

# Turbulence Analysis

July 26, 2022

## 1 Analysis of Turbulence Data

```
[ ]: import numpy as np
import matplotlib.pyplot as plt
import pandas as pd
import statsmodels.api as sm
from scipy import signal
from lmfit import Model
# https://lmfit.github.io/lmfit-py/builtin_models.html
# https://cars9.uchicago.edu/software/python/lmfit/builtin_models.html

import lmfit
label_size = 18
plt.rcParams.update(
    {
        "font.size": label_size,
        "legend.title_fontsize": label_size,
        "legend.fontsize": label_size,
        "axes.labelsize": label_size,
        "xtick.labelsize": label_size,
        "ytick.labelsize": label_size,
        "axes.labelpad": 4,
        # "lines.markersize": 13,
        "lines.linewidth": 2,
    }
)
```

## 2 Exercise 5

```
[ ]: data = np.loadtxt('Data/k10mf_processed_FD2.txt')
```

5.1) Calculate mean  $u_{mean} = \langle u \rangle$  and standard deviation  $u_{std} = \sqrt{\langle u^2 \rangle - (\langle u \rangle)^2}$

```
[ ]: # Sample frequency 20 kHz
f_s = 20_000
T = len(data)/f_s
print(f'The dataset is is {T}s long.')
```

```
# Mean and standard deviation
u_mean = np.mean(data)
u_std = np.std(data)
```

The dataset is 62.88915s long.

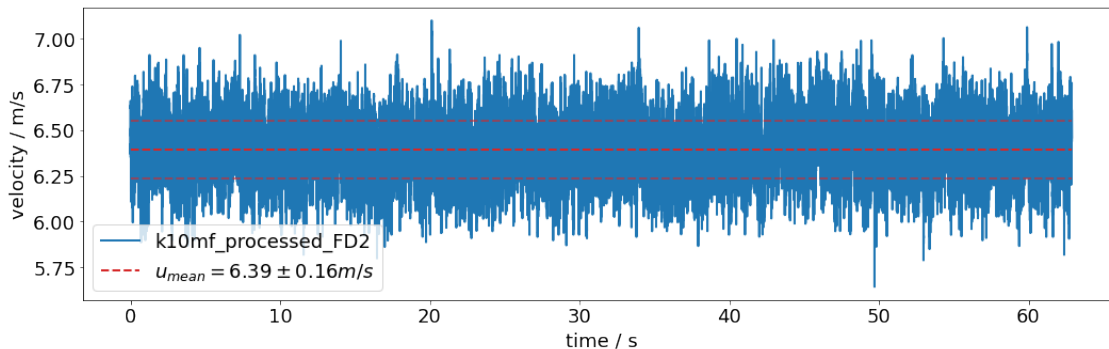
```
[ ]: fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)
time = np.linspace(0,T,len(data))
ax.plot(time, data, label='k10mf_processed_FD2', zorder=1)

# Plot time series and mean +/- std velocity
print(f'The mean velocity is {u_mean:.2f}m/s.\nThe standard deviation is +/-_
↳{u_std:.2f}m/s.')
ax.hlines(u_mean, 0, T, colors='C3', linestyles='--', zorder=2,
↳label=f'$u_{{{mean}}}={{{{{u_mean:.2f}}}} \pm {{{u_std:.2f}}} \text{ m/s}$')
ax.hlines(u_mean+u_std, 0, T, colors='C3', linestyles='--', zorder=2, alpha=0.6)
ax.hlines(u_mean-u_std, 0, T, colors='C3', linestyles='--', zorder=2, alpha=0.6)

ax.set_xlabel('time / s')
ax.set_ylabel('velocity / m/s')
ax.legend()
fig.savefig('Abb/Ex5_Velocity_Time_Series.png')
```

The mean velocity is 6.39m/s.

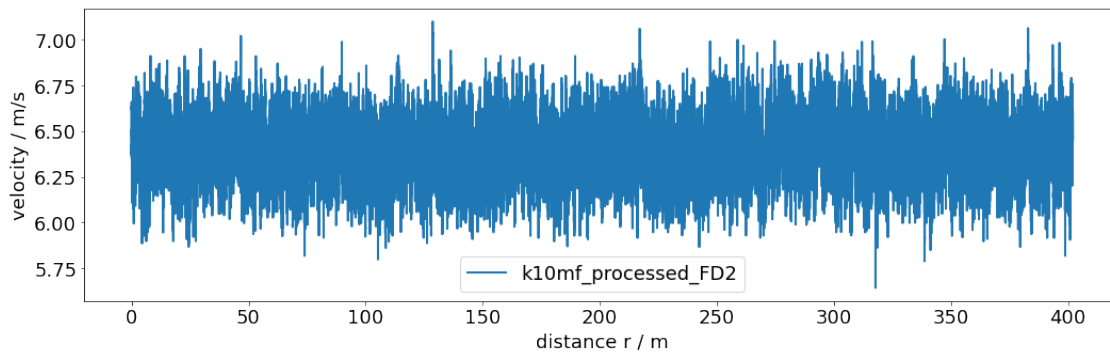
The standard deviation is +/- 0.16m/s.



### 2.0.1 5.2) Use Taylor's hypothesis $x = u_{mean} \cdot t$ to calculate a spatial series from a temporal series

Essentially, you assume that the velocity field you record at one spatial position over time only gets advected but not deformed by the mean velocity.

```
[ ]: fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)
      spatial = time * u_mean
      ax.plot(spatial, data, label='k10mf_processed_FD2', zorder=1)
      ax.set_xlabel('distance r / m')
      ax.set_ylabel('velocity / m/s')
      ax.legend()
      fig.savefig('Abb/Ex5_Velocity_Spatial_Series.png')
```

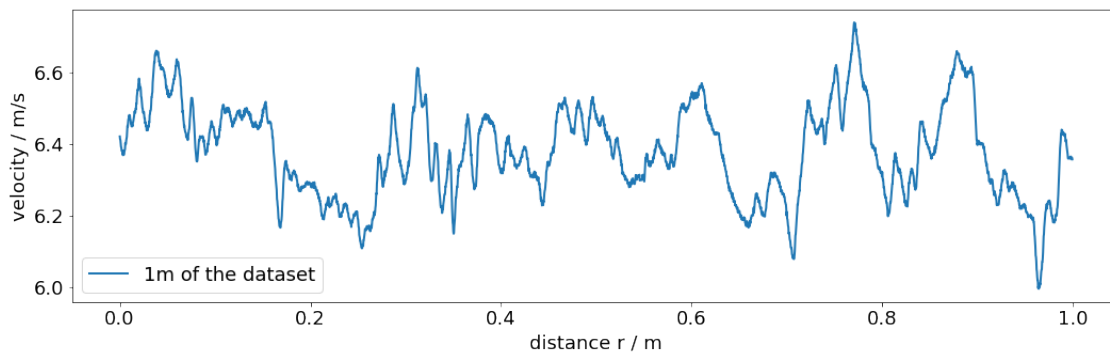
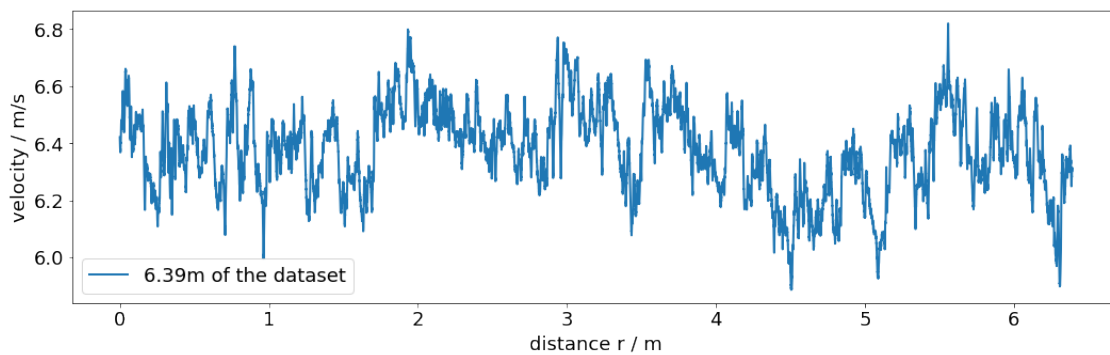
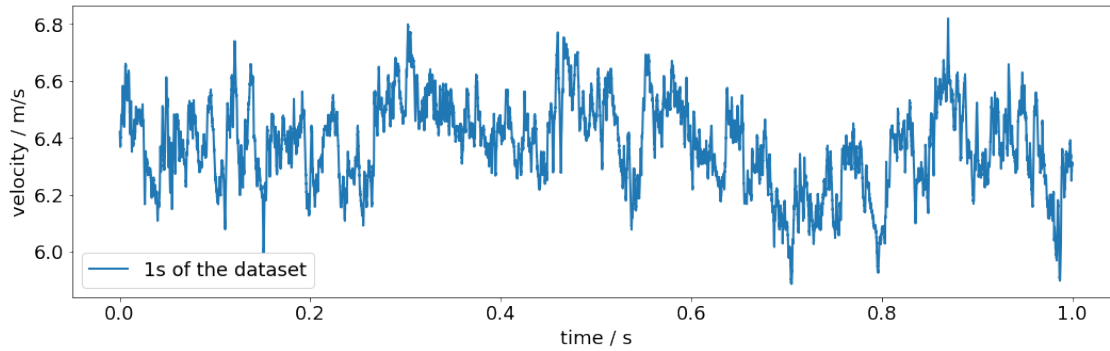


```
[ ]: fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)
      ax.plot(time[:f_s], data[:f_s], label='1s of the dataset', zorder=1)
      ax.set_xlabel('time / s')
      ax.set_ylabel('velocity / m/s')
      ax.legend()

      fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)
      ax.plot(spatial[:f_s], data[:f_s], label=f'{u_mean:.2f}m of the dataset',
      ↪zorder=1)
      ax.set_xlabel('distance r / m')
      ax.set_ylabel('velocity / m/s')
      ax.legend()

      fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)
      ax.plot(spatial[:int(f_s/u_mean)], data[:int(f_s/u_mean)], label='1m of the
      ↪dataset', zorder=1)
      ax.set_xlabel('distance r / m')
      ax.set_ylabel('velocity / m/s')
      ax.legend()
```

```
[ ]: <matplotlib.legend.Legend at 0x1b512501790>
```



