ m} \end{aligned}$$

Problem A roller of radius 2.5 m rolls without slipping along a horizontal plane. The centre O has uniform velocity 25 m/s. Determine the velocity of the points A and B of the roller.

At noslip condn ($\epsilon V_c = 0$)

$$V_c = V_{co} + V_o = -r\omega + V_o = 0 \Rightarrow V_o = r\omega$$

$$\omega = \frac{V_o}{r} = \frac{25}{2.5} = 10 \text{ rad/s}$$

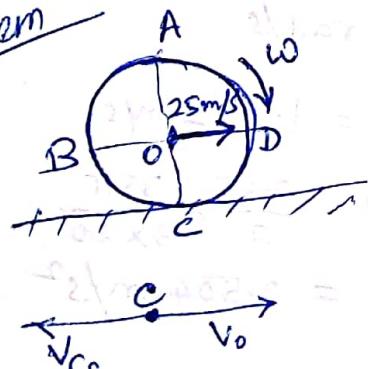
$$V_{B0} = V_{AO} = V_{DO} = \dots = r\omega$$

$$\text{At point B, } V_B = V_{B0} + V_o$$

$$V_B = \sqrt{V_{B0}^2 + V_o^2} = \sqrt{(r\omega)^2 + V_o^2} = \sqrt{(10 \times 2.5)^2 + 25^2} = 35.35 \text{ m/s}$$

$$\theta = \tan^{-1}\left(\frac{25}{10}\right) = 45^\circ$$

$$\text{At pt A, } V_A = V_{AO} + V_o = 2 \times 25 = 50 \text{ m/s.}$$

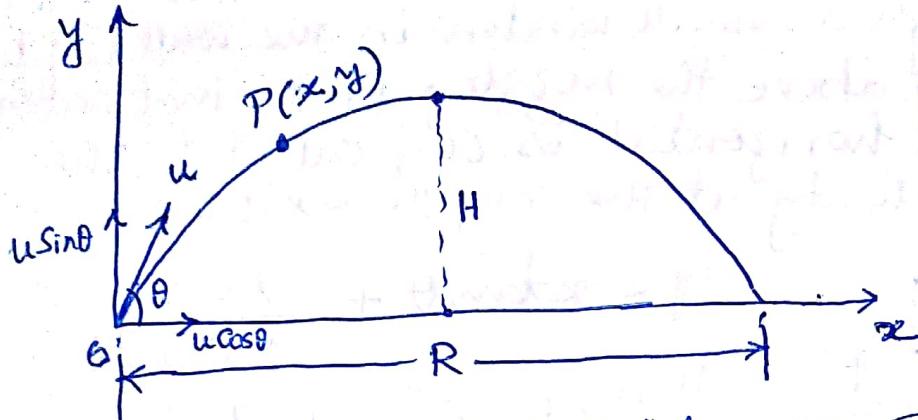


$$(a) V_{co} \quad V_o$$

$$(b) V_{B0} \quad V_B \quad V_o$$

$$(c) V_o \quad V_{AO}$$

Projectile motion :



$$x = u \cos \theta \times t \quad \text{--- (i)}$$

$$y = u \sin \theta \times t - \frac{1}{2} g t^2 \quad \text{--- (ii)}$$

$$\therefore y = (u \sin \theta) \cdot \left(\frac{x}{u \cos \theta} \right) - \frac{1}{2} g \cdot \left(\frac{x}{u \cos \theta} \right)^2$$

$$y = x \tan \theta - \frac{1}{2} g \cdot \frac{x^2}{u^2 \cos^2 \theta} \quad \text{--- (A)}$$

To find the time of flight :

For this, put $y=0$ in eqn (ii)

$$t = \frac{2u \sin \theta}{g}$$

To find the max^m height reached :

$$v^2 = u^2 - 2gH$$

$$0 = (u \sin \theta)^2 - 2gH$$

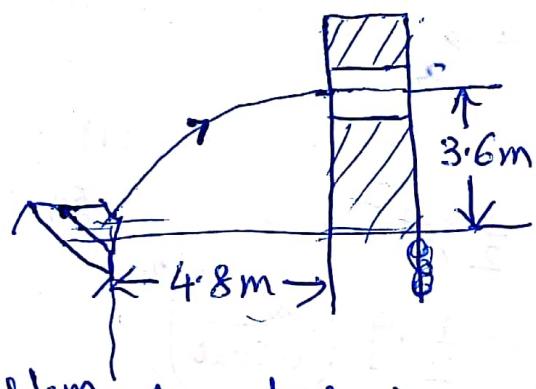
$$H = \frac{u^2 \sin^2 \theta}{2g}$$

To find the max^m horizontal range :

$$R = (u \cos \theta) \times \text{time of flight}$$

$$R = u \cos \theta \times \frac{2u \sin \theta}{g}$$

Problem A fireman holding a nozzle at a horizontal distance of 4.8 m from a vertical wall, wishes to send a jet of water through a small window in the wall located 3.6 m vertically above the nozzle. If the inclination of the jet with horizontal is 60° , calculate the required jet velocity at the nozzle exit.



$$y = xt \tan \theta + \frac{gx^2}{2u^2 \cos^2 \theta}$$

Here, $\theta = 60^\circ$, $y = 3.6 \text{ m}$.

$$x = 4.8 \text{ m}$$

Putting these values :-

$$u = 9.78 \text{ m/s.}$$

Problem A particle is projected at such an angle with the horizontal that the horizontal range is 4-times the greatest height attained by the particle. Find the angle of projection.

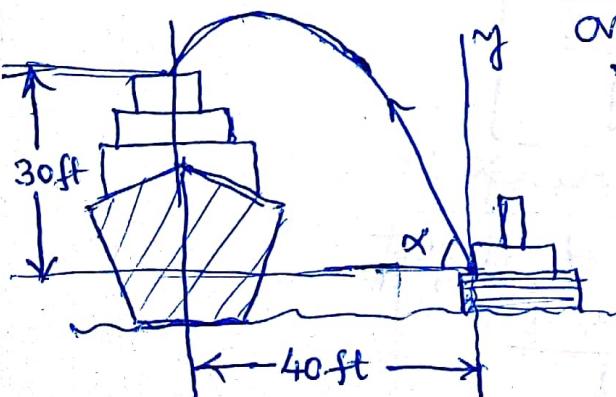
$$\text{Here, } R = 4H_{\max}$$

$$\text{w, } \frac{u^2 \sin 2\theta}{g} = 4 \cdot \frac{u^2 \sin^2 \theta}{2g}$$

$$\text{w, } 2 \sin \theta \cos \theta = 2 \sin^2 \theta$$

$$\text{w, } \tan \theta = 1 \quad \text{w, } \theta = 45^\circ$$

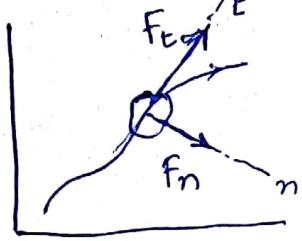
Problem The engine room of a freighter is on fire. A fire-fighter tugboat has drawn alongside and is directing a stream of water to enter the stack of the freighter as shown in figure.



If the initial speed of jet is 70 ft/sec, calculate the angle of projection of jet.

$$[\text{Ans: } \alpha = 81.37^\circ]$$

Particle Dynamics: $\mathbf{F} = F_x \hat{i} + F_y \hat{j} + F_z \hat{k}$



$$\begin{aligned} F_x &= m a_x \\ F_y &= m a_y \\ F_z &= m a_z \end{aligned}$$

$$F_t = m a_t = m \ddot{s} = m r \ddot{\theta} = m \omega^2 r$$

$$F_n = m a_n = m \frac{v^2}{r} = m r \dot{\theta}^2 = m \omega^2 r$$

$$\text{For cylindrical Co-ordinates: } \boxed{\omega^2 = \omega_0^2 + 2\dot{\theta}\theta}$$

$$\mathbf{F} = F_r \hat{e}_r + F_\theta \hat{e}_\theta + F_z \hat{e}_z$$

$$F_r = m \left(\frac{d^2 r}{dt^2} - r \dot{\theta}^2 \right)$$

$$F_\theta = m \left(r \ddot{\theta} + 2 \dot{r} \theta \right)$$

$$F_z = m z \ddot{z}$$

① Force is a function of time or constant: $F_z = m z \ddot{z}$

$$a = \frac{d^2 x}{dt^2}$$

$$F(t) = m \vec{a} = m \frac{d^2 \vec{x}}{dt^2} = m \frac{dV}{dt}$$

$$d \left(\frac{dx}{dt} \right) = F(t) dt$$

$$\text{Integrating: } v = \int F(t) dt + C_1$$

$$\text{Integrating: } x = \int \int F(t) dt + C_1 t + C_2$$

$$\text{Apply B.C.: } t=0, v=v_0 \quad (say)$$

$$t=0, x=x_0$$

② Force is a function of position:

$$m \frac{dV}{dt} = f(x)$$

$$m \frac{dV}{dt} = m \left(\frac{dV}{dx} \right) \cdot \left(\frac{dx}{dt} \right) = F(x)$$

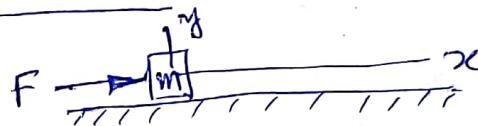
$$m V \frac{dV}{dx} = F(x)$$

Integrating:-

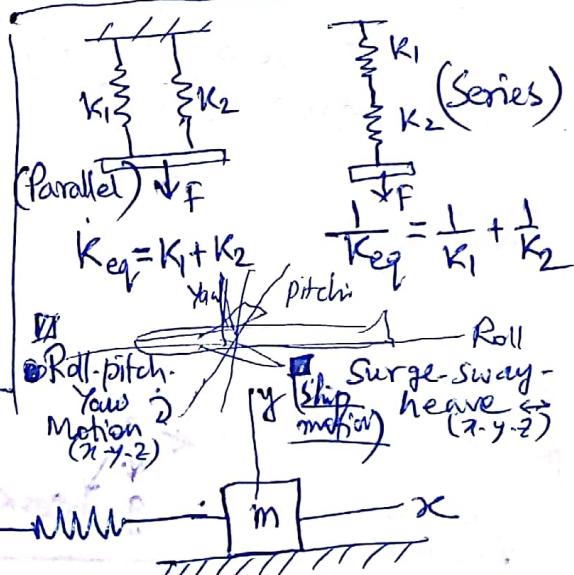
$$m \frac{V^2}{2} = \int F(x) dx + C_1$$

$$m V = \left[\frac{2}{m} \int F(x) dx + C_1 \right]^{\frac{1}{2}}$$

$$\frac{dx}{dt} = \left[\frac{2}{m} \int F(x) dx + C_1 \right]^{\frac{1}{2}}$$



Example - Traction force of automobile w.r.t time



Example - Restoring force of spring-mass system

$$F_{\text{spring}} = - \int k(x) dx$$

$$dt = \frac{dx}{\left[\frac{2}{m} \int F(x) dx + C_1 \right]^{\frac{1}{2}}}$$

$$t = \int \frac{dx}{\left[\frac{2}{m} \int F(x) dx + C_1 \right]^{\frac{1}{2}}} + C_2$$

Now, Apply B.C.

③ Force is a function of Speed: $F(v) = m \cdot \frac{dv}{dt}$

$$\frac{dv}{dt} = \frac{F(v)}{m}$$

$$\frac{dv}{F(v)} = \frac{1}{m} dt$$

$$\int \frac{dv}{F(v)} = \frac{1}{m} t + C_1$$

$$\text{Here, } v = H(t, C_1) \text{ or, } \frac{dx}{dt} = H(t, C_1)$$

Where, H is a function of t & C_1

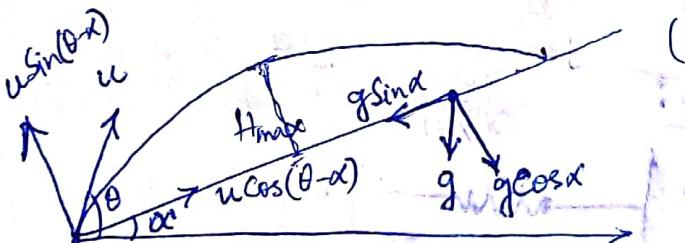
$$x = \int H(t, C_1) dt + C_2$$

Apply B.C. to find the values of C_1 and C_2 .

Lagrange-D'Alembert's Principle: The sum of differences between the forces acting on the system and the time derivative of momentum of the system along ~~with~~ a virtual displacement consistent with the constraints of the system is zero.

$$\sum (F_i - m_i a_i) \delta r_i = 0 \quad \text{or,} \quad \sum (F_i - m_i a_i) = 0$$

Projectile motion on Inclined plane: ie, $F - ma = 0$.



