

# Introduction to Blackbody radiation

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## Introduction

Heat is a form of energy and can be transferred from one point to another by means of three methods namely, conduction, convection, and radiation. Radiation is a process of transference of heat from one point to another without the aid of any intervening medium or without affecting the intervening medium if any. All bodies at all temperatures are capable of emitting heat. The heat emitted by radiation from a body due to its temperature is called thermal radiation. In this document we are going to discuss some key concepts related to blackbody radiation. The key topics that we will be discussing are:

1. Properties of thermal radiation
2. Important definitions needed to understand radiation
3. Laws of radiation
  - (a) Kirchhoff's law
  - (b) Stefan's law
4. Fery's blackbody
5. Energy distribution in blackbody radiation
6. Wien's distribution law
7. Rayleigh-Jean's law
8. Planck's hypothesis and quantum theory
9. To deduce Wien's law from Planck's law
10. To deduce Rayleigh-Jean's law from Planck's law

## 1 Properties of thermal radiation

Thermal radiation exhibits the following properties:

1. They travel in straight lines.
2. They do not require any material medium for their propagation.
3. They travel equally in all directions, in a homogeneous medium.
4. They travel with speed of light.
5. They obey inverse square law.
6. They exhibit the phenomenon of reflection and refraction.
7. They exhibit the phenomenon of interference diffraction and polarization.
8. When they fall on matter, heat is developed.
9. They are electromagnetic waves having wavelength greater than that of visible region.

## 2 Important definitions

1. Monochromatic emissive power: Monochromatic emissive power ( $e_\lambda$ ) of a body at a temperature ( $T$ ) for wavelength ( $\lambda$ ) is defined as the energy radiated, in vacuum, per unit time, per unit area and per unit range around wavelength i.e., lying between  $\lambda - \frac{1}{2}$  to  $\lambda + \frac{1}{2}$ . For a body,  $e_\lambda$  will be different for different values of  $\lambda$  and for different values of  $T$ .
2. Emissive power: Emissive power ( $E$ ) of a body at a temperature  $T$  is defined as the total amount of energy for all wavelengths, radiated per unit time, per unit area of the body. If  $dE$  is the amount of energy radiated per second per unit area for wavelength  $d\lambda$ , then

$$dE = e_\lambda d\lambda \quad (1)$$

The emissive power  $E$ , is given by

$$E = \int_0^\infty e_\lambda d\lambda, \quad (2)$$

measured in the units of  $Jm^{-2}s^{-1}$ .

3. Spectral energy density: Spectral energy density  $u_\lambda$  at any point, is defined as the radiant energy per unit volume, around the point, for wavelengths lying in a unit range around  $\lambda$  i.e., in between  $\lambda - \frac{1}{2}$  to  $\lambda + \frac{1}{2}$ .
4. Total energy density: Total energy density  $u$  at any point is defined as the radiant energy per unit volume, around the point for all wavelengths taken together,

$$u = \int_0^\infty u_\lambda d\lambda \quad (3)$$

5. Monochromatic absorptive power: Monochromatic absorptive power  $a_\lambda$  of a body at temperature  $T$  for a wavelength  $\lambda$  is defined as the ratio of amount of radiation absorbed by the surface in a given interval of time to the total amount of radiation falling upon the surface in that same time for wavelength lying in a unit range around  $\lambda$  i.e., in between  $\lambda - \frac{1}{2}$  to  $\lambda + \frac{1}{2}$ .
6. Absorptive power: Absorptive power ( $a$ ) of a body at temperature  $T$  is defined as the ratio of the amount of radiation absorbed on the surface in a given interval of time to the total amount of the radiation falling on the surface in the same time, for all wavelengths,

$$a = \int_0^\infty a_\lambda d\lambda \quad (4)$$

### 3 Laws of radiation

1. Kirchhoff's law: "The ratio of the emissive power to the absorptive power is the same for all surfaces at the same temperature and is equal to the emissive power of a perfect blackbody at that temperature".

