

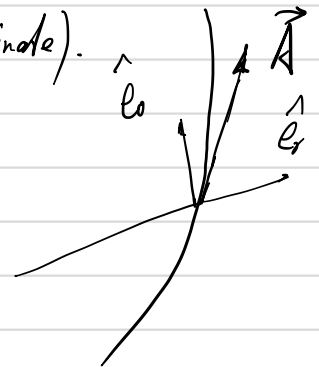
# HW 6 of Plasma

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1. (a) for any tangent vector on line  $r = R \cdot \sin^2 \theta$  (spherical coordinate).

$$\frac{A_r}{A_\theta} = \frac{1}{r} \cdot \frac{dr}{d\theta}, \text{ apply } \vec{B} \text{ into } \vec{A}$$

$$\frac{1}{r} \cdot \frac{dr}{d\theta} = \frac{2\cos\theta}{\sin\theta} \Rightarrow \ln r = 2 \ln \sin\theta + C$$



define  $r(\theta = \pi/2) = R$ , so  $C = R$ ,  $r = R \cdot \sin^2 \theta$  #

$$(b) \quad \vec{b} = \vec{B}/B = \frac{1}{\sqrt{1+3\cos^2\theta}} (2\cos\theta \cdot \hat{e}_r + \sin\theta \cdot \hat{e}_\theta)$$

$$\text{given } \nabla(\vec{f} \cdot \vec{g}) = (\vec{g} \cdot \nabla) \vec{f} + (\vec{f} \cdot \nabla) \vec{g} + \vec{g} \times (\nabla \times \vec{f}) + \vec{f} \times (\nabla \times \vec{g})$$

replace  $\vec{g}, \vec{f}$  by  $\vec{b}$ , and apply  $(\vec{b})^2 = 1$

$$(\vec{b} \cdot \nabla) \vec{b} = \frac{1}{2} \nabla(b^2) - \vec{b} \times (\nabla \times \vec{b}) = -\vec{b} \times (\nabla \times \vec{b})$$

$$b_\phi = 0, \text{ so } \nabla \times \vec{b} = \frac{1}{r} (\partial_r(r b_\theta) - \partial_\theta b_r) \hat{e}_\phi, \quad \nabla \times \vec{b} \perp \vec{b}$$

$$\text{so } |(b \cdot \nabla) b| = |\vec{b}| \cdot |\nabla \times \vec{b}| = |\nabla \times \vec{b}|, \text{ now calculate } \nabla \times \vec{b}$$

$$\partial_r(r b_\theta) = b_\theta, \quad \partial_\theta b_r = \frac{2\sin\theta}{1+3\cos^2\theta} \cdot \frac{1}{\sqrt{1+3\cos^2\theta}}$$

$$|\nabla \times \vec{b}| = \frac{1}{r} \left| \left( \sin\theta + \frac{2\sin\theta}{1+3\cos^2\theta} \right) \cdot \frac{1}{\sqrt{1+3\cos^2\theta}} \right| = \left| \frac{\sin\theta}{r} \cdot \frac{3(1+\cos^2\theta)}{1+3\cos^2\theta} \right|$$

$$\text{at equator, } \theta = \frac{\pi}{2}, \sin\theta = 1, \cos\theta = 0, r = R \sin^2\theta = R$$

$$R_c = 1/|\nabla \times \vec{b}| = R/3 \quad \#$$

(c). Curvature drift  $\vec{v}_R = \frac{m v_{\perp}^2}{q B^2} \cdot \frac{\vec{R}_c \times \vec{B}}{R_c^2}$ , at equator

$$\vec{R}_c = \frac{1}{3} R \cdot \hat{e}_r, \quad \vec{B} = \frac{\mu_0 M}{4\pi} \cdot \frac{1}{R^3} \cdot \hat{e}_\theta$$

so  $\vec{v}_R = \frac{3 m v_{\perp}^2 R^2}{q} \cdot \frac{4\pi}{\mu_0 M} \cdot \hat{e}_\phi$

(d)  $\nabla B$  drift  $\vec{v}_{\nabla B} = \frac{m v_{\perp}^2}{2 q B^2} \cdot \frac{\vec{R}_c \times \vec{B}}{R_c^2}$ , applying (c)

$$\vec{v}_{\nabla B} = \frac{3 m v_{\perp}^2 R^2}{2 q} \cdot \frac{4\pi}{\mu_0 M} \cdot \hat{e}_\phi$$

(e)  $\vec{v}_R$  and  $\vec{v}_{\nabla B}$  have the same direction, while  $|v_R| = 2 |v_{\nabla B}|$

2. Constancy of magnetic momentum indicates the vertical kinetic energy  $E_{\perp}$

has form of:  $\frac{E_{\perp}}{2 B_z} = \frac{E_{\perp}}{2 B_z} \Big|_{z=0}$

so  $v_{\perp}^2 = v_{\perp,0}^2 \left(1 + \left(\frac{z}{L}\right)^2\right)$ , and  $E_{\parallel}$  will reduce to

$$E_{\parallel} = E_{\parallel,0} - (v_{\perp}^2 - v_{\perp,0}^2)$$

$$\Rightarrow v_{\parallel}^2 = v_{\parallel,0}^2 - v_{\perp,0}^2 \cdot \left(\frac{z}{L}\right)^2, \text{ having harmonic oscillation form}$$

on  $z$  direction:  $\frac{1}{2} m \cdot v_{\parallel}^2 + \frac{1}{2} m \cdot \frac{v_{\perp,0}^2}{L^2} \cdot z^2 = \frac{1}{2} m \cdot v_{\parallel,0}^2 \quad (\text{const})$

so bounce frequency  $f = \frac{1}{2\pi} \sqrt{\frac{k}{m}} \quad (k = m \cdot \frac{v_{\perp,0}^2}{L^2})$

as  $k = \frac{2 E \sin^2 \alpha_0}{L^2}$ ,  $f = \frac{\sin \alpha_0}{2\pi L} \sqrt{\frac{2 E}{m}}$

3. (a) normalization parameter  $A$  satisfies  $\int_{-\infty}^{+\infty} f(x) dx = 1$

$$\text{so } \psi_0 \cdot A = 1 \quad A = \psi_0^{-1}$$

$$(b) 1) \langle x \rangle = \int_{-\infty}^{+\infty} x f(x) dx, \text{ since } f \text{ is even function, and } x \text{ is odd}$$
$$x \cdot f \text{ is odd, so } \langle x \rangle = 0$$

$$2) \langle |x| \rangle = 2 \int_0^{+\infty} x f(x) dx = \frac{2}{\psi_0^2} \int_0^1 x(\psi_0 - x) dx = \frac{1}{3} \psi_0$$

$$3) \langle \frac{1}{2} m x^2 \rangle = \frac{1}{2} m \cdot 2 \int_0^{+\infty} x^2 f(x) dx = \frac{m}{\psi_0^2} \int_0^1 x^2(\psi_0 - x) dx = \frac{1}{12} m \cdot \psi_0^2$$