Specific Heat and Calorimetry; Heat engine overview

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Outline

1. Specific heat and calorimetry

2. Specific heat of ideal gas

3. Heat engine overview

1. Specific heat and calorimetry

Specific heat and molar specific heat

- The heat it takes to raise the temperature of 1 kg of a given substance by 1 K is called the **specific heat** c of that substance. In equation:

$$Q = Mc\Delta T$$

- Similarly, the heat it takes to raise the temperature of 1 mole of a substance by 1 K is its **molar specific heat** C. In equation:

$$Q = nC\Delta T$$

Specific heat: units and conventions

- The SI unit for specific heat is J/kg·K

- The SI unit for molar specific heat is J/mol·K

- Since specific heat is concerned only with change in temperature, there is no real difference between measurement in °C and K

Specific heat and thermal energy

- Since **solids** and **liquids** are essentially incompressible, we may assume $W = 0$ when they are heated or cooled. By the first law of thermodynamics, we thus have:

$$\Delta E_{\text{th}} = Q = Mc\Delta T$$

- In contrast, a gas expands when heated under constant pressure, which turns out to contribute $W < 0$. As a result, we generally have $c_P > c_V$ in a gas
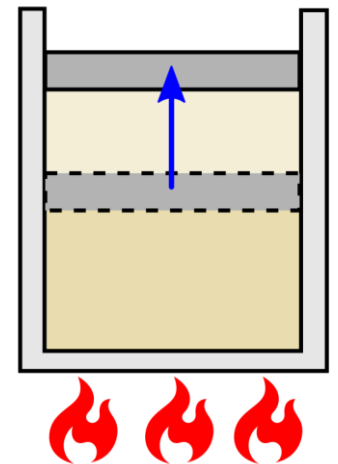
Heat of transformation

- When a substance is undergoing phase change, additional energy is needed to reorganize its microscopic structure

- The **heat of transformation** L is the additional heat needed for this process for each 1 kg of the substance. In equation:

$$Q = ML$$
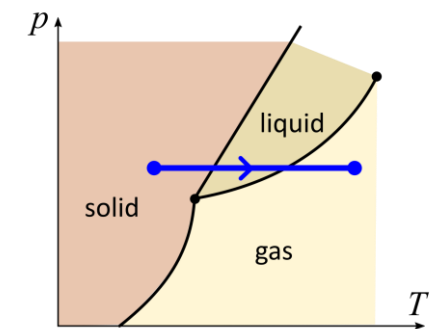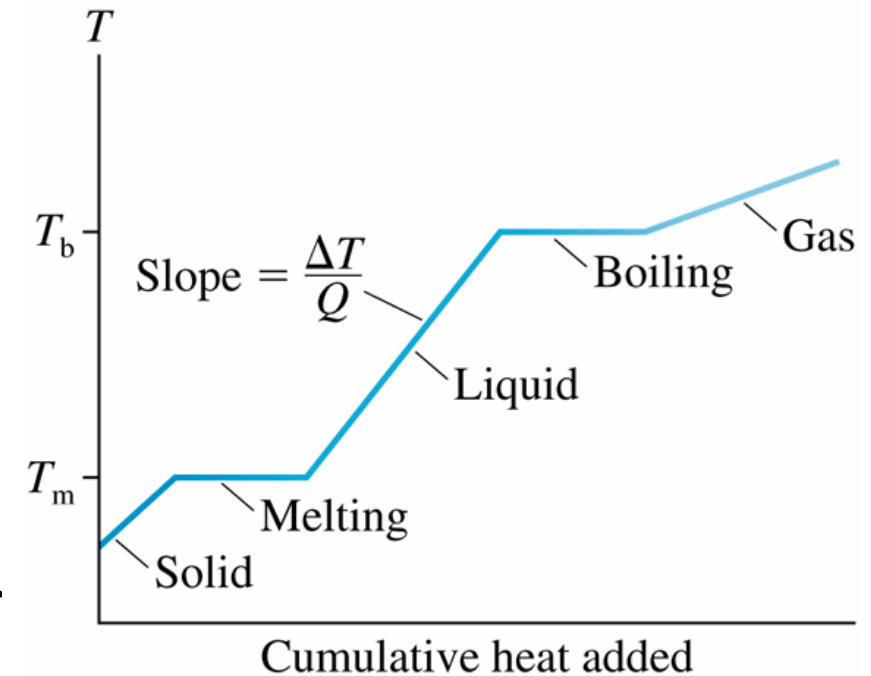
Heat of transformation: signs and conventions

