# Title of My Thesis



SUBMITTED BY

Jan-Philipp Anton Konrad Christ

## Titel meiner Arbeit

## Bachelorarbeit

FAKULTÄT FÜR PHYSIK
QUANTEN VIELTEILCHENSYSTEME/ THEORETISCHE NANOPHYSIK
LUDWIG-MAXIMILIANS-UNIVERSITÄT
MÜNCHEN

VORGELEGT VON

Jan-Philipp Anton Konrad Christ

# Title of My Thesis

## **Bachelor Thesis**

FACULTY OF PHYSICS

QUANTUM MANY-BODY SYSTEMS/ THEORETICAL NANOPHYSICS GROUP

LUDWIG MAXIMILIAN UNIVERSITY

MUNICH

SUBMITTED BY

Jan-Philipp Anton Konrad Christ

First supervisor: Prof. Dr. Fabian Bohrdt, geb. Grusdt

Second supervisor:  $\dots$ 

## NOTATION AND SYMBOLS

Here you can write which notation and conventions you want to use in your formulas.

For example, there are two different sign conventions for the metric tensor  $\eta_{\mu\nu}$  in Minkowski space.

In (general) relativity, it is common to use the sign convention  $\eta_{\mu\nu} = \text{diag}(-,+,+,+)$ , which is called the "mostly plus", "space dominant" or "east coast" sign convention.

Therefore, to lower the components of a contravariant vector  $\mathbf{x} = (ct, x, y, z) =: (x^0, x^1, x^2, x^3)$ , we get

$$x_{0} = \eta_{0\nu}x^{\nu} = -ct \neq x^{0},$$
  

$$x_{1} = \eta_{1\nu}x^{\nu} = +x = x^{1},$$
  

$$x_{2} = \eta_{2\nu}x^{\nu} = +y = x^{2},$$
  

$$x_{3} = \eta_{3\nu}x^{\nu} = +z = x^{3}.$$

On the other hand, in particle physics and quantum field theory, it is common to use the sign convention  $\eta_{\mu\nu} = \text{diag}(+, -, -, -)$ , which is called the "mostly minus", "time dominant" or "west coast" sign convention.

Therefore, to lower the components of a contravariant vector  $\boldsymbol{x}=(ct,x,y,z)=:(x^0,x^1,x^2,x^3),$  we get

$$x_{0} = \eta_{0\nu}x^{\nu} = +ct = x^{0},$$
  

$$x_{1} = \eta_{1\nu}x^{\nu} = -x \neq x^{1},$$
  

$$x_{2} = \eta_{2\nu}x^{\nu} = -y \neq x^{2},$$
  

$$x_{3} = \eta_{3\nu}x^{\nu} = -z \neq x^{3}.$$

Personally, I prefer the "mostly plus" sign convention  $\eta_{\mu\nu} = \text{diag}(-,+,+,+)$ , since it is only necessary to flip the sign of one component by lowering the index instead of three – but that is of course up to you.

Another convention that is often declared at the beginning of a document are the units that are used.

For example, someone might prefer Planck units instead of SI units. Here, I want to show how to convert from SI units to Planck units.

$$[c]_{\rm SI} = 299\,792\,458\,\frac{\rm m}{\rm s} \implies \left[\sqrt{\frac{\hbar G}{c^3}}\right]_{\rm SI} \approx 1.616\times 10^{-35}\,{\rm m} =: l_P$$

$$[\hbar]_{\rm SI} \approx 1.054\times 10^{-34}\,{\rm kg}\frac{{\rm m}^2}{\rm s} \implies \left[\sqrt{\frac{\hbar G}{c^5}}\right]_{\rm SI} \approx 5.391\times 10^{-44}\,{\rm s} =: t_p$$

$$[G]_{\rm SI} \approx 6.674\times 10^{-11}\,\frac{{\rm m}^3}{{\rm kg}\cdot{\rm s}^2} \implies \left[\sqrt{\frac{\hbar c}{G}}\right]_{\rm SI} \approx 2.176\times 10^{-8}\,{\rm kg} =: m_P$$

$$\left[\frac{1}{4\pi\varepsilon_0}\right]_{\rm SI} \approx 8.987\times 10^9\,{\rm kg}\frac{{\rm m}^3}{{\rm s}^2\cdot{\rm C}^2} \implies \left[\sqrt{4\pi\epsilon_0\hbar c}\right]_{\rm SI} \approx 1.875\times 10^{-18}\,{\rm C} =: q_P$$

$$[k_B]_{\rm SI} \approx 1.381\times 10^{-23}\,{\rm kg}\frac{{\rm m}^2}{{\rm s}^2\cdot{\rm K}} \implies \left[\sqrt{\frac{\hbar c^5}{Gk_B^2}}\right]_{\rm SI} \approx 1.416\times 10^{32}\,{\rm K} =: T_P$$