5.3) Calculate correlation function  $C(r) = \langle u'(x+r)u'(x) \rangle$  of the fluctuations  $u' = u - \langle u \rangle$

```
[ ]: # Number of seconds in the dataset to do the computation
sec = 1
data_fluc = data - u_mean

fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)
# Not exactly the same one
```

```

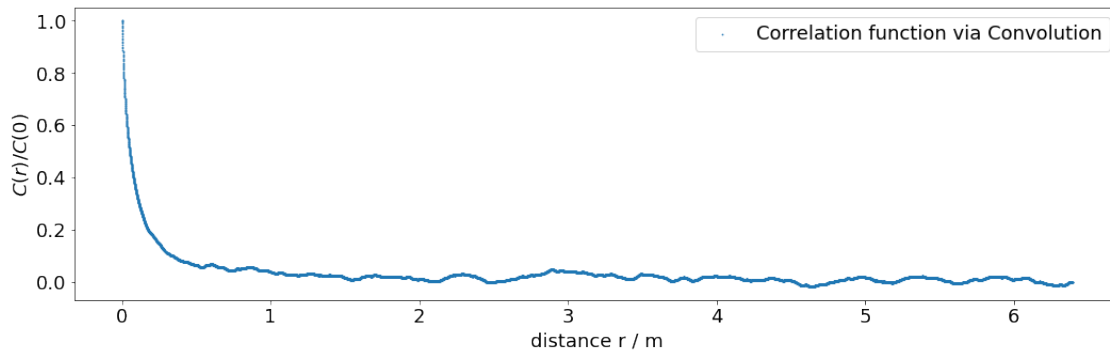
# acf = sm.tsa.acf(data_fluc, nlags=f_s*sec, fft=True)
# ax.plot(spatial[:f_s*sec+1], acf, label='Autocorrelation function', zorder=1)

# Correlation via convolution: Much faster than doing it for each r by my own
# 2s for all r instead of 4 Minutes for only 20.000 r
cf_cov = signal.fftconvolve(data_fluc, data_fluc[::-1], mode='full')
# Devide by how many samples were involved in convolution
xi = np.arange(1, len(data_fluc) + 1)
d = np.hstack((xi, xi[:-1][::-1]))
cf_cov = cf_cov / d

# Only the first half is our correlation
cf_cov = cf_cov[:len(data_fluc)][::-1]
cf_cov_0 = cf_cov[0]
cf_cov = cf_cov/cf_cov_0
ax.scatter(spatial[:f_s*sec+1], cf_cov[:f_s*sec+1], label='Correlation function_
→via Convolution', zorder=1, s=0.5)

ax.set_xlabel('distance r / m')
ax.set_ylabel('$C(r)/C(0)$')
ax.legend()
fig.savefig('Abb/Ex_5Correlation.png')

```



5.4) Determine integral length scale  $L = \int_0^\infty dr \frac{C(r)}{C(0)}$  and by an exponential fit  $C(r) \approx C(0) \cdot e^{-\frac{r}{L}}$

```

[ ]: fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)

# 1) Integral length scale via Integral = Summation of rectangles
step = time[1]-time[0]

# Zero crossing zc
zc = np.argwhere(cf_cov < 0)[0,0]
L_sum = np.sum(cf_cov[:zc])*step

```

```

print(f'Integral length via summation: L = {L_sum:.4f}m')

# 2) Integral length scale via exp-fit
# begin, end of fit in m
begin, end = 0.12, 6
idx_b, idx_e = int(begin/u_mean*f_s) , int(end/u_mean*f_s)

xdat = spatial[idx_b:idx_e]
ydat = cf_cov[idx_b:idx_e]

mod = lmfit.models.ExponentialModel() + lmfit.models.ConstantModel()
# pars = mod.make_params()
result = mod.fit(ydat, x=xdat)
L_fit = result.values['amplitude']
print(f'Integral length via fit: L = {L_fit:.4f}m')

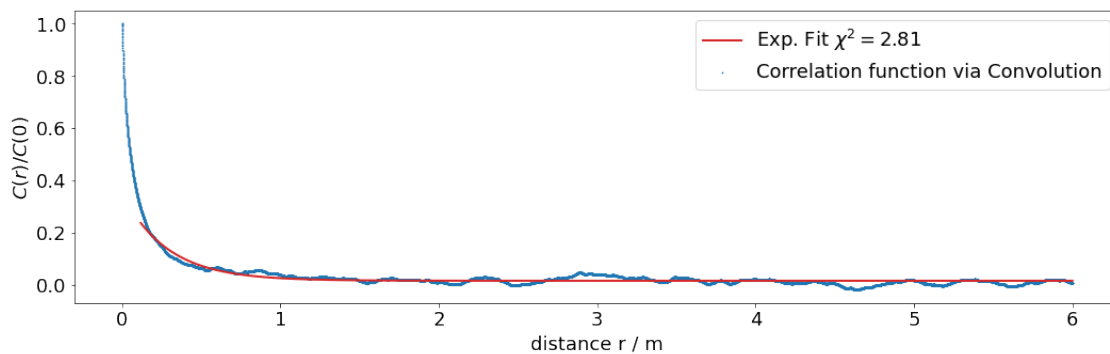
# print(result.fit_report())
plt.plot(xdat, result.eval(result.params, x=xdat), c='C3', label=f'Exp. Fit_
↳ $\chi^2$={{result.chisqr:.2f}}$')
ax.scatter(spatial[:idx_e], cf_cov[:idx_e], label='Correlation function via_
↳ Convolution', zorder=1, s=0.5)

ax.set_xlabel('distance r / m')
ax.set_ylabel('$C(r)/C(0)$')
ax.legend()
fig.savefig('Abb/Ex5_Correlation_Zoomed_Fit.png')

```

Integral length via summation: L = 0.0259m

Integral length via fit: L = 0.3376m



Try out different begin and end points for the fit

```

[ ]: size = 30
L_matrix = np.empty((size,size))

```

```

b_vec = np.linspace(0.1,0.5,size)
e_vec = np.linspace(3,10,size)
for i, begin in enumerate(b_vec):
    for j, end in enumerate(e_vec):
        idx_b, idx_e = int(begin/u_mean*f_s) , int(end/u_mean*f_s)
        xdat = spatial[idx_b:idx_e]
        ydat = cf_cov[idx_b:idx_e]

        mod = lmfit.models.ExponentialModel() + lmfit.models.ConstantModel()
        try:
            result = mod.fit(ydat, x=xdat)
            L = result.values['amplitude']
        except:
            L = None
        L_matrix[i,j] = L

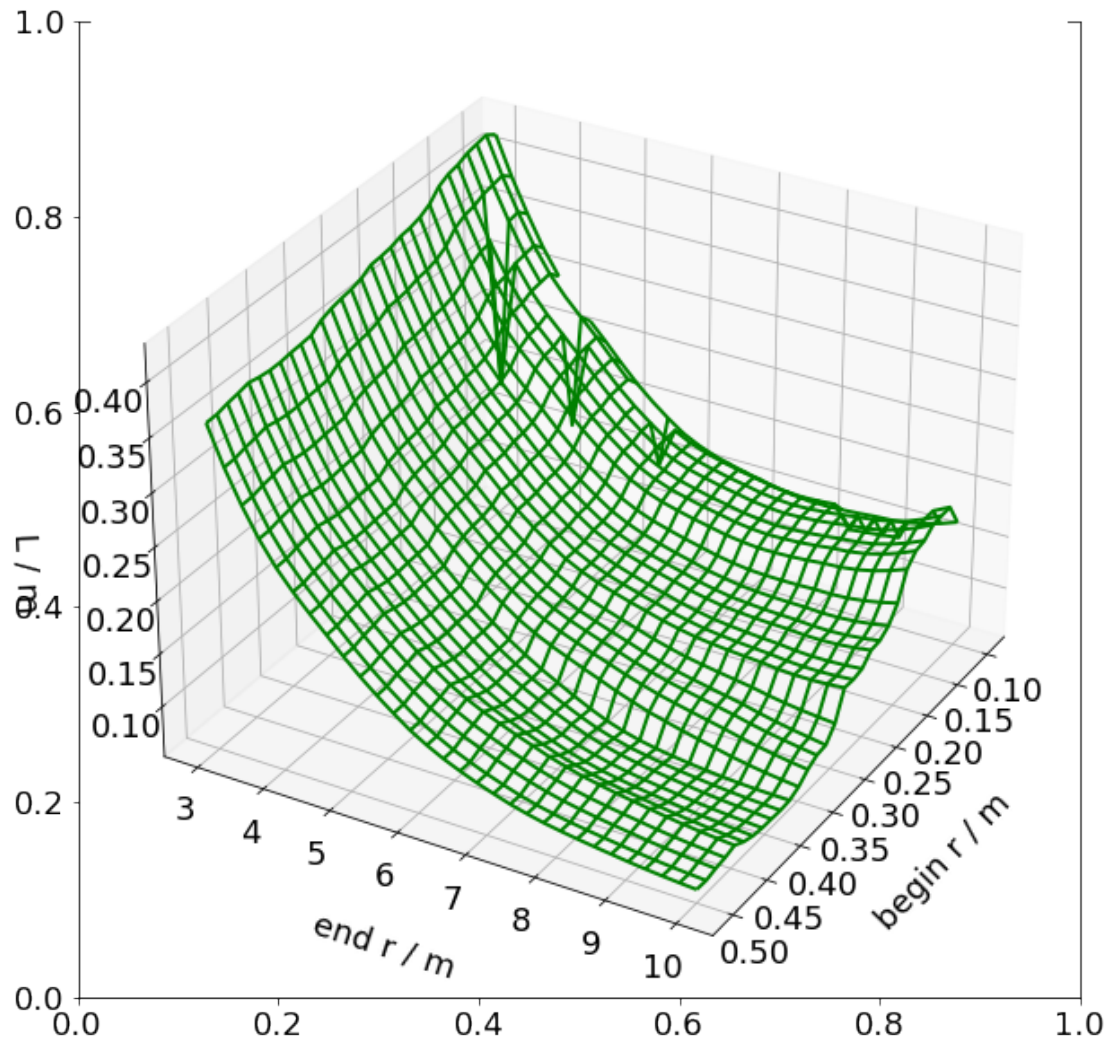
```

The value of L really depends on the start and end point of our fit!

```

[ ]: fig, ax = plt.subplots(figsize=(10,10))
ax = plt.axes(projection = '3d')
X, Y = np.meshgrid(b_vec, e_vec)
# ax.plot3D(b_vec[:len(L_matrix)], e_vec[:len(L_matrix)], L_matrix)
ax.plot_wireframe(X, Y, L_matrix, color = 'green')
ax.set_xlabel('begin r / m')
ax.set_ylabel('end r / m')
ax.set_zlabel('L / m')
ax.view_init(30, 30)
ax.yaxis.labelpad = 20
ax.xaxis.labelpad = 20
ax.zaxis.labelpad = 20
fig.savefig('Abb/Ex5_Different_integral_lengths_L.png')

```



5.5) Determine Taylor length  $\lambda = \sqrt{-\frac{C(0)}{C''(0)}}$

```
[ ]: fig, ax = plt.subplots(figsize=(15,5), tight_layout=True)

# Taylor length scale via polynomial fit
# begin, end of fit in m
dpoints = 50
idx_b, idx_e = 0, 7

xdat = spatial[idx_b:idx_e]
ydat = cf_cov[idx_b:idx_e]

mod = lmfit.models.QuadraticModel()
result = mod.fit(ydat, x=xdat)
C_curvature = result.values['a']
```



```

# print(f'Taylor length lambda via cf_cov[0]= 1 normed = {np.sqrt(-cf_cov[0]/
→C_curvature):.4f}m')
print(f'Taylor length lambda via cf_cov[0] = {np.sqrt(-cf_cov_0/C_curvature):.
→4f}m')

ax.plot(spatial[:dpoints], result.eval(result.params, x=spatial[:dpoints]),
→label=f'Exp. Fit\n$\chi^2={{result.chisqr:.2e}}$')
ax.scatter(spatial[:dpoints], cf_cov[:dpoints], label='Correlation function via
→Convolution', zorder=1)
ax.plot(spatial[:dpoints], cf_cov[:dpoints], alpha=0.5, zorder=1)

ax.set_xlim((-0.001,0.01))
ax.set_ylim((0.85,1.01))
ax.set_xlabel('distance r / m')
ax.set_ylabel('$C(r)/C(0)$')
ax.legend()
fig.savefig('Abb/Ex5_Correlation_Taylor_length.png')

print(result.fit_report())

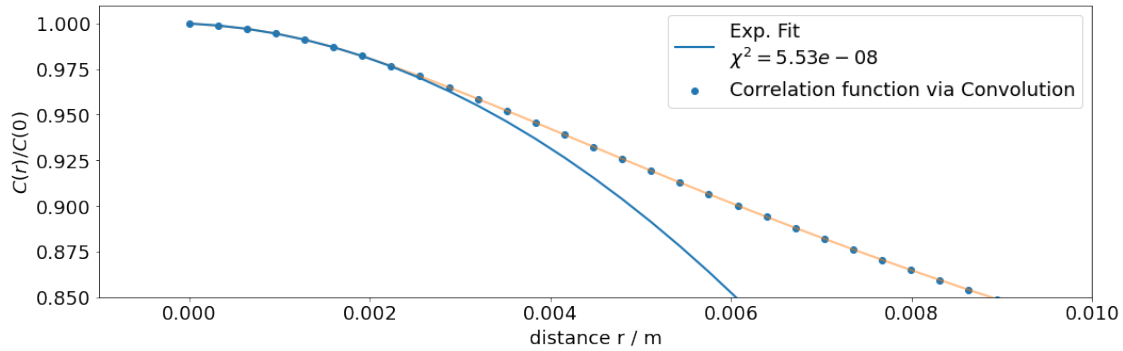
```