- Heat of transformation is also known as **latent heat**

- The SI unit for heat of transformation is is J/kg

- In general, L is different for different phase change. Also, note the sign reverse between heating and cooling. Thus,

$$Q = \pm ML_f \qquad (+ \text{ melt } / - \text{ freeze})$$

$$Q = \pm ML_v \qquad (+ \text{ boil } / - \text{ condense})$$
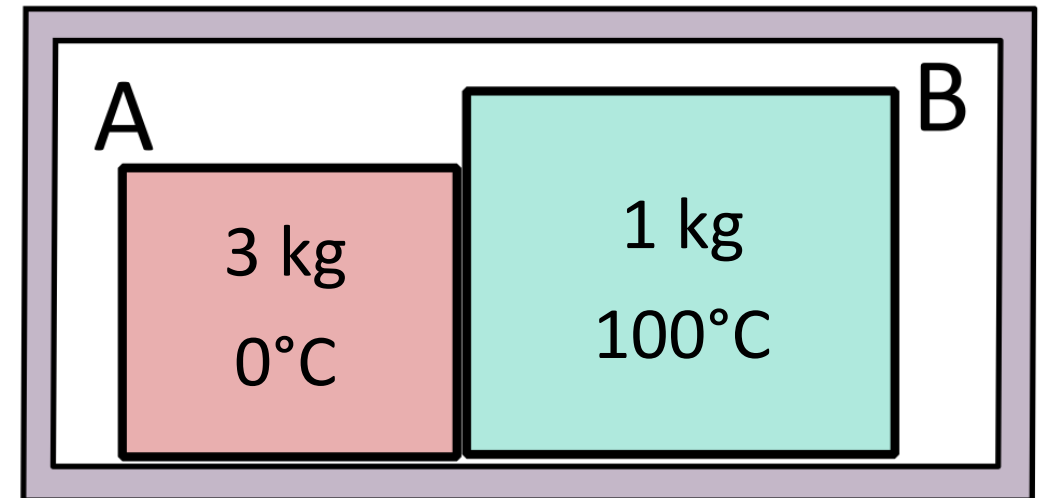
Calorimetry

- When two objects are in thermal contact with each other but isolated from everything else, we must have $|Q_1| = |Q_2|$

- Regardless of the initial temperatures, when thermal equilibrium is reached we must have $T_1 = T_2$

2. Specific heat and calorimetry
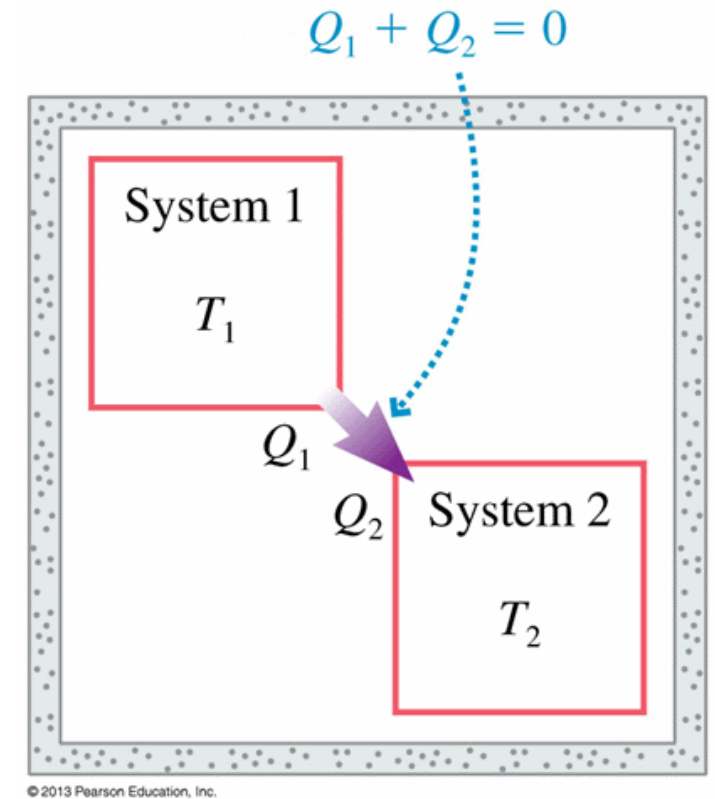
Your turn: heat capacity and heat exchange

Two objects A and B are brought into thermal contact with each other but are otherwise well-isolated from the surrounding. If $c_B = 2c_A$, the final (equilibrium) temperature is:

A. $T_f < 50°C$

B. $T_f = 50°C$

C. $T_f > 50°C$

Numerical practice: calorimetry

A 750 g aluminum pan ($c = 900$ J/kg·K) are removed from the stove and plunged into a sink filled with 10.0 L of water ($c = 4190$ J/kg·K) at 20°C. The water temperature quickly rises before it stabilizes at 24.0°C. What was the initial temperature of the pan in °C?
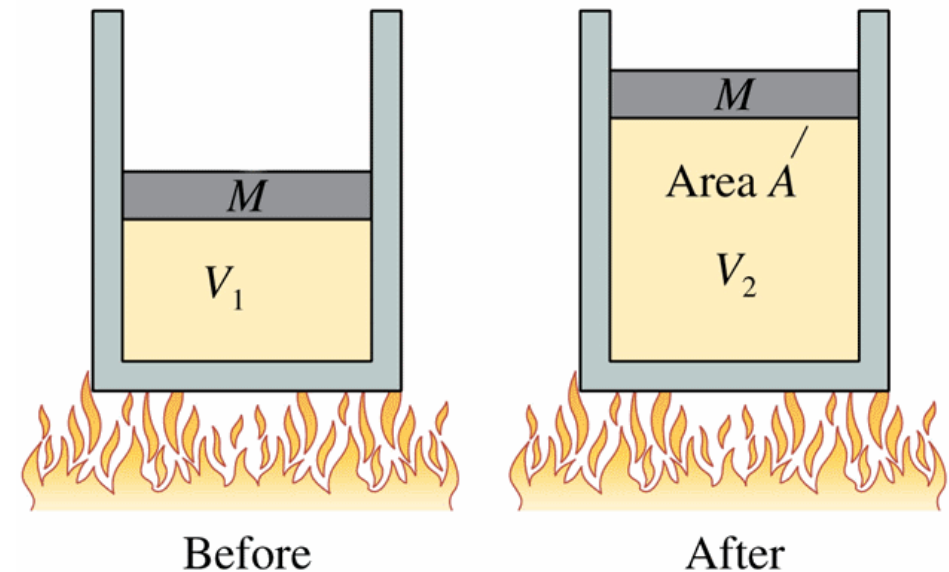
2. Specific heat of ideal gas

Reminder: isochoric vs isobaric processes

- Isochoric ($V = $ const.) process:
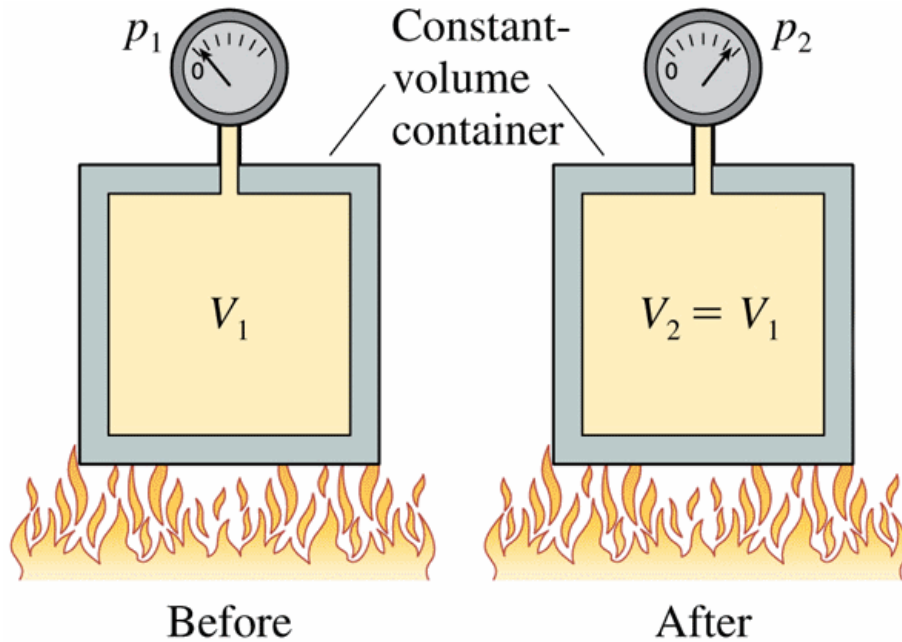
$$Q = \Delta E_{\text{th}} = \frac{3}{2} nR\Delta T$$