(i) $u \cos(\theta-\alpha)$ along inclined plane

(ii) $u \sin(\theta-\alpha)$ along ~~perp~~ to inclined plane

$$s = ut + \frac{1}{2} at^2$$

$$0 = u \sin(\theta-\alpha)t - \frac{1}{2}(g \cos \alpha)t^2$$

$$t = \frac{2u \sin(\theta-\alpha)}{g \cos \alpha} \quad (\text{time of flight})$$

Max height reached w.r.t inclined plane

$$V^2 = u^2 - 2gt$$

$$0 = u^2 - 2(g \cos \alpha) \cdot H_{\max}$$

$$H_{\max} = \frac{u^2 \sin^2(\theta-\alpha)}{2g \cos \alpha}$$

Max horizontal Range

$$S = ut + \frac{1}{2} gt^2$$

$$R = u \cos(\theta-\alpha) \cdot \frac{2u \sin(\theta-\alpha)}{g \cos \alpha}$$

$$R = \frac{2u^2 \sin(\theta-\alpha)}{g \cos \alpha}$$

$$-\frac{1}{2} (g \sin \alpha) \left[\frac{2u \sin(\theta-\alpha)}{g \cos \alpha} \right]^2$$

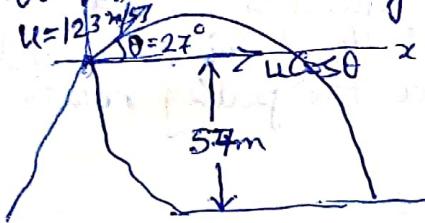
For project down inclined plane $\rightarrow -$

For max range

$$\theta = 45^\circ + \frac{\alpha}{2}$$

$$R_{\max} = \frac{u^2}{g \cos^2 \alpha} (1 - \sin \alpha)$$

Problem A gun is fired from hill top at 27° upward at the speed of 123 m/s. The enemy is 57 m below (vertically). Determine the max^m rise of bullet from horizontal, velocity and time to hit the enemy.



At the instant of hitting the enemy, if the vertical component of velocity

$$v_{\text{hit}}^2 = 55.84^2 + 2 \times 9.8 \times 158.97$$

$$v_{\text{hit}} = 78.96 \text{ m/s}$$

At this instant, final velocity = $\sqrt{109.6^2 + 78.96^2}$
 $= 135.08 \text{ m/s}$

$$u_x = u \cos \theta = 123 \cos 27^\circ = 109.6 \text{ m/s}$$

$$u_y = u \sin \theta = 123 \sin 27^\circ = 55.84 \text{ m/s.}$$

(At max^m height)

$$v_y^2 = u_y^2 - 2gH$$

$$0 = 55.84^2 - 2gH \quad H = 158.97 \text{ m}$$

If the time span required to hit the enemy

$$v = u + gt$$

$$78.96 = 55.84 + 9.8t$$

$$t = 2.358 \text{ sec.}$$

Problem A flywheel of 550 mm dia is brought uniformly from rest up to a speed of 350 rpm in 20 sec. Find the velocity and acceleration of a point on its rim 3 sec after starting from rest.

Initial angular speed $\omega_0 = 0$

$$\text{Speed at } 20 \text{ sec time, } \omega = \frac{2\pi \times 350}{60} = \frac{35\pi}{3} \text{ rad/s.}$$

$$\alpha = \frac{\omega - \omega_0}{t} = \frac{35\pi}{3 \times 20} \text{ rad/s}^2$$

Now, at $t = 3 \text{ sec.}$

$$\omega_t = \omega_0 + \alpha t = 0 + \frac{35\pi}{3 \times 20} \times 3 = \frac{7\pi}{4} \text{ rad/s.}$$

Velocity of the point; $v = r\omega_t \approx \frac{0.55}{2} \times \frac{7\pi}{4} = 1.512 \text{ m/s.}$

Acelⁿ — a) Tangential acclⁿ (a_t) = $r\alpha = \frac{0.55}{2} \times \frac{35\pi}{3 \times 20}$

$$= 0.504 \text{ m/s}^2$$

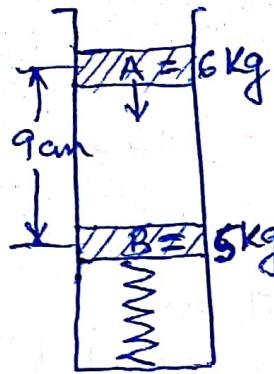
b) Normal acclⁿ (a_n) = $\omega_t^2 r$

$$= \left(\frac{7\pi}{4}\right)^2 \times \frac{0.55}{2} = 8.312 \text{ m/s}^2$$

$$\text{Total accl}^n (a) = \sqrt{a_t^2 + a_n^2}$$

$$= \sqrt{0.504^2 + 8.312^2} = 8.3273 \text{ m/s}^2$$

Problem ..



Block A of wt 6 kg drops on block B of weight 5 kg from a height of 9 cm as shown in figure. Considering the spring constant as $k = 1.85 \text{ N/mm}$, determine the amount of compression of spring over and above the static deformation due to block B.

Solⁿ Velocity of block A just before impact with block B,

$$u_A = \sqrt{2gh} = \sqrt{2 \times 9.81 \times 0.09} \\ = 1.328 \text{ m/s.}$$

After impact, block A & block B both will move down at a common velocity (say V).

Using principle of conservation of momentum,

$$m_A u_A + m_B u_B^1 = (m_A + m_B) V$$

$$\text{or, } 6 \times 1.328 + 0 = (6+5) V$$

$$\text{or, } V = 0.7247 \text{ m/s.}$$

Static deformation of spring due to block B,

$$\delta_{st} = \frac{W_B}{k} = \frac{m_B g}{K} = \frac{5 \times 9.81}{1.85} = 0.0265 \text{ m}$$

If due to impact, the spring gets compressed by an amount of δ , total compression

$$\Delta = \delta_{st} + \delta = (0.0265 + \delta) \text{ m}$$

Just after impact, the total energy stored in the system

$$E_1 = \frac{1}{2} K \delta_{st}^2 + \frac{1}{2} (m_A + m_B) V^2$$

$$= \frac{1}{2} \times 1.85 \times 0.0265^2 + \frac{1}{2} (6+5) \times 0.7247^2 \\ = 3.538 \text{ N-m}$$

When both the blocks come to rest after compressing the spring by δ , total energy stored in the system.

$$E_2 = \frac{1}{2} K \Delta^2 - (m_A + m_B) g \times \delta$$

$$= \frac{1}{2} \times 1.85 \times \Delta^2 - (6+5) \times 9.81 \times \delta \\ = 925 \Delta^2 - 107.878$$

from principle of conservation of energy.

$$E_2 = E_1$$

$$925 \Delta^2 - 107.878 = \cancel{25} 3.538$$

$$\therefore 925(s + 0.0265)^2 - 107.878 = 3.538$$

$$\therefore s^2 - 0.0636s - 3.1226 \times 10^{-3} = 0$$

$$\therefore s = 0.09609 \text{ m} \approx 96.09 \text{ mm}$$

Problem

A spring of unstretched length 0.2 m is connected to blocks A and B having weights of 3 kg and 2 kg, respectively. If initially, the spring is compressed by 5 cm and the system is released from rest, when the spring comes again in its normal length? Take Spring const. $k = 2 \text{ N/mm}$

From principle of

conservation of momentum,



$$m_A u_A + m_B u_B = 0 \quad (\text{as initial momentum of blocks is zero})$$

$$3u_A + 2u_B = 0 \quad \boxed{u_A = -\frac{2}{3}u_B}$$

If the spring expands by an amount s from its compressed position to come back to its original length, the compressive force, $F = k(0.05 - s) = 2 \times 10^3(0.05 - s)$

For displacement of 'ds',

$$\begin{aligned} dW &= F \times ds \\ W &= \int_0^{0.05} dW = \int_0^{0.05} 2 \times 10^3 (0.05 - s) ds \\ &= 2 \times 10^3 \left[0.05s - \frac{s^2}{2} \right]_0^{0.05} \\ &= 2.5 \text{ N-m} \end{aligned}$$

Initial kinetic energy of the block are zero.

$$\text{Final Kinetic energy} = \frac{1}{2} m_A u_A^2 + \frac{1}{2} m_B u_B^2$$

$$\begin{aligned} &= \frac{1}{2} \times 3 \times \left(-\frac{2}{3} u_B \right)^2 + \frac{1}{2} \times 2 u_B^2 \\ &= \frac{5}{3} u_B^2 \end{aligned}$$

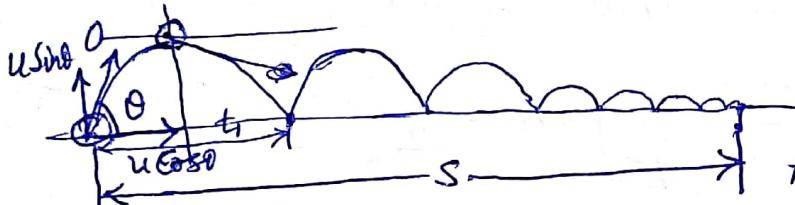
From Work-energy principle,

$$\frac{5}{3} u_B^2 = 2.5$$

$$\therefore u_B = 1.224 \text{ m/s} \quad \therefore u_A = -\frac{2}{3} u_B = -0.816 \text{ m/s.}$$

Problem A glass ball of mass 24 gm is thrown from a horizontal floor at an angle of ' θ ' at velocity u . When the ball stops bouncing after number of impacts, the horizontal distance travelled is 's'. If 'e' is symbolised as the co-efficient of restitution, prove that

$$s = \frac{u^2}{g} \frac{\sin 2\theta}{(1-e)}$$



Time interval: $t_1 = \frac{2u \sin \theta}{g}$

$$t_2 = \frac{2eu \sin \theta}{g}$$

and so on.

$$\text{Total time (T)} = t_1 + t_2 + t_3 + \dots$$

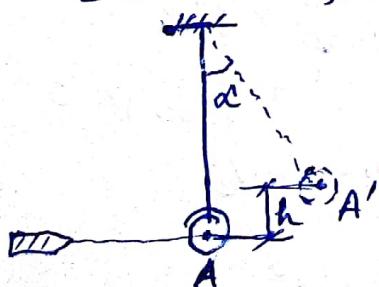
$$\begin{aligned} &= \frac{2u \sin \theta}{g} + \frac{2eu \sin \theta}{g} + \frac{2e^2 u \sin \theta}{g} + \frac{2e^3 u \sin \theta}{g} + \dots \\ &= \frac{2u \sin \theta}{g} (1 + e + e^2 + e^3 + \dots) \\ &= \frac{2u \sin \theta}{g} \times \frac{1}{(1-e)} \quad [\text{from Binomial expansion formulation}] \end{aligned}$$

Total distance travelled

$$\begin{aligned} (S) &= (u \cos \theta) \times \text{Total time (T)} \\ &= u \cos \theta \times \frac{2u \sin \theta}{g} \times \frac{1}{(1-e)} = \frac{u^2 \sin 2\theta}{g(1-e)} \end{aligned}$$

Problem A 0.303 bullet weighing 5.678 gm hits a pendulum of bob 707 gm suspended by an inextensible string of length 1.1 m. If the velocity of strike is 97.3 m/s,

determine the maximum angle of swing (α) when the bullet gets embedded within the bob and also when the bullet escapes the bob at a velocity of 17 m/s. If the bullet rebounds the bob at a velocity of 24.5 m/s, what will be the angle of swing? [$\alpha_1 = 11.26^\circ$, $\alpha_2 = 10.213^\circ$]



(a) When the bullet gets embedded, applying principle of conservation of momentum,

$$\left(\frac{5.678}{1000}\right)V + \left(\frac{707}{1000}\right)V = \left(\frac{5.678}{1000}\right) \times 97.3 + \left(\frac{707}{1000}\right) \times 0$$

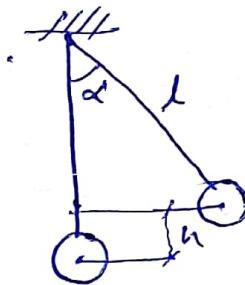
$$V = 0.7752 \text{ m/s. (Combined velocity)}$$

Let, the system gets displaced from position A to A' raising a height 'h'. Now, applying principle of conservation of energy.

$$\textcircled{2} \left(\frac{5.678}{1000} + \frac{707}{1000}\right) \times g \times h = \frac{1}{2} \left(\frac{5.678}{1000} + \frac{707}{1000}\right) V^2$$

or, $h = \frac{V^2}{2g}$.

or, $h = 0.0306 \text{ m.}$



Angle of swing (α)

$$\cos \alpha = \frac{l-h}{l} = \frac{1.1 - 0.0306}{1.1}$$

$$\alpha = 13.546^\circ$$

(b) If escape velocity of bullet $v_1 = 17 \text{ m/s}$ and velocity of bob is ' V_1 ' then using principle's of conservation of momentum,

$$\left(\frac{5.678}{1000}\right) \times 17 + \left(\frac{707}{1000}\right) \times V_1 = \left(\frac{5.678}{1000}\right) \times 97.3$$

$$\text{or, } V_1 = 0.6449 \text{ m/s.} \quad \text{...} + \left(\frac{707}{1000}\right) \times 0$$