Taylor length lambda via cf\_cov[0] = 0.0026m

```

[[Model]]
  Model(parabolic)
[[Fit Statistics]]
  # fitting method   = leastsq
  # function evals   = 13
  # data points      = 7
  # variables        = 3
  chi-square         = 5.5260e-08
  reduced chi-square = 1.3815e-08
  Akaike info crit   = -124.599882
  Bayesian info crit = -124.762151
[[Variables]]
  a: -3738.93693 +/- 125.543927 (3.36%) (init = 0)
  b: -2.12421281 +/- 0.25058120 (11.80%) (init = 0)
  c:  0.99990890 +/- 1.0259e-04 (0.01%) (init = 0)
[[Correlations]] (unreported correlations are < 0.100)
  C(a, b) = -0.961
  C(b, c) = -0.781
  C(a, c) = 0.625

```



Try out different end point of the correlation function and look at Chi-Square for best fit

```
[ ]: num_idx = 15
taylor_lengths = np.empty(num_idx)
chi_sqr = np.empty(num_idx)

indices = np.linspace(1,num_idx,num_idx)

for idx_e in indices:
    idx_e = int(idx_e)
    xdat = spatial[:idx_e]
    ydat = cf_cov[:idx_e]

    mod = lmfit.models.QuadraticModel()
    try:
        result = mod.fit(ydat, x=xdat)
        C_curvature = result.values['a']
        lambda_l = np.sqrt(-cf_cov[0]/C_curvature)
    except:
        lambda_l = None
    taylor_lengths[idx_e-1] = lambda_l
    chi_sqr[idx_e-1] = result.chisqr

fig, ax = plt.subplots(1,2,figsize=(15,7), tight_layout=True)

ax[0].scatter(indices, taylor_lengths, label='Taylor lengths', zorder=1)
ax[0].set_xlabel('end index')
ax[0].set_ylabel('$Taylor length$')

ax[1].scatter(indices, chi_sqr, label='$\chi^2$', zorder=1)
ax[1].set_xlabel('end index')
ax[1].set_ylabel('Chi-Square')
```

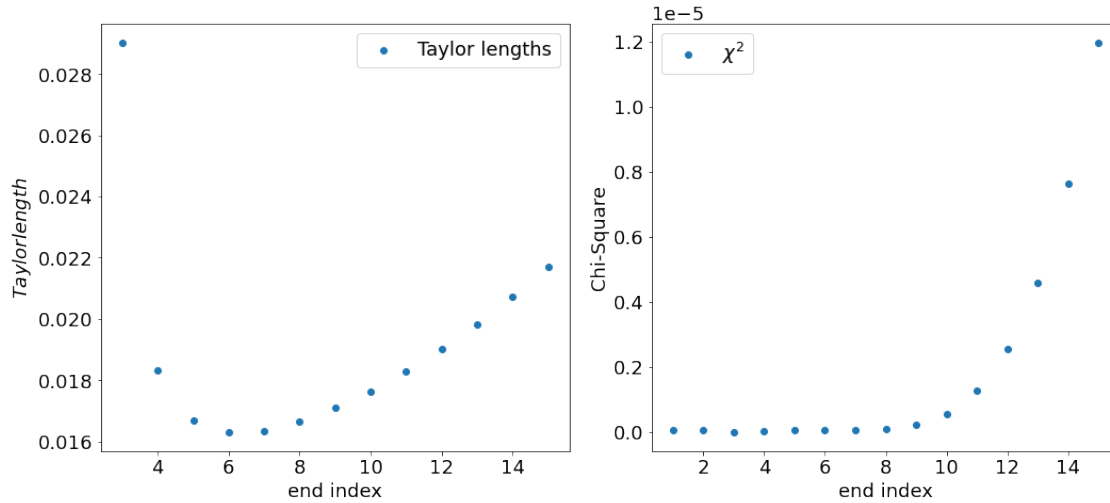
```

ax[0].legend()
ax[1].legend()
print(f'Minimum chi-sqr (best number of indices) at index e_idx = {np.
      ↳argmin(chi_sqr)+1}')

fig.savefig('Abb/Ex5_Correlation_Taylor_length_different_idx.png')

```

Minimum chi-sqr (best number of indices) at index e\_idx = 3



5.6) Calculate Structure Function  $S_n(r) = \langle [u(x+r) - u(x)]^n \rangle$

```

[ ]: dr = u_mean/f_s
r_in_log = np.logspace(-3, 0, 50)
struc = np.empty((7,np.size(r_in_log)))

for n in range(1,7):
    for r in range(np.size(r_in_log)):
        r_idx = int(r_in_log[r]/dr)
        struc[n,r] = np.mean((data[:-r_idx]-data[r_idx:])**n)

```

```

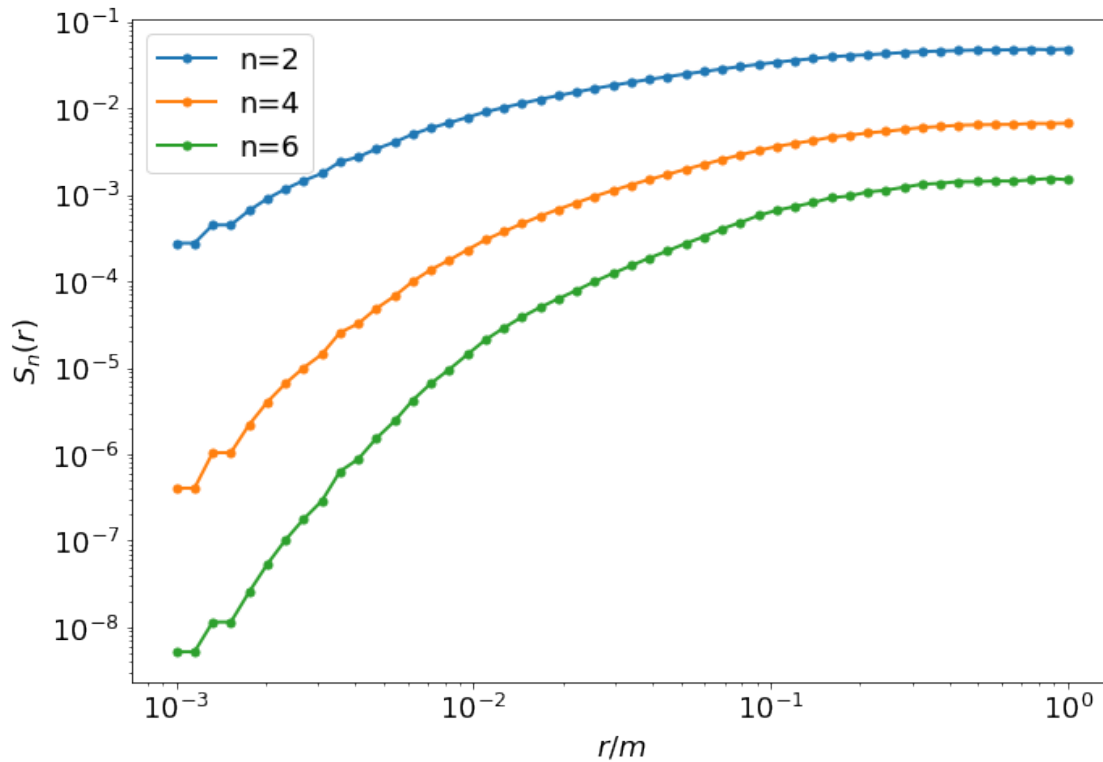
[ ]: # Even structure functions
fig, ax = plt.subplots(figsize=(10,7), tight_layout=True)

dr = u_mean/f_s
ax.loglog(r_in_log, struc[2], label='n=2', marker='.', ms=10)
ax.loglog(r_in_log, struc[4], label='n=4', marker='.', ms=10)
ax.loglog(r_in_log, struc[6], label='n=6', marker='.', ms=10)

ax.set_xlabel('$r / m$')

```

```
ax.set_ylabel('$S_n(r)$')
ax.legend()
plt.savefig('Abb/Ex5_Even_structure_functions')
```



A power-law like  $r^{\zeta_n}$  should look like a linear graph with slope  $\zeta_n$

When  $r$  is large (approximately  $L$ ) then we have Gaussian fluctuations and the model isn't that good anymore

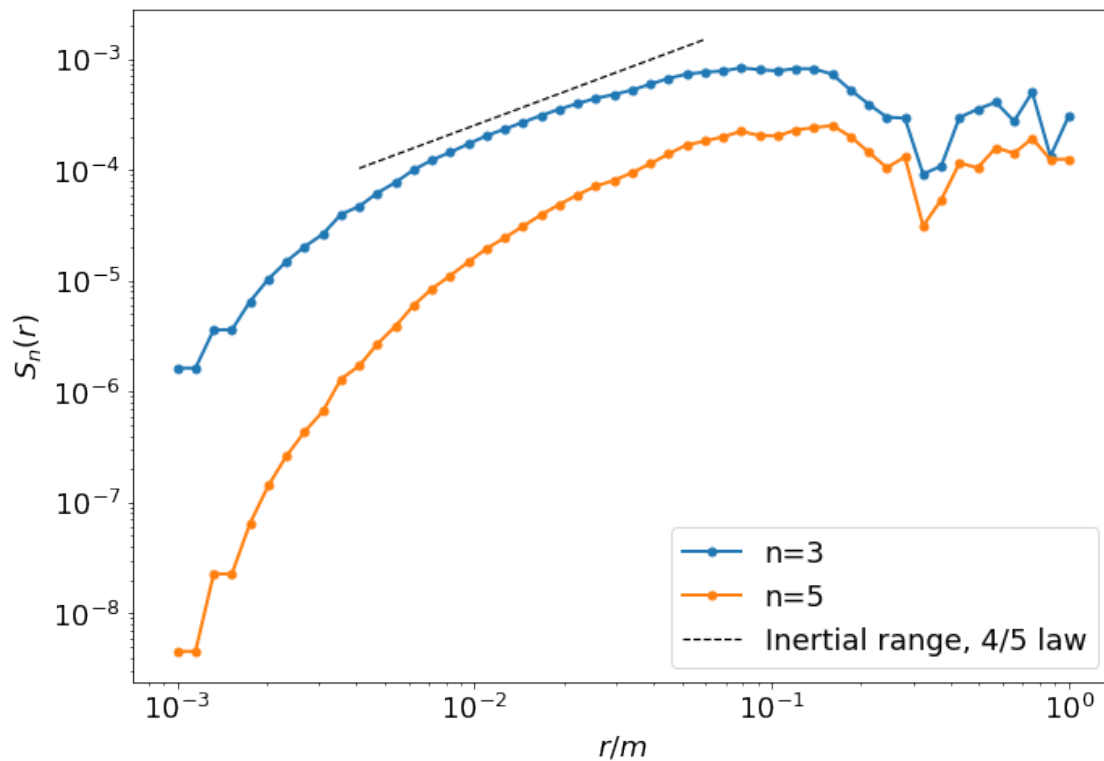
```
[ ]: # Odd structure functions
fig, ax = plt.subplots(figsize=(10,7), tight_layout=True)

ax.loglog(r_in_log, -struc[3], label='n=3', marker='.', ms=10)
ax.loglog(r_in_log, -struc[5], label='n=5', marker='.', ms=10)

# S_3(r,t) = -4/5 <epsilon> r
# Inertial range
y = 4/5*10**(-1.5)*r_in_log[10:30]
ax.plot(r_in_log[10:30], y, label='Inertial range, 4/5 law',
        linestyle='dashed', color='black', lw=1.2)

ax.set_xlabel('$r / m$')
ax.set_ylabel('$S_n(r)$')
ax.legend()
```