- Isobaric ($p = $ const.) process:

$$Q = \Delta E_{\text{th}} - W = \frac{5}{2} nR\Delta T$$

Molar specific heat for monoatomic ideal gas

- Recall the definition of molar specific heat C: $Q = nC\Delta T$

- For constant volume process,

$$Q = \Delta E_{\text{th}} = \frac{3}{2} nR\Delta T$$

- Thus (for monoatomic gas)...

$$C_V = \frac{3}{2} R$$

- For constant pressure process,

$$Q = \Delta E_{\text{th}} = \frac{5}{2} nR\Delta T$$

- Thus (for monoatomic gas)...

$$C_P = \frac{5}{2} R$$

Relationship between C_P and C_V

- From $C_V = \frac{3}{2}R$ and $C_P = \frac{5}{2}R$, we see that $C_P = C_V + R$

- This difference stems from the work done in the constant pressure case, where in $W = -p\Delta V = -nR\Delta T$

- Note that no monoatomic assumption is made. So in general:

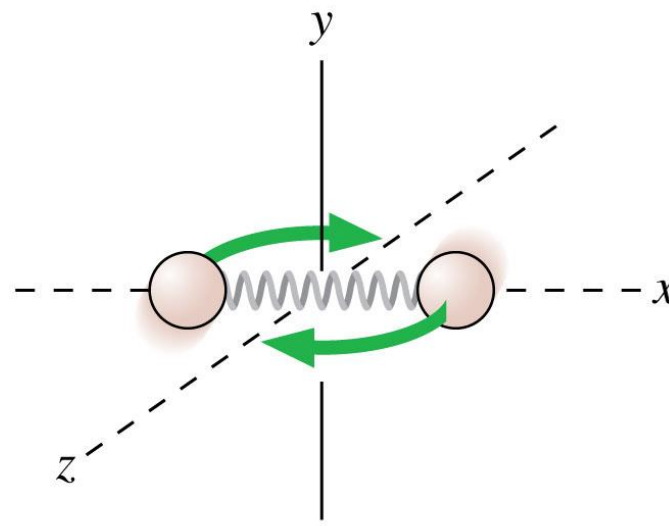$$\boxed{C_P = C_V + R}$$

(**all** ideal gas)

Beyond monoatomic gas: new degrees of freedom

- In addition to an overall translational motion, diatomic molecule can rotate about 2 of its axes and can vibrate about its bond axis
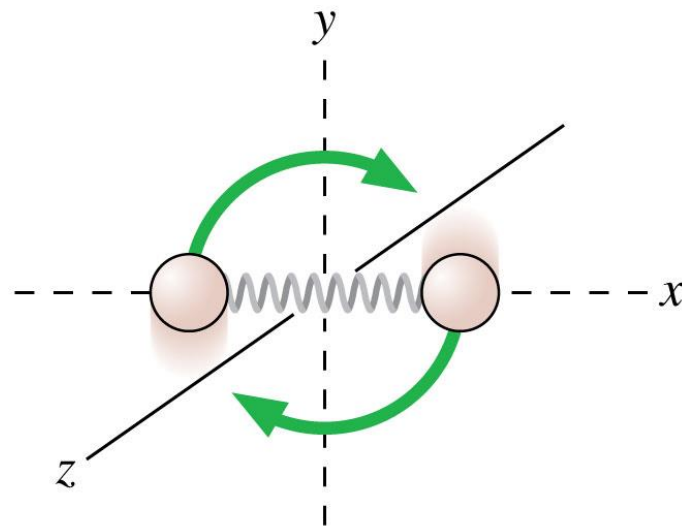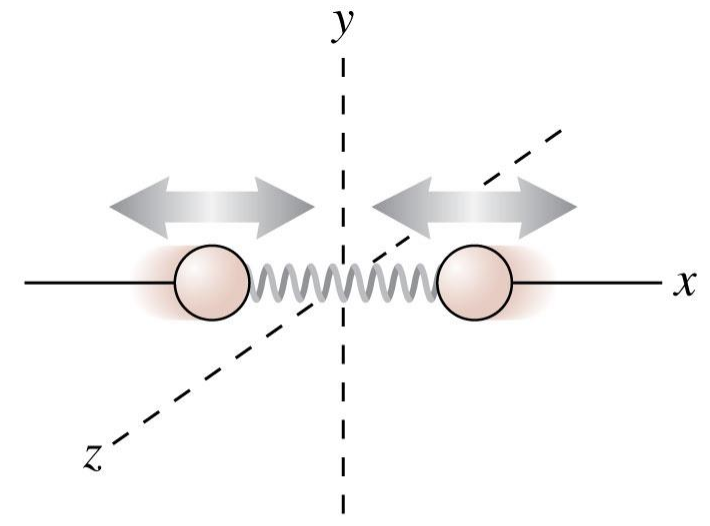