Using principles of conservation of energy, $mgh_1 = \frac{1}{2} m V_1^2$

$$\textcircled{3} \left(\frac{707}{1000}\right) \times g \times h_1 = \frac{1}{2} \left(\frac{707}{1000}\right) \times 0.6449^2$$

$$\text{or, } h_1 = \frac{V_1^2}{2g} = 0.0212 \text{ m}$$

$$\text{so, Angle of swing } (\alpha_1) = \cos^{-1} \left(\frac{1.1 - 0.0212}{1.1} \right) = 11.267^\circ$$

(c) When the bullet rebounds at $v_2 = 24.5 \text{ m/s}$, from principles of conservation of momentum,

$$\left(\frac{5.678}{1000}\right) \times 24.5 + \left(\frac{707}{1000}\right) V_2 = \left(\frac{5.678}{1000}\right) \times 97.3$$

$$V_2 = 0.5847 \text{ m/s.} \quad \text{...} + \left(\frac{707}{1000}\right) \times 0$$

Using principle of conservation of energy, $mgh_2 = \frac{1}{2} m V_2^2$

$$\left(\frac{707}{1000}\right) \times g \cdot h_2 = \frac{1}{2} \left(\frac{707}{1000}\right) \times 0.5847^2$$

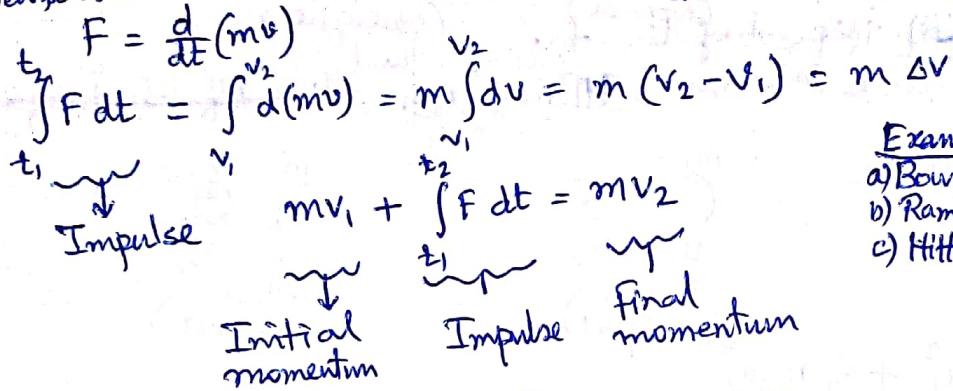
$$\text{or, } h_2 = 0.01743 \text{ m.}$$

$$\therefore \text{Angle of Swing } (\alpha_3) = \cos^{-1} \left[\left(\frac{1.1 - 0.01743}{1.1} \right) \right] = 10.213^\circ$$

Impulse & Momentum

From Newton's 2nd law of motion,

$$\frac{d}{dt}(I_z\omega) = M_2 \quad (\text{Gyroscope rotating about axis of symmetry/spin maintains constant angular velocity of spin})$$



Example of impulse

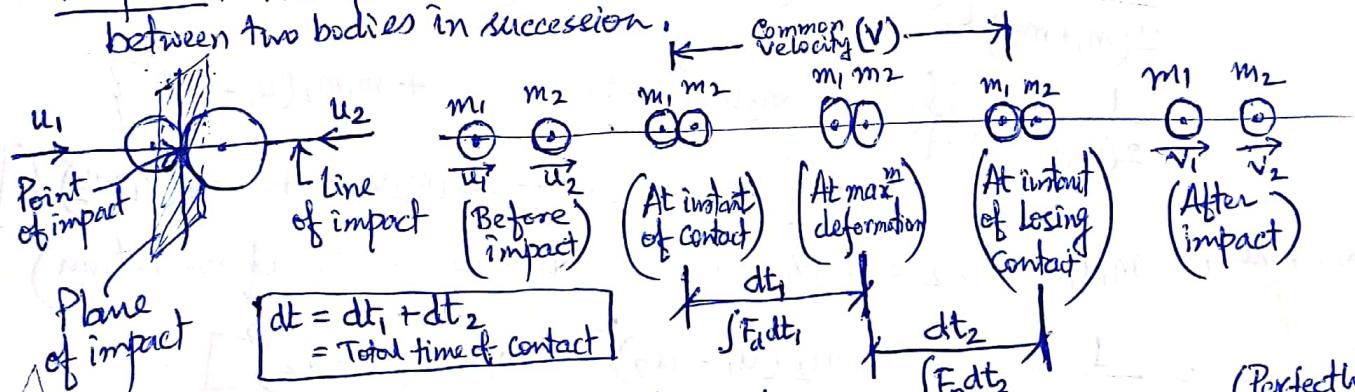
- a) Bouncing of ball
- b) Ramming of pile driver
- c) Hitting a baseball by bat

$$\sum_{i=1}^n (m_i v_{1i})_{x/y/z} + \sum_{i=1}^n \int_{t_1}^{t_2} F_{ix/y/z} dt = \sum_{i=1}^n (m_i v_{2i})_{x/y/z}$$

Principle of conservation of linear momentum: $\sum_{i=1}^n (m_i v_i)_{\text{initial}} = \sum_{i=1}^n (m_i v_i)_{\text{final}}$

Impact: Abrupt short duration encounter

between two bodies in succession.



Geff of restitution: Ratio of magnitude of impulse of force of recovery to impulse of force of deformation.

$$e = \frac{\int F_r dt_2}{\int F_d dt_1}$$

$$e = \frac{m_1 V - m_1 V_1}{m_1 U_1 - m_1 V} = \frac{V - V_1}{U_1 - V} \quad (\text{for Body A}) \quad \textcircled{1}$$

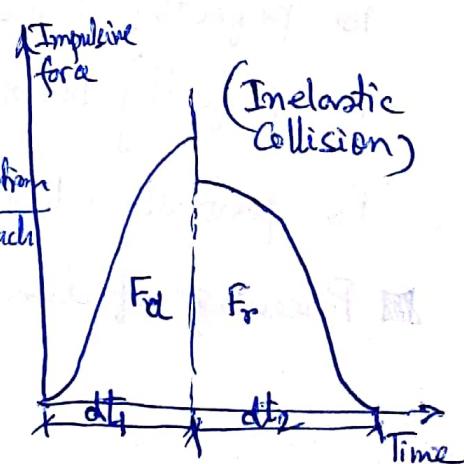
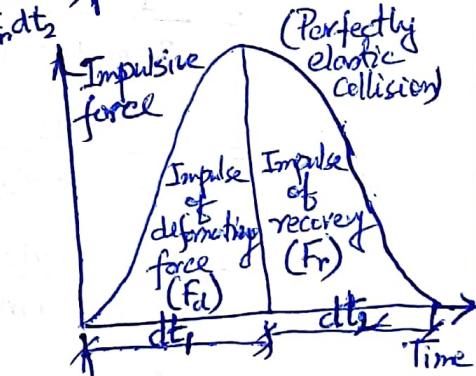
$$e = \frac{m_2 V_2 - m_2 V}{m_2 V - m_2 U_2} = \frac{V_2 - V}{V - U_2} \quad (\text{for body B}) \quad \textcircled{11}$$

From eqn ① & ⑪ applying concept of identity

$$e = \frac{(V - V_1) + (V_2 - V)}{(U_1 - V) + (V - U_2)} = \frac{V_2 - V_1}{U_1 - U_2} = \frac{\text{Vel. of separation}}{\text{Vel. of approach}}$$

If $e = 1 \Rightarrow$ Perfectly elastic collision

If $e = 0 \Rightarrow$ Perfectly plastic collision

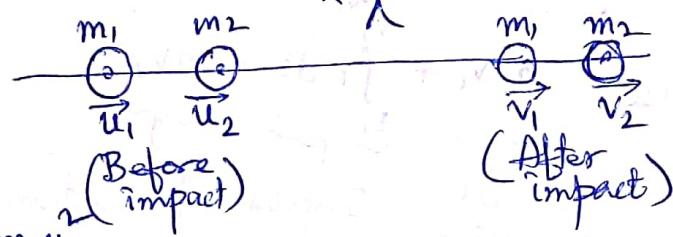


Impact types :-

- (a) Depending on direction of velocity
 - (i) Central (Normal)
 - (ii) Eccentric (Oblique)
- (b) Depending on type of impact
 - billiard ball by golf club, firing shotgun

Example :- Collision betⁿ atoms, Collision of ball by golf club, Firing shotgun

Inelastic collision:



$$KE_{\text{Initial}} = \frac{1}{2}m_1u_1^2 + \frac{1}{2}m_2u_2^2$$

$$KE_{\text{final}} = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2$$

Loss of KE during inelastic collision.

$$\begin{aligned}\Delta KE &= KE_{\text{Initial}} - KE_{\text{final}} = \left(\frac{1}{2}m_1u_1^2 + \frac{1}{2}m_2u_2^2 \right) - \left(\frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2 \right) \\ &= \frac{1}{2(m_1+m_2)} \left[(m_1+m_2)(m_1u_1^2 + m_2u_2^2) - (m_1+m_2)(m_1v_1^2 + m_2v_2^2) \right] \\ &= \frac{1}{2(m_1+m_2)} \left[\{(m_1u_1 + m_2u_2)^2 - 2m_1m_2u_1u_2 + m_1m_2(u_1 + u_2)^2\} \right. \\ &\quad \left. - \{(m_1v_1 + m_2v_2)^2 - 2m_1m_2v_1v_2 + m_1m_2(v_1 + v_2)^2\} \right]\end{aligned}$$

We know, $m_1u_1 + m_2u_2 = m_1v_1 + m_2v_2$ (from conservation of momentum)

$$\begin{aligned}&= \frac{1}{2(m_1+m_2)} \left[m_1m_2(u_1 - u_2)^2 - m_1m_2(v_1 - v_2)^2 \right] \\ &= \frac{m_1m_2}{2(m_1+m_2)} \left[(u_1 - u_2)^2 - e^2(u_1 - u_2)^2 \right] \quad \left[e = \frac{v_2 - v_1}{u_1 - u_2} \right] \\ &= \frac{m_1m_2(u_1 - u_2)^2}{2(m_1+m_2)} (1 - e^2)\end{aligned}$$

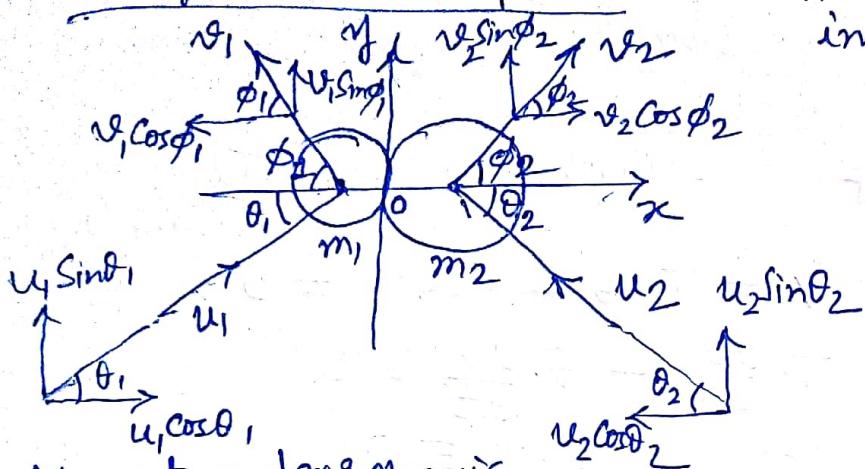
For perfectly elastic collision, $e = 1$. ie, $\Delta KE = 0$.

For perfectly plastic collision, $e = 0$ ie, $\Delta KE = \frac{m_1m_2(u_1 - u_2)^2}{2(m_1+m_2)}$

For partially elastic collision, $0 < e < 1$

Percentage of dissipation of energy = $\frac{\Delta KE}{KE_{\text{Initial}}} \times 100\%$.

Oblique Central Impact:



θ_1, θ_2 are the angles which initial velocities u_1, u_2 make with x -axis.

$\phi_1, \phi_2 \rightarrow$ make with y -axis.

Momentum along y -axis,

$$m_1 u_1 \sin \theta_1 = m_1 v_1 \sin \phi_1$$

$$\text{or, } u_1 \sin \theta_1 = v_1 \sin \phi_1 - \text{(i)}$$

$$m_2 u_2 \sin \theta_2 = m_2 v_2 \sin \phi_2$$

$$\text{or, } u_2 \sin \theta_2 = v_2 \sin \phi_2 - \text{(ii)}$$

From conservation of momentum,

$$m_1 u_1 \cos \theta_1 + m_2 u_2 \cos \theta_2 = m_1 v_1 \cos \theta_1 + m_2 v_2 \cos \theta_2 - \text{(iii)}$$

From Newton's ~~law~~ law of collision of elastic body valid for oblique impacts.