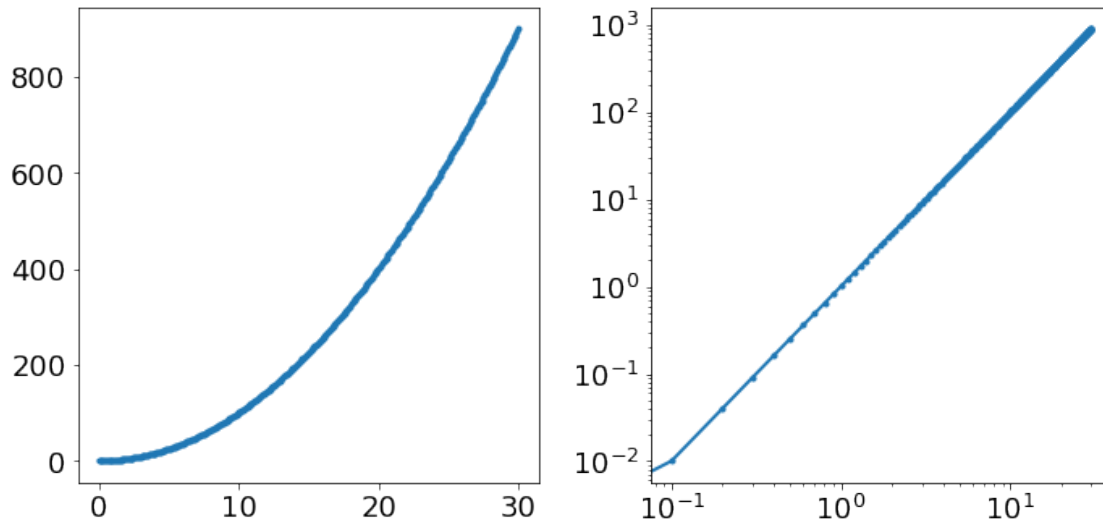
```
plt.savefig('Abb/Ex5_Odd_structure_functions')
```



A power-law like  $r^{\zeta_n}$  should look like a linear graph with slope  $\zeta_n$  if  $\zeta_n$  is only a number. Here for example  $\zeta_n = 2$

```
[ ]: fig, ax = plt.subplots(1,2, figsize=(10,5), tight_layout=True)
r_lin = np.linspace(0,30,300)
y = np.power(r_lin,2)
ax[0].plot(r_lin,y, marker='.')
ax[1].loglog(r_lin, y, marker='.')
```

```
[ ]: [<matplotlib.lines.Line2D at 0x1b513ba1550>]
```



$$S_n(r) = C_n \cdot \langle \epsilon \rangle \cdot r^{\frac{n}{3}} \cdot \left(\frac{L}{r}\right)^{\frac{\mu}{18} n(n-3)}$$

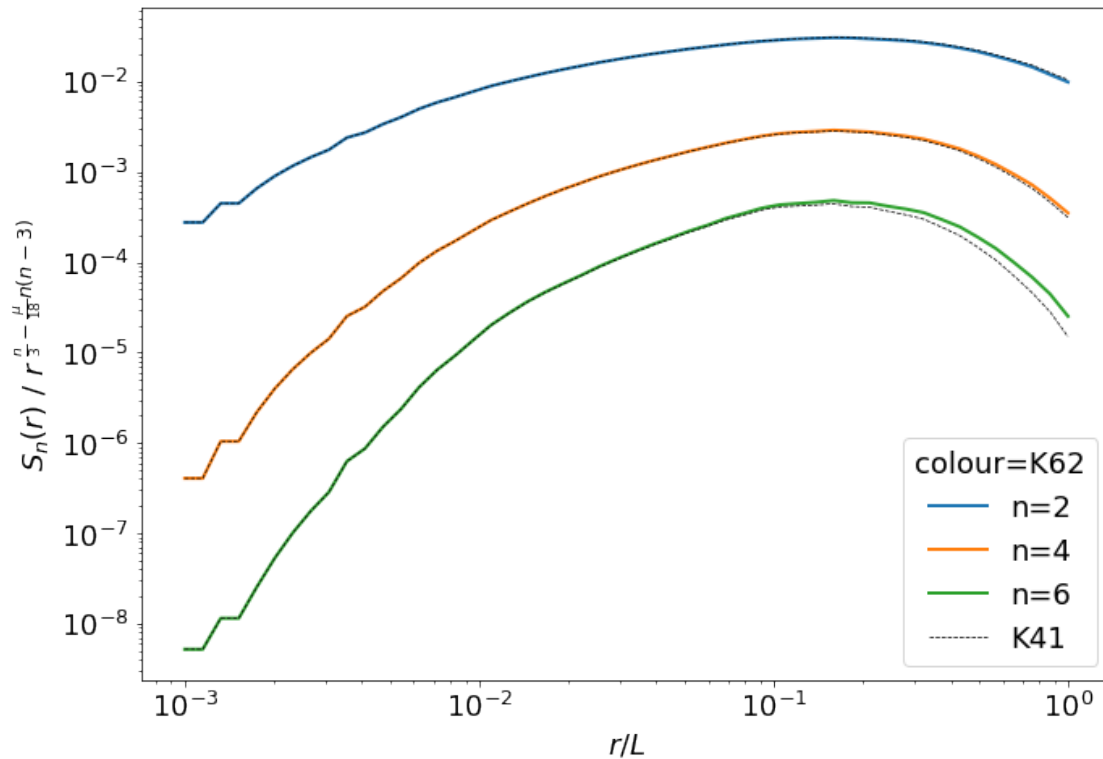
$$S_n(r) = C'_n \cdot \langle \epsilon \rangle \cdot r^{\frac{n}{3} - \frac{\mu}{18} n(n-3)}$$

```
[ ]: # even structure functions
fig, ax = plt.subplots(figsize=(10,7), tight_layout=True)

# From lecture mu is usually between 0.2 < mu < 0.28
# No "real" reason, just from experimental data
r_in_lin = np.power(10,r_in_log)

mu =0.227
ax.loglog(r_in_log, struc[2]/r_in_lin**(2/3+mu/9) , label='n=2')
ax.loglog(r_in_log, struc[2]/r_in_lin**(2/3) , color='black', ls='dashed', lw=0.
↪7)
ax.loglog(r_in_log, struc[4]/r_in_lin**(4/3-2*mu/9), label='n=4')
ax.loglog(r_in_log, struc[4]/r_in_lin**(4/3), color='black', ls='dashed', lw=0.
↪7)
ax.loglog(r_in_log, struc[6]/r_in_lin**(6/3-mu), label='n=6')
ax.loglog(r_in_log, struc[6]/r_in_lin**(6/3), color='black', ls='dashed', lw=0.
↪7, label='K41')

ax.set_xlabel('$r/L$', fontsize=18)
ax.set_ylabel(r'$S_n(r) \propto r^{\frac{n}{3} - \frac{\mu}{18} n(n-3)}$', ↪
↪fontsize=18)
ax.legend(title='colour=K62')
plt.savefig('Abb/Ex5_Even_structure_functions_with_K41')
```



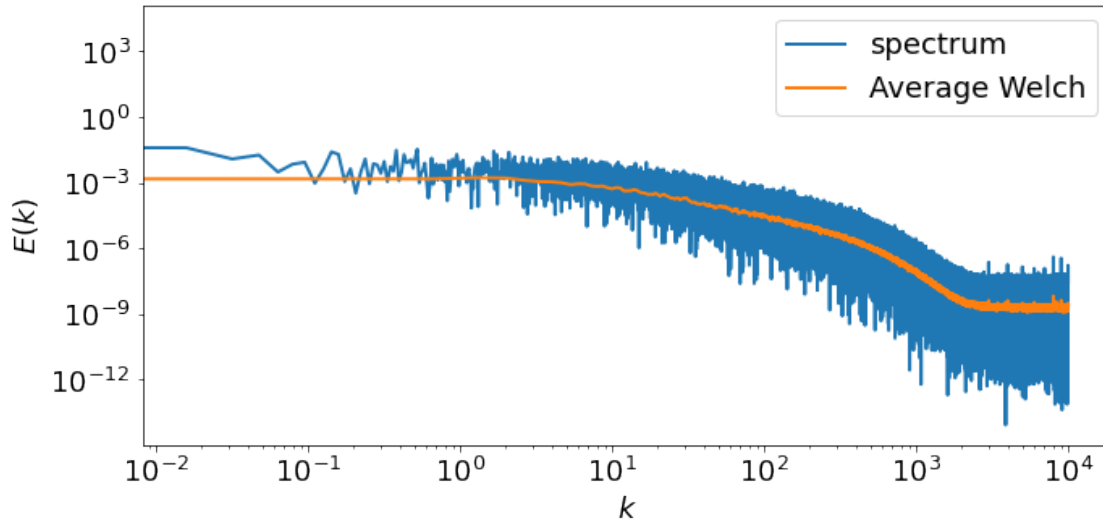
### 3 Exercise 6

6.1) Determine the energy spectrum  $E(k)$  by using Fourier transform of the signal

```
[ ]: data_fourier = np.fft.rfft(data)/np.size(data)*20
k = np.linspace(0, 20*1000/2, np.size(data)//2+1)
spectrum = np.abs(data_fourier**2)

# Estimate power spectral density using Welch's method.
freq_welch, spectrum_average = signal.welch(data, fs=20*1000, nperseg=30000)
# f, Pxx_averaged = signal.welch(data, fs=20*1000, average='median')

[ ]: fig, ax = plt.subplots(figsize=(10,5), tight_layout=True)
ax.loglog(k,spectrum, label='spectrum')
ax.loglog(freq_welch, spectrum_average, label='Average Welch')
ax.set_xlabel('$k$')
ax.set_ylabel('$E(k)$')
ax.legend()
plt.savefig('Abb/Ex6_Spectrum')
```