In Planck units, we have  $[c]_P = 1$ ,  $[\hbar]_P = 1$ ,  $[G]_P = 1$ ,  $[\frac{1}{4\pi\epsilon_0}]_P = 1$ ,  $[k_B]_P = 1$  and therefore  $l_P = 1$ ,  $t_P = 1$ ,  $m_P = 1$ ,  $q_P = 1$ ,  $T_P = 1$ , which is very convienent to express and manipulate equations.

Of course, this "trick" demands to remind ourself, which physical quantity we deal with while manipulating equations.

## \_\_\_\_\_ABSTRACT

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INTRODUCTION

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#### 1.1 Section 01

#### 1.1.1 Subsection 1-01

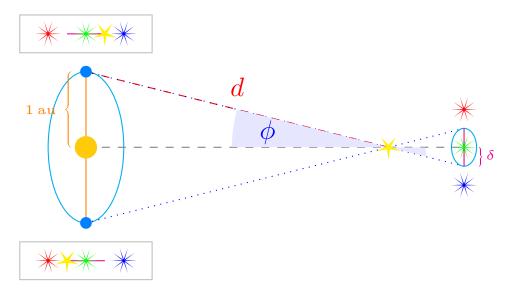


Figure 1.1: Some graphics made with Tikz. Have a look into the directory "figures/tikz/parallax".

## 1.2 Section 02

2 Introduction

#### 1.2.1 Subsection 2-01

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#### 1.2.2 Subsection 2-02

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#### 1.2.3 Subsection 3-03

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## 1.3 Section 03

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#### 1.3.1 Subsection 3-01

**1.3 Section 03** 3

#### 1.3.2 Subsection 3-02

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#### 1.3.3 Subsection 3-03

4 Introduction

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## 2.1 Second Quantization

#### 2.1.1 Occupation number representation and the Fock Space

I can cite stuff [Leavitt].

#### 2.1.2 Bosonic creation and annihilation operators

Hello, here is some text without a meaning. This text should show what a printed text will look like at this place. If you read this text, you will get no information. Really? Is there no information? Is there a difference between this text and some nonsense like "Huardest gefburn"? Kjift – not at all! A blind text like this gives you information about the selected font, how the letters are written and an impression of the look. This text should contain all letters of the alphabet and it should be written in of the original language. There is no need for special content, but the length of words should match the language.

## 2.1.3 One-Body-Operators

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## 2.1.4 Two-Body-Operators

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#### 2.1.5 Change of basis

#### 2.2 Section 02

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#### 2.2.1 Subsection 2-01

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#### 2.2.2 Subsection 2-02

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#### 2.2.3 Subsection 3-03

**2.3 Section 03** 7

### 2.3 Section 03

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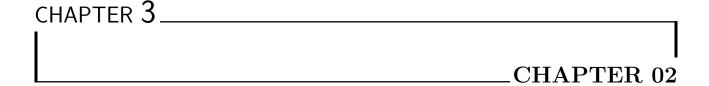
#### 2.3.1 Subsection 3-01

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#### 2.3.2 Subsection 3-02

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#### 2.3.3 Subsection 3-03



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## 3.1 Section 01

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#### 3.1.1 Subsection 1-01

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#### 3.1.2 Subsection 1-02

10 **Chapter 02** 

#### 3.1.3 Subsection 1-03

Hello, here is some text without a meaning. This text should show what a printed text will look like at this place. If you read this text, you will get no information. Really? Is there no information? Is there a difference between this text and some nonsense like "Huardest gefburn"? Kjift – not at all! A blind text like this gives you information about the selected font, how the letters are written and an impression of the look. This text should contain all letters of the alphabet and it should be written in of the original language. There is no need for special content, but the length of words should match the language.