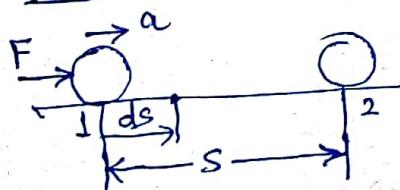
$$\text{Co-eff of restitution (e)} = \frac{v_2 \cos \phi_2 - v_1 \cos \phi_1}{u_2 \cos \theta_2 - u_1 \cos \theta_1} - \text{(iv)}$$

For perfectly elastic collision, $e = 1$

$$\text{or, } v_2 \cos \phi_2 - v_1 \cos \phi_1 = u_1 \cos \theta_1 - u_2 \cos \theta_2 - \text{(iva)}$$

Using eqn(i) to (iva), solving eqn(iv) we can get four unknowns such as, v_1, v_2, ϕ_1 and ϕ_2 , considering u_1, u_2, θ_1 & θ_2 are known.

Work-Energy Principle: If a body is subjected to force 'F' to travel elementary 'ds' distance in time 'dt' then Workdone by the force,



$$dW = F \cdot ds$$

From Newton's 2nd law of motion, $F = m \cdot a$

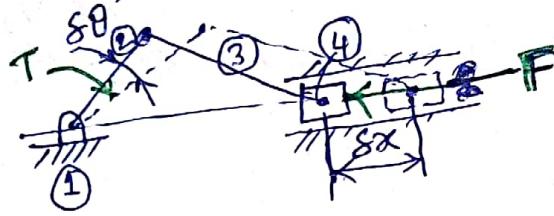
$$\text{Where, } a = \frac{dv}{dt} = \frac{dv}{ds} \cdot \frac{ds}{dt} = v \frac{dv}{ds}$$

$$\therefore dW = (m \cdot a) \cdot ds = m \cdot v \frac{dv}{ds} \cdot ds = mv \cdot dv$$

Integrating both sides:-

$$W_{1-2} = \int dW = m \int v dv = \frac{1}{2} m (v_2^2 - v_1^2) = KE_2 - KE_1 \text{ ie, workdone on the object equals to change in K.E. of the object. This is called Work-energy principle.}$$

Principles of Virtual Work:
The sum of workdone during a virtual small displacement from the equilibrium is equal to zero.



Total Workdone

= Workdone by Torque (T)

+ Workdone by force (F)

$$W = T \cdot \delta\theta + F \cdot \delta x = 0$$

$$\therefore T \cdot \frac{d\theta}{dt} + F \cdot \frac{dx}{dt} = 0$$

$$\therefore T \cdot \omega + F \cdot v = 0$$

$$\therefore T = - \frac{F \cdot v}{\omega}$$

D'Alembert's Principle:

Inertia forces and couples + External forces and torques on a body together provide the static equilibrium.

Mathematically,

$$\begin{aligned} \sum F + F_{\text{inertia}} &= 0 \\ \sum T + C_{\text{inertia}} &= 0 \end{aligned}$$

For static equilibrium,

$$F_{\text{inertia}} = 0 \quad \text{and} \quad C_{\text{inertia}} = 0.$$

then,

$$\begin{aligned} \sum F &= 0 \\ \sum T &= 0 \end{aligned}$$

D'Alembert's principle states:-

The system of forces acting on a body in motion is in dynamic equilibrium with the ~~inertia forces~~ of the body & ~~inertia couples~~.

where,

$$F_{\text{inertia}} = -m f_g$$

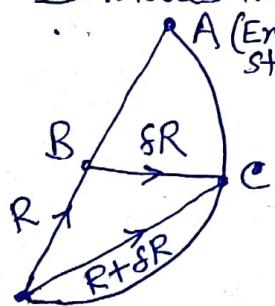
$$C_{\text{inertia}} = -I g \alpha$$

$$\sum F = f_1 + f_2 + f_3 + \dots + f_n \quad (\text{External forces})$$

$$\sum T = T_1 + T_2 + T_3 + \dots + T_n \quad (\text{External torques})$$

Hamilton's Principle:

It states that — Time integral of the difference between the kinetic and potential energies of dynamical system is ~~not~~ stationary.



From Newton's 2nd Law, $\mathbf{F} = m \frac{d^2 \mathbf{R}}{dt^2}$

Workdone during virtual displacement $\delta R'$,

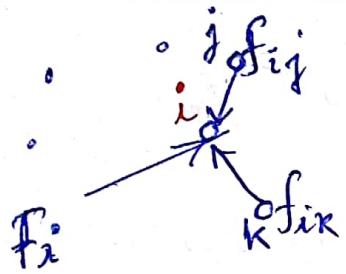
$$\delta W = \mathbf{F} \cdot \delta \mathbf{R} = m \frac{d^2 \mathbf{R}}{dt^2} \cdot \delta \mathbf{R} \quad \text{--- (i)}$$

Kinetic energy, $KE = \frac{1}{2} m \left(\frac{d\mathbf{R}}{dt} \right)^2$ or, $\delta KE = m \left(\frac{d\mathbf{R}}{dt} \right) \cdot \delta \left(\frac{d\mathbf{R}}{dt} \right)$

$$\delta W + \delta KE = m \frac{d}{dt} \left(\frac{d\mathbf{R}}{dt} \cdot \delta \mathbf{R} \right)$$

$$\int (\delta W + \delta KE) dt = \left[m \frac{d\mathbf{R}}{dt} \cdot \delta \mathbf{R} \right]_{t_0}^{t_1} + \int (KE - PE) dt = 0$$

General motion of a system of particle :



F_i = Applied force

f_{ij} = Interactive force between particles

from Newton's 2nd law :

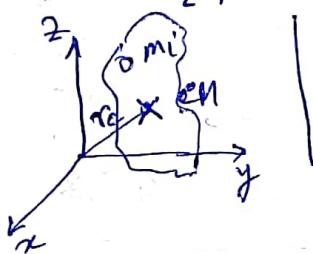
$$m_i \frac{d^2 r_i}{dt^2} = F_i + \sum_{j=1}^n f_{ij} \quad (i \neq j)$$

$$\sum_{i=1}^n m_i \frac{d^2 r_i}{dt^2} = \sum_{i=1}^n F_i + \sum_{i=1}^n \sum_{j=1}^n f_{ij}$$

$$F = \sum_{i=1}^n m_i \frac{d^2 r_i}{dt^2} = \frac{d^2}{dt^2} \sum_{i=1}^n m_i r_i$$

$$F = M \frac{d^2 r_c}{dt^2} \quad (M r_c)$$

$$= M \frac{d^2 r_c}{dt^2}$$



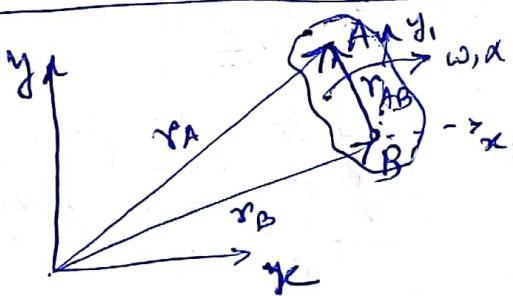
Since, $f_{ij} = -f_{ji}$

First moment vector

$$r_c \sum_{i=1}^n m_i = \sum_{i=1}^n m_i r_i \equiv \sum m_i r_i$$

$$r_c = \frac{\sum m_i r_i}{\sum m_i} = \frac{\sum m_i r_i}{M}$$

Governing kinetic eqn of Rigid body :



$$r_A = r_B + r_{AB}$$

$$v_A = v_B + v_{AB}$$

$$v_A = v_B + \omega \times r_{AB} \quad \text{where, } v_{AB} = \frac{dr_{AB}}{dt} = \omega \times r_{AB}$$

$$v_A = v_B + \underbrace{\omega \times r_{AB}}_{\text{Rotational component}}$$

$$v_A = \underbrace{v_B}_{\text{Translation component}} + \underbrace{\omega \times r_{AB}}_{\text{Rotational component}}$$

$$\text{Now, } a_A = a_B + a_{AB}$$

$$a_{AB} = \frac{d}{dt} (\omega \times r_{AB})$$

$$= \frac{d\omega}{dt} \times r_{AB} + \omega \times \frac{dr_{AB}}{dt}$$

$$= \alpha \times r_{AB} + \omega \times (\omega \times r_{AB})$$

$$a_A = a_B + \underbrace{\alpha \times r_{AB}}_{\text{Translation part}} + \underbrace{\omega \times (\omega \times r_{AB})}_{\text{Rotational part}} = a_B + (a_{AB})_T + (a_{AB})_R$$

$$\begin{cases} (a_{AB})_T = \alpha \times r_{AB} \\ (a_{AB})_R = \omega^2 \times r_{AB} \end{cases}$$



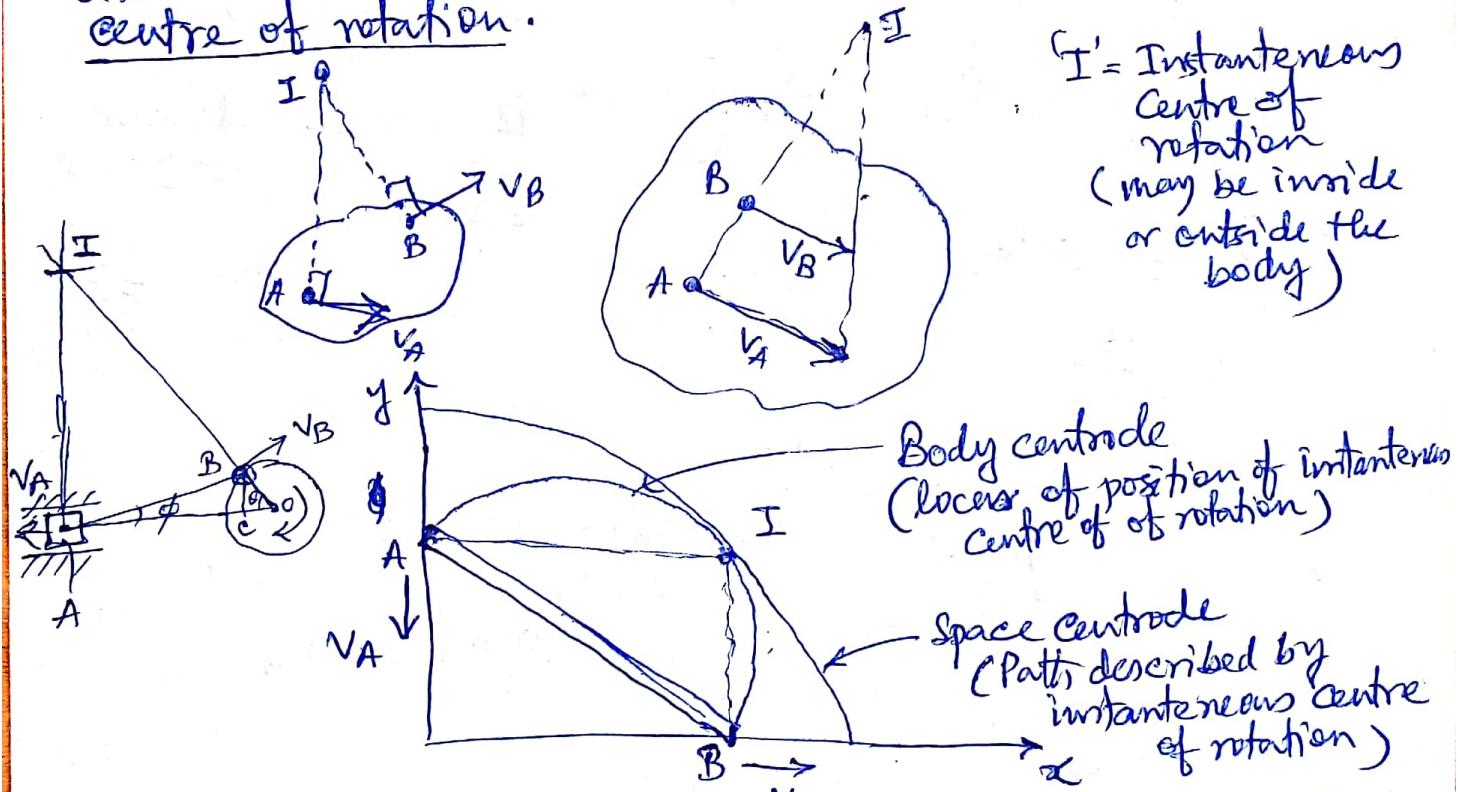
$$\begin{aligned}v_A &= v_B \\r_A \omega_A &= r_B \omega_B \\ \frac{\omega_A}{\omega_B} &= \frac{r_B}{r_A}\end{aligned}$$

$$r_A \alpha_A = -r_B \alpha_B$$

$$\frac{\alpha_A}{\alpha_B} = -\frac{r_B}{r_A}$$

Instantaneous Centre of Rotation:

At a given instant, the velocities of various particles of a body could ~~not~~ be expressed as the result of a pure rotation whose axis is perpendicular to the plane. This axis intersects the plane at a point called instantaneous centre of rotation.