Rotation end-over-end about the z-axis

Rotation end-over-end about the y-axis

Vibration back and forth along the x-axis

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Beyond monoatomic gas: counting degrees of freedom

- Each additional "degree of freedom" can store energy and thus contributes additionally to E_th and hence C_V

- For diatomic molecule, the counting goes:

  3 translation + 2 rotation + 2 vibration = 7 degrees of freedom

  1 potential + 1 kinetic

- For translation, we get $\frac{1}{2} k_B T$ per degree of freedom per molecule, how about the other new degrees of freedom?

Beyond monoatomic gas: the equipartition theorem

- The **equipartition theorem** from *classical* statistical mechanics says:

  The thermal energy of a system of particles is equally divided among all degrees of freedom. More specifically, the energy stored in each degree of freedom is $\frac{1}{2} k_B T$ per particle

- However, because of *quantum* mechanics, the contributions from rotation and vibration are "activated" only at higher temperatures
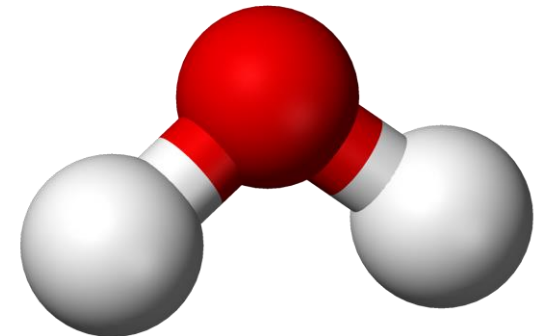
Beyond the monoatomic gas: putting all together
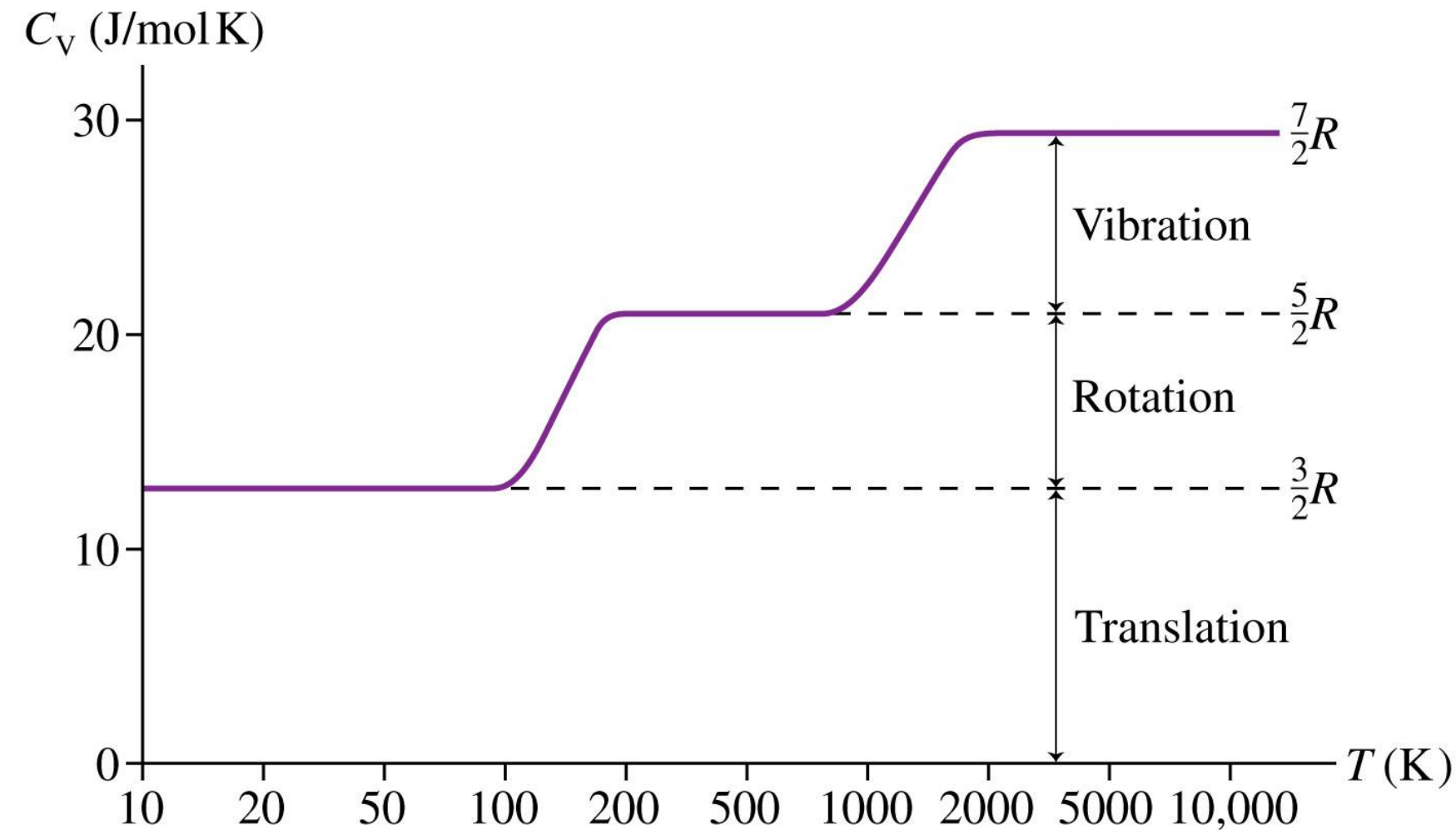
Upshot: use $C_V = \dfrac{5}{2}R$

