### Clustering of Mg II absorbers

Dual degree project: 2017-18

H. S. Sunil Simha PH13B011

Indian Institute of Technology, Madras Guided by

Dr. L Sriramkumar

Dr. R Srianand, IUCAA, Pune

26th November 2017

#### Contents

- Introduction
  - Mg II absorbers along QSO sightlines
- 2 Theory
  - Cosmological perturbation theory
  - Non-linear collapse
  - Press-Schechter formalism
- Oata Analysis
  - Observed clustering

# Mg II absorbers along QSO sightlines

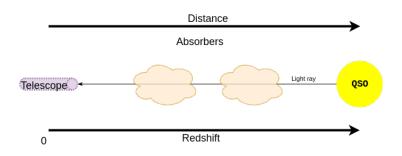


Figure: The absorption system

- Mg II ions have been known to be tracers of neutral gas.
- One of the strongest lines and thus easily spotted.
- Specifically looking at the doublet 2796-2803 Å

## Goals of this project

- To study the distribution of Mg II absorbers in redshift space
- To model this distribution and hence make inferences on the distribution of dark matter halos

Understanding the relationship between cold gas and dark matter halos would help understand structure formation better.

#### Work to be done

Explaining the distribution of these absorbers using a dark matter halo model and a gas distribution.

### Background universe

#### The cosmological principle

The universe is homogeneous and isotropic.

- True at large enough scales: Of the order of 100 MPc.
- Evolution governed by the FLRW equations.

$$H^{2} = H_{0}^{2} \left[ \frac{\Omega_{m0}}{a^{3}} + \frac{\Omega_{r0}}{a^{4}} + \Omega_{\Lambda 0} + \frac{1 - \Omega_{0}}{a^{2}} \right]$$

$$\frac{\ddot{a}}{a} = -\frac{H_{0}^{2}}{2} \left[ \frac{3P}{\rho_{cr_{0}}} + \Omega \right]$$
(1)

- The early universe being largely homogeneous and isotropic is reflected in the CMB.
- $\bullet$  Temperature fluctuations are of the order of  $10^{-5}$  of the average.

□ > <回 > < 亘 > < 亘 > < 亘 < り へ ⊙ </p>

# Cosmological perturbation theory

- Small fluctuations allow for a perturbative treatment.
- Since I am only interested in structures much smaller than the Hubble radius, I can use Newtonian theory.

#### Newtonian limit

The gravitational field obeys Poisson's equation. In terms of co-moving coordinates:

$$\nabla_{\mathbf{x}}^{2}\phi = 4\pi G a^{2}\rho + 3a\ddot{a} \tag{2}$$

Assuming the universe to be filled with fluid,

$$\frac{\partial \rho}{\partial t_x} + 3H\rho + \frac{1}{a} \nabla_x (\rho \mathbf{v}) = 0$$
 (continuity eqn.)

$$\frac{\partial \mathbf{v}}{\partial t_x} + H\mathbf{v} + \frac{(\mathbf{v} \cdot \nabla_{\mathbf{x}})\mathbf{v}}{a} = -\left(\frac{\nabla_{\mathbf{x}}P}{\rho a} + \frac{\nabla_{\mathbf{x}}\phi}{a}\right)$$

(Euler eqn.)

## Linear perturbation theory

Defining the density contrast as  $\delta = \rho/\rho_b - 1$ , and combining the Euler and continuity equations in a matter dominated universe, we get:

$$\partial_t^2 \delta + 2H \partial_t \delta = \frac{\nabla^2 P}{\rho_b a^2} + \frac{1}{a^2} \nabla \cdot (1+\delta) \nabla \phi + \frac{1}{a^2} \partial_i \partial_j [(1+\delta) v^i v^j]$$
 (3)

Linear order: The Meszaros equation

$$\partial_t^2 \delta + 2H \partial_t \delta = \frac{1}{a^2} \left( \frac{\nabla^2 P}{\rho_b} + 4\pi G \rho_b \delta \right) \tag{4}$$

Jean's length:

$$\lambda_J = \sqrt{\frac{\pi}{G\rho_b}} c_s \tag{5}$$

Perturbations of wavelength greater than this grow while the others die out. 4 D F 4 P F F F F F F F

### Non-linear theory: Spherical collapse I

A spherically symmetric, uniformly overdense region is considered. By conservation of energy,

$$\frac{\dot{r}^2}{2} - \frac{GM}{r} = E \tag{6}$$

Starting from a point where  $\dot{r}$  was nearly Hr, the region behaves as if it were a closed universe by itself.

$$r = X(1 - \cos\Theta), t + T = Y(\Theta - \sin\Theta), X^3 = GMY^2$$

Equation of a cycloid!



## Non-linear theory: Spherical collapse II

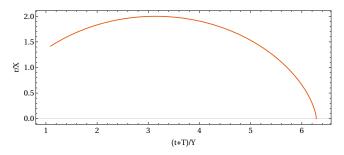


Figure: Evolution of a spherically overdense region.

Of course, collapse stops before r=0 because of pressure generated by fluid. The system **virializes** and comes to a halt at  $r=r_{max}/2$  with a density of  $170\rho_b(t_{coll})$  in a matter dominated universe.

4 D > 4 D > 4 E > 4 E > E 99 C

#### Press-Schechter formalism

The PS formalism estimates the number of objects collapsed within a mass range of  $[M, M + \delta M]$  at a given redshift.

#### Press-schechter mass distribution

$$\frac{dn}{dM} = -\sqrt{\frac{2}{\pi}} \frac{\rho_m}{M} \frac{d\sigma}{dM} \frac{\delta_c}{\sigma^2} e^{\delta_c^2/2\sigma^2}$$
 (7)

- $\rho_m/M$  is the average number density of objects of mass M.
- $\bullet$   $\sigma$  is the variance of the linear power spectrum of density perturbations smoothed by a window function.
- $\delta_c$  is the initial critical over-density above which non-linear collapse happens. This is generally taken to be 1.686



# Observed clustering

- Data of over 30,000 QSO sightlines taken from SDSS.
- Counted pairs along LoS in redshift space as a function of velocity separation.
- Calculated the expected histogram from current model of redshift space distribution.

#### Completeness correction

- Surveys tend to report less objects than there are simply because of sensitivity limitations.
- Used data from MC simulations by G.B.
   Zhu to correct for this.

## Observed clustering

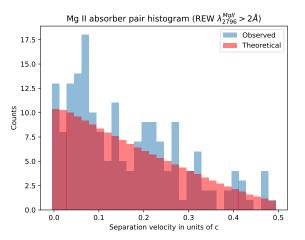


Figure: A comparison of theoretical and observed histograms of pairs of Mg II absorbers. The theoretical estimate has been obtained from the empirical distribution of absorbers in Zhu & Menard, The Astrophysical Journal, 770:130 (15pp), 2013 June 20

12 / 13

### References

- Zhu, Menard, The JHU-SDSS Metal Absorption Line Catalogue, The Astrophysical Journal, 770:130 (15pp), 2013 June 20
- Nestor, Turnshek & Rao, Mg II absorption systems in SDSS QSO spectra, The Astrophysical Journal, 628:637654, 2005 August 1
- **T. Padmanabhan**, Structure formation in the universe, Cambridge 1993
- Jim Peebles, Large Scale Structure of the Universe, Princeton 1992
- Peter Schneider, Extragalactic Astronomy and Cosmology, Springer 2006