6.2) Compute histogram of the velocity increments  $u(x+r) - u(x)$

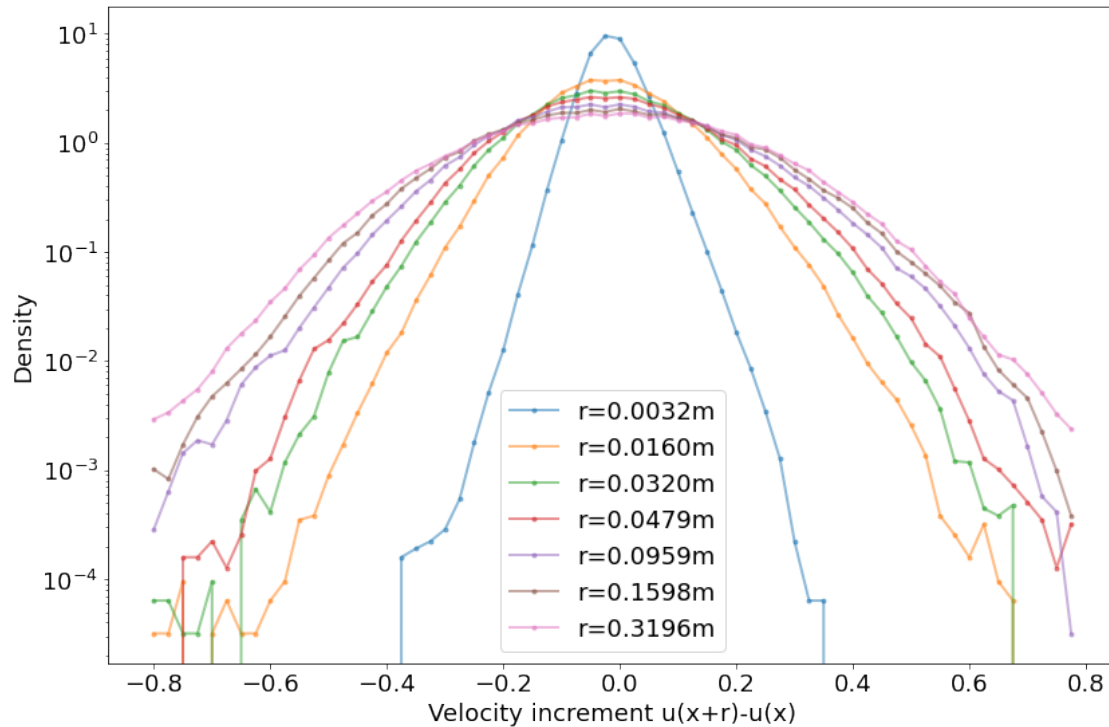
```
[ ]: fig, ax = plt.subplots(figsize=(12,8), tight_layout=True)
dr = u_mean/f_s

r_vec_idx = [10,50,100,150,300,500,1000]
for r in r_vec_idx:
    velocity_increments = data[r:] - data[:-r]

    hist, bins = np.histogram(velocity_increments, bins=64, density=True,
    ↪ range=[-0.8,0.8])
    std_increment = np.std(velocity_increments)
    ax.semilogy(bins[:-1], hist, alpha=0.5, marker='.', label=f'r={ (dr*r):.
    ↪ 4f}m')

ax.set_xlabel('Velocity increment u(x+r)-u(x)')
ax.set_ylabel('Density')
ax.legend()
plt.savefig('Abb/Ex6_Density')
```





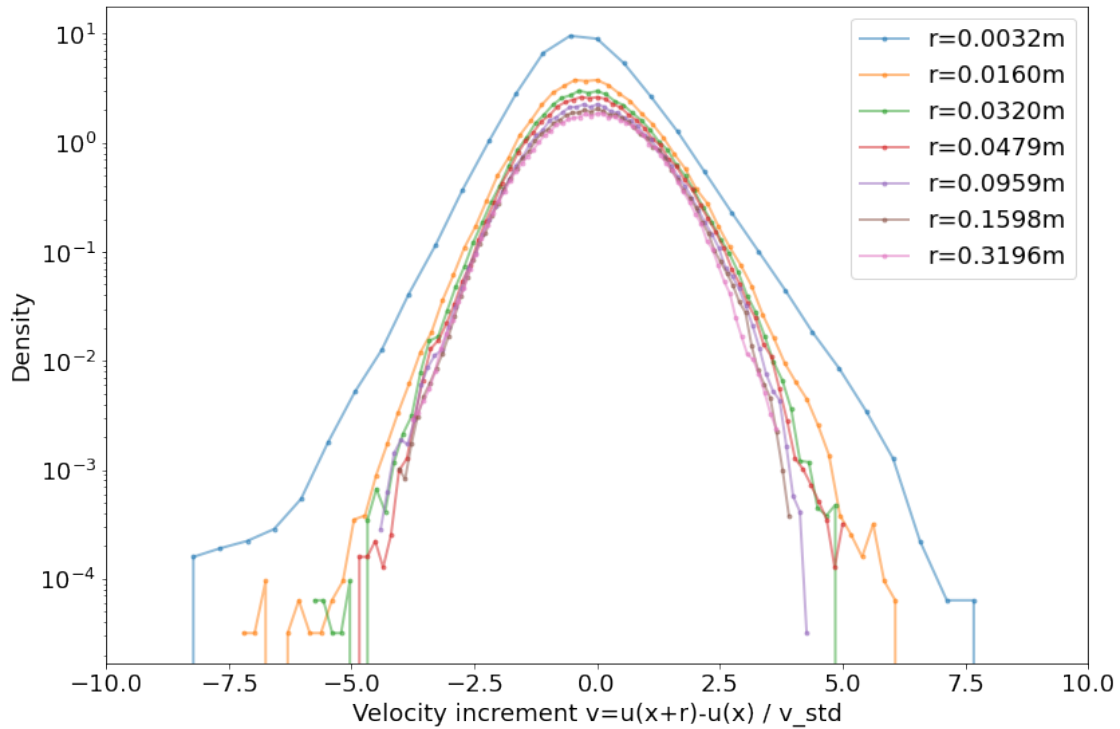
For a better comparison, plot x-axis as multiple of standard deviation

```
[ ]: fig, ax = plt.subplots(figsize=(12,8), tight_layout=True)
dr = u_mean/f_s

r_vec_idx = [10,50,100,150,300,500,1000]
for r in r_vec_idx:
    velocity_increments = data[r:]-data[:-r]

    hist, bins = np.histogram(velocity_increments, bins=64, density=True,
    ↪range=[-0.8,0.8])
    std_increment = np.std(velocity_increments)
    ax.semilogy(bins[:-1]/std_increment, hist, alpha=0.5, marker='.',
    ↪label=f'r={dr*r:.4f}m')

ax.set_xlabel('Velocity increment v=u(x+r)-u(x) / v_std')
ax.set_ylabel('Density')
ax.set_xlim(-10,10)
ax.legend()
plt.savefig('Abb/Ex6_Density_in_std')
```



6.3) Now we multiply with a factor of  $10^{-i}$  so the histograms are not on top of each other. We can clearly see that the probability density functions are NOT self-similar. For small  $r$  so in the region of the Kolmogorov microscale we see a Non-Gaussian pdf and on large scale a Gaussian pdf. For self-similarity the pdf's must be the whole time Gaussian or Non-Gaussian, so independent of scale  $r$ .

The pdf of small  $r$  could be described by a quasi-normal Ansatz  $\rightarrow$  superstatistics with summation of wheighted Gaussians with different variances

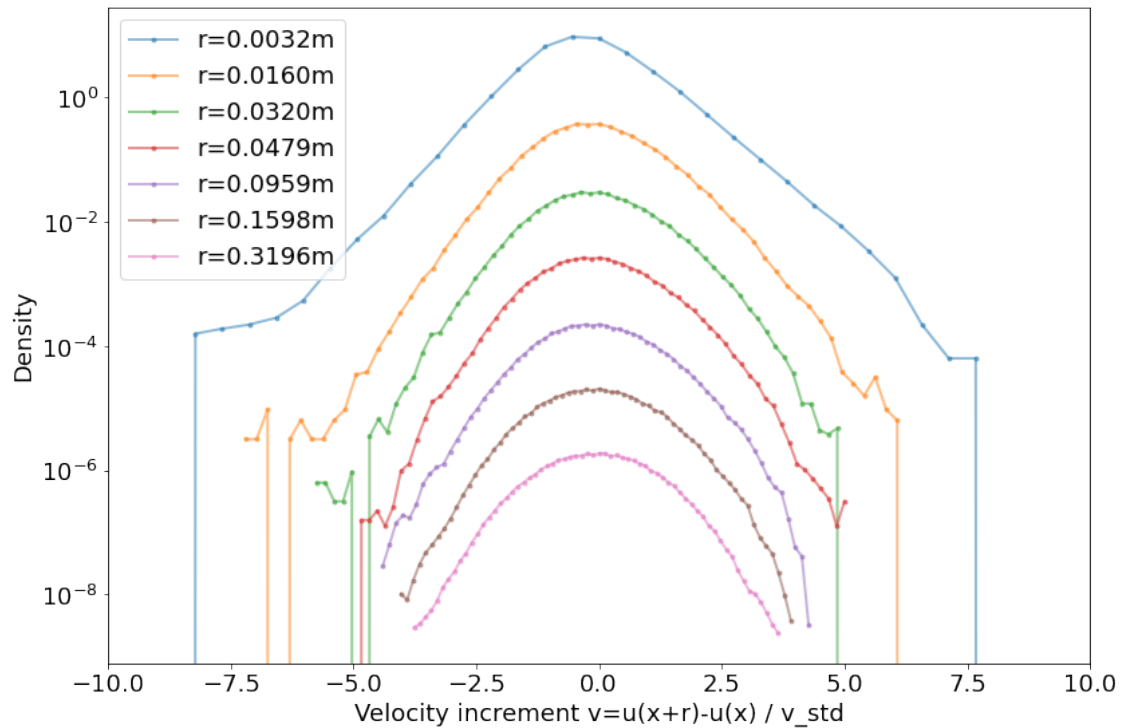
```
[ ]: fig, ax = plt.subplots(figsize=(12,8), tight_layout=True)
dr = u_mean/f_s

r_vec_idx = [10,50,100,150,300,500,1000]
for i, r in enumerate(r_vec_idx):
    velocity_increments = data[r:] - data[:-r]

    hist, bins = np.histogram(velocity_increments, bins=64, density=True,
    range=[-0.8,0.8])
    std_increment = np.std(velocity_increments)
    ax.semilogy(bins[:-1]/std_increment, hist*10**(-i), alpha=0.5, marker='.',
    label=f'r={dr*r:.4f}m')

ax.set_xlabel('Velocity increment v=u(x+r)-u(x) / v_std')
ax.set_ylabel('Density')
```

```
ax.set_xlim(-10,10)
ax.legend()
plt.savefig('Abb/Ex6_Density_in_std_moved')
```



## 4 Exercise 7

7.1) Calculate local energy dissipation rate  $\epsilon(x) = 2\nu(\partial_x u(x))^2$ , kinematic viscosity air  $\nu \approx 0.15\text{cm}^2/\text{s}$