for diatomic gas at "typical" temperatures 200 K $\lesssim T \lesssim$ 800K

Your turn: counting degrees of freedom

Consider the water molecule (see figure). How many **rotational** degrees of freedom does it has?

A. 2 degrees of freedom

B. 3 degrees of freedom

C. 4 degrees of freedom

D. 6 degrees of freedom

3. Heat engine overview

Basic properties of a heat engine

- **Heat engine** converts heat into useful work

- Conceptually, car engines, steam generators, etc. are all heat engine of some kind

- A heat engine operates in **cycles**, i.e., it returns to its initial state periodically

- A heat engine general operates between a **hot reservoir** and a **cold reservoir**

Energy-transfer diagram and relevant quantities

- T_H = temperature of hot reservoir

- T_C = temperature of cold reservoir

- Q_H = heat absorbed from hot reservoir

- Q_C = heat released to cold reservoir

- W_out = useful work output

*Q_H, Q_C, and W_out are values **per cycle**, and are all taken to be **positive**
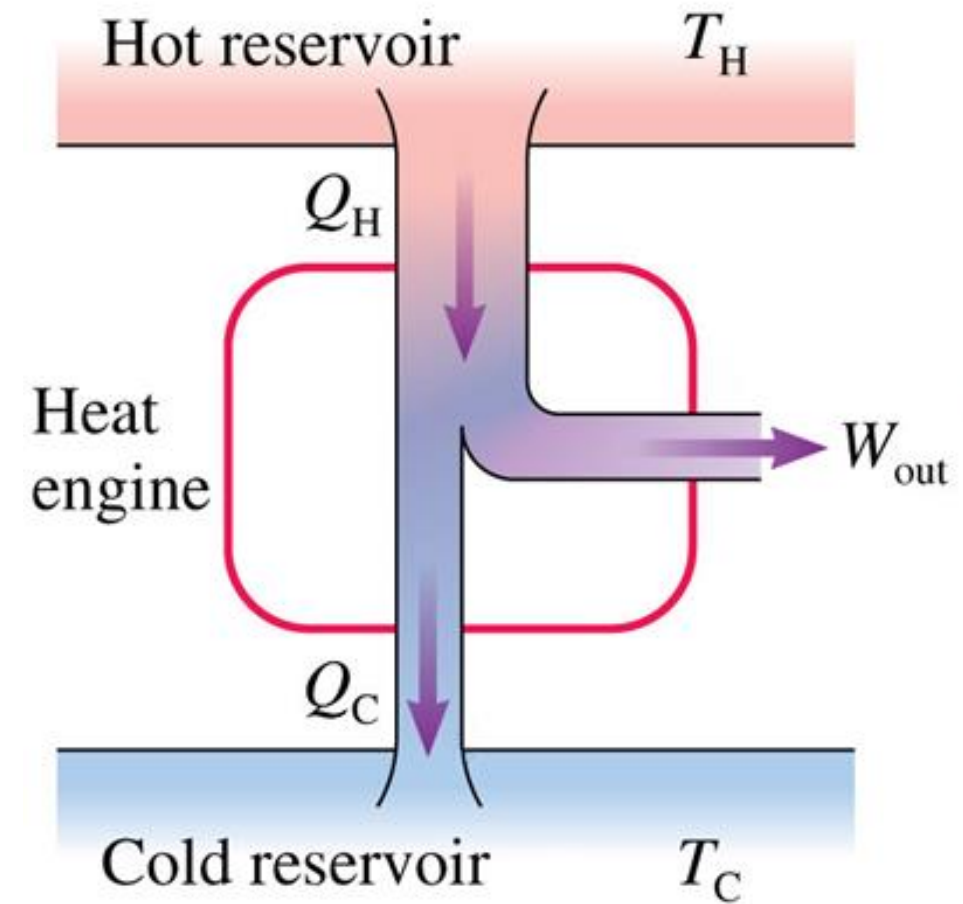
First law of thermodynamics on heat engine