The axis passing through the instantaneous centre point and at right angles to the plane of motion is called instantaneous axis of rotation. The instantaneous centre changes its position at every moment and its locus is called centrode. The surface generated by the instantaneous axis is called axode.

Important points on instantaneous centre:-

- It is the point about which the body appears to rotate
- It is not a fixed point but changes from one instant to another as the body rotates.
- The velocity at an instantaneous centre is zero.
- Instantaneous centre of a body may be inside or outside the body
- The dirⁿ of velocity at a point on a body is normal to the line joining the point and the instantaneous centre.
- I.C. lies at infinite distance when velocities at two points are equal, parallel & in same dirⁿ.

Coriolis' Law :

If a point moves along a path that has a rotational motion, the absolute acceleration of the point is given by the vector sum of three accelerations:—

(a) acceleration of the coincident point relative to which the point under consideration is moving

(b) acceleration of the point relative to the coincident point and

(c) Coriolis component of acceleration.

Mathematically,

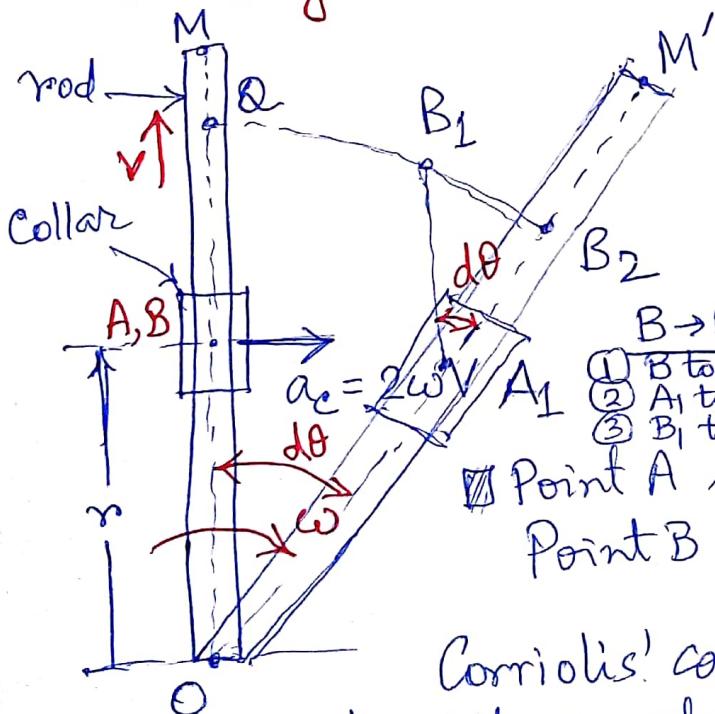
$$\vec{a}_A = \vec{a}_B + \vec{a}_{Aw} + \vec{a}_C$$

Absolute accⁿ of particle at point A

Accⁿ of point B of moving frame 'w' coinciding with A

Accⁿ of 'A' relative to moving frame 'w'

Coriolis' accⁿ.



Consider a collar slides at const velocity along the rod OM towards point 'Q'. Simultaneously the rod rotates at const angular velocity 'ω' about 'O'. When the collar moves on the rod to point 'Q' in time 'dt', the rod rotates through an angle $d\theta$. So, finally $A \rightarrow A_1$, $B \rightarrow B_2$

- ① $B \rightarrow B_2$ due to rotation of rod.
- ② $A_1 \rightarrow B_1$ due to outward velocity of collar on the rod.
- ③ $B_1 \rightarrow B_2$ due to Coriolis component of accⁿ of the rod.

Point A is on the rod.
Point B is on the collar.

Coriolis' component acts in the direction of relative velocity vector for the two coincident point (A and B) rotated by 90° in the direction of angular velocity of rod and it is expressed as,

$$\vec{a}_c = 2\omega V_{ba}$$

$$\text{Coriolis force} = m \vec{\omega} \times \vec{v}_c \quad \text{Where}$$

m = mass of the
collar

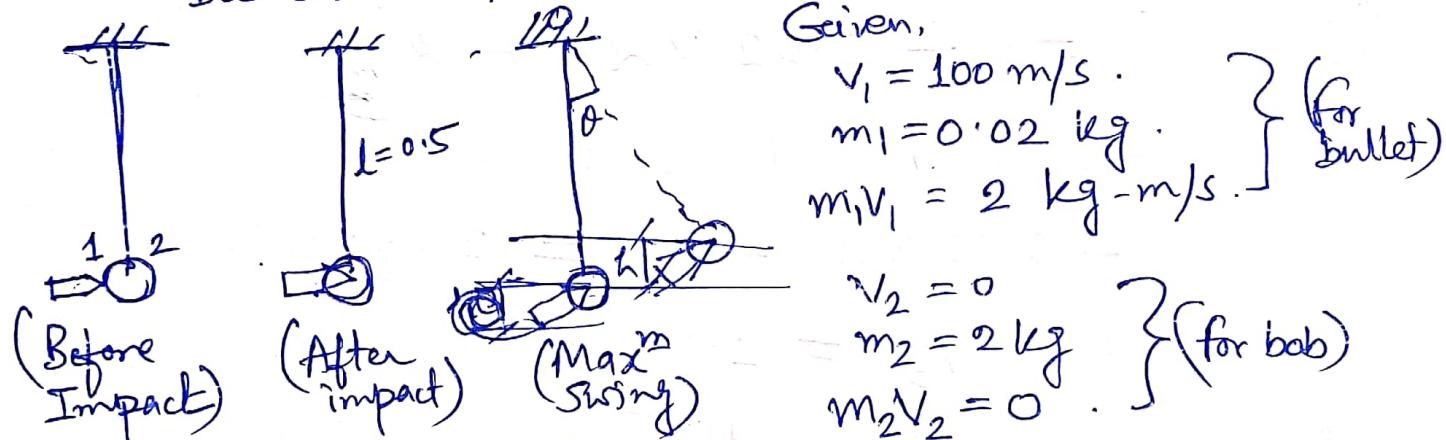
The concept of Coriolis acceleration is very useful in the study of long range projectiles and the bodies whose motions are appreciably affected by the rotation of earth.

Chasle's theorem:

It states that any general displacement of a rigid body can be represented by a combination of translatory motion and rotational motion.

Problem A bullet of mass 20 g moving with a velocity of 100 m/s. hits a 2 kg bob of a simple pendulum horizontally as shown in figure. Determine the max^m angle through which the pendulum string 0.5 m long may swing if:-

- (a) the bullet get embedded in the bob.
- (b) the bullet escapes from other end at 20 m/s.
- (c) the bullet is rebounded from the surface of the bob at 20 m/s.



Now, as per conservation of momentum,

$$\text{So, } m_1 v_1 + m_2 v_2' = 2 \text{ kg m/s } \quad \text{(Initial momentum)}$$

Case (a): Just after impact, the bullet and bob becomes united to travel at a velocity of V' m/s and their momentum = $(2 + 0.02)V' = 2.02V' \text{ kg-m/s}$.

Equating initial and final momentum,

$$2.02V' = 2 \quad \text{or, } V' = 0.99 \text{ m/s.}$$

After the impact,

the bob must rise by a distance h according to principle of energy conservation,

$$\frac{1}{2} m V'^2 = mgh$$

$$\text{or, } h = \frac{\frac{1}{2} m V'^2}{mg} = 0.05 \text{ m.}$$

$$\therefore \theta = \cos^{-1} \left(\frac{0.5 - 0.05}{0.5} \right) = 25.85^\circ$$

Case (b): If the bullet escaped from other end of bob at a velocity of 20 m/s, the final momentum would be given by, $= (0.02 \times 20) + 2 \times V'_2$

Equating with initial momentum, we get,

$$0.02 \times 20 + 2 V'_2 = 2 \text{ kg-m/s} \quad \text{or, } V'_2 = 0.8 \text{ m/s.}$$

$$So, h = \frac{V_2'^2}{2g} = \frac{0.8^2}{2 \times 9.81} = 0.0326 \text{ m.}$$

$$\theta = \cos^{-1} \left(\frac{l-h}{l} \right) = \cos^{-1} \left(\frac{0.5 - 0.0326}{0.5} \right) = 20.8^\circ$$

Case(e): If instead the bullet rebounded at 20 m/s, the final momentum would be given by,

$$= -0.02 \times 20 + 2V_3'$$

Equating with initial momentum,

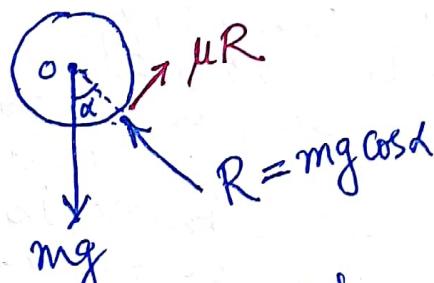
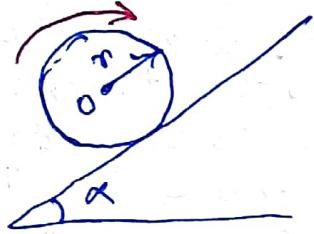
$$(-0.02 \times 20) + 2V_3' = 2 \text{ kg-m/s.}$$

$$\therefore V_3' = 1.2 \text{ m/s.}$$

$$\therefore h = \frac{V_3'^2}{2g} = \frac{1.2^2}{2 \times 9.81} = 0.0734 \text{ m}$$

$$\theta = \cos^{-1} \left(\frac{l-h}{l} \right) = \cos^{-1} \left(\frac{0.5 - 0.0734}{0.5} \right) = 31.4^\circ.$$

Problem A uniform circular cylinder of mass 'm' and radius 'r' is given an initial angular velocity ' ω_0 ' and no initial translational velocity. It is placed in contact with the plane inclined at ' α ' to the horizontal and it moves in upward direction. If ' μ ' = co-eff of friction for sliding between cylinder and the plane, find the distance the cylinder moves before sliding stops. Assume, $\mu > \tan \alpha$



Using angular momentum eqn to determine the time 't' required to decrease the velocity from ' ω_0 ' to ' ω '

$$-\mu R r t = I(\omega - \omega_0) = \frac{mr^2}{2} (\omega - \omega_0)$$

$$\text{or, } -\mu (mg \cos \alpha) \cdot r \cdot t = \frac{mr^2}{2} (\omega - \omega_0)$$

$$\text{or, } t = \frac{r(\omega_0 - \omega)}{2\mu g \cos \alpha} \quad \dots \dots \dots \quad (1)$$

The velocity of the centre 'O' parallel to the plane is obtained by applying linear momentum equation,

$$(umg \cos \alpha - mg \sin \alpha) \times t = m(v - 0) \quad \dots \dots \quad (2)$$

(with positive dirⁿ
up the plane)

From eqⁿ (1) & (2)

Note:

At the instant sliding stops
and pure rolling begins.

$$v = r\omega$$

~~$$1) mg(\mu \cos \alpha - \sin \alpha) \times t = mr\omega$$~~

~~$$\text{or, } \frac{rg(\mu \cos \alpha - \sin \alpha) \times r(\omega_0 - \omega)}{2\mu g \cos \alpha} = \omega r \beta$$~~

$$\omega = \frac{\mu \cos \alpha - \sin \alpha}{3\mu \cos \alpha - \sin \alpha} \cdot \omega_0 \quad \text{and} \quad t = \frac{r\omega_0}{g(3\mu \cos \alpha - \sin \alpha)}$$

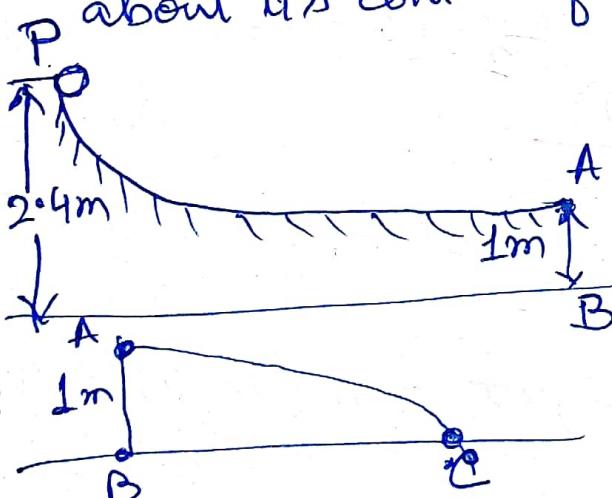
The forces acting on cylinder remain constant during sliding and hence, acceleration is constant.