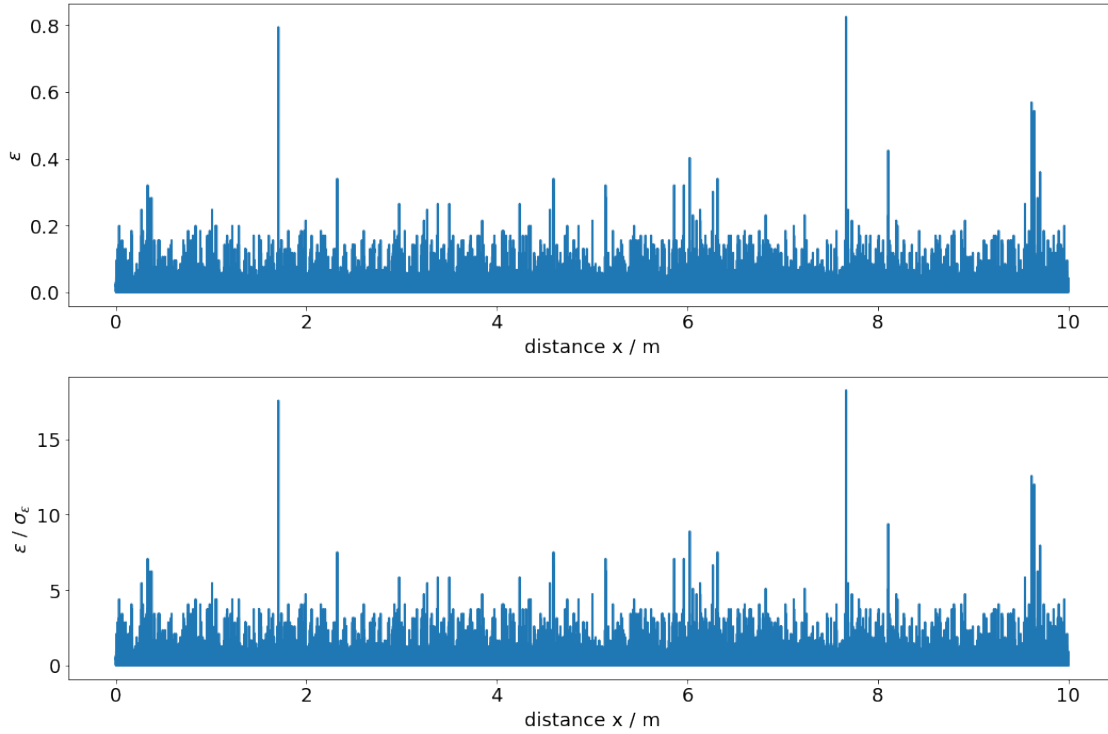
```
[ ]: fig, ax = plt.subplots(2,1,figsize=(15,10), tight_layout=True)

nu = 0.15 * 10**(-4)
du = data[1:] - data[:-1]
epsilon = 2 * nu * (du/dr)**2
epsilon_std = np.std(epsilon)

# end in meter
end = 10
idx_e = int(end/u_mean*f_s)

ax[0].plot(spatial[:idx_e], epsilon[:idx_e])
ax[0].set_xlabel('distance x / m')
ax[0].set_ylabel('$\epsilon$')
```

```
ax[1].plot(spatial[:idx_e], epsilon[:idx_e]/epsilon_std)
ax[1].set_xlabel('distance x / m')
ax[1].set_ylabel('$\epsilon$ / $\sigma_\epsilon$')
fig.savefig('Abb/Ex7_Epsilon.png')
```



Taylor length  $\lambda = \sqrt{15 \cdot \frac{\nu}{\langle \epsilon \rangle}} \cdot u_{rms}$  with  $\vec{u}_{rms} = \sqrt{\frac{\langle \vec{u}^2 \rangle}{N}}$

Kolmogorov microscale  $\eta = (\frac{\nu^3}{\langle \epsilon \rangle})^{1/4}$  from dimensional analysis  $[\langle \epsilon \rangle] = m^2/s^3$  and  $[\nu] = m^2/s$  see also  $\pi$ -theorem

```
[ ]: taylor_l_eps = np.sqrt(15*nu/np.mean(epsilon)) * np.sqrt(np.mean(np.
    ↪square(data_fluc)))
kolmog_microscale = (nu**3/np.mean(epsilon))**0.25
print(f'Taylor length with local energy dissipation rate: {taylor_l_eps}m')
print(f'The Kolmogorov microscale is: {kolmog_microscale}m')
print(f'Average dissipation rate: {np.mean(epsilon)}m^2/s^3')
```

Taylor length with local energy dissipation rate: 0.016887069520983355m

The Kolmogorov microscale is: 0.0006420463044883548m

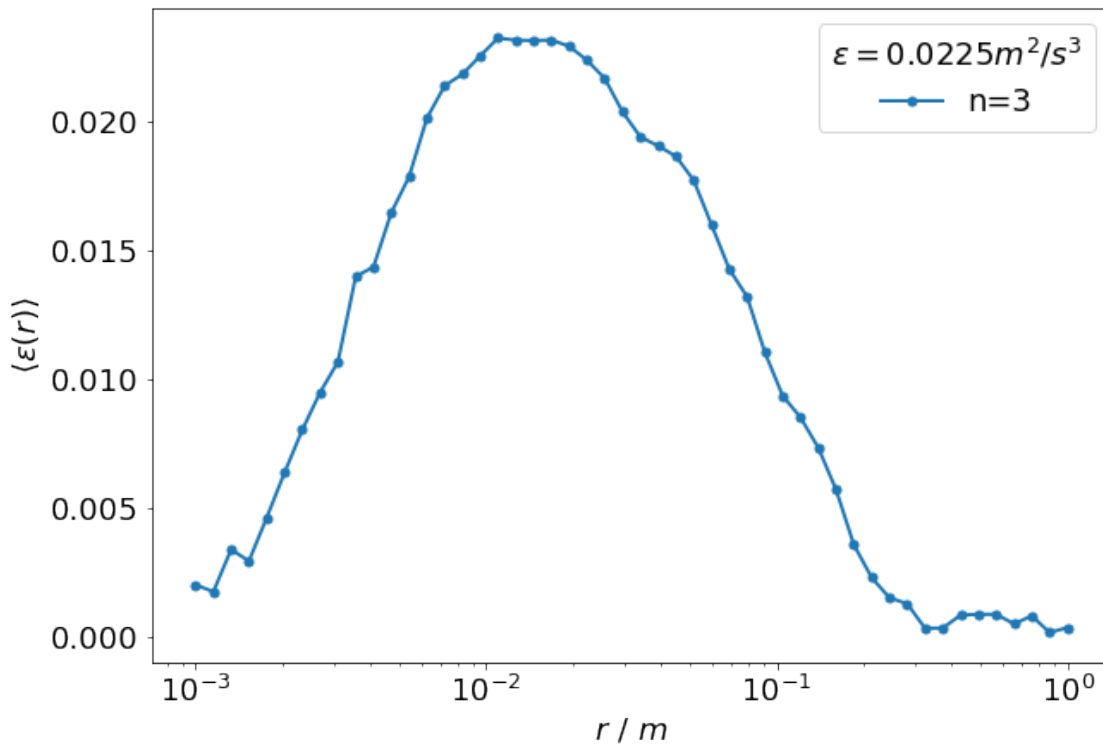
Average dissipation rate: 0.01986133216996104m<sup>2</sup>/s<sup>3</sup>

7.1) From 4/5 law  $S_3(r, t) = -\frac{4}{5}\langle \epsilon \rangle r \Rightarrow \langle \epsilon \rangle = S_3(r) \cdot -\frac{5}{4} \cdot \frac{1}{r}$

```
[ ]: fig, ax = plt.subplots(figsize=(10,7))
eps_avg = -5/4*struc[3]/r_in_log
ax.semilogx(r_in_log,eps_avg, label='n=3', marker='.', ms=10)

arg_max = np.argmax(eps_avg)
eps_avg_struct_func = np.mean(eps_avg[arg_max-3:arg_max+3])
print(f'Average dissipation rate by structure function = {eps_avg_struct_func:.
↪4f}m^2/s^3')
ax.set_xlabel('$r \setminus m$')
ax.set_ylabel(r'$\langle \epsilon(r) \rangle$')
ax.legend(title = f'$\langle \epsilon \rangle = {eps_avg_struct_func:.4f} m^2/s^3$')
plt.savefig('Abb/Ex7_Epsilon_from_4_5_law')
```

Average dissipation rate by structure function = 0.0225m<sup>2</sup>/s<sup>3</sup>



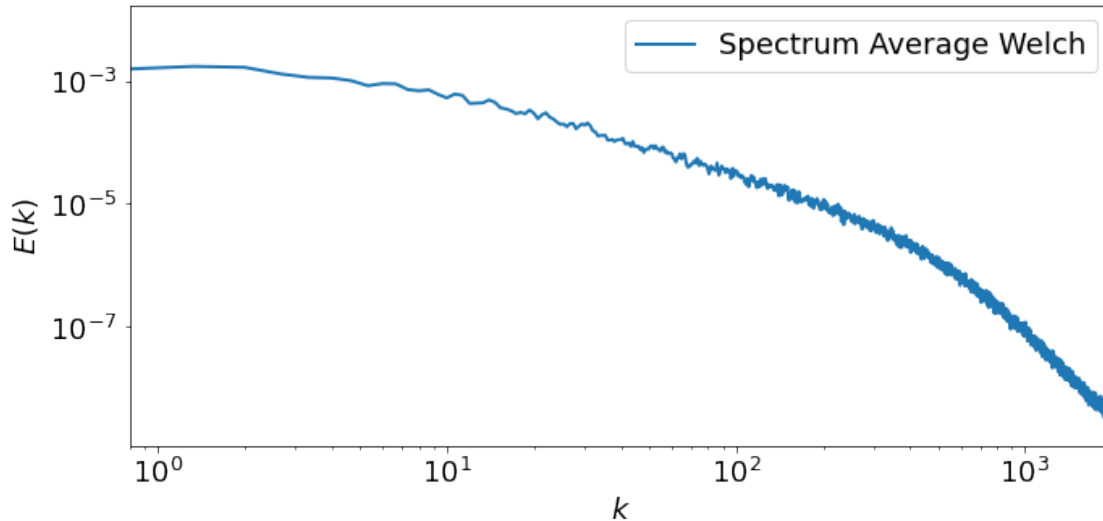
7.1) From  $E(k)$ , K41 phenomenology. In the inertial range we have  $E(k) = C_k \langle \epsilon \rangle^{2/3} \cdot k^{-5/3}$   
 We have three ranges: forcing, inertial range, dissipation range

```
[ ]: fig, ax = plt.subplots(figsize=(10,5), tight_layout=True)
ax.loglog(freq_welch, spectrum_average, label='Spectrum Average Welch')

ax.set_xlim(0.8,2000)
ax.set_ylim(np.min(spectrum_average), np.max(spectrum_average)*10)
```

```
ax.set_xlabel('$k$')
ax.set_ylabel('$E(k)$')
ax.legend()
```

```
[ ]: <matplotlib.legend.Legend at 0x1b5125e2670>
```



```
[ ]: fig, ax = plt.subplots(figsize=(10,5), tight_layout=True)
eps_avg = (spectrum_average * freq_welch**(5/3))**(3/2)

ax.loglog(freq_welch, eps_avg, label='Spectrum Average Welch')
ax.set_xlim(0.8,2000)

arg_max = np.argmax(eps_avg)

N_l, N_r = 250, 150
ax.loglog(freq_welch[arg_max-N_l:arg_max+N_r], eps_avg[arg_max-N_l:
    ↪arg_max+N_r], label='Spectrum Average Welch')
eps_avg_struct_func = np.mean(eps_avg[arg_max-N_l:arg_max+N_r])

# Fit Linear Model to the orange data points
mod = lmfit.models.LinearModel()
xdat = freq_welch[arg_max-N_l:arg_max+N_r]
ydat = eps_avg[arg_max-N_l:arg_max+N_r]

result = mod.fit(ydat, x=xdat)
eps_avg_spectrum = result.values['intercept']
print(f'Average dissipation rate by structure function = {eps_avg_spectrum:.
    ↪4f}m^2/s^3')
```