If  $e_\lambda$  and  $a_\lambda$  represents the emissive power and absorptive power of a given surface,  $E_\lambda$  and  $A_\lambda$  the corresponding values for perfect black surface at the same temperature, then according to the law,  $\frac{e_\lambda}{a_\lambda} = \frac{E_\lambda}{A_\lambda}$ .

But  $A_\lambda$  for a perfectly blackbody is unity. Hence,  $\frac{e_\lambda}{a_\lambda} = E_\lambda$  where,  $E_\lambda$  is some function of  $\lambda$  and  $T$ .

#### 3.1 Note:

- (a) Though it is proved for a body inside an enclosure, it is valid for all bodies under all conditions for pure temperature radiations, since the absorptive and emissive powers of a body depends only on the nature of the surface and is independent of the surroundings.
  - (b)  $\frac{e_\lambda}{a_\lambda}$  is a constant for all kinds of surfaces for given temperature and wavelength.
  - (c) Since  $\frac{e_\lambda}{E_\lambda} = a_\lambda$ , it follows that, emissivity of a surface is equal; to its absorptive power.
  - (d) Since,  $E_\lambda$  for a given temperature is constant,  $e_\lambda = a_\lambda$ , i.e., emissive power varies as the absorptive power and hence, good emitters must also be a good absorbers.
  - (e) If  $a_\lambda$  is a constant for all wavelengths then  $\frac{e_\lambda}{E_\lambda}$  is a constant, then the body is called a gray body.
  - (f) A small hole in the wall of a constant temperature enclosure acts as a perfect absorber as well as a good emitter of radiation.
2. Stefan's law: "The amount of radiant energy emitted per second per unit area of a surface of a blackbody is directly proportional to the fourth power of its absolute temperature."

$$E \propto T^4 \quad (5)$$

, where,  $E$  is the radiant energy emitted per second per unit area of the blackbody, and  $T$  is the absolute temperature of the body.

The above equation can be written of the form,

$$E = \sigma T^4 \quad (6)$$

where  $\sigma = 5.678 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$ .

The proportionality constant introduced in the above equation is called stefan's constant. Later by combining with theoretical aspect, the law was modified by Boltzmann.

If a blackbody  $B_1$  at temperature  $T_1$  is surrounded by another blackbody  $B_2$  at temperature  $T_2$ , where  $T_1 > T_2$ , then the net amount of energy emitted per second per unit area of a blackbody  $B_1$  is given by

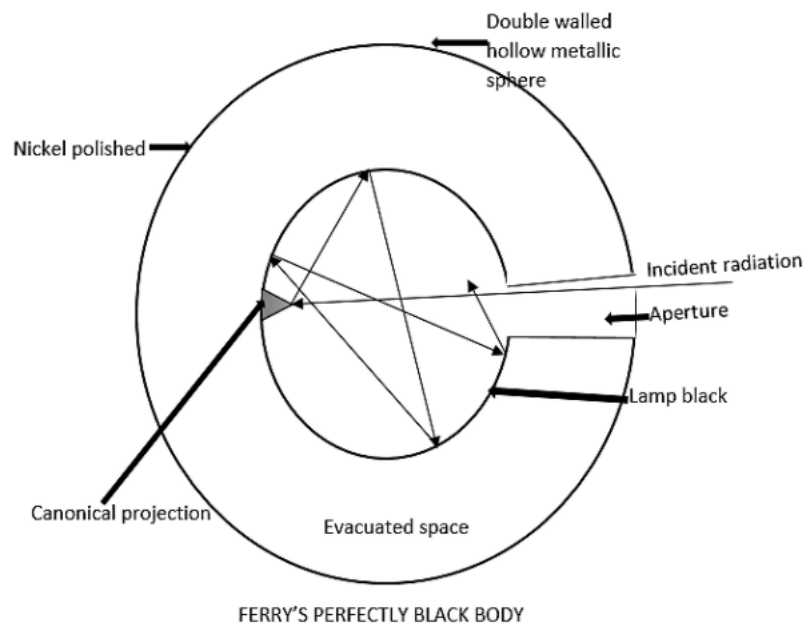
$$E = \sigma(T_1^4 - T_2^4) \quad (7)$$

This is known as Stefan-Boltzmann's law.