The distance travelled by sliding

$$d = \frac{1}{2} vt = \frac{1}{2} \omega r \cdot t$$

$$= \frac{\pi^2 \omega_0^2 (\mu \cos \alpha - \sin \alpha)}{2g (3\mu \cos \alpha - \sin \alpha)^2}$$

Problem A small sphere rolls down without slipping from the top of a track in a vertical plane. The track has an elevated section and a horizontal part as shown in Figure. The horizontal part is 1 m above the ground level and the top of the track is 2.4 m above the ground. Find the distance on the ground with respect to point B' (which is vertically below the end of the track as shown) where the sphere lands. During its flight projectile, does the sphere continue to rotate about its centre of mass?



At Point P,

$$KE_1 = 0$$

$$PE_1 = 2.4mg$$

At point A,

$$KE_2 = \frac{1}{2} mv^2 + \frac{1}{2} I \omega^2$$

$$= \frac{1}{2} mv^2 + \frac{1}{2} \left(\frac{2}{5} mR^2\right) \frac{v^2}{R^2}$$

$$= \frac{1}{2} mv^2 + \frac{1}{5} mv^2$$

$$= \frac{7}{10} mv^2$$

$$PE_2 = 1mg$$

$$Z = -\frac{gx^2}{2V_0^2}$$

$$\approx x = \sqrt{2 \times 2g \times \frac{1}{g}}$$

$$= 2 \text{ m.}$$

By conservation of energy

$$0 + 2.4mg = \frac{7}{10} mv^2 + 1mg$$

$$1.4mg = 0.7mv^2$$

$$\therefore v = \sqrt{2g}$$

After leaving A, the sphere would continue to rotate at the same rotational speed in absence of any external moment acting on it.

Problem The rectilinear motion of a particle is defined by the time-displacement eqⁿ: $x = x_0 (2e^{-kt} - e^{-2kt})$, in which x_0 is the initial displacement, e is natural logarithmic base. Find the max^m velocity of the particle.

$$x = x_0 (2e^{-kt} - e^{-2kt})$$

$$\dot{x} = 2x_0 (-k) e^{-kt} - x_0 (-2k) e^{-2kt}$$

$$= -2kx_0 [e^{-kt} - e^{-2kt}]$$

For attaining max^m velocity.

$$\ddot{x} = 0$$

$$e^{-kt} - e^{-2kt} = 0$$

$$\text{or } e^{kt} = 2$$

$$\text{or, } t = \frac{1}{k} \ln 2$$

$$\therefore \text{Max}^m \text{ velocity} = \dot{x}|_{\text{max}} = \dot{x}|_{t=\frac{1}{k} \ln 2} = -\frac{kx_0}{2}.$$

Problem An automobile starting from rest increases its speed from '0' to 'v' with constant acceleration ' α ', runs at this speed for a time and finally comes to rest with constant acceleration ' β '. If the total distance travelled is ' S ', find the total time(t) required.

$$\text{while moving with accel}^m \alpha, \text{time reqd } (t_1) = \frac{v-0}{\alpha} = \frac{v}{\alpha}$$

$$\text{while moving with decel}^m \beta, \text{time reqd } (t_3) = \frac{v-0}{\beta} = \frac{v}{\beta}$$

$$\text{Distance travelled during } t_1, S_1 = \frac{1}{2} \alpha t_1^2 = \frac{1}{2} (\alpha t_1) \cdot t_1$$

$$= \frac{1}{2} vt_1$$

$$\text{Distance travelled during } t_3, S_3 = \frac{1}{2} \beta t_3^2 = \frac{1}{2} (\beta t_3) \cdot t_3$$

$$= \frac{1}{2} vt_3$$

$$\text{Distance travelled during } t_2, S_2 = vt_2$$

$$\text{Now, } S = S_1 + S_2 + S_3 = \frac{1}{2} vt_1 + vt_2 + \frac{1}{2} vt_3$$

$$v, vt_2 = s - \frac{1}{2} v(t_1 + t_3)$$

$$v, t_2 = \frac{s}{v} - \frac{1}{2}(t_1 + t_3)$$

Total time of travel,

$$t = t_1 + t_2 + t_3$$

$$= \frac{v}{\alpha} + \left[\frac{s}{v} - \frac{1}{2}(t_1 + t_3) \right] + \frac{v}{\beta}$$

$$= \frac{v}{\alpha} + \left[\frac{s}{v} - \frac{1}{2} \frac{v}{\alpha} - \frac{1}{2} \frac{v}{\beta} \right] + \frac{v}{\beta}$$

$$= \frac{s}{v} + \frac{v}{2} \left(\frac{1}{\alpha} + \frac{1}{\beta} \right)$$

If, $\alpha = \beta$ then

$$t = \frac{s}{v} + \frac{v}{\alpha}$$

Problem The acceleration of a particle along a straight line is given by the eqⁿ; $a = 4 - \frac{t^2}{9}$. If the particle starts with zero velocity from a position $x=0$, find:-

(i) its velocity after 6 sec. (ii) distance travelled in 6 sec.

$$\text{Here, } a = 4 - \frac{t^2}{9} = \frac{dv}{dt}$$

$$v, dv = \left(4 - \frac{t^2}{9} \right) dt$$

Integrating:-

$$\int_0^t dv = \int_0^t \left(4 - \frac{t^2}{9} \right) dt$$

$$v = 4t - \frac{t^3}{27}$$

$$\frac{dx}{dt} = \left(4t - \frac{t^3}{27} \right)$$

$$dx = \left(4t - \frac{t^3}{27} \right) dt$$

Integrating:-

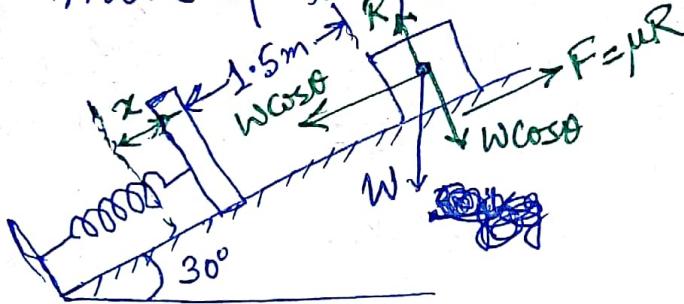
$$\int_0^x dx = \int_0^t \left(4t - \frac{t^3}{27} \right) dt$$

$$x = 2t^2 - \frac{t^4}{108}$$

At $t = 6$ sec,

$$v = 16 \text{ m/s} \text{ and } x = 60 \text{ m}$$

Problem A block of 50 kg mass is released from rest and slides down a 30° inclined plane as shown. After sliding 1.5 m down the plane, the block hits spring of spring constant (K) = 25 N/mm. If the co-eff of friction between the block and the plane is 0.2, make calculations for maximum deflection of the spring, and the max^{ns} velocity of the block. Also find the distance the block will move up the plane due to rebound.



Let, x is the deformation of the spring.

$$\therefore \text{Distance moved by the block} = (1.5 + x)$$

$$\therefore \text{Workdone by the block} = \text{Net force on block} \times \text{distance moved by the block}$$

$$= 160.29 \times (1.5 + x)$$

$$\text{Now, Workdone by the spring} = -\frac{1}{2} Kx^2 = -\frac{1}{2} (25 \times 10^3) x^2$$

$$= -12.5 \times 10^3 x^2$$

$$\therefore \text{Total Workdone} = W_B + W_S$$

$$= 160.29 \times (1.5 + x) - 12.5 \times 10^3 x^2$$

The block starts from rest and finally comes to rest when there is max^{ns} deflection of spring. Obviously the change in K.E. is zero.

Applying Work-energy principle,

$$\text{Workdone} = \text{Change in kinetic energy} = 0$$

$$160.29 \times (1.5 + x) - 12.5 \times 10^3 x^2 = 0$$

$$\text{or}, 12.5 x^2 - 0.16 x - 0.24 = 0$$

$$\text{or}, x = 0.145 \text{ m or } -0.132 \text{ m}$$

Hence acceptable
 $x = 0.145 \text{ m}$

Net force on the block when it starts moving downward

$$= W \sin \theta - \mu R$$

$$= mg \sin \theta - \mu w \cos \theta$$

$$= mg \sin \theta - \mu mg \cos \theta$$

$$= mg (\sin \theta - \mu \cos \theta)$$

$$= 160.29 \text{ N.}$$

(i) According to Newton's Second Law, $F = ma$

$$160 \cdot 29 = 50 a$$

$$\therefore a = 3.21 \text{ m/s}^2$$

From kinematic relations,

$$v^2 - u^2 = 2as$$

$$v^2 = 2 \times a \times s = 2 \times 3.21 \times 1.5$$

$$v = 3.10 \text{ m/s. (Max^m vels. of block)}$$

(ii) When the block moves up the inclined plane due to rebound, both the components of weight of the block along the plane and the frictional forces act in the same direction.

$$\begin{aligned}\therefore \text{Force on the block} &= W \sin \theta + \mu mg \cos \theta \\ &= mg \sin \theta + \mu mg \cos \theta \\ &= 330.21 \text{ N.}\end{aligned}$$

If 's' be the distance

which is moved up by the block due to rebound,

Work done by forces acting on the block

$$= 330.21 \times s$$

This workdone equals to energy stored in spring

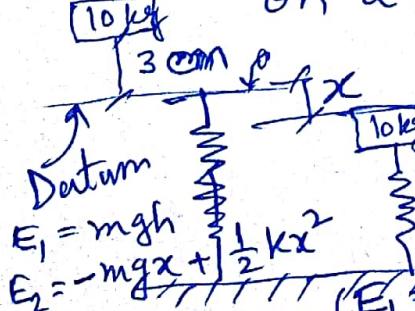
$$\therefore \frac{1}{2} k x^2 = 330.21 \times s.$$

$$\therefore \frac{1}{2} \times (25 \times 10^3) \times (0.145)^2 = 330.21 \times s$$

$$\therefore s = 0.796 \text{ m.}$$

Ex Problem

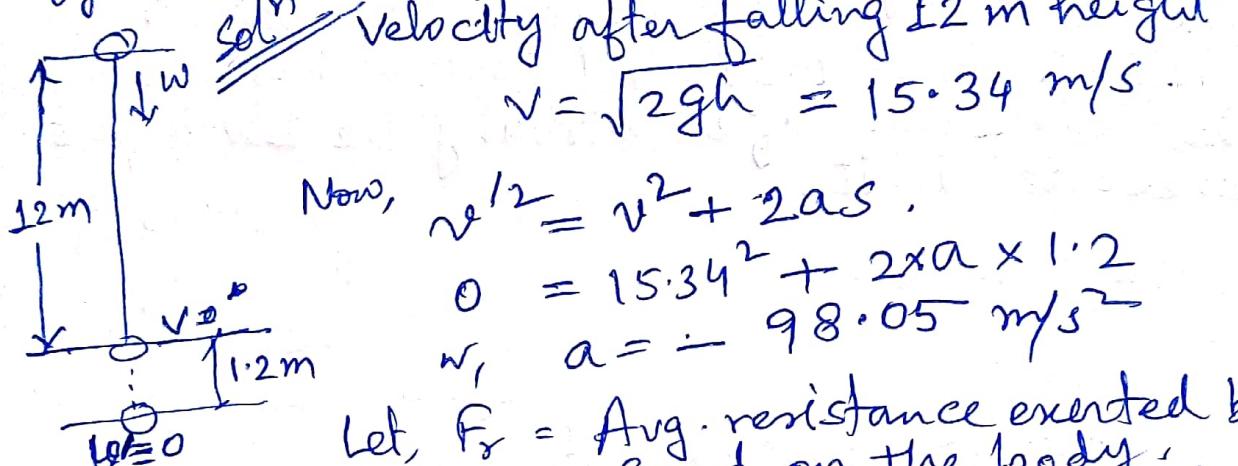
A body of mass 10 kg is made to fall 3 cm height on a spring of stiffness 120 N/cm. Find the



displacement of spring using the concept that total energy of mass-spring system remains constant.

$$[An x = 3.178 \text{ cm}]$$

Problem A body of weight 100N falls from a height of 12 m on a sand bed. It is estimated that the body penetrates 1.2 m into the sand before coming to rest. Determine the average thrust exerted by the sand on the body.



Let, F_r = Avg. resistance exerted by Sand on the body.

∴ Net force acting on body in upward direction

$$F_n = F_r - W$$

$$\text{w, } \frac{100}{9.81} \times 98.05 = F_r - 100 \quad [\text{As, } F_n = ma]$$

$$\text{w, } F_r = 1100 \text{ N}$$

Problem A train weighing 4000 kN has a frictional resistance of 5 N/kN of weight. Determine the steady pull which the locomotive must exert if the speed of the train is to be increased from 30 km/hr to 60 km/hr within a period of 1.5 min.

