

Brachistochrone

Sylvain Vanneste

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Problem set up

On an other planet, similar to ours, some cities are built up in the sky. The house and small neighborhood are isolated on floating rocks.

Akara lives up a cliff, and she would like to send gifts via a roller-coaster to her friend Buni which lives just a little further down in the valley. She built a tubular structure with super-conductive magnets, such that there is no air friction, as well as no contact friction (since the wagons are levitating). There is no engine, and the packages are simply dropped with no initial speed on the rail and left to the gravitational pull of earth.

However, she would like to know which shape to give to its roller-coaster in order to minimise the time a package takes to arrive to her friend.

In order to picture the problem, we will consider the origin of the referential placed at Akara's house, and Buni's coordinates are noted (x_B, y_B) . For the problem to be possible, we consider that Buni's house is lower than Akara's. That is to say, $y_B < 0$

Historical introduction

Johann Bernoulli posed the same problem in June, 1696. Well, without the roller-coaster and superconductive magnets part :

"Given two points A and B in a vertical plane, what is the curve traced out by a point acted on only by gravity, which starts at A and reaches B in the shortest time.

At 4 p.m. on 29 January 1697 when he arrived home from the Royal Mint, Isaac Newton found the challenge in a letter from Johann Bernoulli. Newton stayed up all night to solve it and mailed the solution anonymously.

In an attempt to outdo his brother, Jakob Bernoulli created a harder version of the brachistochrone problem. In solving it, he developed new methods that were refined by Leonhard Euler into what the latter called (in 1766) the calculus of variations. Joseph-Louis Lagrange did further work that resulted in modern infinitesimal calculus.

Personal notes

We want to minimise the time it take for an object to goes from A to B , t_{AB} . We define v as the speed of the object, g the gravitational constant, and m the object's mass. The trajectory of the object can be defined by relation between its horizontal and vertical position, $y = f(x)$. For simplicity, we will write $y = y(x) = f(x)$.

Using energy conservation, we have

$$\begin{aligned} E_p + E_c &= E_{pi} + E_{ci} \\ \Leftrightarrow mgy + \frac{1}{2}mv^2 &= mgy_A + 0 \\ \Leftrightarrow v(y) &= \sqrt{2g(y_A - y)}, \end{aligned} \quad (1)$$

where $v(y)$ is a function of the object heigh y .

The time t_{AP} for the object to go from point A to an arbitrary position P , with $P \in [A, B]$, is written

$$t_{AP} = \int \frac{ds}{v(y)}, \quad (2)$$

with ds the infinitesimal distance increment of the object. It can be expressed using Pythagorean law

$$\begin{aligned} ds^2 &= dx^2 + dy^2 \\ \Leftrightarrow ds &= \sqrt{1 + \frac{dy^2}{dx^2}} dx \\ &= \sqrt{1 + y'^2} dx, \end{aligned} \quad (3)$$

where the derivative of the function is written $y' = dy/dx = f'(x)$.

Using Eqs. (1) and (3), we can now simplify the time Eq. (2),

$$\begin{aligned} t_{AP} &= \int_A^P \frac{ds}{v(y)}, \\ &= \int_{x_A}^{x_P} \frac{\sqrt{1 + y'^2(x)}}{\sqrt{2g(y_A - y(x))}} dx, \\ &= \frac{1}{\sqrt{2g}} \int_{x_A}^{x_P} \frac{\sqrt{1 + y'^2(x)}}{\sqrt{(y_A - y(x))}} dx, \end{aligned} \quad (4)$$

Therefore, the time t_{AB} the object takes to go from point A to B is depends on the trajectory $y(x)$. Seeking a minimising this time is equivalent to seek for a function $y(x)$ which minimise the Eq. (4) (where $x_P = x_B$).

Before seeking for such solution, we first compare the physical behaviour for some known functions $y(x)$, namely : the straight line, the reciprocal function, and a second order polynomial function.

The straight line

The straight line between point A and B can be expressed as

$$y = \alpha x + \beta, \quad (5)$$

where $\alpha = (y_B - y_A)/(x_B - x_A)$ and $\beta = y_A - \alpha x_A$. The function derivative is therefore $y' = \alpha$, and the time for an object to reach a point P is written

$$\begin{aligned} t(x) &= \frac{1}{\sqrt{2g}} \int_{x_A}^{x_P} \frac{\sqrt{1 + y'^2(x)}}{\sqrt{(y_A - y(x))}} dx, \\ &= \frac{1}{\sqrt{2g}} \int_{x_A}^{x_P} \frac{\sqrt{1 + \alpha^2}}{\sqrt{y_A - \alpha x - \beta}} dx, \\ &= \sqrt{\frac{1 + \alpha^2}{2g}} \int_{x_A}^{x_P} \frac{1}{\sqrt{y_A - \alpha x - \beta}} dx, \\ &= -\sqrt{\frac{1 + \alpha^2}{2g}} \frac{2}{\alpha} \left[\sqrt{y_A - \alpha x - \beta} \right]_{x_A}^{x_P}, \\ &= -\sqrt{\frac{1 + \alpha^2}{2g}} \frac{2}{\alpha} \left[\sqrt{y_A - \alpha x_P - \beta} - \sqrt{y_A - \alpha x_A - \beta} \right] \end{aligned} \quad (6)$$

Conversely, the position x_P can be written as a function of the time,

$$x_P(t) = -\frac{1}{\alpha} \left[\left(-t \frac{\alpha}{2} \sqrt{\frac{2g}{1 + \alpha^2}} + \sqrt{-\alpha x_A - \beta + y_A} \right)^2 - \beta + y_A \right]. \quad (7)$$

The reciprocal function

We now choose the reciprocal function

$$y(x) = \frac{1}{x + \gamma} + \eta, \quad (8)$$

therefore $y'(x) = -(x + \gamma)^{-2}$.

Since $x_A = 0$ and $y_A = 0$, we have $\eta = -1/\gamma$, and $\gamma = \pm \frac{\sqrt{x_B} \sqrt{-x_B y_B + 4}}{2\sqrt{x_B}} - \frac{x_B}{2}$. We choose the $+$ sign for the second condition in order for $y(x)$ to be convex.

The time for an object to reach a point P is computed by numerical integration, following

$$t(x) = \frac{1}{\sqrt{2g}} \int_{x_A}^{x_P} \frac{\sqrt{1 + y'^2(x)}}{\sqrt{(y_A - y(x))}} dx. \quad (9)$$

$$(10)$$

2nd order polynomial function

We now choose a 2nd order polynomial function of the form

$$y(x) = ax^2 + bx + c, \tag{11}$$

with the condition $y(x) \leq y_A$. Therefore $y'(x) = 2a + b$. We also choose $y(x)$ to be convex, therefore $y(x)'' \geq 0$, which leads to $a \geq 0$. Since $x_A = 0$ and $y_A = 0$, we have $c = 0$. Finally, $b = y_B - ax_B^2$.

Refs