## 4 Blackbody

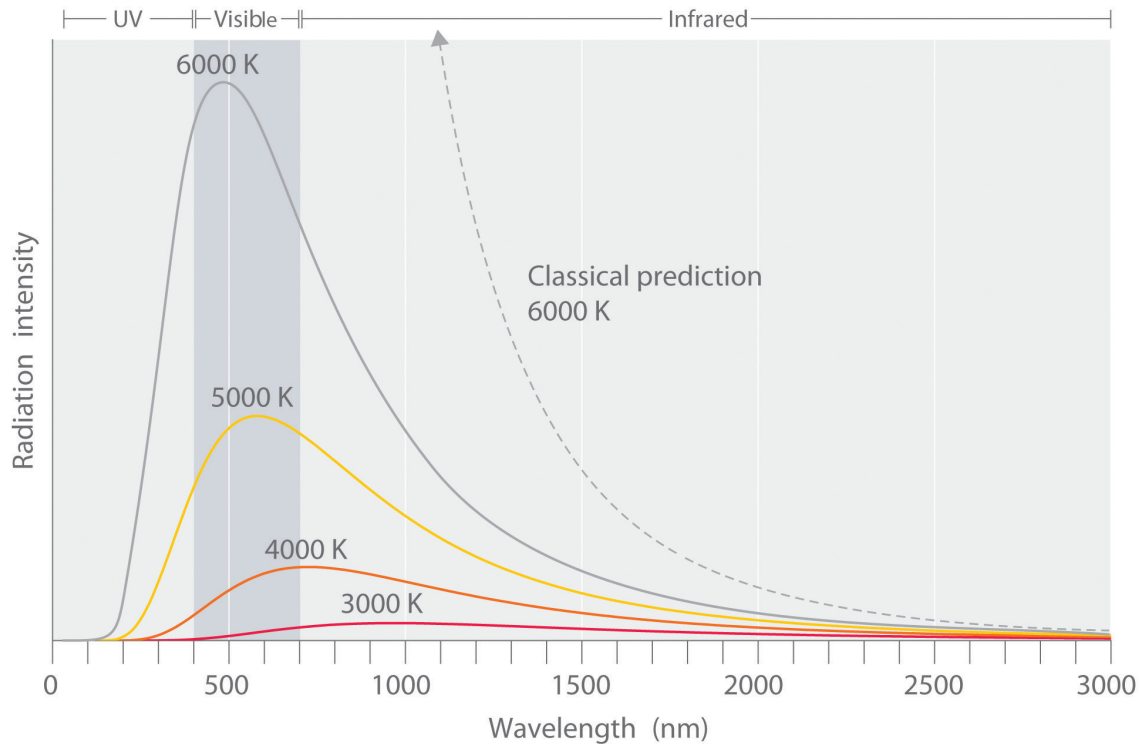
A blackbody is one which absorbs all the radiation (corresponding to all wavelengths) that falls on it. It also emits all radiations when maintained at constant temperature. Such radiations are called blackbody radiations.

In practice, a perfect blackbody is not available. the type of perfect blackbody devised by Fery is shown. It consists of a hallow, double walled copper sphere, coated with lamp black on its inner surface. It has a fine hole on one side and a conical projection just opposite the hole on the other side. When radiations enter the hole, they suffer multiple reflections and are completely absorbed. This body acts as blackbody absorber. When this body is placed in a heat bath at a constant temperature then the heat radiations will come out of the hole. Then the hole acts like a blackbody radiator. It should be remembers that only the hole and not the walls of the body acts as a blackbody radiator.



## 5 Energy distribution curve

Blackbody emits all radiations when maintained at constant temperature. The distribution of energy in a blackbody radiation for different wavelengths and at various temperatures was determined experimentally by Lummer and Pringsheim in 1899. A graph of intensity  $E_\lambda$  and wavelength  $\lambda$  was plotted at different temperatures.