Solution Initial velocity (u) = 30 km/hr = 8.33 m/s.
Final " " (v) = 60 km/hr = 16.66 m/s.

$$v = u + at$$

$$\text{w, } a = \frac{v-u}{t} = \frac{16.66-8.33}{1.5 \times 60} = 0.092 \text{ m/s}^2$$

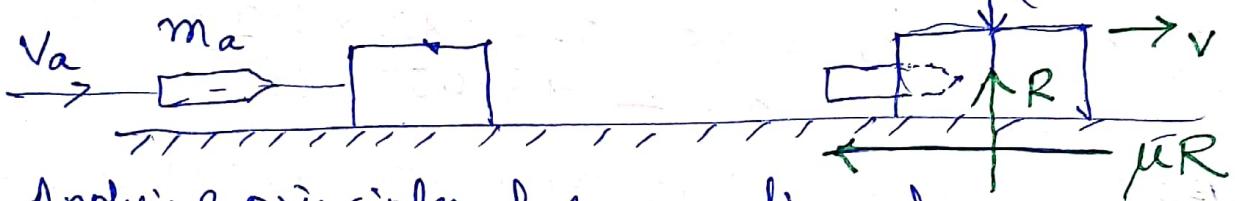
frictional resistance (F_f) = 5 N per KN of train weight
 $= 5 \times 4000 = 20,000 \text{ N}$.

Net force (F) = Frictional force (F_f) - Frictional resistance (F_f)
or steady pull

$$\text{w, } \frac{4000 \times 10^3}{9.81} \times 0.092 = F_t - 20,000 \quad [\text{where } F = ma]$$

$$F_t = 57.512 \text{ KN}$$

Problem A wooden block of weight 40 N rests on a rough horizontal plane having friction co-eff $\mu = 0.35$. The block is struck by a bullet travelling horizontally with a velocity of 750 m/s and weighing 0.25 N. Work out the distance by which the block is displaced from its initial position considering that the bullet after striking the block gets embedded in it.



Applying principles of conservation of momentum,

$$m_a V_a + m_b V_b = (m_a + m_b) V$$

$$\left(\frac{0.25}{9.81}\right) \times 750 + \left(\frac{40}{9.81}\right) \times 0 = \left(\frac{0.25 + 40}{9.81}\right) \times V$$

$$\text{or } V = 4.657 \text{ m/s.}$$

$$\text{force of friction } (F_f) = \mu R = \mu (m_a + m_b)g$$

$$= 0.35 \left[\frac{0.25 + 40}{9.81} \right] \times 9.81$$

$$= 14.087 \text{ N.}$$

Applying Work-energy Correlation

Workdone to overcome the frictional force
= kinetic energy lost by the block with bullet embedded

$$\text{or, } F_f \times s = \frac{1}{2} (m_a + m_b) \times V^2$$

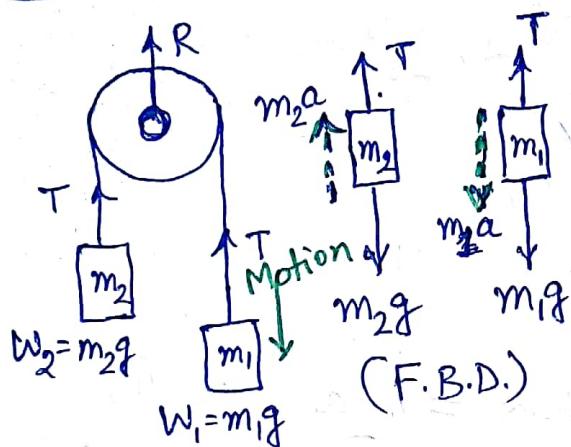
$$\text{or, } 14.087 \times s = \frac{1}{2} \times \left(\frac{0.25 + 40}{9.81} \right) \times (4.657)^2$$

$$\text{or, } s = 3.097 \text{ m.}$$

Motion of Connected Bodies:

Motion of two bodies connected each other when are subjected to some external forces, both bodies would be moving with same velocity and acceleration.

Case-1: Two bodies connected by a string passing over smooth pulley



For mass m_1 :

$$m_1g - T = m_1a \dots \dots \text{(i)}$$

For mass m_2 :

$$T - m_2g = m_2a \dots \dots \text{(ii)}$$

Adding eqn (i) & (ii) we get:-

$$m_1g - m_2g = (m_1 + m_2)a$$

Putting the value of 'a' in eqn (i)

$$a = \left(\frac{m_1 - m_2}{m_1 + m_2} \right) g = \left(\frac{W_1 - W_2}{W_1 + W_2} \right) g$$

$$m_1g - T = m_1a = m_1 \left(\frac{m_1 - m_2}{m_1 + m_2} \right) g \quad [m_1 = \frac{W_1}{g} \text{ & } m_2 = \frac{W_2}{g}]$$

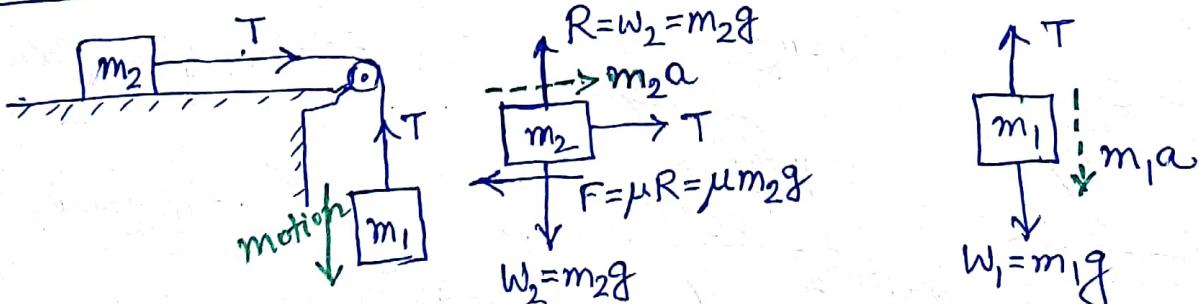
$$\therefore T = m_1g \left[1 - \frac{m_1 - m_2}{m_1 + m_2} \right] = \left(\frac{2m_1m_2}{m_1 + m_2} \right) g = \left(\frac{2W_1W_2}{W_1 + W_2} \right) g$$

Pressure on the pulley

$$R = \sqrt{P^2 + Q^2 + 2PQ \cos\theta} = \sqrt{T^2 + T^2 + 2T \cdot T \cdot \cos 0^\circ} = \sqrt{4T^2} = 2T$$

$$\therefore R = 2T = \frac{4W_1W_2}{W_1 + W_2}$$

Case-2: Two bodies connected at the edge of a horizontal surface



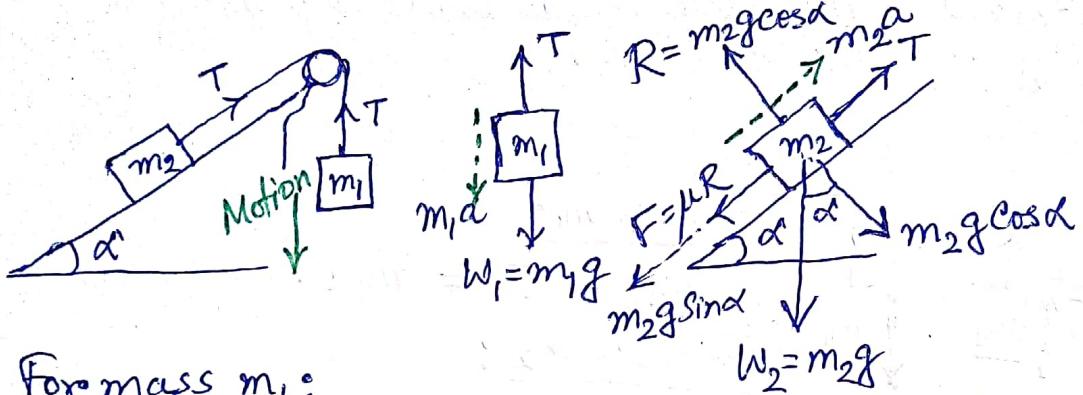
For mass m_1 : $T - \mu m_2 g = m_1 a \dots \dots \text{(i)}$

For mass m_2 : $m_2g - T = m_2 a \dots \dots \text{(ii)}$

Solving eqn (i) & (ii) we get,

$$a = \frac{m_1g - \mu m_2 g}{m_1 + m_2} = \frac{W_1 - \mu W_2}{W_1 + W_2} \quad \text{and} \quad T = \frac{m_1m_2(1+\mu)}{m_1 + m_2} \cdot g \\ = \frac{W_1W_2(1+\mu)}{W_1 + W_2}$$

Case-3: Two bodies connected by a string one hangs free and the other on rough inclined plane



For mass m_1 :

$$m_1 g - T = m_1 a \quad \dots \dots \text{(i)}$$

For mass m_2 :

$$T - m_2 g \sin \alpha - \mu m_2 g \cos \alpha = m_2 a \quad \dots \dots \text{(ii)}$$

Solving eqn (i) & (ii), we get,

$$a = \frac{m_1 g - m_2 g \sin \alpha - \mu m_2 g \cos \alpha}{m_1 + m_2} = \frac{(W_1 - W_2 \sin \alpha - \mu W_2 \cos \alpha)}{W_1 + W_2}$$

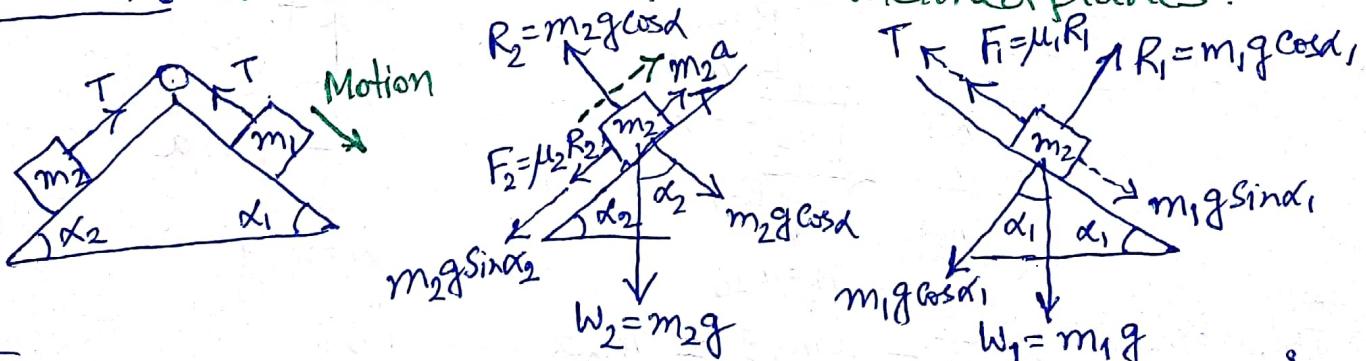
$$T = \frac{m_1 m_2 (1 + \sin \alpha + \mu \cos \alpha) g}{m_1 + m_2} = \frac{W_1 W_2 (1 + \sin \alpha + \mu \cos \alpha)}{W_1 + W_2}$$

If, $\mu = 0$, $a = \frac{(W_1 - W_2 \sin \alpha)}{W_1 + W_2} \times g$ and $T = \frac{W_1 W_2 (1 + \sin \alpha)}{W_1 + W_2}$

Pressure on pulley

$$R = \sqrt{P^2 + Q^2 + 2PQ \cos \theta} = \sqrt{T^2 + T^2 + 2T \cdot T \cos(90^\circ - \alpha)} = \sqrt{2T^2(1 + \sin \alpha)} \\ = T \sqrt{2(1 + \sin \alpha)}$$

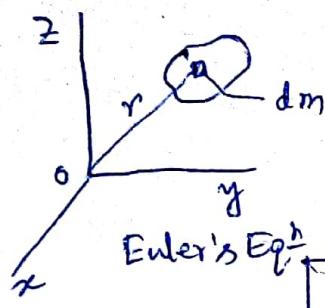
Case-4: Two bodies connected over inclined planes.