#### 3.2 Section 02

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#### 3.2.1 Subsection 2-01

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#### 3.2.2 Subsection 2-02

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#### 3.2.3 Subsection 2-03

**3.3 Section 03** 

#### 3.3 Section 03

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#### 3.3.1 Subsection 3-01

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#### 3.3.2 Subsection 3-02

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#### 3.3.3 Subsection 3-03

12 Chapter 02

CHAPTER 4	
1	
	CONCLUSION

14 Conclusion

## DETAILED CALCULATIONS

## A.1 Deriving the flow equations in the case of no ndependence

First the canonical generator  $\hat{\eta}$  has to be evaluated:

$$\hat{\eta} := \hat{\eta}(\lambda) := \left[\hat{\mathcal{H}}_{0}, \hat{\mathcal{H}}_{int}\right] = \left[\sum_{k} \omega_{k} \hat{a}_{k}^{\dagger} \hat{a}_{k}, \sum_{q \neq q'} V_{q,q'} \hat{a}_{q}^{\dagger} \hat{a}_{q'} + \sum_{p,p'} \left(W_{p,p'} \hat{a}_{p}^{\dagger} \hat{a}_{p'}^{\dagger} + W_{p,p'}^{*} \hat{a}_{p} \hat{a}_{p'}\right)\right]$$

$$= \sum_{k} \sum_{q,q'} \omega_{k} V_{q,q'} \left[\hat{a}_{k}^{\dagger} \hat{a}_{k}, \hat{a}_{q}^{\dagger} \hat{a}_{q'}\right] + \sum_{k} \sum_{p,p'} \left(\omega_{k} W_{p,p'} \left[\hat{a}_{k}^{\dagger} \hat{a}_{k}, \hat{a}_{p}^{\dagger} \hat{a}_{p'}^{\dagger}\right] + \omega_{k} W_{p,p'}^{*} \left[\hat{a}_{k}^{\dagger} \hat{a}_{k}, \hat{a}_{p} \hat{a}_{p'}\right]\right)$$

$$= \sum_{k} \sum_{q,q'} \omega_{k} V_{q,q'} \left(\hat{a}_{k}^{\dagger} \hat{a}_{q'} \delta_{k,q} - \hat{a}_{q}^{\dagger} \hat{a}_{k} \delta_{k,q'}\right)$$

$$+ \sum_{k} \sum_{p,p'} \left(\omega_{k} W_{p,p'} \left(\hat{a}_{k}^{\dagger} \hat{a}_{p}^{\dagger} \delta_{k,p'} + \hat{a}_{k}^{\dagger} \hat{a}_{p'}^{\dagger} \delta_{k,p}\right) - \omega_{k} W_{p,p'}^{*} \left(\hat{a}_{p} \hat{a}_{k} \delta_{k,p'} + \hat{a}_{p'} + \hat{a}_{p'} \hat{a}_{k} \delta_{k,p}\right)\right)$$

$$= \sum_{q \neq q'} V_{q,q'} (\omega_{q} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} + \sum_{p,p'} \left(W_{p,p'} (\omega_{p} + \omega_{p'}) \hat{a}_{p}^{\dagger} \hat{a}_{p'}^{\dagger} - W_{p,p'}^{*} (\omega_{p} + \omega_{p'}) \hat{a}_{p} \hat{a}_{p'}\right)$$

$$(A.2)$$

Since  $\hat{\eta}$  has the same form as  $\hat{\mathcal{H}}_{int}$ ,  $\left[\hat{\eta}, \hat{\mathcal{H}}_{0}\right]$  follows by inspection of A.2:

$$\left[\hat{\eta}, \hat{\mathcal{H}}_{0}\right] = -\sum_{q \neq q'} V_{q,q'} (\omega_{q} - \omega_{q'})^{2} \hat{a}_{q}^{\dagger} \hat{a}_{q'}$$

$$-\sum_{p,p'} \left( W_{p,p'} (\omega_{p} + \omega_{p'})^{2} \hat{a}_{p}^{\dagger} \hat{a}_{p'}^{\dagger} + W_{p,p'}^{*} (\omega_{p} + \omega_{p'})^{2} \hat{a}_{p} \hat{a}_{p'} \right)$$
(A.3)

16 Detailed Calculations

The commutator of the generator and  $\hat{\mathcal{H}}_{int}$  needs more work:

In the following, A.5-A.8 will be evaluated separately:

A.5:

$$\begin{split} & \left[ \sum_{q \neq q'} V_{q,q'}(\omega_{q} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{q'}, \sum_{\tilde{q} \neq \tilde{q}'} V_{\tilde{q},\tilde{q}'} \hat{a}_{\tilde{q}}^{\dagger} \hat{a}_{\tilde{q}'} \right] \\ & = \sum_{q \neq q'} \sum_{\tilde{q} \neq \tilde{q}'} V_{\tilde{q},\tilde{q}'} V_{q,q'}(\omega_{q} - \omega_{q'}) \left[ \hat{a}_{q}^{\dagger} \hat{a}_{q'}, \hat{a}_{\tilde{q}}^{\dagger} \hat{a}_{\tilde{q}'} \right] \\ & = \sum_{q \neq q'} \sum_{\tilde{q} \neq \tilde{q}'} V_{\tilde{q},\tilde{q}'} V_{q,q'}(\omega_{q} - \omega_{q'}) \left( \hat{a}_{q}^{\dagger} \hat{a}_{\tilde{q}'} \delta_{q',\tilde{q}} - \hat{a}_{\tilde{q}}^{\dagger} \hat{a}_{q'} \delta_{q,\tilde{q}'} \right) \\ & = \sum_{q \neq q'} \sum_{\tilde{q}'} V_{q',\tilde{q}'} V_{q,q'}(\omega_{q} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{\tilde{q}'} - \sum_{q \neq q'} \sum_{\tilde{q}} V_{\tilde{q},q} V_{q,q'}(\omega_{q} - \omega_{q'}) \hat{a}_{\tilde{q}}^{\dagger} \hat{a}_{q'} \\ & = \sum_{q,q'} \sum_{\tilde{q}'} V_{\tilde{q},q'} V_{q,q'}(\omega_{q} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{\tilde{q}'} - \sum_{q,q'} \sum_{\tilde{q}} V_{\tilde{q},q} V_{q,q'}(\omega_{q} - \omega_{q'}) \hat{a}_{\tilde{q}}^{\dagger} \hat{a}_{q'} \\ & = \sum_{q,q'} \sum_{\tilde{q}} V_{\tilde{q},q'} V_{q,\tilde{q}}(\omega_{q} - \omega_{\tilde{q}}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} - \sum_{q,q'} \sum_{\tilde{q}} V_{q,\tilde{q}} V_{\tilde{q},q'}(\omega_{\tilde{q}} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} \\ & = \sum_{q \neq q'} \sum_{\tilde{q}} V_{\tilde{q},q'} V_{q,\tilde{q}}(\omega_{q} - \omega_{\tilde{q}}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} - \sum_{q \neq q'} \sum_{\tilde{q}} V_{q,\tilde{q}} V_{\tilde{q},q'}(\omega_{\tilde{q}} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} \\ & + \sum_{k} \sum_{\tilde{q}} V_{\tilde{q},k} V_{k,\tilde{q}}(\omega_{k} - \omega_{\tilde{q}}) \hat{a}_{k}^{\dagger} \hat{a}_{k} - \sum_{k} \sum_{\tilde{q}} V_{q,\tilde{q}} V_{\tilde{q},q'}(\omega_{\tilde{q}} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} \\ & + \sum_{k} \sum_{\tilde{q}} V_{\tilde{q},q'} V_{q,\tilde{q}}(\omega_{q} - \omega_{\tilde{q}}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} - \sum_{q \neq q'} \sum_{\tilde{q}} V_{q,\tilde{q}} V_{\tilde{q},q'}(\omega_{\tilde{q}} - \omega_{q'}) \hat{a}_{q}^{\dagger} \hat{a}_{q'} \\ & + \sum_{k} \sum_{\tilde{q}} V_{\tilde{q},k} V_{k,\tilde{q}}(\omega_{k} - \omega_{\tilde{q}}) \hat{a}_{q}^{\dagger} \hat{a}_{k} \end{aligned} \tag{A.9}$$

**A.6**:

$$\left[ \sum_{q \neq q'} V_{q,q'}(\omega_q - \omega_{q'}) \hat{a}_q^{\dagger} \hat{a}_{q'}, \sum_{\tilde{p},\tilde{p}'} \left( W_{\tilde{p},\tilde{p}'} \hat{a}_{\tilde{p}}^{\dagger} \hat{a}_{\tilde{p}'} + W_{\tilde{p},\tilde{p}'}^* \hat{a}_{\tilde{p}} \hat{a}_{\tilde{p}'} \right) \right]$$

$$= \sum_{q \neq q'} \sum_{\tilde{p},\tilde{p}'} V_{q,q'}(\omega_q - \omega_{q'}) \left( W_{\tilde{p},\tilde{p}'} \left[ \hat{a}_q^{\dagger} \hat{a}_{q'}, \hat{a}_{\tilde{p}}^{\dagger} \hat{a}_{\tilde{p}'}^{\dagger} \right] + W_{\tilde{p},\tilde{p}'}^* \hat{a}_{\tilde{p}} \left[ \hat{a}_q^{\dagger} \hat{a}_{q'}, \hat{a}_{\tilde{p}'} \right] \right)$$

$$(A.10)$$

APPENDIX B	
I	
	THE SECOND APPENDIX

Here comes the second appendix.

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Munich, 22.06.2023 .....