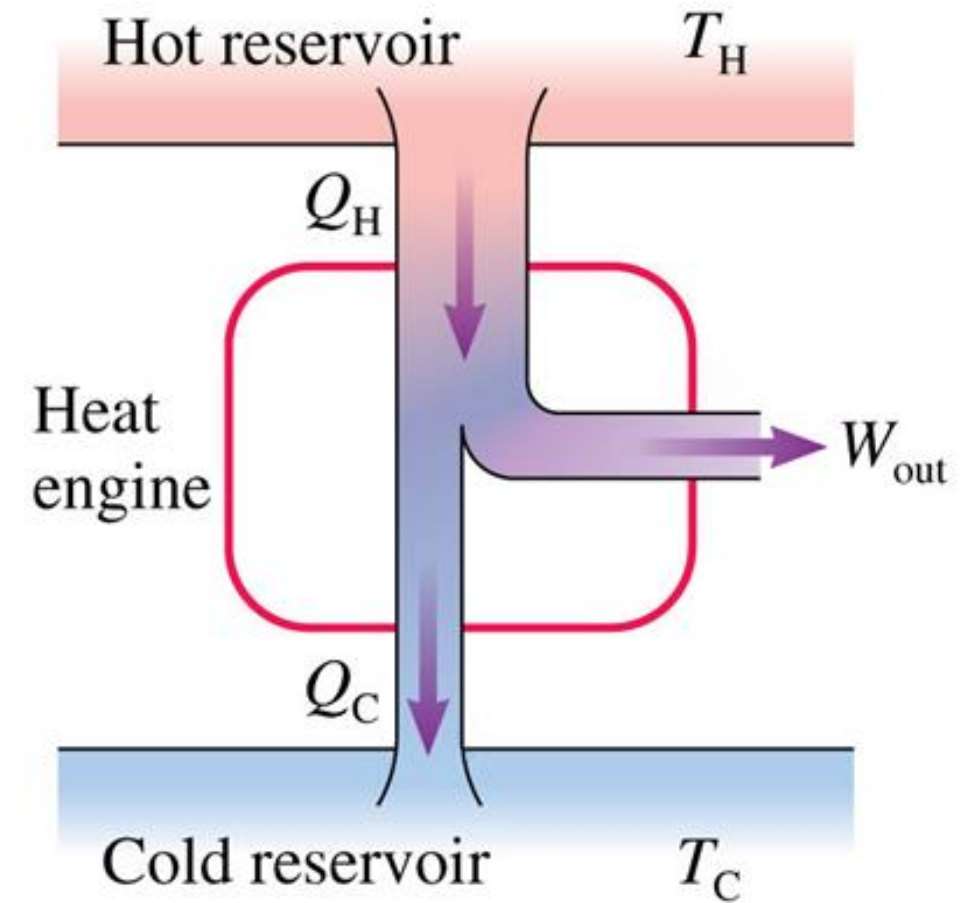- Relating our definitions,

$$Q = Q_H - Q_C \quad \text{(per cycle)}$$

$$W = -W_{\text{out}}$$

- Since the heat engine is cyclic,

$\Delta E_{\text{th}} = 0$ after each cycle

- Thus, 1st law implies:

$$\boxed{W_{\text{out}} = Q_H - Q_C}$$

Hot reservoir $\qquad T_{\text{H}}$

Q_{H}

Heat engine

W_{out}

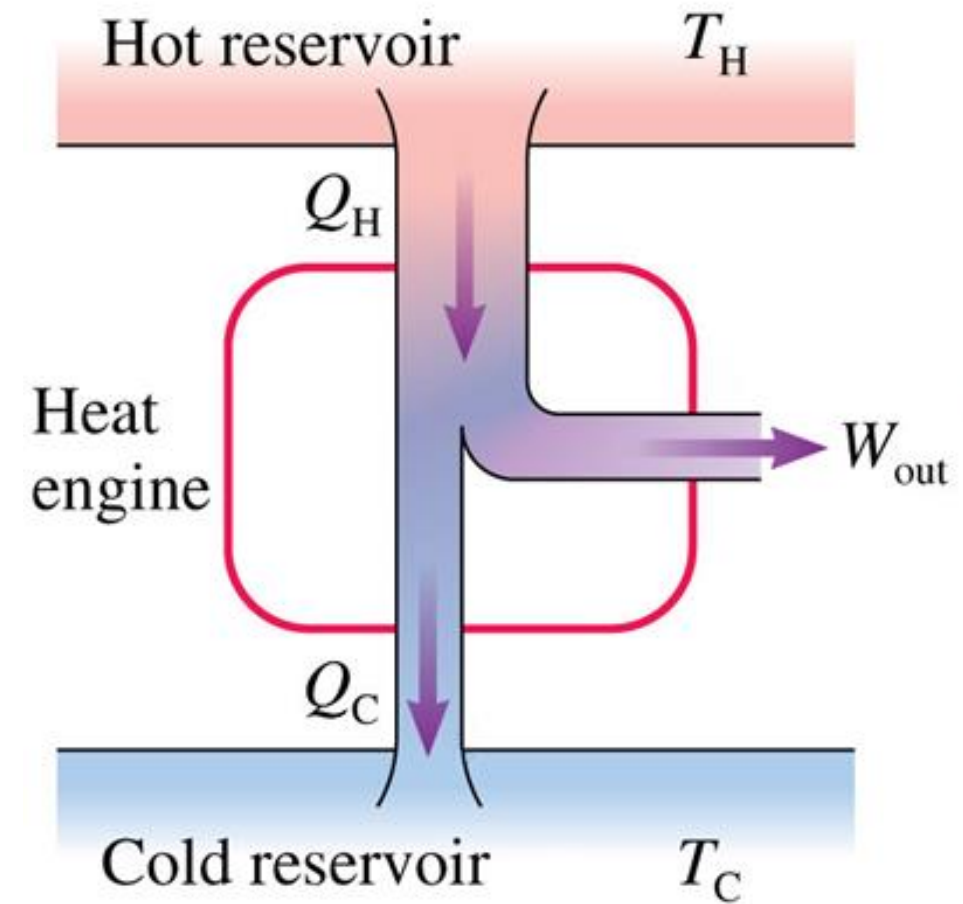Q_{C}

Cold reservoir $\qquad T_{\text{C}}$

Heat engine efficiency

- The **thermal efficiency** of a heat engine, η, is defined as:

$$\eta = \frac{W_{\text{out}}}{Q_H}$$

- One important question for heat engine is whether η is limited

Your turn: heat engine efficiency

Rank the following heat engines by thermal efficiency, from largest to smallest

A. (II) > (III) = (IV) > (I)

B. (II) = (IV) > (I) = (III)

C. (III) > (I) = (II) > (IV)

D. (IV) > (II) > (I) > (III)