For mass m_1 : $m_1 g \sin \alpha_1 - T - \mu_1 m_1 g \cos \alpha_1 = m_1 a \quad \dots \dots \text{(i)}$

For mass m_2 : $T - m_2 g \sin \alpha_2 - \mu_2 m_2 g \cos \alpha_2 = m_2 a \quad \dots \dots \text{(ii)}$

Momentum and Angular momentum



Momentum of elementary mass = $dm \times v = dm \times \omega r$
 Moment of momentum of " " = momentum \times radius
 For entire mass, moment of momentum = $dm \times \omega r \times r$
 $= \int dm \cdot \omega r^2$
 $= \omega \int dm r^2$
 $= \omega I_c$ [for circular section,
 $I = \frac{\pi}{64} d^4$]

Impulse-momentum:

From Newton's 2nd law: $F = ma = m \frac{dv}{dt}$
 $\int F dt = m \int v_2 - v_1 dv = m(v_2 - v_1)$

From Euler's eqⁿ: $M_c = I_c \alpha = I_c \cdot \frac{d\omega}{dt}$
 $\int M_c dt = I_c \int_{\omega_1}^{\omega_2} d\omega = I_c(\omega_2 - \omega_1)$

Linear momentum Conservation principle: $m_1 v_1 + m_2 v_2 = m_1 v'_1 + m_2 v'_2$
 Angular " " " " " " : $I_a \omega_1 + I_b \omega_2 = I_a \omega'_1 + I_b \omega'_2$

Work-Energy formulation:

Linear: $dW = F \cdot dr$

$$W = \int dW = \int_{r_1}^{r_2} F \cdot dr$$

$$PE = mgf \theta h \text{ and } KE = \frac{1}{2} mv^2 + \frac{1}{2} I_c \omega^2$$

$$W_{\text{linear}} = \int_{r_1}^{r_2} F \cdot dr = \frac{1}{2} m(v_2^2 - v_1^2) = KE_{2\text{lin}} - KE_{1\text{lin}}$$

$$W_{\text{Angular}} = \int_{\theta_1}^{\theta_2} M \cdot d\theta = \frac{1}{2} I_c(\omega_2^2 - \omega_1^2) = KE_{2\text{Angular}} - KE_{1\text{Angular}}$$

$$W_{\text{Total}} = W_{\text{Linear}} + W_{\text{Angular}} = \int_{r_1}^{r_2} F \cdot dr + \int_{\theta_1}^{\theta_2} M \cdot d\theta$$

$$= \frac{1}{2} m(v_2^2 - v_1^2) + \frac{1}{2} I_c(\omega_2^2 - \omega_1^2)$$

$$= \left[\frac{1}{2} m v_2^2 + \frac{1}{2} I_c \omega_2^2 \right] - \left[\frac{1}{2} m v_1^2 + \frac{1}{2} I_c \omega_1^2 \right]$$

$$(PE_1 - PE_2)_{\text{Total}} = (KE_2 - KE_1)_{\text{Total}} = KE_2 - KE_1$$

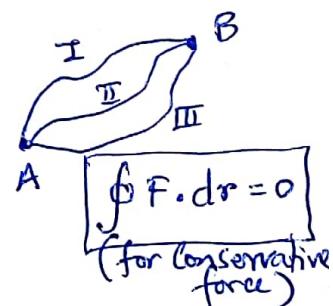
$$(KE_p + PE_1) = (KE_2 + PE_2)_T$$

(Principle of conservation of energy)

Angular: $dW = M \cdot d\theta$

$$W = \int dW = \int_{\theta_1}^{\theta_2} M \cdot d\theta$$

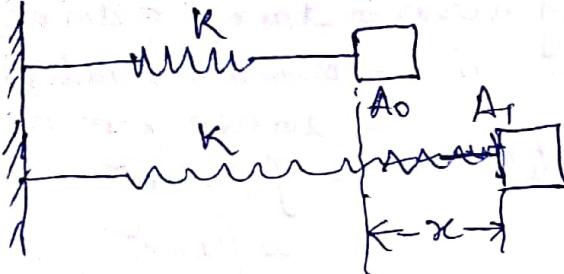
where,
 M = Moment



$$\int F \cdot dr = 0$$

(for conservative force)

Workdone by Spring:



Spring force \propto displacement

$$F \propto x$$

$$\therefore F = -Kx \quad K = \text{Spring stiffness}$$

Elementary workdone
by spring = $F dx = -Kx dx$

$$W = - \int_{x_1}^{x_2} Kx dx = -\frac{1}{2} Kx^2$$

$$W = \int_{x_1}^{x_2} Kx dx = -\frac{1}{2} (x_2^2 - x_1^2) = \frac{1}{2} (x_1^2 - x_2^2)$$

Graphically, Again, $W = \frac{1}{2} \times (F_1 + F_2) \times (x_2 - x_1)$

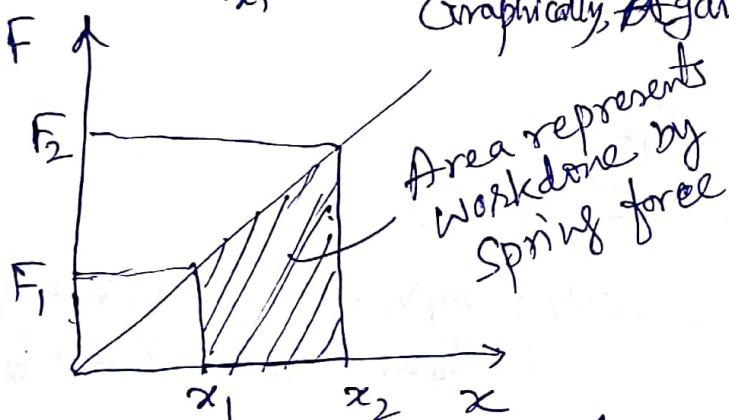
= Area of trapezium

$$= \frac{1}{2} \{(-Kx_1) + (-Kx_2)\} \times (x_2 - x_1)$$

$$= -\frac{1}{2} K (x_2 + x_1)(x_2 - x_1)$$

$$= -\frac{1}{2} K (x_2^2 - x_1^2)$$

$$= \frac{1}{2} K (x_1^2 - x_2^2)$$

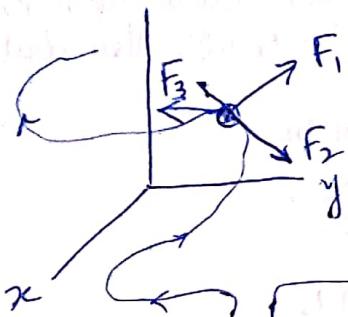


Workdone by the spring force

is positive when $x_2 < x_1$

i.e., when the spring is returning to its original position

Kinetics of Particle:



$$\sum \mathbf{F} = m\mathbf{a}$$

$$\sum F_x = m a_x$$

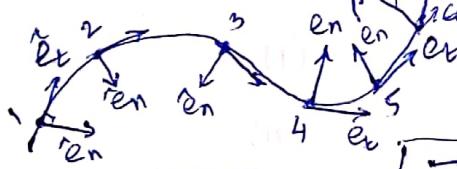
$$\sum F_y = m a_y$$

$$\sum F_z = m a_z$$

$$a_n = \frac{v^2}{P}$$

$$\text{where, } P = \frac{(1 + y_1^2)^{3/2}}{y_2}$$

$\hat{e}_t \rightarrow$ tangential component of unit vector
(Dirⁿ of velocity or motion)



$$v = v \hat{e}_t$$

$$\vec{a} = a_t \hat{e}_t + a_n \hat{e}_n$$

Changes speed
(By brake or accelerator)

Changes dirⁿ
(By steering wheel)

$$\bar{a} = \frac{d\bar{v}}{dt} = \frac{d}{dt}(v \hat{e}_t)$$

$$= v \left(\frac{d\hat{e}_t}{dt} \right) + \frac{dv}{dt} \hat{e}_t$$

$$\text{Now, } \frac{d\hat{e}_t}{dt} = \left(\frac{d\hat{e}_t}{d\theta} \right) \left(\frac{d\theta}{dt} \right) = \hat{e}_n \cdot \omega$$

$$= v \omega \hat{e}_n + \frac{dv}{dt} \hat{e}_t$$

$$\text{Now, } v = \frac{ds}{dt} = P \frac{d\theta}{dt} = P\omega$$

$$= P \cdot \omega \cdot \hat{e}_n + \frac{dv}{dt} \hat{e}_t$$

$$= P\omega^2 \hat{e}_n + \frac{dv}{dt} \hat{e}_t$$

$$= \frac{(P\omega)^2}{P} \cdot \hat{e}_n + \left(\frac{dv}{dt} \right) \cdot \hat{e}_t$$

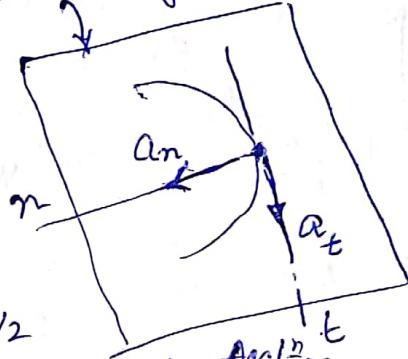
$$= \frac{v^2}{P} \cdot \hat{e}_n + \left(\frac{dv}{dt} \right) \cdot \hat{e}_t$$

$$= v \hat{e}_t + \left(\frac{v^2}{P} \right) \hat{e}_n$$

a_t or
tangential accⁿ

Centripetal
accⁿ (a_n)

motion only occurs
in Osculating plane



Note: Accⁿ or
force along
bi-normal dirⁿ is zero
 $a_t \text{ or } F_b = 0$ ($a_t \times a_n$)

$\hat{e}_n \rightarrow$ normal component
of unit vector
along radius of
curvature.
(Dirⁿ of turning)

Straight line motion:

$$P \rightarrow \infty, \quad a_n = \frac{v^2}{P} = 0, \quad a = a_t = v^2$$

Constant Speed along the path

$$a_t = v^2 = 0, \quad a = a_n = \frac{v^2}{P}$$

Constant tangential accⁿ.

$$a_t = (a_t)_c \quad v = (v_0) + (a_t)_c t$$

$$s - s_0 = (v_0)t + \frac{1}{2}(a_t)_c t^2$$

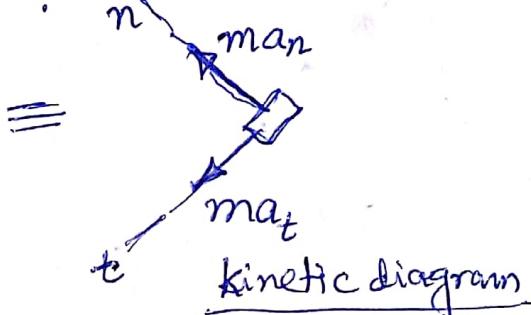
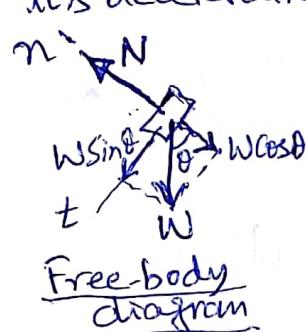
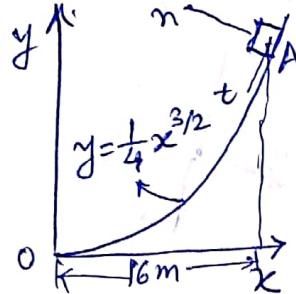
$$v^2 - v_0^2 = 2(a_t)_c (s - s_0)$$

Path: $y = f(x)$

$$P = \frac{\left[1 + \left(\frac{dy}{dx} \right)^2 \right]^{3/2}}{\left(\frac{d^2y}{dx^2} \right)}$$



Problem A 10 kg crate travels along a smooth slope as shown. If at the point 'A', its speed is 20 m/s, determine the normal force exerted by the slope to the crate. Also what is its acceleration?



By force balancing,

$$\sum F_n = N - W \cos \theta = m a_n \quad \dots \dots \quad (1)$$

$$\sum F_t = W \sin \theta = m a_t \quad \dots \dots \quad (2)$$

$$\text{Now, } y = \frac{1}{4} x^{3/2} \text{, so, } \frac{dy}{dx} = \text{Slope} = \frac{3}{8} x^{\frac{1}{2}}$$

$$\text{At } x=16 \text{ m, Slope} = \frac{3}{8} \times 16^{\frac{1}{2}} = 1.5 = \tan \theta$$

$$\text{So, } \theta = \tan^{-1}(1.5) = 56.3^\circ$$

$$\text{We know, } a_n = \frac{v^2}{r} \text{, Here, } v = 20 \text{ m/s.}$$

$$r = \frac{[1 + (\frac{dy}{dx})^2]^{3/2}}{\frac{d^2y}{dx^2}}, \text{ Now, } \frac{d^2y}{dx^2} = \frac{3}{8} \frac{x^{-1/2}}{x^2} = \frac{3}{16} x^{-3/2}$$

$$\left| \frac{d^2y}{dx^2} \right|_{x=16} = \frac{3}{64}$$

$$r = 125 \text{ m}$$

$$\text{So, } a_n = \frac{v^2}{r} = \frac{20^2}{125} = 3.2 \text{ m/s}^2$$

From eqn (1) & (2),

$$N - 10 \times 9.81 \cos 56.3^\circ = 10 \times 3.2$$

$$10 \times 9.81 \sin 56.3^\circ = 10 \times a_t$$

Solving the above eqn we get

$$a_t = 8.16 \text{ m/s}^2$$

$$N = 86.4 \text{ N}$$

$$a = \sqrt{a_n^2 + a_t^2} = 8.77 \text{ m/s}^2$$