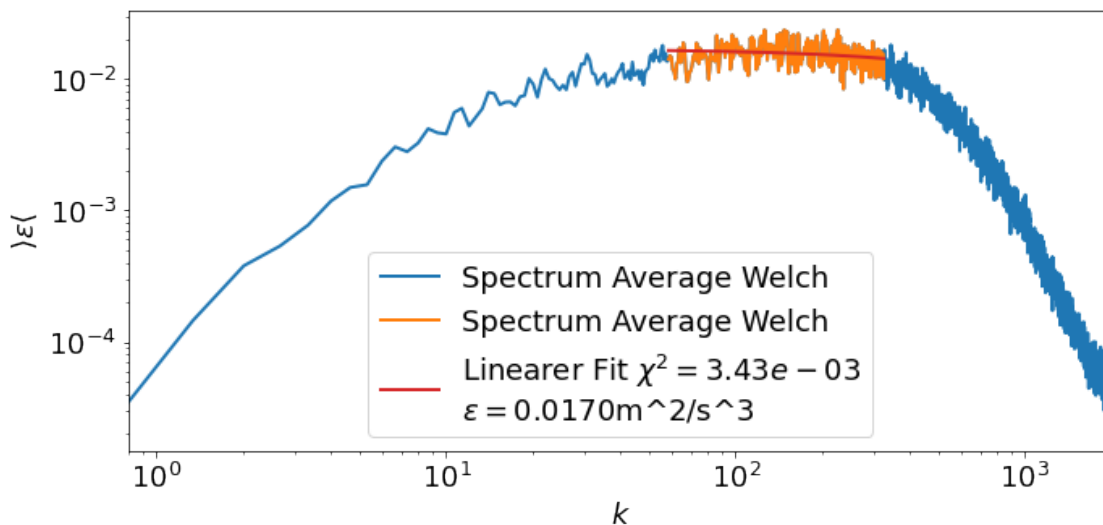
```

ax.plot(xdat, result.eval(result.params, x=xdat), c='C3', label=f'Linearer Fit_
↪ $\chi^2={{result.chisqr:.2e}}$ \n $\epsilon = {{eps_avg_spectrum:.4f}}$ m^2/
↪ s^3')

ax.set_xlabel('$k$')
ax.set_ylabel(r'$\angle \epsilon \angle$')
ax.legend()
plt.savefig('Abb/Ex7_Epsilon_from_K41_spectrum')

```

Average dissipation rate by structure function =  $0.0170 \text{ m}^2/\text{s}^3$



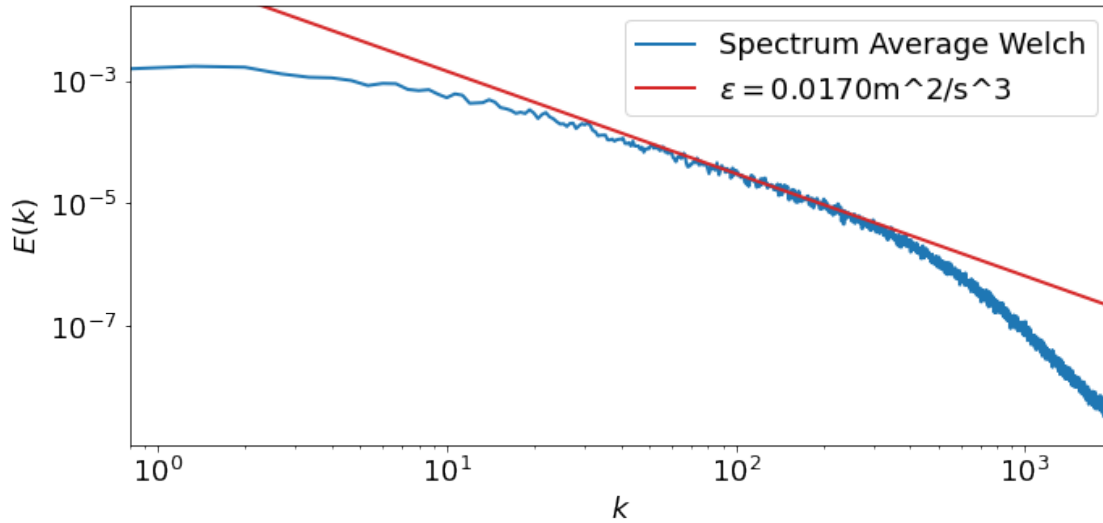
Plot of  $E(k)$  but with the determined epsilon

```

[ ]: fig, ax = plt.subplots(figsize=(10,5), tight_layout=True)
ax.loglog(freq_welch, spectrum_average, label='Spectrum Average Welch')
E_k = eps_avg_spectrum**(2/3)*freq_welch**(-5/3)
ax.plot(freq_welch, E_k, c='C3', label=f'$\epsilon = {{eps_avg_spectrum:.
↪ 4f}}$ m^2/s^3')

ax.set_xlim(0.8,2000)
ax.set_ylim(np.min(spectrum_average), np.max(spectrum_average)*10)
ax.set_xlabel('$k$')
ax.set_ylabel('$E(k)$')
ax.legend()
plt.savefig('Abb/Ex7_Spectrum_wwith_epsilon')

```



7.2) Calculate the scale resolved energy dissipation rate  $\epsilon_r(x) = \int_x^{x+r} dx' \epsilon(x')$ .

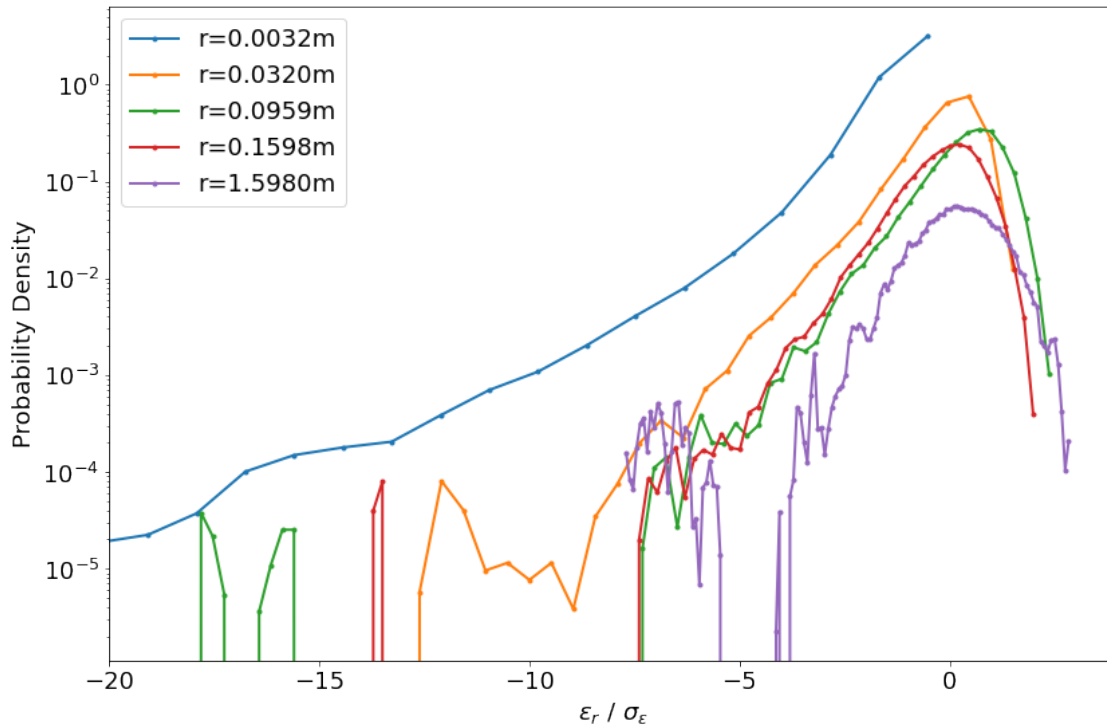
For large scale  $r$  we have a more Gaussian curve and for very small scales we have a different probability function (log-normal).

```
[ ]: fig, ax = plt.subplots(1,1,figsize=(12,8), tight_layout=True)
r_vec_idx = [10,100,300,500,5000]
std_eps_vec = np.empty(len(r_vec_idx))

for i, r in enumerate(r_vec_idx):
    epsilon_r = np.cumsum(np.roll(epsilon,r)[r:]) - np.cumsum(epsilon[r:])
    hist, bins = np.histogram(epsilon_r, density=True, bins=128)
    std_epsilon = np.std(epsilon_r)
    std_eps_vec[i] = std_epsilon
    ax.semilogy(bins[:-1]/std_epsilon, hist, marker='.', label=f'r={dr*r} :.
    ↪4f}m')

ax.set_xlabel('$\epsilon_r \ / \ \sigma_{\epsilon}$')
ax.set_ylabel('Probability Density')
plt.xlim(-20, 4)
ax.legend()
plt.savefig('Abb/Ex7_Resolved_energy_dissipation')
```





6.5) Check whether the histogram follows a lognormal distribution and try to determine the intermittency coefficient from  $\sigma(r)^2 = A + \mu \ln(\frac{L}{r})$ . For  $n = 3$  we had from lecture 15.06  $\langle \epsilon \rangle = ae^{\sigma^2/2}, \sigma^2 = A + \mu \ln \frac{L}{r}$

