ECE 697 Modeling and High-Performance Control of Electric Machines HW 3 Solutions Spring 2022

Problem 16 A Two-Phase Four-Pole Generator

(a) Using an outward normal for the flux surface $(\hat{\mathbf{n}} = \hat{\mathbf{r}})$, the flux linkage in stator loop a - a' due to the rotor's magnetic field is

$$\begin{split} \lambda_{a} &= \int_{0}^{\ell_{1}} \int_{-\pi/2}^{0} B_{R \max} \cos \left(n_{p}(\theta - \theta_{R})\right) \hat{\mathbf{r}} \cdot \left(r_{S} d\theta d\ell \hat{\mathbf{r}}\right) + \int_{0}^{\ell_{1}} \int_{\pi/2}^{\pi} B_{R \max} \cos \left(n_{p}(\theta - \theta_{R})\right) \hat{\mathbf{r}} \cdot \left(r_{S} d\theta d\ell \hat{\mathbf{r}}\right) \\ &= \left. \frac{r_{S} \ell_{1} B_{R \max}}{n_{p}} \sin \left(n_{p}(\theta - \theta_{R})\right) \right|_{-\pi/2}^{0} + \left. \frac{r_{S} \ell_{1} B_{R \max}}{n_{p}} \sin \left(n_{p}(\theta - \theta_{R})\right) \right|_{\pi/2}^{\pi} \\ &= \left. \frac{r_{S} \ell_{1} B_{R \max}}{n_{p}} \left(-\sin \left(n_{p} \theta_{R}\right) + \sin \left(n_{p} \left(\pi/2 + \theta_{R}\right)\right) + \sin \left(n_{p} \left(\pi - \theta_{R}\right)\right) - \sin \left(n_{p} \left(\pi/2 - \theta_{R}\right)\right)\right) \right) \\ &= \frac{r_{S} \ell_{1} B_{R \max}}{2} \left(-\sin \left(2\theta_{R}\right) + \sin \left(2\left(\pi/2 + \theta_{R}\right)\right) + \sin \left(2\left(\pi - \theta_{R}\right)\right) - \sin \left(2\left(\pi/2 - \theta_{R}\right)\right)\right) \\ &= \frac{r_{S} \ell_{1} B_{R \max}}{2} \left(-4\sin(2\theta_{R})\right) \\ &= -2r_{S} \ell_{1} B_{R \max} \sin(2\theta_{R}). \end{split}$$

(b) Does the positive direction of travel around each of the two flux surfaces of phase a coincide with the positive direction of current? Yes

(c)

$$\xi_{a-a'} = -\frac{d\lambda_a}{dt} = 4r_S \ell_1 B_{R\max} \omega_R \cos(2\theta_R) = 4r_S \ell_1 B_{R\max} \omega_R \cos(2\omega_R t)$$

(d) Similarly $(n_p = 2)$

$$\begin{split} \lambda_b(\theta_R) &= \int_0^{\ell_1} \int_{-\pi/4}^{+\pi/4} B_{R \max} \cos \left(n_p(\theta - \theta_R) \right) \hat{\mathbf{r}} \cdot \left(r_S d\theta d\ell \hat{\mathbf{r}} \right) + \int_0^{\ell_1} \int_{3\pi/4}^{5\pi/4} B_{R \max} \cos \left(n_p(\theta - \theta_R) \right) \hat{\mathbf{r}} \cdot \left(r_S d\theta d\ell \hat{\mathbf{r}} \right) \\ &= \left. \frac{r_S \ell_1 B_{R \max}}{n_p} \sin \left(n_p(\theta - \theta_R) \right) \right|_{-\pi/4}^{+\pi/4} + \left. \frac{r_S \ell_1 B_{R \max}}{n_p} \sin \left(n_p(\theta - \theta_R) \right) \right|_{3\pi/4}^{5\pi/4} \\ &= \left. \frac{r_S \ell_1 B_{R \max}}{n_p} \left(\sin \left(2(\pi/4) - 2\theta_R \right) \right) - \sin(2\left(-\pi/4 \right) - 2\theta_R) \right) + \sin(2\left(5\pi/4 \right) - 2\theta_R) - \sin(2\left(3\pi/4 \right) - 2\theta_R) \right) \\ &= \frac{r_S \ell_1 B_{R \max}}{2} \left(\cos \left(2\theta_R \right) + \cos(2\theta_R) + \cos(2\theta_R) + \cos(2\theta_R) \right) \\ &= \frac{r_S \ell_1 B_{R \max}}{2} \left(4\cos(2\theta_R) \right) \\ &= 2r_S \ell_1 B_{R \max} \cos(2\theta_R). \end{split}$$

Note that

$$\lambda_b(\theta_R) = \lambda_a(\theta_R - \pi/4) = -2r_S\ell_1 B_{R\max} \sin(2(\theta_R - \pi/4)) = 2r_S\ell_1 B_{R\max} \cos(2\theta_R).$$

The induced voltage in phase b - b' is

$$\xi_{b-b'} = -d\lambda_b/dt = 4r_S \ell_1 B_{R\max} \omega_R \sin(2\theta_R) = 4r_S \ell_1 B_{R\max} \omega_R \sin(2\omega_R t).$$

(e) With the electric field given by

$$\vec{\mathbf{E}}(\theta - \theta_R) = \omega_R B_{R \max} r_S \cos\left(n_p(\theta - \theta_R)\right) \hat{\mathbf{z}}$$

the voltage in phase a - a' is computed as

$$\begin{split} \xi_{a-a'} &= \int_{a'}^{a} \vec{\mathbf{E}}(\theta - \theta_R) \cdot d\vec{\boldsymbol{\ell}} \\ &= \int_{side\ a'_1} \left(\omega_R B_{R\max} r_S \cos\left(n_p \left(-\pi/2 - \theta_R \right) \right) \hat{\mathbf{z}} \right) \cdot \left(-d\ell \hat{\mathbf{z}} \right) + \int_{side\ a_1} \left(\omega_R B_{R\max} r_S \cos\left(n_p \left(0 - \theta_R \right) \right) \hat{\mathbf{z}} \right) \cdot \left(d\ell \hat{\mathbf{z}} \right) \\ &+ \int_{side\ a'_2} \left(\omega_R B_{R\max} r_S \cos\left(n_p \left(\pi/2 - \theta_R \right) \right) \hat{\mathbf{z}} \right) \cdot \left(-d\ell \hat{\mathbf{z}} \right) + \int_{side\ a_2} \left(\omega_R B_{R\max} r_S \cos\left(n_p \left(\pi - \theta_R \right) \right) \hat{\mathbf{z}} \right) \cdot \left(d\ell \hat{\mathbf{z}} \right) \\ &= \omega_R B_{R\max} r_S \ell_1 \left(-\cos\left(2 \left(-\pi/2 - \theta_R \right) \right) + \cos\left(2\theta_R \right) - \cos\left(2 \left(\pi/2 - \theta_R \right) \right) + \cos\left(2 \left(\pi - \theta_R \right) \right) \right) \\ &= \omega_R B_{R\max} r_S \ell_1 \left(-\cos\left(\pi + 2\theta_R \right) + \cos\left(2\theta_R \right) - \cos\left(\pi - 2\theta_R \right) + \cos\left(2\theta_R \right) \right) \\ &= 4r_S \ell_1 B_{R\max} \omega_R \cos\left(2\theta_R \right) \\ &= 4r_S \ell_1 B_{R\max} \omega_R \cos\left(2\omega_R t \right) \end{split}$$

which is the same result as using Faraday's law in part (c).