### 5.1 Observations

1. Blackbody emits all types of radiations ranging from lower wavelength to higher wavelength.
2. The energy is not distributed uniformly among the wavelengths.
3. At a given temperature the intensity of radiation increase with increase in wavelength and becomes maximum at a particular wavelength ( $\lambda_m$ ). By farther increasing the wavelength, the intensity of radiation decreases.
4. The wavelength ( $\lambda_m$ ) corresponding to maximum intensity get shifted towards lower wavelengths with increase of temperature.
5. Energy emitted due to all wavelengths increasing with increase in temperature.  
(This is in accordance with Stefan-Boltzmann's law)

## 6 Wien's distribution law

### 6.1 Wien's displacement law

"The wavelength  $\lambda_m$  corresponding to maximum intensity of emission of blackbody radiation is inversely proportional to the absolute temperature  $T$  of the body emitting radiation."

$$\lambda_m \propto \frac{1}{T}$$

or

$$\lambda_m = \frac{b}{T} \quad (8)$$

where,  $b = 2.898 \times 10^3 mK$ . This is called the displacement law because the maximum intensity of radiation emitted by a blackbody gets shifted towards shorter wavelength side with increase in temperature of the body.

### 6.2 Fifth power law

"The maximum energy  $E_m$  of the peak emission is directly proportional to the fifth power of absolute temperature of the blackbody."

$$E_m \propto T^5$$

or

$$E_m = (constant)T^5 \quad (9)$$

Combining the two laws, the energy density of radiation in the wavelength range  $\lambda$  and  $\lambda+d\lambda$  as

$$E_\lambda = (constant)\lambda^5 f(\lambda, T) \quad (10)$$

This is called Wien's distribution law. It explains only the increasing part of the curve, in the shorter wavelength region but failed to explain the decreasing part of the curve, the higher wavelength region.

## 7 Rayleigh-Jean's law

Wien's distribution law failed in explaining the decreasing part of the energy distribution curve, Lord Rayleigh, therefore, sought to determine its form on the basis of classical ideas and was later developed by Jean. According to them the energy density of the blackbody radiation within the wavelength range  $\lambda$  and  $\lambda + d\lambda$  is

$$E_\lambda = 8\pi kT\lambda^{-4}d\lambda \quad (11)$$

This is called Rayleigh-Jean's law of energy distribution of blackbody radiation. This agrees well with the experimental results in the longer wavelength region but fails in the shorter wavelength region.

### 7.1 The Ultraviolet catastrophe

According to the Rayleigh-Jean's law, the energy density is given by,

$$E_\lambda = 8\pi kT\lambda^{-4}d\lambda$$

The energy density of blackbody radiation will continuously increase with the decrease of wavelength and approaches infinity with the  $\lambda \rightarrow 0$ . This shows that all the thermal energy in the universe will manifest itself in the form of short wavelengths. This is against to the experimental observations. This fatal objection to the Rayleigh-Jean's law at shorter wavelength is known as the ultraviolet catastrophe. The concept led classical physics to fail.

## 8 Planck's hypothesis and quantum theory

In order to explain the distribution of energy in the spectrum of blackbody over the entire range of wavelengths, Max Planck proposed a quantum theory of radiation in 1901. It is based on these following assumptions:

1. The chamber which emits radiation is assumed to consists of large number of oscillators, each having its own natural frequency.
2. These oscillators cannot emit or absorb energy continuously.
3. Radiation is absorbed or emitted by oscillators in discrete packets of energy called quanta. Energy of each quantum is  $E = h\nu$ , where  $h = 6.626 \times 10^{-34} Js$  is called the planck's constant and  $\nu$  is the frequency of the radiation emitted or absorbed. Thus, the energy of an oscillator may be  $h\nu, 2h\nu, 3h\nu \dots$  or  $nh\nu$ , where  $n$  is an integer.

The Planck's law of radiation.

$$E_{\lambda}d\lambda = \frac{8\pi hc}{\lambda^5} \left( \frac{1}{e^{\frac{hc}{kT\lambda}} - 1} \right) d\lambda \quad (12)$$