For large  $r$  the fit is not good because it is more a normal Gaussian

```
[ ]: fig, ax = plt.subplots(1,1,figsize=(15,10), tight_layout=True)
r_vec_idx = [10, 100, 300, 500, 5000]
std_eps_vec = np.empty(len(r_vec_idx))
sigma_vec = np.empty(len(r_vec_idx))
xdat_smooth = np.linspace(0,20, 100)
for i, r in enumerate(r_vec_idx):

    epsilon_r = np.cumsum(np.roll(epsilon,r)[r:]) - np.cumsum(epsilon[r:])
    hist, bins = np.histogram(epsilon_r, density=True, bins=128)
    std_epsilon = np.std(epsilon_r)

    arg_max = np.argmax(hist)
    N_l, N_r = 20, 20
    # Keep the index in bounds
    if arg_max + N_l >= 128:
        N_l = 128 - arg_max

    xdat = (bins[arg_max-N_r:arg_max+N_l])
```

```

xdat = -(xdat - np.max(xdat))
xdat = xdat /std_epsilon
ydat = hist[arg_max-N_r:arg_max+N_l]

ax.semilogy(xdat, ydat, marker='.', label=f'r={{(dr*r):.4f}}m')

# Fit Log Normal model to the data points
mod = lmfit.models.LognormalModel()
result = mod.fit(ydat, x=xdat)
sigma = result.values['sigma']
sigma_vec[i] = sigma
print(f'sigma: {sigma}')

ax.plot(xdat_smooth, result.eval(result.params, x=xdat_smooth), c='k',
ls='--', alpha=0.8, label=f'LogNorm Fit  $\chi^2={{result.chisqr:.2e}}$ ,
 $\sigma={{sigma:.3f}}$ ')

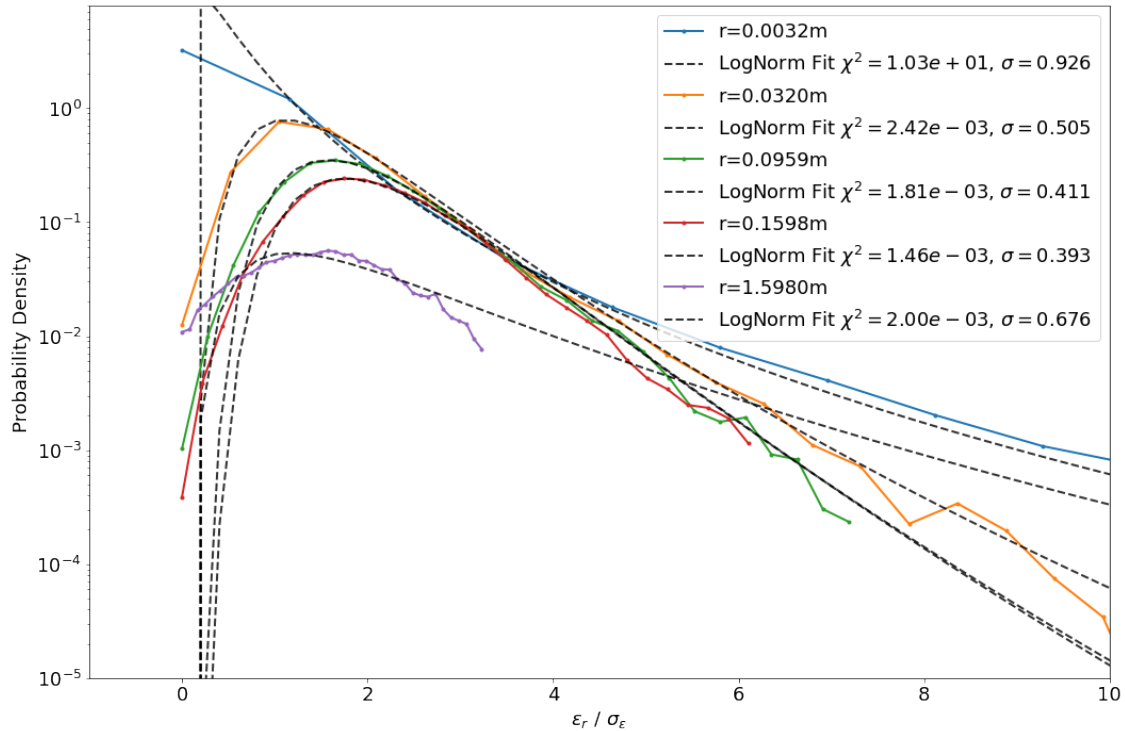
ax.set_xlabel('$\epsilon_r \ / \ \sigma_{\epsilon}$')
ax.set_ylabel('Probability Density')
ax.set_ylim(10**(-5),8)
ax.set_xlim(-1,10)
ax.legend()
plt.savefig('Abb/Ex7_Resolved_energy_dissipation_sigma_fit')

```

```

sigma: 0.9263541881786528
sigma: 0.505437918046481
sigma: 0.4113939820145063
sigma: 0.39252788746521516
sigma: 0.675599521066637

```



Plot  $\sigma^2$  over  $1/r$

```
[ ]: fig, ax = plt.subplots(1,1,figsize=(8,6), tight_layout=True)
r = np.array([spatial[i] for i in r_vec_idx])
ax.plot(np.log(L_fit/r), sigma_vec**2, marker='.', ms=10, label='$\sigma^2$')

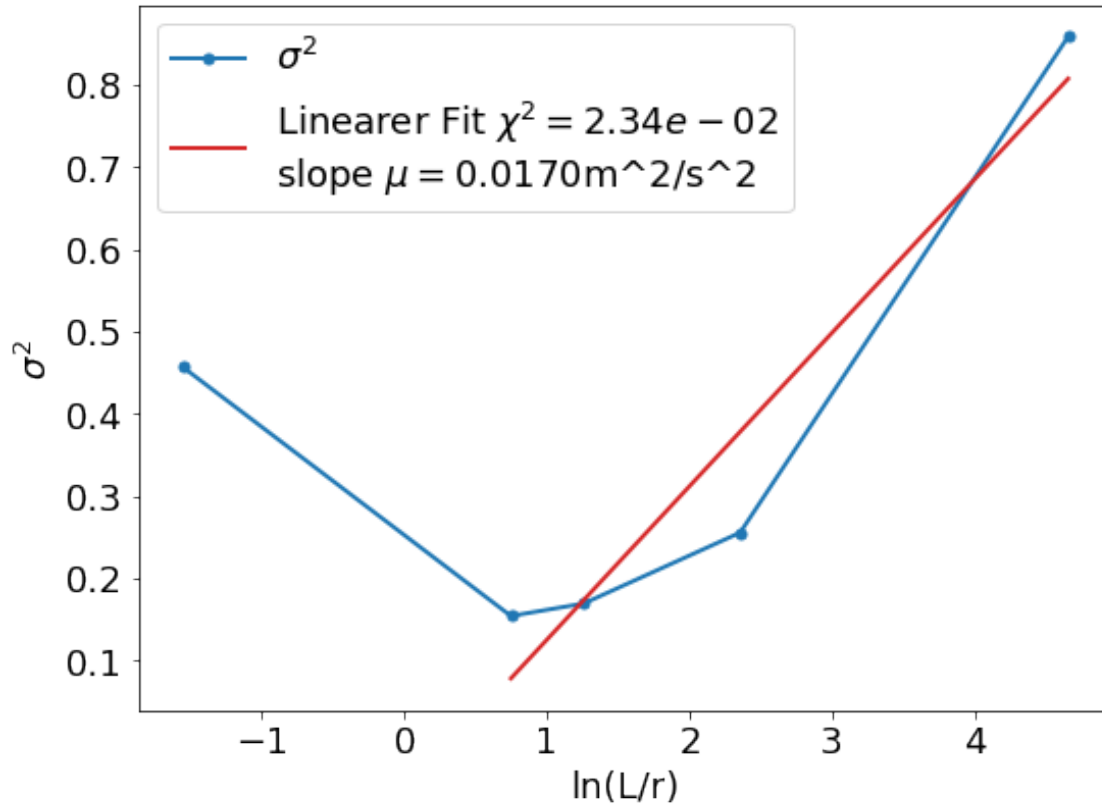
# Fit Linear Model to the orange data points
mod = lmfit.models.LinearModel()
xdat = np.log(L_fit/r)[: -1]
ydat = sigma_vec[: -1]**2

result = mod.fit(ydat, x=xdat)
mu_computed = result.values['slope']
print(f'Average dissipation rate by structure function = {mu_computed:.4f}m^2/
      ↪ s^3')
ax.plot(xdat, result.eval(result.params, x=xdat), c='C3', label=f'Linearer Fit
      ↪ $\chi^2={{result.chisqr:.2e}}$nslope $\mu = {{eps_avg_spectrum:.
      ↪ 4f}}$m^2/s^2')

ax.set_xlabel('ln(L/r)')
ax.set_ylabel('$\sigma^2$')
```

```
ax.legend()
plt.savefig('Abb/Ex7_intermittency_via_sigma_squared')
```

Average dissipation rate by structure function =  $0.1863\text{m}^2/\text{s}^3$



## 5 Exercise 8

8.1) Calculate flatness  $S_4(r)/3S_2(r)^2$ ,  $S_n(r) = r^{\zeta_n}$  with  $\zeta_n = \frac{n}{3} - \frac{18}{\mu}n(n-3)$  so we get as flatness  $= \frac{1}{3}r^{-\frac{4\mu}{9}}$ . Flatness with  $S_2(r)^2$  NOT 3 as on the exercise sheet. The flatness for a Gaussian would be one because the quadratic  $n^2$  term was only there due to the Non Gaussianity.

The  $\mu$  via flatness is in a good and plausible range

```
[ ]: fig, ax = plt.subplots(1,1,figsize=(10,7), tight_layout=True)
ax.semilogx(r_in_log, struc[4]/(3*struc[2]**2), label='Data of flatness',
            marker='.', ms=10)
idx_b, idx_e = 13,28

xdat = r_in_log[idx_b:idx_e]
ydat = (struc[4]/(3*struc[2]**2))[idx_b:idx_e]
ax.semilogx(xdat, ydat, label='Data for Fit, Inertial range', marker='.', ms=10)
```

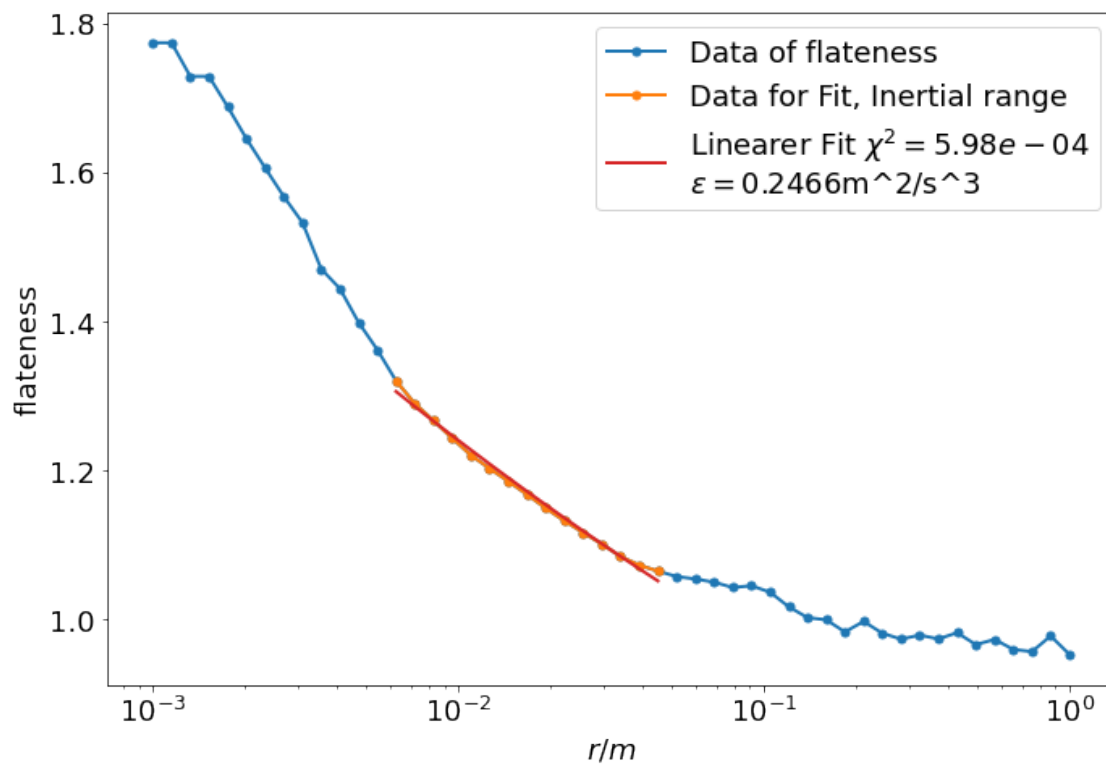
```

# Fit Linear Model to the orange data points
mod = lmfit.models.PowerLawModel()
# mod.set_param_hint('amplitude', vary=True)
# pars = mod.make_params(amplitude=1/3)
result = mod.fit(ydat, x=xdat)
mu_flatness = result.values['exponent']
print(f'Mu via flatness = {mu_flatness*(-9/4):.4f}')
ax.plot(xdat, result.eval(result.params, x=xdat), c='C3', label=f'Linearer Fit_')
↳ $\chi^2 = \{ \{ result.chisqr : .2e \} \} \$ \n \$ \epsilon = \{ \{ mu\_flatness * (-9/4) : .
↳ 4f \} \} \$ m^2 / s^3 '$

ax.set_xlabel('$r / m$')
ax.set_ylabel('flatness')
ax.legend()
plt.savefig('Abb/Ex8_flatness_and_mu')

```

Mu via flatness = 0.2466



[ ]: