

**Search for Long-Lived Neutral Particles decaying to
Photons and Missing Transverse Energy in pp Collisions at
 $\sqrt{S} = 8$ TeV**

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Acknowledgements

Dedication

What ever it is that doesnt workout, get up...I think youve got to realized that there is no personal ambition you have which can be extinguished by anybody else...Only YOU by giving up your dreams can extinguish them. And if your dreams doesnt work out EXACTLY like you intended, it will still take you someplace interesting and you will make a DIFFERENCE. So, DONT be AFRAID TO FAIL because you probably will whether you are afraid to or not and it's scary. You just gotta get out. The world belongs to tomorrow not yesterday. Dont give anybody else permission to take your life away, just keep living and keep giving and never make the perfect enemy the good.Never think that what Im doing is too little to make a difference, THAT'S NOT TRUE!, THAT'S NOT TRUE! Do something everyday. Someday, for all of us, it will be our last day and what would matter will be all the steps we took along the way and what they amounted to, NOT the home run we hit on day X. I wish you well.

US Pres Bill Clinton, Dallas, TX

No, I have nothing against Mr. Einstein. He is a kind person and has done many good things, some of which will become part of the music. I will write to him and try to explain that the ether exists, and that its particles are what keep the Universe in harmony, and the life in eternity.

Nikola Tesla

The Universe is unbiased, she chooses to reveal herself to whomever she wants.

Tambe E. Norbert

Abstract

We have performed a search for events with delayed photons and large missing transverse momentum from the decay of neutral long-lived particles produced in proton-proton collisions with a center of mass energy, $\sqrt{S} = 8$ TeV at the LHC. Capitalizing on the excellent timing resolution of the CMS electromagnetic calorimeter, we searched for events with delayed photons using the measured photon arrival time at the electromagnetic calorimeter. We observed a single event, consistent with our background expectations from the standard model and proceed to set limits on the cross section, $\sigma_{\tilde{\chi}_1^0} > 0.02$ pb, for the production and decay of the lightest neutralino ($\tilde{\chi}_1^0$) with mass, $m_{\tilde{\chi}_1^0} \geq 235$ GeV/ c^2 , and lifetime, $\tau_{\tilde{\chi}_1^0} \geq 35$ ns, according to a R-parity conserving Gauge Mediated Supersymmetry Breaking model. We show, for the first time, that using only timing information from the CMS electromagnetic calorimeter, the CMS detector is sensitive to neutralinos with lifetimes up to 40 ns and masses up to 260 GeV/ c^2 . A description of timing measurement by the electromagnetic calorimeter and its performance is presented.

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Chapter 1

Introduction

The search for Long-Lived Neutral Particles (LLNP) produced in pp collisions at $\sqrt{S} = 8$ TeV by the Large Hadron Collider (LHC) on data recorded by the Compact Muon Solenoid (CMS) detector is presented in this thesis.

The Standard Model (SM) describes the interactions of fundamental particles with unmatched precision and has been very successful in terms of all its predictions of the properties of particles agreeing with experiments. However, there are strong indications that the SM is incomplete. For example, the SM does not include gravity in its formulation and cannot describe Dark Matter (DM) [1, 2, 3], which is about five times more abundant in the universe.

There are many extensions to the SM where DM and other new phenomena are described. For example, many models Beyond the Standard Model (BSM) predict that LLNPs, which are not described by the SM, can be produced at the LHC. These LLNPs can decay into stable, neutral particles which are good candidates for DM particles. The signature for the decay of a LLNP, is the presence of Missing Transverse Energy (MET), which indicates the escape of the undetected candidate DM particle, and electromagnetic particles, like photons, which are detected late with respect to LHC pp collision times, in the same event. The combination of a late photon and a large MET are the central idea in the search presented in this thesis.

The LHC completed a successful first long operation in 2012. Data from pp collisions during this period of operation (LHC RUN 1) was recorded by the CMS detector.

The CMS detector, including the Electromagnetic Calorimeter (ECAL), was fully operational and recorded the arrival time and energy of electromagnetic particles produced from pp collisions with excellent resolution. This resolution was maintained throughout the entire LHC RUN 1 by continuous energy and timing performance monitoring and calibration. The data collected during this period is what we used in the search presented here.

Understanding the signature of a LLNP in CMS detector is provided using Monte Carlo simulated signal samples of the production and decay of LLNPs. A brief introduction of the SM and Supersymmetry (SUSY) and the description of GMSB models,[16], with LLNPs is presented in Chapters 2.

In Chapter 3, the LHC and CMS particle detector are described with an emphasis on the sub-detectors used in this search for LLNPs.

A description of the photon timing measurements is discussed in Chapter 4, while the reconstruction of events and their constituent particles is described in Chapter 5, where the definition of quantities like MET and jets is presented.

Chapter 6 describes the search method in detail including the data samples and triggers used, trigger efficiency, search observable, event selection requirements and background estimation techniques. The systematics arising from signal efficiency and acceptance in our signal event selection are also discussed in this chapter. In Chapter 7, we describe the statistical methods used to estimate the precision of the measurement.

Limits on production cross sections times branching ratio for different lifetimes and mass of a LLNP, within the context of our benchmark GMSB model is discussed in Chapter 8. The interpretation of our results in terms of the possible extensions to the SM by our analysis is also presented in this chapter and in Chapter 9 the conclusions are presented.

Chapter 2

Phenomenology of Long-Lived Particles

2.1 The Standard Model

The Standard Model (SM) is a mathematical description of the fundamental components of matter and their interactions. Fermions are the building blocks of matter and their interaction is mediated by vector bosons. Both fermions and bosons get their mass by interacting with a scalar boson in a process of spontaneous gauge symmetry breaking,[6].

2.1.1 Fermions

Fermions are half-integer spin ($\frac{1}{2}\hbar$) particles and exists in nature as either, together with their anti-particle (has an opposite charge to its corresponding particle), as either leptons (ℓ) or quarks (q). Both leptons and quarks come in 3 *generations* or *flavors* and according to the SM we don't know why this is the case. The 3 generations of quarks and leptons are arranged in a mass hierarchy with the third generation being the heaviest. The second and third generations are meta-stable and do disintegrate (decay) into the first generations. The lifetime of the heaviest SM fermion, top quark (t), is a few orders of *yoctosecond* (ys) (10^{-24} seconds), according to the SM.

Quarks carry fractional electric charges and participate in weak and electromagnetic interactions and also in strong interactions since in addition to the electric charge, they

also carry a *color* charge. Each generation of quarks is a pair of an “*up-type*” and a “*down-type*” quark. The *Up-type* quarks (**up** (u), **charm** (c), **top** (t)) have a charge of $+\frac{2}{3}$, while the *down-type* quarks (**down** (d), **strange** (s), **bottom** (b)) have a charge of $-\frac{1}{3}$. Quarks do not exist as free particles but are found in nature bound together in the form of composite particles called *Hadrons*.

Leptons on the other hand exist free in nature and come in pairs consisting of a given lepton flavor and its corresponding neutrino. Leptons from all 3 generations (**electron** (e), **muon** (μ), **tau** (τ)) have a charge of -1 , while their corresponding neutrinos (**electron neutrino** (ν_e), **muon neutrino** (ν_μ), **tau neutrino** (ν_τ)) have a charge of 0 . Leptons can participate only in weak and electromagnetic interactions. In Table 2.1, we present a summary of the fermions in the SM.

2.1.2 Bosons

Bosons have integer spins. According to experiments, two types of bosons exists: scalar ($0\hbar$) and vector ($1\hbar$) bosons. Vector bosons are responsible for mediating the interactions between fermions and the only scalar boson (Higgs boson), discovered so far, provide mass to both fermions and bosons.

The vector bosons W^\pm (charged) and Z (neutral) are massive and mediate the *weak* interactions, while the massless vector bosons: *photon* (γ) and 8 *gluons* (g) mediate the *electromagnetic* and *strong* interactions, respectively [6, 7]. We provide in Table 2.2, the vector bosons and the interaction they mediate in the SM.

2.1.3 Spontaneous Symmetry Breaking

In the SM, the weak, electromagnetic and strong interactions can be described using *gauge* symmetries and a *conserved quantum quantum* associated to the symmetry. The gauge symmetry of the SM is a combination of 3 different gauge symmetries

$$SU(3)_C \otimes SU(2)_L \otimes U(1)_Y. \quad (2.1)$$

where, $SU(3)_C$ is the gauge symmetry describing strong interactions with *color* (C) charge as its conserved quantum quantity. Quarks participate in strong interactions mediated by 8 gluons because of the color charge

$SU(2)_L \otimes U(1)_Y$ is the gauge symmetry describing the weak and electromagnetic interactions also known as the *electro-weak* interaction and the conserved quantum number derived from a combination of conserved quantum numbers: *isospin* (T_3) of $SU(2)_L$ and *hypercharge* (Y) of $U(1)_Y$ symmetry. After the spontaneous breaking of this electro-weak symmetry ($SU(2)_L \otimes U(1)_Y \rightarrow U(1)_Q$), the resulting conserved quantum number describing the electromagnetic interaction which is based on the $U(1)_Q$ symmetry, is the electric charge, Q . The electric charge is related to T_3 and Y through the relation

$$Q = T_3 + \frac{Y}{2}. \quad (2.2)$$

Corresponding to the $SU(2) \otimes U(1)$ symmetry, are 4 massless vector bosons or *eigen* states, $W_\mu^{1,2,3}, B_\mu$, which combine to form the physical electro-weak vector bosons or *mass* states, W^\pm, Z, γ . These mass states are related to the eigen states through a rotation which is expressed as

$$W^\pm = \frac{W_\mu^1 \pm iW_\mu^2}{\sqrt{2}}, \quad \begin{pmatrix} \gamma \\ Z \end{pmatrix} = \begin{pmatrix} \cos \theta_w & \sin \theta_w \\ -\sin \theta_w & \cos \theta_w \end{pmatrix} \begin{pmatrix} B_\mu \\ W_\mu^3 \end{pmatrix} \quad (2.3)$$

The angle, θ_w , is known as the *Weinberg angle* and relates the physical electric charge, e , to the coupling strengths, g and g' of $SU(2)_L$ and $U(1)_Y$ gauge symmetries, respectively, in a relation given as $e = g \sin \theta_w = g' \cos \theta_w$. This angle is one of the parameters of the SM and has been measured from experiments to be equal to 28.726 degrees, which is, surprisingly, equal to the arc cosine of the ratio of the mass of the W boson to that of the Z boson.

According to the SM, the only way to realize the massive vector bosons and explain how fermions get their mass is through the *Higgs mechanism*, [8], and involves spontaneously breaking the gauge symmetry of the SM, expressed as

$$SU(3)_C \otimes SU(2)_L \otimes U(1)_Y \xrightarrow{\text{SSB into}} SU(3)_C \otimes U(1)_Q. \quad (2.4)$$

In the SM, spontaneous symmetry breaking is realized through the potential of a spin-0 complex scalar field called the *Higgs boson*. The Higgs boson takes the representation of an weak isospin *doublet*. Its potential, $V(\phi^*, \phi) = \mu^2(\phi^*\phi) + \lambda(\phi^*\phi)^2$, is written in terms of the Higgs field (ϕ) and its complex conjugate (ϕ^*), where μ and λ

are parameters. Among the many possible choices of configuration for this potential, only one particular configuration, represented by the choice of the parameters: $\mu^2 < 0$ and $\lambda > 0$, such that the minimum value of this potential or the *Vacuum Expectation Value* (VEV) of the Higgs field, $|\phi_0| = \sqrt{\frac{-\mu^2}{\lambda}} = \nu \approx 246 \text{ GeV}$, spontaneously breaks the $SU(2)_L \otimes U(1)_Y$ symmetry into a $U(1)$ symmetry. This particular configuration of the potential of the Higgs field is shown in Figure 2.1.

Fermions get their mass through their interaction with the Higgs field and the amount of the fermion's mass, m_f , is proportional to the strength of its interaction (Yukawa coupling λ_f) with the Higgs field, given in Equation 2.5.

The massive vector bosons, Z and W^\pm , get their mass, $m_Z \approx 80 \text{ GeV}/c^2$ and $m_{W^\pm} \approx 91 \text{ GeV}/c^2$, respectively, by engulfing or “*eating*” the available massless components (*Nambu-Goldstone bosons*) of the complex Higgs doublet. Out of the four scalar fields, only one physically massive scalar particle *Higgs boson* remains. The masses of the charge and neutral vector bosons, given in Equation 2.5, are connected through the cosine of the *Weinberg angle*.

$$m_f = \lambda_f \frac{\nu}{\sqrt{2}}, \quad \frac{m_{W^\pm}}{m_Z} = \frac{\frac{1}{2}\nu g}{\frac{1}{2}\nu \sqrt{g^2 + g'^2}} = \cos \theta_w \quad (2.5)$$

It took nearly 40 years for the Higgs boson to be discovered after its prediction in the 1960s. The Higgs boson was discovered at CERN on July 04, 2012, through its decay into two photons, $H \rightarrow \gamma\gamma$, and a pair of Z bosons, $H \rightarrow ZZ$, channels. Its mass has been measured to be, $m_H = 125 \pm 0.21 \text{ GeV}/c^2$. As shown in Figure 2.2, every SM particle interacts with the Higgs boson either directly or indirectly as in the case for photons and gluons.

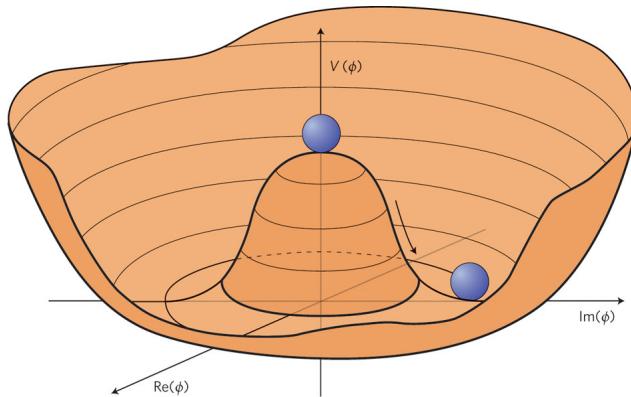


Figure 2.1: Higgs boson “Mexican hat” potential, $V(\phi^*\phi) = \mu^2(\phi^*\phi) + \lambda(\phi^*\phi)^2$, which leads to spontaneous symmetry breaking with choice of parameters $\mu^2 < 0$, $\lambda > 0$.

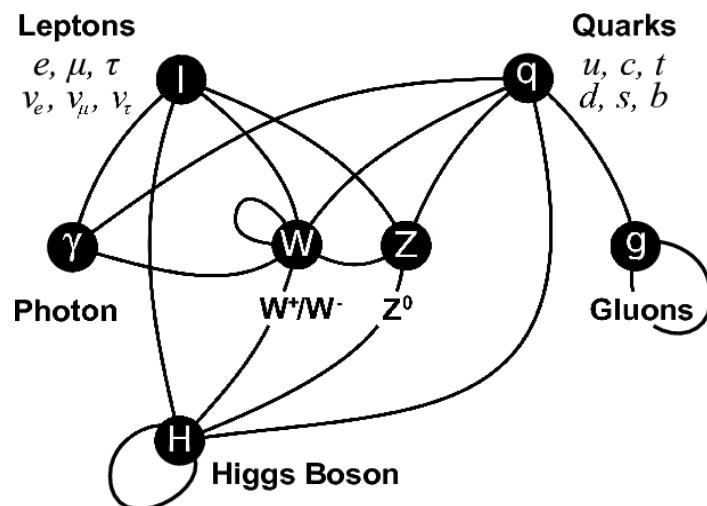


Figure 2.2: Fermions and bosons and their interactions in the SM. All bosons (except the Higgs) mediate interactions between fermions.

Fermions	Generation			Charge
	First	Second	Third	
Leptons	$\begin{pmatrix} e^- \\ \nu_e \end{pmatrix}$	$\begin{pmatrix} \mu^- \\ \nu_\mu \end{pmatrix}$	$\begin{pmatrix} \tau^- \\ \nu_\tau \end{pmatrix}$	$\begin{pmatrix} -1 \\ 0 \end{pmatrix}$
Quarks	$\begin{pmatrix} u \\ d \end{pmatrix}$	$\begin{pmatrix} c \\ s \end{pmatrix}$	$\begin{pmatrix} t \\ b \end{pmatrix}$	$\begin{pmatrix} +\frac{2}{3} \\ -\frac{1}{3} \end{pmatrix}$

Table 2.1: Fermions of the SM. Particle symbols are explained in text.

Bosons	Interaction	Symmetry	Characteristic Lifetime
W^\pm, Z	Weak	$SU(2)_L \otimes U(1)_Y$	10^{-8} to 10^{-13} seconds
γ	Electromagnetic	$U(1)_Q$	10^{-14} to 10^{-20} seconds
g	Strong	$SU(3)_C$	$< 10^{-22}$ seconds

Table 2.2: Interaction mediating vector bosons in the SM and the characteristic lifetime of the interaction.

2.1.4 Beyond Standard Model Physics

It is obvious that the Higgs boson is a crucial component in the formulation and experimental success of the SM, yet the SM itself provides no explanation for the above choice of parameters of the Higg's potential or whether there is only one type of Higgs field to which all fermions couple with to get their mass. Some models Beyond the SM (BSM), such as *Supersymmetry*, allows for the possibility of more than one Higgs field and predicts the existence of additional particles to those in the SM. The following are some of the reasons for the interest in BSM physics, particularly Supersymmetry:

Explaining the Stability Of the Higgs Bosons's Mass

Predictions of the Higgs boson's mass, m_H^2 , by the SM, include corrections to the Higgs boson's mass, δm_f^2 , from its coupling with fermion. These corrections from first order interaction (1-loop level) of the Higgs boson with fermions are computed using the Feynman diagram shown in Figure 2.3(a) and given as

$$\delta m_f^2 = \frac{1}{16\pi^2} |\lambda_f|^2 \left(-2\Lambda^2 + 6m_f^2 \ln \left(\frac{\Lambda}{m_f} \right) + \dots \right) \quad (2.6)$$

where, λ_f is the coupling parameter of the Higgs boson to fermions with the interaction Lagrangian density written as $\lambda_f H \bar{f} f$. Λ represents an arbitrary high energy scale (of the order 10^{18} GeV) known as the *cut-off* energy scale. Because $\Lambda = 10^{18}$ GeV is very large, we would expect the correction to the Higgs boson's mass to also be very large eventually making the Higgs boson's mass, m_H^2 , to be very large. However, the Higgs boson's mass measured from experiment is about 125 GeV/ c^2 , which is not very large as expected. A dilemma thus arises, if one is to trust SM predictions of the Higgs boson's mass, why are all these large corrections to the Higgs boson's mass not observed in experiments? The SM cannot provide an explanation for this. This problem is described as the stability of the Higgs mass.

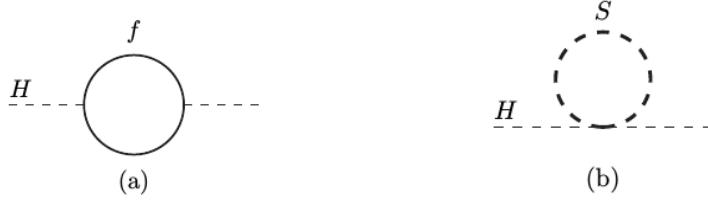


Figure 2.3: Higgs mass contributions arising from the Higgs field coupling to fermions (a) and scalar (b) fields.

On the other hand, SUSY, provides a plausible explanation as to why these corrections are not observed in experiments. According to SUSY, there are in addition to the scalar Higgs boson new supersymmetric scalar particles, which are yet to be observed, that can also couple to the Higgs field as shown in the 1-loop Feynman diagram in Figure 2.3(b). The contributions to the Higgs boson's mass from these new scalar supersymmetric particles, δm_S^2 , is computed at the order of 1-loop and given as

$$\delta m_S^2 = \frac{1}{16\pi^2} |\lambda_S|^2 \left(\Lambda^2 - 2m_S^2 \ln \left(\frac{\Lambda}{m_S} \right) + \dots \right) \quad (2.7)$$

where λ_S is also the coupling constant of the Higgs field to the new scalar particle given as, $\lambda_S H S \bar{S}$. An interesting observation comparing Equations 2.6 to Equation 2.7, up to the Λ^2 -terms is that their signs are different. The opposite sign corrections to the Higgs boson's mass cancel each other and as a result their net contribution to the Higgs mass, m_H^2 is zero. This cancellation which happens at all orders of corrections is the reason why one does not observe very large Higgs boson mass from experiment.

Dark Matter and Long-Lived Particles

There is enough evidence [1], that dark matter (matter which does not interact with light) exists in the universe and if this form of matter is made of particles, then, the lifetime of a dark matter particle must be comparable to the age of the universe it must also be neutral. Such neutral and stable particles do not exist in the SM. Therefore, dark matter particles must certainly be interacting with a new type of interaction which, obviously, cannot be described by the SM.

Long-Lived Neutral Particles (LLNP) which are meta-stable (particles that are

stable for a period of time and then decay) and can decay into possible dark matter particles through new kind of interactions also do not exist in the SM. These LLNPs are predicted, by models BSM, to have lifetimes beyond the characteristic lifetimes of particles and interactions in the SM seen in Table 2.2. In addition to Supersymmetry providing a possible explanation for stability of the Higgs boson's mass, it also predicts the existence of LLNPs which can decay into dark matter candidate particles with lifetime beyond those of the SM,[2, 3, 16, 4, 5, 18]. These LLNPs decay into dark matter candidate particles through new kind of interactions. Such predictions from Supersymmetry has motivated many of the studies of Supersymmetry as an interesting candidate among many models beyond the SM, in both theory and experiment. It is certainly a motivation for the study presented in this thesis.

2.2 Supersymmetry

Supersymmetry (SUSY) relates space-time symmetries (rotation and translation) to gauge symmetries ($SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$) [10, 11]. The generators of SUSY, Q , transform fermions to bosons and bosons to fermions;

$$Q|\text{Fermion}\rangle = |\text{Boson}\rangle, \quad Q|\text{Boson}\rangle = |\text{Fermion}\rangle, \quad (2.8)$$

and as a consequence, the fermions and their corresponding bosons have the same mass.

2.2.1 The Minimal Supersymmetric Standard Model

The *Minimal Supersymmetric Standard Model* (MSSM) uses only one SUSY generator, Q , in its formulation as an extension of the SM, to include the supersymmetric partners of the particles in the SM. The number of fundamental particles in the MSSM is obviously doubled, however, the gauge symmetries are the same. Supersymmetric particles and their SM counterparts differ by a half-integer in spin (s). Bosons (fermions) in the SM have superpartners which are fermions (bosons).

Two Higgs complex scalar fields are required to give mass to the **up-typed** and **down-typed** quarks and leptons, [12, 13], and their superpartners are fermions called *higgsinos*. The names of supersymmetric particles are derived from their SM counterparts by adding an “*s*” in front of the SM particle name. For example, a *selectron* is the supersymmetric partner of the electron, *squarks* are the supersymmetric partners of SM quarks. The symbol of SUSY particles carry the “~” sign above the symbol. For example if q is a SM particle then its supersymmetric partner is written as \tilde{q} . Superpartners of bosons are indicated with the suffix “-ino”, as in *wino*, *Higgsinos*, *binos*, *neutralinos*, *charginos* and *gluinos*.

The superpartners of SM fermions are scalars called *sfermions* (\tilde{l}), sneutrinos ($\tilde{\nu}$) and squarks (\tilde{q}) while *gluinos* (\tilde{g}) are the superpartners of the massless gauge bosons of strong interaction, gluons. The vector bosons of electro-weak interaction have fermionic superpartners *Winos* and *Binos*. These can mix to form a pair of mass eigenstates called *charginos* ($\tilde{\chi}_j^\pm, j = 1, 2$), i.e., $\tilde{\chi}_{1,2}^\pm$ are mixtures of $\tilde{W}^+, \tilde{W}^-, \tilde{H}^+, \tilde{H}^-$ and a quartet of mass eigenstates called *neutralinos* ($\tilde{\chi}_i^0, i = 1, \dots, 4$), i.e., $\tilde{\chi}_{1-4}^0$ are mixtures of

$\tilde{B}^0, \tilde{W}^0, \tilde{h}^0, \tilde{H}^0$. We summarize in Tables 2.3 and 2.4 particles in the MSSM and how they relate to their SM partners. The gauge and mass eigen mass states of fermionic supersymmetric particles in the MSSM is also shown.

Standard Model			Supersymmetry		
Particle	Symbol	Spin	Particle	Symbol	Spin
quark	q	$\frac{1}{2}$	squark	\tilde{q}	0
lepton	ℓ	$\frac{1}{2}$	slepton	$\tilde{\ell}$	0
W Bosons	W^\pm, W^0	1	Wino	$\tilde{W}^\pm, \tilde{W}^0$	$\frac{1}{2}$
B Boson	B	1	Bino	\tilde{B}	$\frac{1}{2}$
gluon	g	1	gluino	\tilde{g}	$\frac{1}{2}$
Higgs Bosons	$H \times 4$	0	higgsino	$\tilde{H} \times 4$	$\frac{1}{2}$
Graviton	G	2	gravitino	\tilde{G}	$\frac{3}{2}$

Table 2.3: Particles in the MSSM. SUSY particles (sparticles) have a “~” on the symbol

Particle Names	Gauge eigen states	Mass eigen states
squark	\tilde{q}	\tilde{q}
slepton	$\tilde{\ell}$	$\tilde{\ell}$
Neutralinos	$\tilde{W}^0, \tilde{B}^0, \tilde{H}_1^0, \tilde{H}_2^0$	$\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$
Charginos	$\tilde{W}^+, \tilde{W}^-, \tilde{H}^+, \tilde{H}^-$	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$
Higgs bosons	$H_1^0, H_2^0, H_1^-, H_2^+$	h^0, H^0, A^0, H^\pm
gluino	\tilde{g}	\tilde{g}
Gravitino	\tilde{G}	\tilde{G}

Table 2.4: Gauge and mass eigenstates of SUSY particles in the Minimal Supersymmetric SM (MSSM).

Supersymmetry Breaking

A consequence of SUSY is that supersymmetric particles and their SM partners must have the same mass. However, because no supersymmetric particle have been observed with the same mass as SM particles, it means SUSY is not an exact symmetry of nature and must be broken. SUSY breaking should happen such that the mass of supersymmetric particles are higher than their standard model counterparts. The two ways of breaking SUSY are explicit and spontaneous breaking. Similar to the motivation of spontaneous symmetry breaking through the Higgs mechanism in the SM, spontaneous SUSY breaking is theoretically favored because it preserves renormalisability also in the MSSM. SUSY breaking is realized by introducing SUSY-breaking terms known as *soft breaking terms* to the Lagrangian density, which ensures that coupling constants which may lead to quadratic divergence are not introduced and the mass of the supersymmetric particles can be at about the TeV energy scale [15, 14, ?]. In theory SUSY breaking happens in the so-called *Hidden Sector* and can be mediated to the MSSM sector through interactions which do not change the flavor of the particles called *flavor blind interactions* e.g. gauge interactions [14, 16]. SUSY models based on gauge interactions mediating SUSY breaking from hidden sector to MSSM sector are called *Gauge Mediated Supersymmetry Breaking* (GMSB) models. GMSB models are interesting at an experimental level because they predict the existence of long-lived particles which can decay into light gravitino, a spin-3/2 superpartner of the graviton field, which is a good candidate dark matter particle since it is stable and neutral.

Double Higgs

Two complex Higgs doublets are required in the MSSM, $\mathbf{H}_d = (H_1^0, H_1^-)$ and $\mathbf{H}_u = (H_2^+, H_2^0)$ to give mass to the down-typed quarks (and leptons) and to the up-typed quarks, respectively. One SUSY partner is introduced for each of the four Higgs particles. Particles interact with the Higgs bosons through a superpotential whose Lagrangian density can be written as

$$W_{\text{MSSM}} = \bar{u}\mathbf{y}_u Q H_u - \bar{d}\mathbf{y}_d Q H_d - \bar{e}\mathbf{y}_e L H_d - \mu H_d H_u \quad (2.9)$$

to obtain their masses. H_u , H_d , Q , L , \bar{u} , \bar{d} , \bar{e} are chiral superfields of the chiral supermultiplets (spin-1/2 and spin-0 fields),[13, 15]. The dimensionless couplings $\mathbf{y_u}, \mathbf{y_d}$, and $\mathbf{y_e}$ are 3×3 matrices of the Yukawa couplings. The superpartners of these Higgs bosons are fermions called *higgsinos* ($\tilde{H}_1^0, \tilde{H}_1^-, \tilde{H}_2^0, \tilde{H}_2^+$) and those of the gauge bosons are called *gauginos* ($\tilde{B}^0, \tilde{W}^0, \tilde{W}^-, \tilde{W}^+$). The neutral gauginos and higgsinos ($\tilde{B}^0, \tilde{W}^0, \tilde{H}_1^0, \tilde{H}_2^0$) mix together to produce four neutral fermions called *Neutralinos* ($\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$) while the charged gauginos and higgsinos ($\tilde{W}^-, \tilde{W}^+, \tilde{H}_1^-, \tilde{H}_2^+$) mix together to produce four charged fermions called *Charginos* ($\tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm$), presented in Table 2.4. The mass of the neutralinos and charginos depend on the mixing parameters of these gauginos and higgsinos and is model dependent.

R-Parity

In supersymmetry, it is, in principle, acceptable to include terms in the MSSM Lagrangian density which may not be conserved quantum numbers like the *baryon* (B) and *lepton* (L) numbers. For example, a consequence of adding such a term is that predictions of the proton's lifetime is much shorter than measured in experiments. Since no such phenomenon which violate these quantum numbers have been observed in experiments, SUSY models are constructed with the introduction of an additional matter symmetry called *R-Parity*, which relates the quarks to leptons through the baryon and lepton numbers. R-parity is defined as

$$R_P = (-1)^{3(B-L)+2S} \quad (2.10)$$

where, S is the particle's spin. SM particles like quarks have an *even* R-parity, $R_P = 1$, while supersymmetric particles like squarks have odd parity $R_P = -1$.

The phenomenological consequence of R-parity are the following: first, in the decay of supersymmetric particles, the lightest SUSY particle (LSP) have odd parity, $R_P = -1$, it is electrically neutral and considered to be absolutely stable. This makes the LSP a good dark matter candidate particle [2, 5], second, every supersymmetric particle produced will eventually decay into an odd number of LSPs, third, supersymmetric particles can only be produced in pairs. This phenomenological consequences make *R-parity Conserving* (RPC) SUSY models very attractive for experimental studies.

Cross Sections, Decay Rate and Branching Ratio

Cross Section

The cross section of producing a particle at a collider, like the LHC, is the probability that the proton beams will collide and interact in a certain way to produce the particle. The cross sections provides a way of counting the number of the particle produced when these proton beams collide. The cross section of producing a supersymmetric particle at the LHC depends on the following: the available energy of the proton beams compared to the mass of the particle, the type of interaction or size of the couplings during collision and the flux of the proton beams. Since the predicted mass of supersymmetric particles is higher than that of SM particles, we expect very few supersymmetric particles, if any, to be produced at the LHC, and even fewer supersymmetric particles for massive particles. The typical cross section of producing a supersymmetry particle at the LHC is of the order of $1 \text{ pb} = 10^{-12} \times 10^{-24} \text{ cm}^2$ or at times $1 \text{ fb} = 10^{-15} \times 10^{-24} \text{ cm}^2$ for extremely rare SUSY processes which can be compared to that for a standard model process, like, the production of the Z or W^\pm bosons which is of the order of a few $nb = 10^{-9} \times 10^{-24} \text{ cm}^2$.

Supersymmetric particles are most likely to be produced at the LHC, through the process of strong interactions ($pp \rightarrow \tilde{g}\tilde{g}, \tilde{q}\tilde{q}$) than electro-weak interactions processes ($pp \rightarrow \tilde{\chi}^\pm\tilde{\chi}^\mp, \tilde{\chi}^0\tilde{\chi}^\pm$), since the cross section for strong interaction processes is larger. Figure 2.4, shows the cross section, on the vertical y-axis, against the mass of the supersymmetric particle, on the horizontal x-axis. With 20 fb^{-1} of LHC integrated luminosity (relates the cross section and the rate of particle produced) of data from proton-proton (pp) collisions at $\sqrt{S} = 8 \text{ TeV}$, we expect to see more events (vertical z-axis) with supersymmetric particles produced at the LHC through strong interactions than electro-weak interactions for a given SUSY sparticle mass.

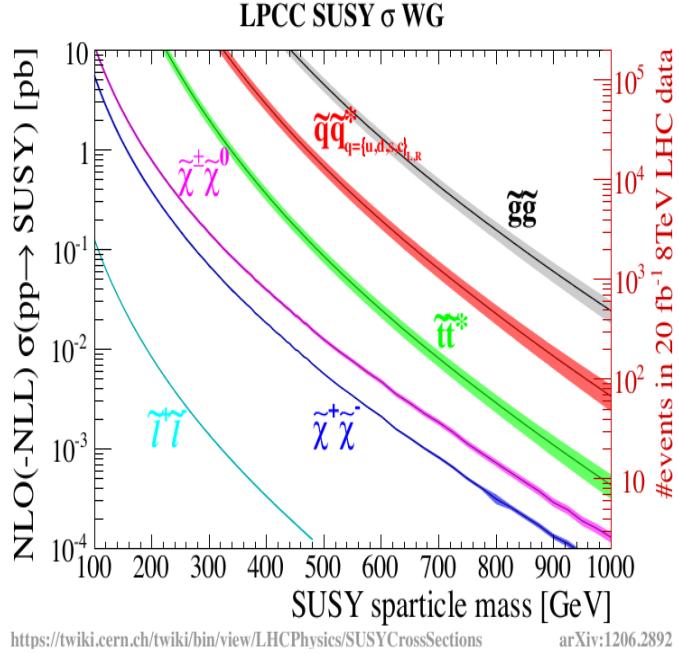


Figure 2.4: Cross section against SUSY sparticle mass for producing supersymmetric particles in different modes of interactions at the LHC. More supersymmetric particles are produced through strong interaction, $pp \rightarrow \tilde{g}\tilde{g}$, processes than the any other process.

Decay Rate and Branching Ratio

After the production of a SUSY sparticle, the particle may decay either immediately after it was produced to live for a while before it decays. The rate at which the SUSY sparticle decays is called the *decay width* (Γ) and it is related to the mean lifetime (τ) through the relation

$$\tau = \frac{\hbar}{\Gamma}. \quad (2.11)$$

Since the decay width depends on the coupling or the mass difference between the SUSY sparticle and its decay products, SUSY particles with small mass differences of large coupling to the daughter particles has smaller decay width and therefore have long lifetimes. The final products of the decay depends on the branching ratio for a given decay channel. In R-parity conserving models, SUSY sparticles are produced in pairs and each decay through a cascade decay process into the LSP and its SM partner.

A SUSY sparticle may decay through strong interactions in which case its mean lifetime

is about 10^{-17} to 10^{-25} seconds, through electromagnetic interactions in which case its mean lifetime maybe between 10^{-20} seconds to about 10^{-14} seconds and through weak interactions in which case its mean lifetime maybe vary from 10^{-13} seconds to 10^{-8} seconds.

By Long-lived particles, we are interested in particles decays through none of the SM interactions and has a typical mean lifetime ranging from a few nanoseconds to hundreds of a nanosecond. These are the kind of long-lived particles we study in this thesis.

2.2.2 Gauge Mediated Supersymmetry Breaking Models

Gauge Mediated Supersymmetry Breaking (GMSB) models are based on the idea that *Messenger* particles through gauge interactions are responsible for mediating SUSY breaking from the *hidden sector* to the visible or MSSM sector. The energy scale of the hidden sector, which is the Vacuum Expectation Value (VEV) of scalar fields responsible for SUSY breaking (some kind of *superHiggs mechanism*) is denoted by a fundamental parameter, \mathbf{F} , also known as the *Fundamental SUSY Breaking Parameter*. The mass of the messenger particle, \mathbf{M}_{mess} , is related to \mathbf{F} through an *effective SUSY breaking* energy scale parameter, Λ , at the MSSM level. This relationship can be written as $\mathbf{F}/\Lambda = \mathbf{M}_{\text{mess}}$ and Λ can be interpreted as the energy scale at which SUSY breaking is felt by the MSSM particles.

For a given choice of the value of \mathbf{F} , the mass spectrum of all supersymmetric particles including the Lightest SUSY Particle (LSP) is well defined. Since the value of \mathbf{F} is not known and can vary, the mass of the LSP can also vary. This variation of the fundamental SUSY breaking scale or equivalently the mass of the LSP is parametrized by a scaling parameter, C_{grav} .

In GMSB models and in SUSY models in general, the value of the VEVs, v_u and v_d , of the two Higgs doublets in the MSSM is not known and is therefore an additional parameter, $\tan(\beta)$, in most models. This parameter relates the VEVs; $\tan(\beta) = \frac{v_u}{v_d}$. Another parameter, $\text{sgn}(\mu)$, defines the sign, which is arbitrary, of the parameter μ in the Higgs Boson's potential.

Thus, in GMSB models there are only (compared to the SM which has at least 19) 5

input parameters given as

$$\{\mathbf{A}, \quad \mathbf{M}_{\text{mess}}, \quad \mathbf{N}_5, \quad \tan(\beta), \quad sgn(\mu), \quad C_{\text{grav}}\} \quad (2.12)$$

where, \mathbf{N}_5 is the number of messenger particles mediating SUSY breaking. Different choices for the values of these parameters will represent different scenarios for GMSB models.

A unique feature in GMSB models is that the mass of the Lightest SUSY Particle (LSP) can be made (through C_{grav}) to be as light as possible in order for a model prediction of the mass of the dark matter candidate particle to agree with bounds from cosmology. In GMSB models, the LSP is most certainly the gravitino (\tilde{G}) as long as, \mathbf{M}_{mess} is very small compared to M_{Pl} , the energy scale where gravitational interactions become relevant. The gravitino's mass, $m_{\tilde{G}}$, can be expressed in terms of the parameter C_{grav} and M_{Pl} as

$$m_{\tilde{G}} = C_{\text{grav}} \cdot \frac{\mathbf{A} \mathbf{M}_{\text{mess}}}{\sqrt{3} M_{pl}} \quad (2.13)$$

where, $M_{Pl} = 1.3 \times 10^{19}$ GeV/c². By varying the scaling parameter, C_{grav} , the mass of the gravitino will also vary and the decay rate of the Next-to-Lightest-Supersymmetric Particle (NLSP) to the gravitino will also vary. Varying the decay rate will vary the lifetime of the NLSP from decaying immediately or promptly to being long-lived. The decay rate of the NLSP to the gravitino depends on the mass difference between the NLSP and the gravitino and also on the coupling or strength of its interaction with the gravitino, [14, 16, 17].

2.3 Long-Lived Neutral Particles in GMSB Models

In GMSB, *gravitino-scalar-chiral fermion* ((ψ, \tilde{G}, ϕ)) and *gravitino-gaugino-gauge boson* ((λ, \tilde{G}, A)) interactions, between the gravitino and its partners in the same supermultiplet, shown in Figure 2.5, are present. These interactions allow for any supersymmetric particle, \tilde{P} , to decay to its SM partner, P , and the gravitino (\tilde{G}) with the decay rate suppressed due to the nature of the interaction coupling strength which goes like $\frac{1}{\sqrt{F}}$.

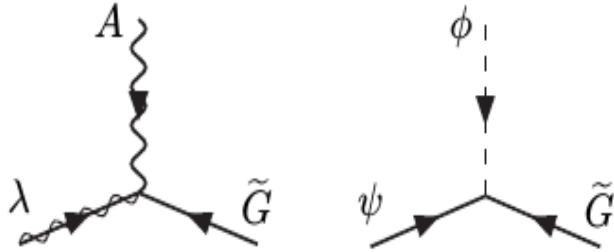


Figure 2.5: Feynman diagrams of Gravitino-Gaugino ((λ, A)) (left) and Gravitino-Scalar ((ψ, ϕ)) (right) couplings. The coupling strength goes like $\frac{1}{\sqrt{\mathbf{F}}}$

The decay rate for $\tilde{P} \rightarrow P + \tilde{G}$ can be computed as

$$\Gamma(\tilde{P} \rightarrow P + \tilde{G}) = \frac{m_{\tilde{P}}^5}{16\pi \mathbf{F}^2} \left(1 - \frac{m_P^2}{m_{\tilde{P}}^2}\right)^2 \quad (2.14)$$

where, the terms in the the decay rate are due to the nature of the interaction and also from the kinematic phase space integral where the mass of the gravitino can be neglected in the computation.

The decay width is larger for smaller values of fundamental SUSY breaking parameter, \mathbf{F} or equivalently for smaller gravitino masses, $m_{\tilde{G}}$. If we re-write the fundamental SUSY breaking scale as $\mathbf{F} = C_{grav} \cdot m_{\tilde{G}} \sqrt{3} M_{pl}$, we see that for different values of C_{grav} , which signifies changing the mass of the gravitino or the fundamental SUSY breaking scale, the decay width can also be vary. Thus, by varying the decay width, we can vary the lifetime of the supersymmetric particle \tilde{P} , for fixed mass, from being prompt to being long-lived.

In a scenario where the, SUSY particle \tilde{P} is not the Next-To-Lightest Supersymmetric Particle (NLSP), the decay, $\tilde{P} \rightarrow P + \tilde{G}$, suffers from competition with other probably favorable decay channels and this process is not competitive enough to happen at a collider experiment. However, if \tilde{P} is the NLSP, then there is no competition and this decay is always sure to happen at a collider.

2.3.1 Benchmark Scenario

In this thesis, we focus on a benchmark working model of the R-parity conserving GMSB models called the *Snowmass Point and Slopes* (SPS8),[18]. In the SPS8 model, the following choice of GMSB parameters are made:

$$M_{\text{mess}} = 2\Lambda \quad \tan(\beta) = 15 \quad N_5 = 1 \quad (2.15)$$

where Λ and C_{grav} are the free to vary.

The LSP is the gravitino (\tilde{G}) and the lightest neutralino ($\tilde{\chi}_1^0$) is the NLSP. Since the $\tilde{\chi}_1^0$ is a mixture of mixture of the supersymmetric particles Bino (\tilde{B}^0), Wino (\tilde{W}^0), higgsino ($\tilde{H}_u^0, \tilde{H}_d^0$), the SM partner it decays to together with the gravitino will depend on the choice of parameters which affect the mixing. Since the parameters: Λ , $\tan \beta$, and $\text{sgn}(\mu)$ do affect the mixing, they have been chosen in the SPS8 model to maximize the branching ratio involving the decay of the $\tilde{\chi}_1^0$ to a photon (γ) and a gravitino over the other possible SM partners: Z boson, Z' and a higgs boson (h).

Thus, this thesis is about the search for decay of the lightest neutralino which is our long-lived neutral particle to a photon and the gravitino. By varying Λ and C_{grav} , we study the production of $\tilde{\chi}_1^0$ for different masses and lifetimes to photon and gravitino.

Particle Mass Spectra

Figure 2.6 presents the mass spectrum for supersymmetric particle as predicted by the SPS8 model. The mass difference is determined by the supersymmetry breaking energy scale represented by Λ . As shown in Figure 2.6, the mass of the SUSY particles, including the lightest neutralino, $\tilde{\chi}_1^0$, is larger for the case where the effective SUSY breaking scale, $\Lambda = 180 \text{ TeV}$ (right plot) compared to the case where $\Lambda = 100 \text{ TeV}$ (left plot). This means that, SUSY sparticle's mass gets heavier with increase value of $\Lambda = 100$.

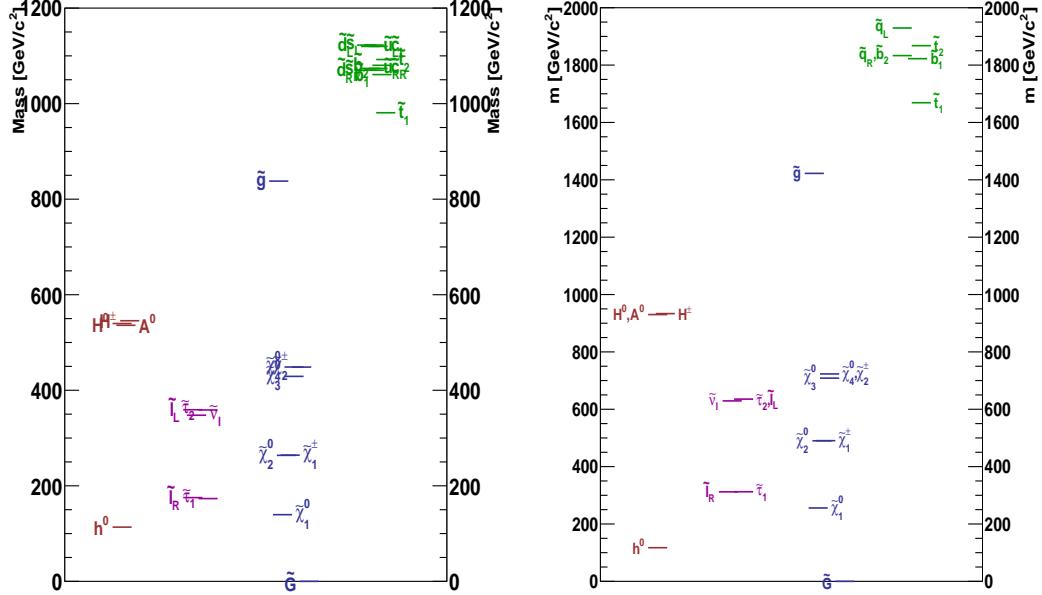


Figure 2.6: SUSY mass spectra for benchmark SPS8 model: $\Lambda = 100 \text{ TeV}$, $C_{grav} = 200$ (left) and $\Lambda = 180 \text{ TeV}$, $C_{grav} = 200$ (right). SUSY mass spectra including the $\tilde{\chi}_1^0$ increased for increase value of Λ

Lightest Neutralino as Long-Lived Neutral Particle

The neutralino being the next-to-lightest-supersymmetry particle (NLSP), can decay into a photon (γ), Higgssino (\tilde{H}^0), Z boson and gravitino (\tilde{G}) [17, 16].

We will concentrate on the production of neutralinos from processes like $pp \rightarrow \tilde{g}\tilde{g}$, $\tilde{q}\tilde{q}$, as these processes have a higher production cross section at the LHC. We mentioned earlier that a probable manner in which neutralinos can be produced is from the production and subsequent decay of higher massive supersymmetric particles. Some of these higher massive supersymmetric particles include squarks (\tilde{q}), excited squarks (\tilde{q}^*) and gluinos (\tilde{g}). In this scenario, the neutralino is produced *indirectly* or as we say from the *cascade decay* of higher massive supersymmetric particles. The Feynmann diagram for these production processes, $pp \rightarrow \tilde{g}\tilde{g}$, $\tilde{q}\tilde{q}^*$, are given in figure 2.7. Squarks and gluinos do not directly decay into gravitinos but through neutralinos and eventually to gravitinos because their coupling to the gravitinos is not possible. The reason for

this is that, in GMSB models, there are no gravitino-gluino-gauge boson or gravitino-squark-gauge boson couplings but rather gravitino-gaugino-gauge boson or gravitino-scalar-chiral fermion couplings.

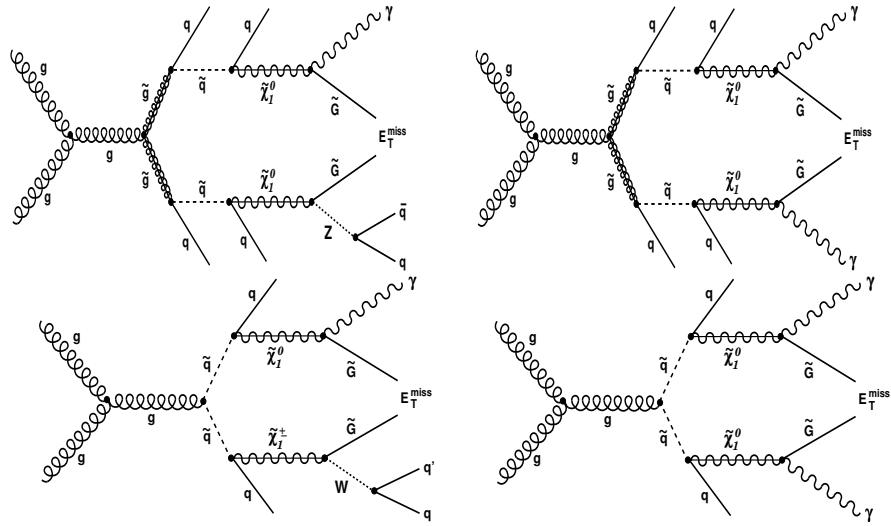


Figure 2.7: Feynmann diagrams for neutralino production from the cascade decay of a produced gluino (top) and squark (bottom). The final event has a single (left diagrams) or double photons (right diagrams) neutralino decay at LHC.

The probability for a neutralino ($\tilde{\chi}_1^0$), produced with energy $E_{\tilde{\chi}_1^0}$ and mass $m_{\tilde{\chi}_1^0}$ to travel a distance x before decaying to a photon and gravitino in the laboratory frame can be expressed as $\mathcal{P}(x) = 1 - \exp\left(-\frac{x}{L}\right)$, where the distance traveled in a particle detector by the neutralino is given by Equation 2.16.

$$L = \left(c\tau_{\tilde{\chi}_1^0} \right) \cdot (\beta\gamma)_{\tilde{\chi}_1^0} \quad (2.16)$$

From Equation 2.16, it is clear that this distance depends on two main factors. The boost factor, $(\beta\gamma)_{\tilde{\chi}_1^0} = \frac{|\vec{p}_{\tilde{\chi}_1^0}|}{m_{\tilde{\chi}_1^0}} = \sqrt{\left(\frac{E_{\tilde{\chi}_1^0}}{m_{\tilde{\chi}_1^0}}\right)^2 - 1}$, which indicates how fast the neutralino is traveling before it decays. For slow moving neutralino, $(\beta\gamma)_{\tilde{\chi}_1^0} \ll 1$.

This means that the momentum ($p_{\tilde{\chi}_1^0}$) of the neutralino during production from gluino or squarks decays must be much smaller than the its mass, $m_{\tilde{\chi}_1^0}$. Neutralinos, produced, with their mass and momentum satisfying the slow condition, $p_{\tilde{\chi}_1^0}/m_{\tilde{\chi}_1^0} \ll 1$,

are definitely good candidates for detectable long-lived neutralinos. The other factor is the inherent long lifetime, $c\tau_{\tilde{\chi}_1^0}$, of the neutralino. Neutralinos with $c\tau_{\tilde{\chi}_1^0} > 1$ are long-lived and would make good candidates for detectable long-lived neutralinos. This inherent neutralino lifetime can be expressed as given in equation 2.17.

$$c\tau_{\tilde{\chi}_1^0} \approx \left(\frac{m_{\tilde{\chi}_1^0}}{\text{GeV}} \right)^{-5} \left(\frac{\sqrt{\mathbf{F}}}{\text{TeV}} \right)^4 \quad (2.17)$$

$$c\tau_{\tilde{\chi}_1^0} \approx C_{grav}^2 \left(\frac{m_{\tilde{\chi}_1^0}}{\text{GeV}} \right)^{-5} \left(\frac{\sqrt{\mathbf{\Lambda} \cdot \mathbf{M}_{mess}}}{\text{TeV}} \right)^4 \quad (2.18)$$

It is important to note that by changing the supersymmetry breaking scale, \mathbf{F} , the lifetime of the neutralino also changes. In the SPS8 model, the parameter C_{grav} is used to adjust the inherent lifetime of the neutralino. Thus, we re-write the neutralino lifetime as given in equation 2.18. This equation is used to simulate physics events with the production and decay of neutralino in CMS detector using Monte Carlo (MC) simulations. The supersymmetry breaking scale, $\mathbf{\Lambda}$, determines the mass of gluino (\tilde{g}), squarks (\tilde{q}) which decay to the neutralino. Therefore the neutralino momentum, $p_{\tilde{\chi}_1^0}$ is determined by the masses of gluino and squarks. If the gluino or squark decays to the neutralino in association with a many gluons and quarks seen in the detector as *jets*, then the neutralino momentum is small with the ratio $p_{\tilde{\chi}_1^0}/m_{\tilde{\chi}_1^0} \ll 1$, this means the neutralino is slow and therefore long-lived. However, if the gluino or squark is decays with less number of jets, then the neutralino momentum is not so small and the neutralino is not very long-lived. It is worth noting that kinematic properties of the neutralino like momentum, arrival time at the CMS ECAL and the number of associated jets can be influenced by the gluino or squark decay properties. The gluino or squark is produced during proton-proton (pp) collisions at the LHC and the data recorded by the CMS detector is analyzed to search for events with neutralino decay.

The produced $\tilde{\chi}_1^0$ can decay to γ, Z bosons, *Higgs*, e^+e^- and the \tilde{G} . However, its decay to a γ and \tilde{G} is the dominant mode of decay. With the *BR* for single photon decay of the $\tilde{\chi}_1^0$ to a $\gamma + \tilde{G}$ about 83 to 94%. We observed that 97 to 99% of all our simulated signal events contain at least a single photon.

2.4 Previous Search Experiments

The have been previously other search experiments for neutral long-live particles decaying to photons. Obviously, negative findings from these experiments led to possible search exclusion regions in terms of the lifetime, mass and cross section of the existence of supersymmetry particles in different supersymmetry models. Results from experiments(DO, CDF, CMS and ATLAS) [19, 20, 21, 22, 23] of the search for Neutralino NLSP decaying to photon and gravitino interpreted using the SPS8 benchmark model is shown in 2.8. These results show that within the SPS8 model, neutralinos with mass $m_{\tilde{\chi}_1^0} \leq 245$ GeV and proper decay length $c\tau_{\tilde{\chi}_1^0} \leq 6000$ mm have not been found at hadron colliders. The diagram on the left of figure 2.8 are exclusion results in the neutralino mass or supersymmetry breaking scale Λ on the horizontal axis and the nuetralino lifetime, $c\tau_{\tilde{\chi}_1^0}$ on the vertical exis from the 7 TeV search analysis by the ATLAS experiment while the diagram on the right is that for CMS experiment. The colored regions on the plots shows the parameter space where these searches have been performed and the findings were negative.

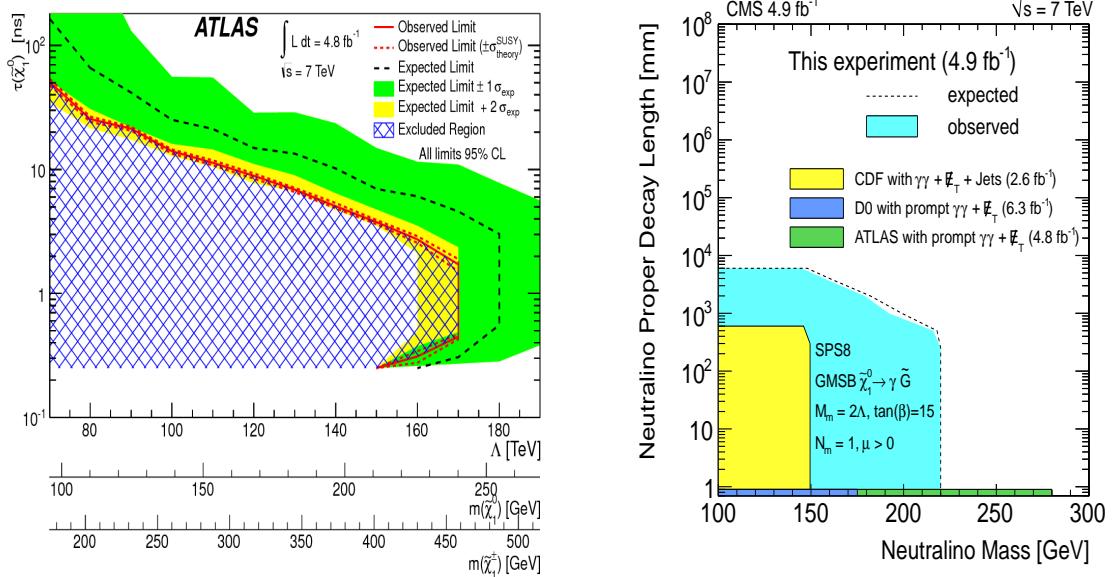


Figure 2.8: Neutralino lifetime and mass upper limit from ATLAS(left) and CMS(right) 7 TeV analysis with non-pointing photons and MET.

Chapter 3

Hadron Collider and Detector

This section describes the particle collider and detector. The first section describes the particle accelerator which is the Large Hadron Collider (LHC) and the next section describes the Compact Muon Solenoid (CMS) detector with emphasis to those sections directly relevant to this analysis. A detailed description of the LHC and CMS detector is found in [24] and [25].

3.1 Large Hadron Collider

3.1.1 Overview

The LHC is a proton-proton and heavy ion collider designed to achieve a center of mass \sqrt{S} energy of 14 TeV. It is hosted by the European Organisation for Nuclear Research (CERN). Unlike linear colliders, the LHC is a circular collider with nearly 27 km in circumference located at the border between France and Switzerland. It is designed to smash protons and ions against each other controlled by powerful magnets at four main points. The Compact Muon Solenoid (CMS) is one of the multi-purpose particle detectors at each collision point. Fig. 3.1 shows the LHC and the different stages before particle collision.

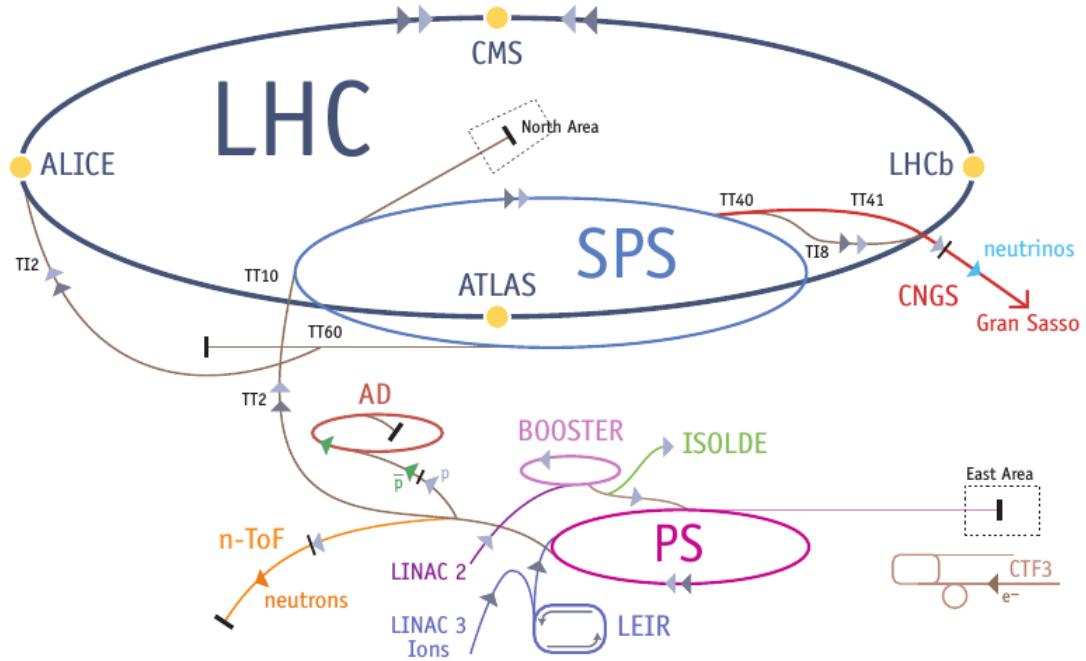


Figure 3.1: Schematic diagram showing the full Large hadron Collider.

3.1.2 Colliding Energy

Hydrogen ions also known as protons from hydrogen gas where the orbiting electron has been stripped away is inserted into a linear accelerator called *Linac 2*. Using electromagnetic fields in Radio Frequency (RF) cavities, these protons are accelerated to an energy of 50 MeV creating a stream of particles called *particle beams* arranged in packets known as *bunches*. Protons from the Linac2 are injected into the circular synchrotron Booster (PSB). The circular synchrotron accelerator ensures that the protons pass many times through a cavity with their energy slowly increasing each time to reach the design energy. The PSB accelerates the protons up to 1.4 GeV and inject them into the Proton Synchrotron (PS) which increases their energy to 25 GeV. These protons traveling at 99.93% the speed of light are sent to the Super Proton Synchrotron (SPS) and accelerated to an energy of 450 GeV. They are finally transferred into the LHC ring (accelerating in a clockwise and anti-clockwise direction) and accelerated for about 20 minutes to their nominal energy of 7 TeV. By now the protons are traveling with the

speed of 99.9999% the speed of light. Powerful bending magnets are used to keep the beams traveling in the circular LHC ring. The advantage of circular particle colliders over fix target is that, the energy available to make new particles called the *center of mass* (COM) energy, denoted as \sqrt{S} is simply the sum of the energy of the two beams i.e. $\sqrt{S} = E_{\text{beam1}} + E_{\text{beam2}}$ compared to $\sqrt{E_{\text{beam}}}$ for fix target experiments. For the LHC, each beam is designed to have energy of 7 TeV making $\sqrt{S} = 14$ TeV. In circular colliders, synchrotron radiation (is inversely proportional to the mass of the charge particle to the fourth power) by an accelerating charge particle contributing to loss its energy. This would require a continuous addition of energy after each turn to maintain the beam energy to a stable value. However, since the proton's mass is about 0.938 GeV, and are to be accelerated to about 7 TeV, this loss of energy is not very significant unlike for electrons whose mass is about 0.000511 GeV and the energy loss is more. Thus protons are preferred to electrons for a circular collider. Then again, the debris of particles produced when electrons collide is much less compared to that of protons making analysis in a hadron collider very challenging.

3.1.3 Luminosity

In colliding beams experiments, the center of mass energy available for the production of new effects is very important. However, the number of useful interactions producing effects (events) is equally important, especially in cases where the probability (also known as cross section, σ) of producing rare events is very small. The quantity which measures the ability of a particle accelerator to produce the events from the required number of interactions is called *luminosity*. The luminosity is also the proportionality factor between the number of events per second and the cross section. Luminosity (\mathcal{L}) is therefore a measure of the number of collisions that can be produced in a collider per squared area per second. The cross section is calculated from theory while the luminosity depends on factors ranging from the flux i.e. number of particles per second of the beams, the beam sizes at collision, and the frequency of collision. For physics experiments, the integrated luminosity which is total luminosity over a given period of time usually a year gives the amount of data that has been recorded by a given detector.

Using the luminosity (\mathcal{L}) and the cross section (σ_p) of a given process, we can calculate event rate (\mathcal{R}) or the number of events per second produced in proton collisions

by the given interaction process. By calculating the event rate, we are measuring a given cross section (σ_p) through $(\sigma_p = \frac{\mathcal{R}}{\mathcal{L}}$ in order to prove or disprove theories which make prediction on σ_p .

$$\mathcal{R} = \mathcal{L} \cdot \sigma_p \quad (3.1)$$

In CMS we have a “recorded” and “delivered” luminosity. Delivered luminosity refers to the luminosity delivered by LHC to CMS and one would expect this to be equal to the amount recorded. However, there are instances where the CMS detector is unable to take data either because the data acquisition chain (DAC) or one of the CMS sub-detectors is temporarily down and also trigger dead time. Figure 3.2 shows the total integrated luminosity delivered by LHC and recorded using the CMS detector during the 8 TeV proton-proton collision by the LHC.

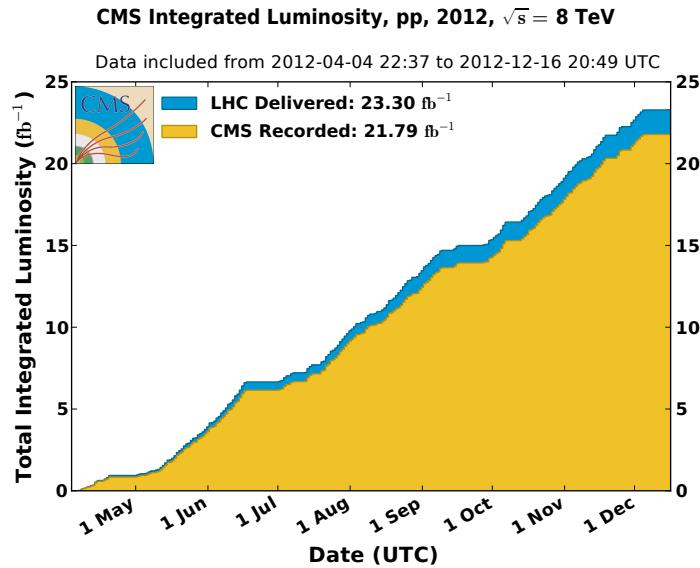


Figure 3.2: Cumulative luminosity versus day delivered to (blue), and recorded by CMS (orange) during stable beams and for p-p collisions at 8 TeV center-of-mass energy in 2012.

3.1.4 LHC Bunch Structure

Each LHC proton bunch is placed inside an RF buckets during beam filling. The filling scheme is such that not all RF buckets have proton bunches. Those empty buckets or beam gaps are necessary to avoid parasitic collision near IP and to make room for beam dump and beam halos known as beam cleaning. The separation in time between two

buckets/bunches filled or unfilled is approximately 2.5 ns. There are about 10^{11} protons per bunch during filling and acceleration. However, during filling and eventual bunch splitting at PS, it is possible that some empty buckets are filled with a much smaller proton population compared to the main bunch. These buckets with few protons can be either trailing the main bunch by $\Delta t = 2.5, 5.0, 7.5, \dots$ ns, or leading the main bunch by $\Delta t = -2.5, -5.0, -7.5, \dots$ ns. If these less populated bunches are 2.5 to 3.0 ns spaced in time from each other, they are referred to as *satellite* bunches and if 5.0 ns, they are referred to as *ghost* bunches. Figure 3.3 shows ghost and satellite bunches and a main proton bunch.

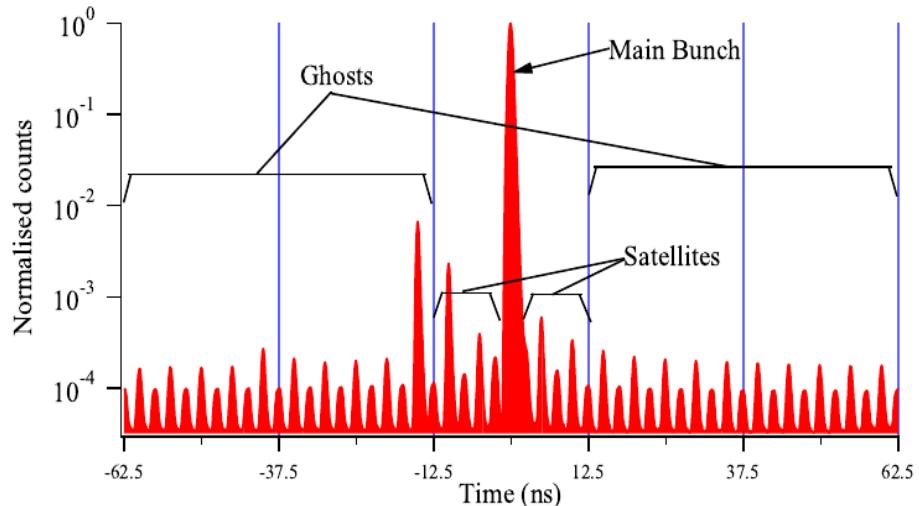


Figure 3.3: Longitudinal Profile taken with Longitudinal Density Monitor (LDM) detector showing definition of Ghost/Satellite bunches with respect to main bunches.

The presence of ghost/satellite bunches increases the uncertainty in LHC luminosity measurements and can also generate proton-proton interactions near but not at the collision region. Effects on ghost/satellite bunches on instantaneous luminosity measurements have been studied using CMS, ATLAS and ALICE detectors. Their results showing clear observation of physics events produced from ghost and satellite bunch collisions is shown in figure 3.4. CMS uses energy deposits in the endcap calorimeters to observe time space which is consistent with the expectation from ghost/satellite bunches while in ATLAS uses the Longitudinal Density Monitor (LDM) detector to study ghost/satellite bunches.

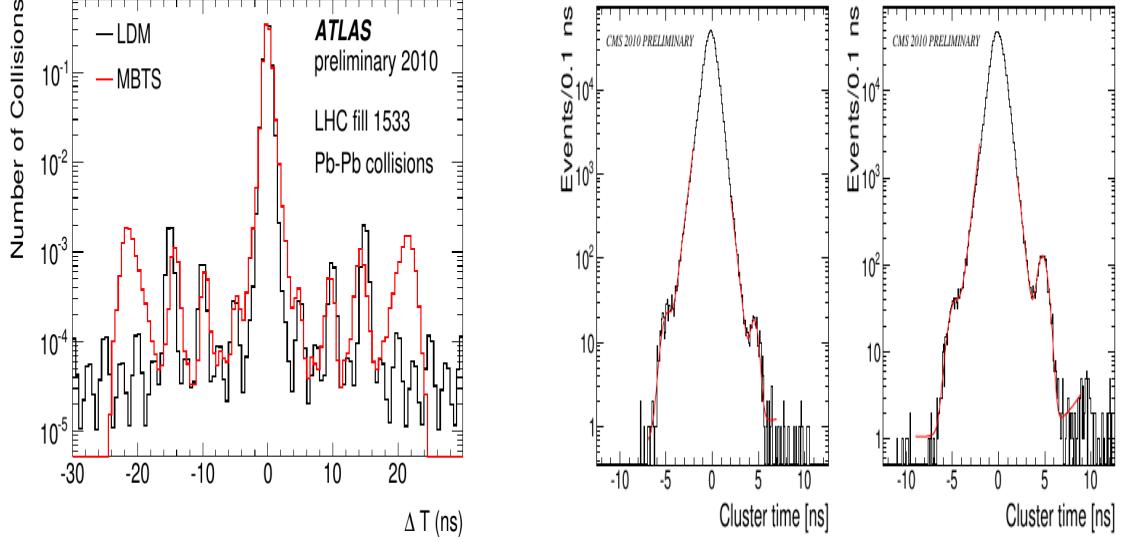


Figure 3.4: (left) Arrival time distribution(red) of ATLAS MBTS for LHC fill 1533 during 2010 Pb-Pb run and LDM profile(black) for Beam2(same for Beam1).
 (Right) Timing of Clusters in the CMS endcap calorimeters for fill 1089:Left: EEP detector(left side of IP $z > 0$) Right: EEM detector(right side of IP, $z < 0$). Plots from ATLAS [33] and CMS, [34]

Table 3.1 gives a summary of the LHC design conditions compared to the conditions used during the LHC RUN 1 operation.

LHC Operation Parameters 2010-2013

Parameter	2010 value	2011 Value	2012/13 Value	Design Value
Beam energy[TeV]	3.5	3.5	4.0	7
β^* in IP 5[m]	3.5	1.0	0.6	0.55
Bunch spacing [ns]	150	75/50	50	25
Number of bunches	368	1380	1380	2808
Protons/bunch	1.2×10^{11}	1.45×10^{11}	1.7×10^{11}	1.15×10^{11}
Normalised emittance[mm.rad]	≈ 2.0	≈ 2.4	≈ 2.5	3.75
Peak luminosity[$cm^{-2}s^{-1}$]	2.1×10^{32}	3.7×10^{33}	3.7×10^{33}	1×10^{34}
Evts/bunch crossing	4	17	37	19
Stored Beam energy(MJ)	≈ 28	≈ 110	≈ 140	≈ 362
Int. Luminosity by CMS[pb^{-1}]				-
Circumference[km]	26.659	26.659	26.659	26.659
Dipole Magnet B[T]	8.33	8.33	8.33	8.33

Table 3.1: LHC operation parameter conditions during RUN 1, 2010-2013

3.2 Compact Muon Solenoid

3.2.1 Overview

The Compact Muon Detector (CMS) is a modern particle detector design for many different particle detection capability. It is one of the general purpose detectors located at one of the proton-proton collision points along the 27 km LHC ring. Its main feature is the presence of a superconducting solenoid of 6 m internal diameter providing a field of 3.8 T for measuring a charge particle's momentum as the particle bends under the influence of this field traveling in the detector. This magnetic field encloses an entirely silicon pixel and strip tracker detector used for vertex finding and for detecting and reconstructing the tracks of charged particles, a lead-tungstate scintillating-crystals electromagnetic calorimeter (ECAL) and a brass-scintillating sampling hadron calorimeter (HCAL). Very long lived particles like muons are measured in gas-ionization detectors embedded in the flux-return iron-yoke located at the outermost section of the detector. It has a simple cylindrical structure consisting of barrel and endcap detectors and an extensive forward calorimetry and detectors to provide a near 4π solid angle coverage assuring good hermetic particle detection. The CMS apparatus has an overall length of 21.6 m, a diameter of 14.6 m, and weighs 12,500 tons. Figure 3.5 shows the CMS detector indicating the different sub-detectors and their material design type. We provide a performance summary and material type of each sub-detector in Table 3.2 of the CMS detector. The CMS uses a coordinate system with the origin coinciding with the center of the detector where proton-proton or nominal collision occurs. This point is commonly referred to as the *interaction point* (IP). The direction of x , y , and z -axes are as shown in figure 3.6. However, for particle identification, CMS uses a more convenient coordinate system based on the polar coordinates. In this polar coordinate system, the azimuthal angle, ϕ , is measured in the $x - y$ plane, where $\phi = 0$, is the x -axis and $\phi = \pi/2$, the y -axis. The radial distance in this plane is denoted R and the polar angle θ measured from the z -axis is related to *pseudo-rapidity*, η , through the relation; $\eta = -\ln \tan(\frac{\theta}{2})$. The coordinate system (η, ϕ) and its radial distance R identifies a point in the cylindrical volume of the CMS detector. In the coming sections, we describe the geometry, material characteristics and functionality of the CMS subdetectors used in our analysis.

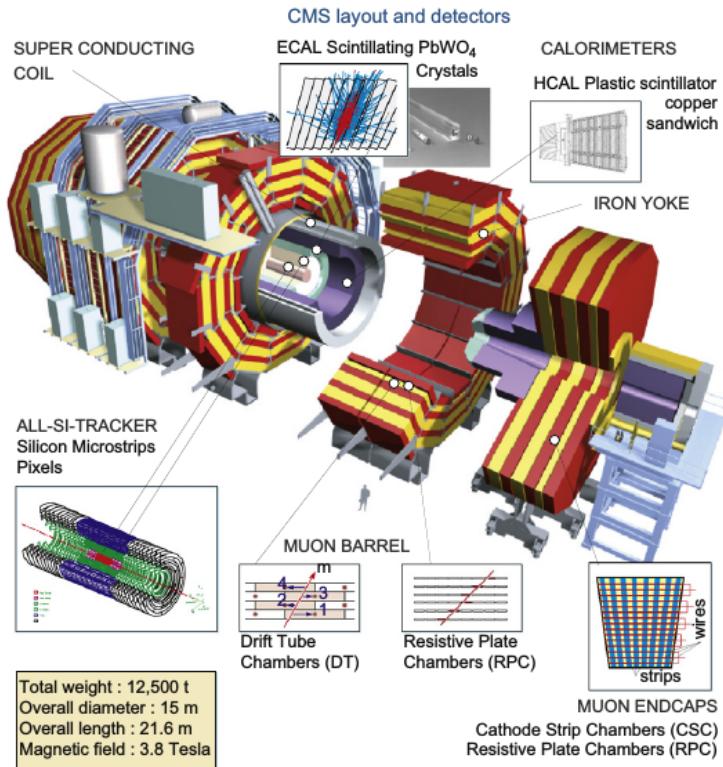


Figure 3.5: CMS Detector showing the different subdetectors and their material.

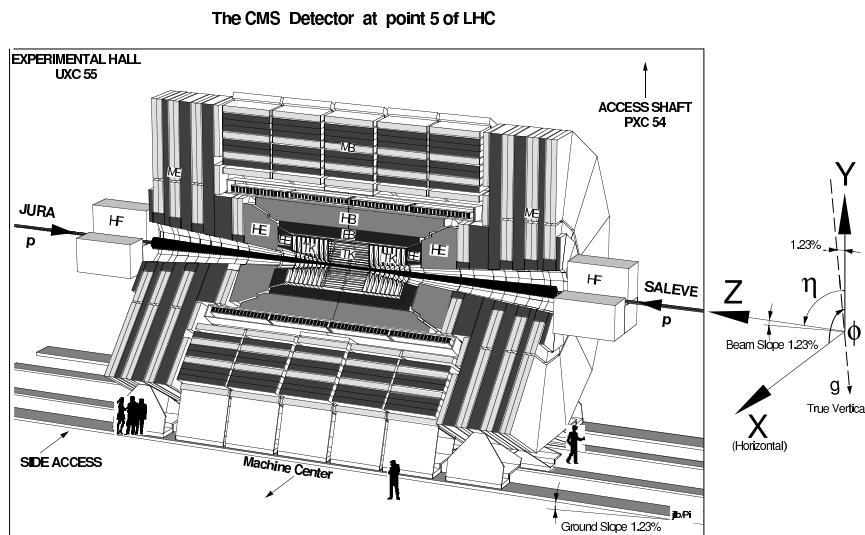


Figure 3.6: CMS detector schematic view with definition of $x - y - z$ coordinates.

CMS Detector and Resolution

Subdetector	Quantity	Resolution	Uses
Tracker	Momentum[GeV/c]	$\sigma_T/p_T \approx 1.5 \times 10^{-4} p_T + 0.005$	Silicon Pixels and Strips
ECAL	Energy[GeV]	$\sigma/E \approx 3\%/E + 0.003$	PbWO ₄ Crystals
ECAL	Time[ns]	$\sigma(\Delta t) = \frac{N}{A_{eff}/\sigma_n} \oplus \sqrt{2}\bar{C}$	PbWO ₄ Crystals
HCAL	Energy[GeV]	$\sigma/E \approx 100\%/E + 0.05$	Brass + Scintillator
Muon Chambers	Momentum[GeV/c]	$\sigma_T/p_T \approx 1\% \quad 50 \text{ GeV to } 10\% \quad 1 \text{ TeV}$	inner tracker + Muon Systems
Magnetic field	B-field strength[T]	3.8 T + 2 T	Solenoid + Return Yoke
Triggers	On/Off-line	Levels	L1(On-line) +HLT(Off-line)(L2+L3)

Table 3.2: CMS detector material, [25], and resolution(Time resolution: $N \approx 35 \text{ ns}$, $\bar{C} \approx 0.070 \text{ ns}$ [42])

3.2.2 Calorimeter

A CMS calorimeter absorbs a good fraction of energy of an incident particle and produces a signal with an amplitude proportional to the energy absorbed. This absorption is through the cascade production of secondary particles with energy of the incident particle directly proportional to the number of secondary particles produced. There are two types of calorimeters choices used in the CMS detector; the *Electromagnetic calorimeter* (ECAL); for absorbing the energy of electromagnetic particles such as photons and electrons and a *Hadronic calorimeter* (HCAL) made of more than one type of material for stopping and absorbing the energy of hadrons such as kaons and pions through strong interactions. The combined calorimeter detectors of CMS covers a region in $|\eta| < 5$ making it nearly hermetic for good missing energy measurements. The ECAL

and HCAL are arranged in a nested fashion shown in figure 3.7 so that electromagnetic particles can be distinguished from hadronic particles by comparing the depth of the particle shower penetration in both calorimeters.

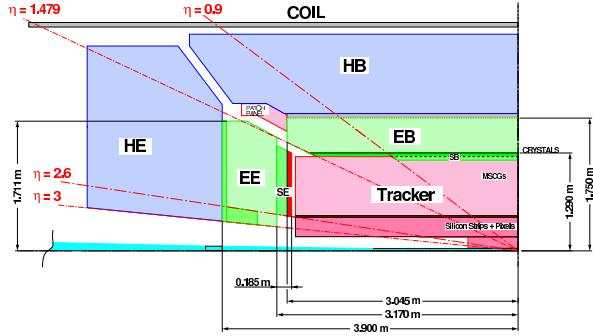


Figure 3.7: Schematic diagram of CMS calorimetry system with HCAL enclosing ECAL in the Barrel and Endcap regions.

Electromagnetic Calorimeter

The electromagnetic calorimeter (ECAL) detects photons and electrons. High energy photons and electrons are detected through their interaction with the lead tungstate (PbWO_4) crystals. During this interaction which happens either through electromagnetic showering or electron-positron pair production (*bremsstrahlung*), the incoming photon or electron deposit practically almost all of its energy. There are 75848 crystals in total mounted in a cylindrical geometry, with a barrel (EB) and an endcap (EE) structure. The choice of PbWO_4 crystals as calorimetry material by CMS for operation in the LHC environment is because of its a high density (8.28 g/cm^3), short radiation length ($X_0=0.89 \text{ cm}$) and a small Molière radius (22 cm). In a high radiation dose and fast timing (25 ns proton bunch spacing) environment like the LHC, PbWO_4 crystals is preferred to other crystal materials for its high radiation resistance and a scintillation decay time which is comparable to the LHC bunch crossing interval of 25 ns and about 80% of the light is emitted in 25 ns. The probability of an electromagnetic object with high energy to interact either through *Bremsstrahlung* or *pair production* with the material in ECAL is proportional to the nuclear charge, Z , of the material. PbWO_4 is a high Z material and this makes it once more the preferred material choice for electromagnetic calorimetry by CMS. The small Molière radius ensures that on average about

95 % of the electromagnetic shower energy is contained within a crystal crystal volume of about 9 crystals. This reduces the transverse spread of the electromagnetic cascade from multiple scattering of electrons and helps improve on the estimation of the transverse position of impact of an incident particle. It also provides a fine granularity for measuring the particle's energy by providing fewer overlap of particle signals. Its dense nature also allows for the electromagnetic shower to develop early and therefore likely to be fully contained within a compact device like CMS.

The EB section of the ECAL covers a pseudo-rapidity of $|\eta| < 1.479$. It has 61,200 crystals providing a granularity of 360 degree fold in ϕ and (2×85) -fold in η . The crystals are mounted in a quasi-projective geometry so that their axes make an angle of 3% with respect to a line vector from the nominal interaction vertex in η and ϕ directions. This avoids cracks aligned with a particle's trajectory. A crystal in EB is approximately 0.0174×0.0174 in $\eta - \phi$ or $22 \times 22 \text{ mm}^2$ at its front face and $26 \times 26 \text{ mm}^2$ at its rear face. Each crystal is 230 mm long corresponding to about $25.8 X_0$ radiation lengths. The crystal's radial distance measuring from the center of the face of the crystal to the beam line is 1.29 m. A number of crystals are placed in a thin-walled alveolar structure made with aluminum forming a *submodule*. Each submodule is arranged into 4 modules of different types according to their η position. There are about 400 to 500 crystals in each module and these 4 combined make one *supermodule* containing 1700 crystals. To reduce crystal reflective lost, the aluminum surface is coated to avoid oxidation leading to coloration. On the rear end of each EB crystal, two *Avalanche Photodiodes* (APD) is glued to collect the scintillating light from the crystals converting light into charge current which is further collected by the read-out electronics.

The endcap sector covers a pseudo-rapidity region of $1.479 < |\eta| < 3.0$ with a Preshower (ES) detector made of silicon strip sensors interleaved with lead placed immediately in front of it. The purpose of the preshower is to identify photons from the decay of neutral pion, $\pi^0 \rightarrow \gamma\gamma$ and also to help separate photons producing electrons through pair production from photons not producing electrons before their arrival at the EE. The endcap located on the $+z$ side of the nominal interaction is denoted EE+ while the other located on the $-z$ side is denoted as EE-. The longitudinal distance between the IP and the center of the surface of the EE crystals is 3.154 cm. Each endcap is divided into two halves called *Dees* with each Dee holding 3662 crystals. Crystals in EE with identical shape

are grouped into 5×5 units called *supercrystals* (SC). The crystals in the SC form an $x - y$ grid. Each crystal is 220 mm ($24.7 X_o$) in length and has a front face and rear cross section of 28.62×28.62 square mm and 30×30 square mm, respectively. Vacuum Phototriodes (VPT) instead of APDs is glued on the rear face of each crystal for scintillating light conversion into electrical signals. The VPT is used in the EE because of its high resistance to radiation and smooth operation in a strong magnetic field environment. These APDs and VPTs are used because of their high gain relative to regular photodiodes with no gain and the fact that they are not affected by the high magnetic field. Although the light yield for PbWO_4 crystals is rather low (≈ 70 photons/ MeV), these photo-detectors have internal gain (50 for APDs and 10 for VPTs) and quantum efficiency of 75 % for APDs and 20 % for VPTs of the emission wavelength. This makes it possible that signals from incident particles with energies of a few to high GeV longer than noise.

The signals from the APDs and VPTs are digitized by voltage-sensitive analogue-to-digital converters and through fibre-optic links transported as light signals to the counting room located adjacent to the experimental cavern.

The energy resolution and geometry structure of the ECAL ensures that the photon or electron's arrival energy, time, position and even the direction through the shape of its electromagnetic shower in the crystals can be identified and measured with good precision.

Hadronic Calorimeter

The CMS Hadron Calorimeter (HCAL) is comprised of four distinct subdetectors: the Barrel (HB), the Endcap (HE), the Outer Barrel (HO), and the Forward (HF). Unlike the ECAL, the HB, HE and HO subdetectors are scintillator-sampling calorimeters with embedded wavelength shifting fibers (WLS). HB, HE and HO uses brass plates as the inactive material and plastic scintillator with WLS as the active material. The brass plate is used for absorbing the hadronic shower which comprise of an *electromagnetic*(particles like π^0 s, η s and other mesons generated in the absorption process and decay to γ s which develop electromagnetic (em) showers) and *non-electromagnetic* components. The plastic scintilator is divided into 16 η sectors resulting in segmentation of $\Delta\eta \times \Delta\phi = 0.087 \times 0.087$. It was chosen for its long-term stability and moderate

radiation hardness. energy. The scintillating light through the WLS brings the light to hybrid photodiodes (HPDs) in the HB and HE. HPDs which have high electrical noise and will be replaced with silicon photon multipliers (SiPM) which have low noise during the current CMS detector upgrade. The HB and HE combine cover a region in pseudo-rapidity of $|\eta| < 3$. The HB covering the region $|\eta| < 1.3$, is divided into two-half barrel (HB+ and HB-) sections each composed of 18 identical 20° wedges in ϕ . Each wedge is made of flat brass alloy and steel(only front and back plates) absorber plate. HE covers $1.3 < \eta < 3.1$ and has plastic scintillation tiles with granularity of $\Delta\eta \times \Delta\phi = 0.087 \times 0.087$ for $|\eta| < 1.6$ and $\Delta\eta \times \Delta\phi = 0.17 \times 0.17$ for $|\eta| > 1.6$. The HO is an extension of HB outside the solenoid and thus utilizes the solenoid coil as an additional absorber. It is used to identify the starting shower and to measure the shower energy deposited after HB. The first active layer of the scintillating tiles is situated directly behind the ECAL in order to actively sample low energy showering particles from the support material between the ECAL and HCAL.

The HF occupies a pseudo-rapidity region of $3 < |\eta| < 5$. Its purpose is to provide a closer to 4π hermetic phase space coverage required for missing transverse energy calculation or MET. MET is the established signal for very weakly interacting particles like neutrino and supersymmetric particles like gravitino which travel through the detector undetected. HF consists of radiation hard quartz fibers embedded in steel absorbers running parallel to the beam axis. The signal from Cherenkov light emitted in the quartz fibers in response to charged particles makes it possible to detect all charge particles in the forward region. The HF calorimeter has long and short fibers for better sampling and to distinguish showers generated by electrons and photons from those generated by hadrons. The choice of quartz fibers is because of its high resistance to the high radiation in the forward detectors and its fast production of light through Cherenkov process.

For $|\eta| < 1.48$, the HCAL cells map on to 5×5 ECAL crystal arrays to form calorimeter towers projecting outwards from near the nominal interaction point. In each tower, the energy in ECAL and HCAL cells is summed to define the calorimeter energy tower. The energy ratio of an HCAL tower to an ECAL in a calorimeter energy tower can be used to improve photons and electron identification.

3.2.3 Muon Chambers

Muons unlike electrons and hadrons do not deposit most of their energy in the calorimeters. They are capable of traveling across the entire CMS detector into the muon chambers. Muons produce tracks which run across the CMS detector starting from the silicon pixel and strip subdetector closest to the IP called the *Tracker* and depositing very little fraction of their energy in the calorimeters unto the muon chambers. The muon chambers use the process of ionization and a 2 T magnetic field from the return iron yokes (bending the tracks of charge particles) to measure the momentum of charged particles. The three different types of muon chambers used by the CMS are: the drift tubes (DT) chambers in the barrel, cathode strip chambers (CSC) in the endcaps and resistive plate chambers (RPC) glued to the DT and CSC chambers. Four layers or stations of DT/RPC and CSC/RPC are embedded in an interleaved style with the iron yoke for track reconstruction and triggering. Figure 3.9 is a longitudinal view of the CMS detector showing the position of the muon stations. The DT and CSC record track segments characterized by the position of the track and the bending angle. This information is used to determine the precise transverse momentum and charge of particles during particle reconstruction. The RPCs(DTs and CSC will also be used after the current detector upgrade) are dedicated L1 trigger chambers used to determine the candidate muon's approximate transverse momentum and proton bunch crossing number. The RPC has a timing resolution of about 3 ns.

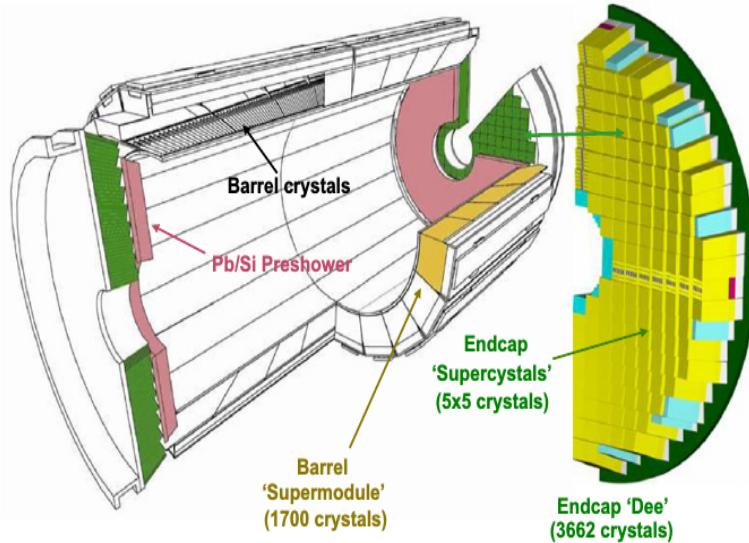


Figure 3.8: Layout of the CMS electromagnetic calorimeter showing the arrangement of crystal modules, supermodules in the barrel with the preshower in front of endcap with supercrystals.

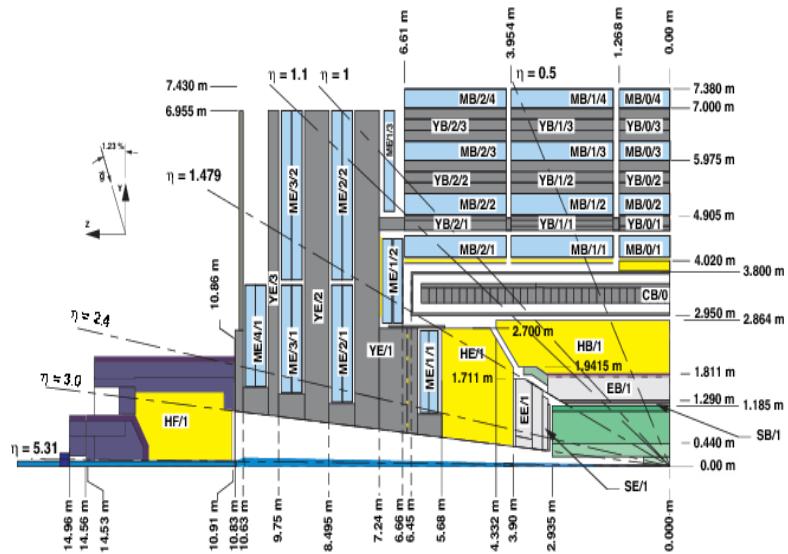


Figure 3.9: Cross section view showing the coverage range of CMS sub-detectors and their longitudinal distance from the IP.

3.2.4 Triggering

In CMS, there are a billion interactions including *pile up* (PU) happening each second. This means data from each 25 ns proton-proton collision has to be processed and stored before the next collision happens. Also, since not all these collisions produce interesting physics events, we have to be capable of selecting only interesting physics events produced from proton-proton collisions with sufficient energy. The process of selecting such interesting events is called *triggering*. CMS uses a two level triggering system for selecting interesting events produced with enough energy from collisions. The comprise of the *Level-1* (L1) and *High Level Triggers* (HLT) triggers.

The L1 triggers is a hardware designed electronics system implemented in FPGA and ASIC technology and uses information from the calorimeter, muon trigger and a global trigger board. The global trigger board makes the final decision based on the calorimeter and muon triggers to reject or keep an event for further processing at the HLT trigger. The L1 trigger is responsible for selecting the best 100,000 events/second from the initial 1 billions events/second produced.

The HLT is a software comprised of implemented selection algorithms running on a farm of more than 1000 standard computers. These complex algorithms include instructions like, match tracks to hits from the muon chambers, select energy deposits above a certain threshold in the calorimeters with no tracks for electromagnetic objects, and begins the first step of event selection. Just like the L1 trigger, the HLT uses assimilated and synchronized information from different parts of the CMS detector to create the entire event. By the time this selection process is complete, there are now only 100 events/second with the remaining 99,900 thrown away. Taking an average event size to be 1 Megabyte, in a stable and effective LHC proton collision period of a year or 10^7 seconds, CMS produces about a Petabyte of data which is stored and used later for offline physics analysis.

Chapter 4

Time Reconstruction and Resolution

ECAL Time Overview

The ECAL was designed to precisely measure the energy of electrons and photons produced in proton-proton collisions, with a target barrel resolution of 0.5% for photons with energies larger than 50 GeV. In addition to energy measurements, the combination of fast scintillation for PbWO₄ crystals, the electronic pulse shaping and sampling rate of 40 MHz allow for excellent time measurements to be made with ECAL. A pulse shape matching of less than 1% non-linearity by the pulse shaping chip in each channel, ensures a linear and uniform pulse height from among all the channels in the ECAL. The time measured by each channel is reconstructed from 10 discrete samples of the digitized analog pulse height of the electronics, using a time reconstruction algorithm.

4.1 Electromagnetic Calorimeter Readout Electronics

The scintillating light from the PbWO₄ crystals is readout and converted into electric current using Avalanche Photo-Diodes (APD) and Vacuum Photo-Triodes (VPT) in the barrel and endcap detector regions, respectively.

The ECAL electronics readout and control system comprise of the Front End (FE) electronics (mainly consisting of radiation hard ASIC chips) installed inside the detector

volume at the back of ECAL and the Off Detector (OD) electronics in the upper counting room. the FE electronics is connected to the OD electronics by 100 m radiation hard Gigabit Optical fiber Links (GOL) for transporting the optical data stream.

The FE has a crystal arrangement in matrixes of 5×5 crystals called *Trigger Towers* in barrel, which is based on the electronics arrangement and a 5×5 mechanical arrangement of crystals called *Super-Crystals* in the endcaps. Each FE board hosts five Very Front End (VFE) boards with each holding five analog channels shown in right schematic diagram of Figure 4.1. Each channel, shown in the left schematic diagram of Figure 4.1, consists of a Multi-Gain Pre-Amplifier (MGPA) and a 12 bit Analog-to-Digital Converted (ADC) used to amplify, shape and digitize the signal coming from the Photo-Diodes.

The MGPA chip uses 3 gain ranges with gain ratios of 1,6 and 12 to span the overall dynamic range of the signal. Equipped with a CR-RC filter with a pulse shaping time of 40 ns and less than 1% of non-linearity, the MGPA ensures a linearity and pulse shape matching across all three ranges which allows for precise pulse shape reconstruction. The readout and precision performance requires a multi range 12-bit ADC chip with a sampling frequency of 40 MHz to digitize the analog pulse signal of the highest unsaturated range into 10 discrete samples, with an electronic noise of about 40 MeV.

The FE board also is also equipped with a Clock and Control Unit (CCU) chip which enables the exchange of control information between the OD and the FE electronics through a Clock and Control link.

The digitized data are stored in pipeline buffers until a Trigger Level 1 (LV1) accept decision is made. At the same time the digitized data are sent to a digital signal processing FERNIX chip which performs Trigger Primitive Generation (TPG) for the LV1 trigger system. The TPs are sent to the Trigger Concentration Card (TCC) on the OD electronics at a rate of 1 word/25 ns. Upon a positive trigger LV1 decision, the data samples corresponding to a single event are transferred to primary event buffers which are eventually merged and sent to the Data Concentration Card (DCC) through Readout links at 800 Mbps. The crystal readout data from the FE board is collected by the OD electronics DCC.

The OD board hosts the DCC, TCC and a Clock and Control System card (CCS). The TCC completes TPG process by combining the TPs generated and transmit them

to the regional Calorimeter Trigger system while the CCS is primarily task with distributing fast timing signals (LHC 40 MHz clock and trigger commands) to the FE and OD electronics and to synchronize their operation. A common clock distribution serving 5×5 readout unit of the FE helps minimize channel-to-channel synchronization.

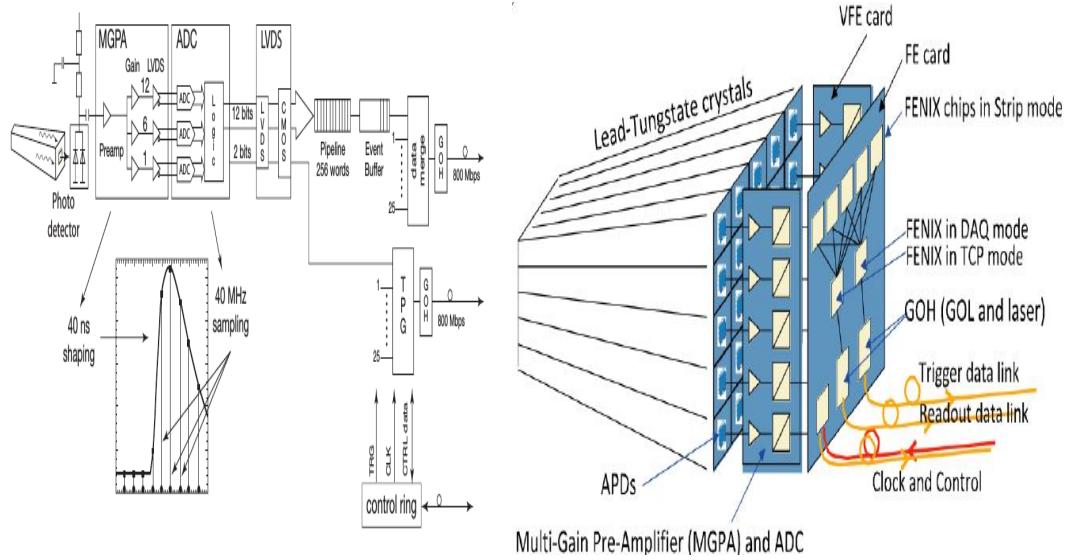


Figure 4.1: Schematic diagram of the CMS ECAL electronics readout for a single channel(Left) and the entire Trigger Tower or Front End (FE) Card(Right).

4.2 Time Reconstruction

An analog pulse shape from a single channel is shown in Figure 4.2(a). Overlaying the pulse shape are typical 10 digitized samples in red. The first three samples are taken in the absence of a signal and correspond to the pedestal. The ADC chip responsible for the digitization, has a sampling frequency of 40 MHz, i.e. one sample is made every 25 ns, which is the same rate as the LHC proton-proton bunch collision frequency of one bunch crossing every 25 ns. In addition, the timing phase between each sample and the next is adjusted so that the maximum of the signal pulse shape corresponds to one of the samplings within 1 ns. A time reconstruction algorithm uses these 10 digitized samples to measure the time of a single channel by finding the precise time corresponding to the maximum of the pulse shape, T_{\max} .

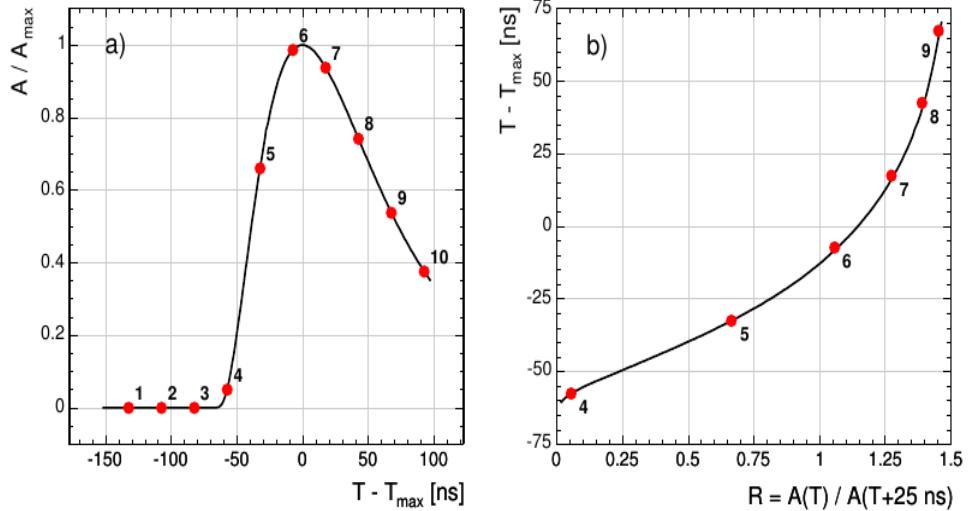


Figure 4.2: (a) A measured ECAL pulse shape for a single channel. (b) $\mathbf{T} - \mathbf{T}_{\max}$ Vs $R(T)$ relationship whose inverse is $T(R)$. Solid line is pulse shape from test beam while dots are typical 10 discrete samples corresponding to signal from proton-proton collision.

Precise measurement of \mathbf{T}_{\max} relies on good knowledge of the pulse shape. To find \mathbf{T}_{\max} using the pulse shape, a ratio algorithm uses the ratio

$$R_i = A(T_i)/A(T_{i+1}), \quad (4.1)$$

where T_i is the time of the time of the beam crossing that the i^{th} pulse height sample is taken. Each pulse height ratio is uniquely determined from the pulse shape as shown in Figure 4.2(b) which shows the the pulse height ratio, R_i , as a function of the time that the pulse height of the first of a pair of samples is taken relative to the time of the maximum height of the pulse, $\mathbf{T}_i - \mathbf{T}_{\max}$. Using the inverse of this function, $T(R)$, we can obtained $\mathbf{T}_{\max} = T_i - T(R_i)$ and the uncertainty, σ_i , for each R_i point.

The uncertainty on each measurement is a product of the derivative of the function $T(R)$ and the uncertainty on the value of R_i which itself depends on three separate uncertainties: the noise fluctuation (σ_n) of each sample, the uncertainty in the estimation of the pedestal value which is always subtracted from the measured value and the truncation during 12-bit digitization. These uncertainties are uncorrelated and can be added in quadrature,[38].

A more precise value of T_{max} is obtained by using the ratios R_4 through R_7 . The other ratios are not used for the estimation of T_{max} because their associated uncertainties are large due to the large slope of the $T(R)$ function. A weighted average of each T_{max} and σ obtained from ratios R_4 through R_7 as expressed in Equation 4.2 gives the best estimate of T_{max} and uncertainty.

$$T_{max} = \frac{\sum \frac{T_{max,i}}{\sigma_i^2}}{\sum \frac{1}{\sigma_i^2}} , \quad \frac{1}{\sigma_T^2} = \sum \frac{1}{\sigma_i^2} \quad (4.2)$$

The sum is from $i = 4, \dots, 7$.

4.3 ECAL Time Performance from Test Beam

4.3.1 ECAL Time Resolution

The intrinsic time resolution of ECAL, $\sigma(t)$, measured during test beam, consist of three main contributions which can be summed in quadrature since they are uncorrelated. These three contributions are the noise, stochastic and constant terms. The *Noise* (N) term arises from the electronic noise, coherent movement of the baseline and effects overlapping hits some of which come from other soft proton-proton collisions called *pile up* (PU) events. The *Stochastic* term (S) is from fluctuations in the number of photons collected during the sample times. Lastly, the *Constant* term (C), whose contribution is independent of the energy deposited and arises from variations in the point of shower initiation within the crystal and variations in the pulse shape for each channel. The full expression for time resolution with all three contributions is given in Equation 4.3,

$$\sigma^2(t) = \left(\frac{N}{A/\sigma_n} \right)^2 + \left(\frac{S}{\sqrt{A}} \right)^2 + C^2, \quad (4.3)$$

where A is the measured amplitude corresponding to the energy deposited and σ_n is the intrinsic noise in the amplitude for individual channel. σ_n has a value of 42 MeV and 140 MeV in the barrel and endcap, respectively. $N = 33$ ns has been estimated from Monte Carlo (MC) simulation studies. Contribution from the stochastic term, (S) is small, with a value of $S < 7.9$ ns·MeV^{1/2}.

To measure the intrinsic time resolution of ECAL and study the pulse shape, H2 and

H4 test beam facilities at CERN were used. Prior to installation at the CMS detector, about 25% of the barrel and endcap crystals were exposed to electron beams with energy between 15 GeV and 250 GeV. The time resolution was obtained from the Gaussian distribution of the difference in time of two crystals sharing energy and belonging to the same electromagnetic shower. Note that the contribution from crystal-to-crystal synchronization does not contribute to the constant term in this study. Neglecting the stochastic term because its contribution is negligible, Equation 4.3 reduces to:

$$\sigma^2(t_1 - t_2) = \left(\frac{N}{A_{eff}/\sigma_n} \right)^2 + 2\bar{C}^2 \quad (4.4)$$

where $A_{eff} = A_1 A_2 / \sqrt{A_1^2 + A_2^2}$, while $t_{1,2}$ and $A_{1,2}$ are the times and amplitudes of the two crystals. \bar{C} is their residual constant term contribution. The time resolution is measured from the standard deviation of a Gaussian fit to the time distribution from each slice of A_{eff}/σ_n of the A_{eff}/σ_n distribution. The resulting distribution of $\sigma(t_1 - t_2)$ of these standard deviations plotted against A_{eff}/σ_n is used to extract the noise and residual constant terms. The result presented in Figure 4.3 gives a noise factor $N = (35.1 \pm 0.2)$ ns and $\bar{C} = (20 \pm 4)$ ns obtained from test beam experiment.

The value of the N in this study agrees with our Monte Carlo estimate to within 6%. Other studies of the timing resolutions using the time difference between two basic clusters of separate electromagnetic shower was used and the results gave similar values of N .

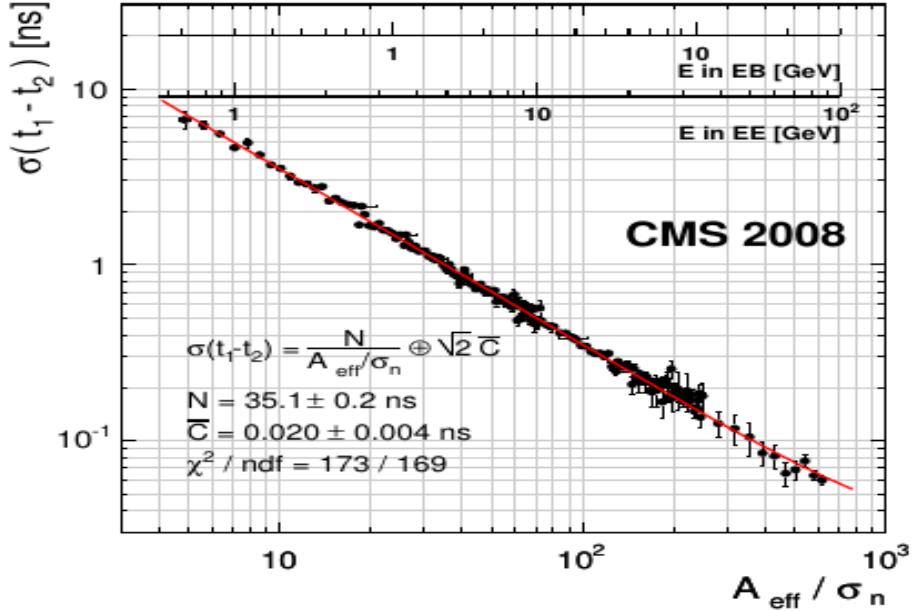


Figure 4.3: Difference in the time measurements as a function of A_{eff}/σ_n of two crystals sharing an energy and belonging to the same electromagnetic shower obtained during electron testbeam measurements. The single crystal energy scales for barrel (EB) and end-cap (EE) is overlaid. The fitted results give $N = (35.1 \pm 0.2)$ ns and $\bar{C} = (20 \pm 4)$ ns.

An intrinsic time resolution better than 0.1 ns for energy $E > 20$ GeV in EB is realized from test beam. Knowing that it takes on average 4.3 ns for a photon to reach the ECAL surface produced from proton-proton interaction point, this demonstrates an intrinsic time resolution of 2% for photons with energy $E > 20$ GeV in the barrel.

4.4 ECAL Time Performance from Collision

Ofcourse, we do not expect the ECAL time resolution in the LHC proton-proton collision environment to be as precised as test beams studies show. Therefore, it is imperative to determine the ECAL time resolution in an LHC environment as knowing the ECAL time resolution in LHC operation can be used for other purposes like eliminating particles such as cosmic rays, beam halo muons, electronic noise and out-of-time proton-proton interactions with broad time distributions. There are factors like clock time variations

over extended time, timing bias with energy and losses in crystal transparency due to radiation which would lead to different pulse shapes which worsen the ECAL time resolution. Thus, a robust crystal time calibration procedure is required if we are to match the intrinsic time resolution realized in test beam studies that we mentioned above. In addition, precise ECAL time during collision can be used in the search for photons with late arrival time produced from the decay of slow moving particles with $\beta \ll 1$, photons traveling along a path which significantly deviates from the obvious straight path from the proton interaction point.

4.4.1 Time Calibrations

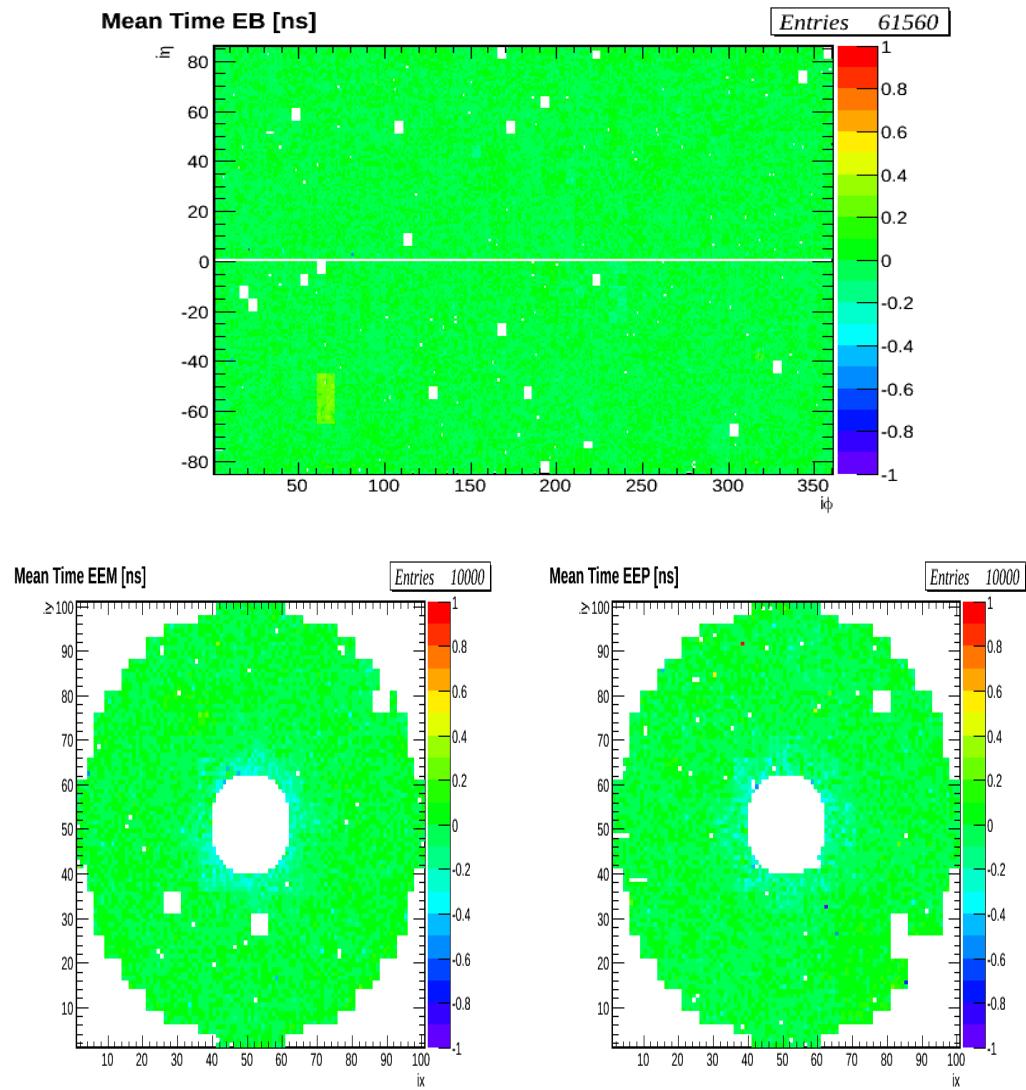
Variations in the crystal time maybe caused by the non-uniform response to electromagnetic showers of each crystal causing difference in pulse shapes or by clock time shifts introduced by interventions during technical shutdowns for machine repairs. These variations can be about 1 to 3 ns on average. To reduce these contributions, we continuously time align all 75,848 PbWO₄ crystals in ECAL by performing crystal *time calibration* once or twice every month. Time calibrating these crystals ensures that there is a uniform response by all the crystals to photons produced from proton-proton collisions, at interaction point, and traveling along a straight path with speeds close to the speed of light to ECAL. It equally guarantees the synchronization of most particles of an event and assigns each event to the correct LHC proton bunch crossing.

The presence of the “ T_{max} Phase”, the difference in pulse shape between each crystal, variation in time of flight by a few nanosecond (ns) and the different intrinsic delays in each channel motivates a time calibration at two separate levels. At the level of the front end electronics (FE) consisting of 5×5 crystals, We perform an initial internal time synchronization by adjusting the time measured by each trigger tower in steps of 1.04 ns among all the trigger towers. Determining the trigger tower time constants needed for time aligning these trigger towers takes place during the process of *Hardware Synchronization*. Offline, we assign time constants to each crystal, during hit reconstruction using the CMS reconstruction software during the transition from uncalibrated hits to calibrated hits.

Offline Time Calibration

The purpose of the offline calibration is to provide time calibration constants for each crystal during event reconstruction. These constants are derived from data recorded during proton-proton collisions and are used to adjust for any clock time shift. A crystal time shift of about 1 ns is caused by the relative time shift between the CMS and the LHC clocks and time shift of about 3 ns to 5 ns is caused by time offsets in the local CMS clock introduced during hardware interventions. The local time offsets is specific only to certain regions of the ECAL detector. The calibration constants for each crystal is the reverse sign of the average time of all the reconstructed energy deposits (rechits) in the crystal. New sets of calibration constants are produced each month throughout the entire LHC run period. Each set of time constants produced has an *interval of validity* (IOV). A total of 17 IOV time constants were produced for the entire LHC run in 2011 and 44 IOV time constants for 2012.

The data sample use for producing the calibration time constants contains mostly electromagnetic particles of mostly loosely triggered photon, electron and hadrons with large electromagnetic shower which travel with nearly the speed of light and largely deposit most of their energy at ECAL. Events of these data samples undergo a rigorous event selection process in order to minimize contamination from events not produced from proton-proton collisions like cosmic muons, beam halos and electronic noise. The produced constants are validated through a full cycle hit reconstruction-time adjustment-hit re-reconstruction procedure of a chosen sub data sample. The validate constants are then used during full event reconstruction. The maps in Figure 4.4, show two dimensional distribution maps of the average time (time calibration constants) for each of the 61,200 crystals in EB and 14648 crystals in EE. The figure shows each crystal average time before (*top 3 plots*) and after (*bottom 3 plots*) calibration. More information about crystal time calibration for the entire LHC Run 1 is found in [39].



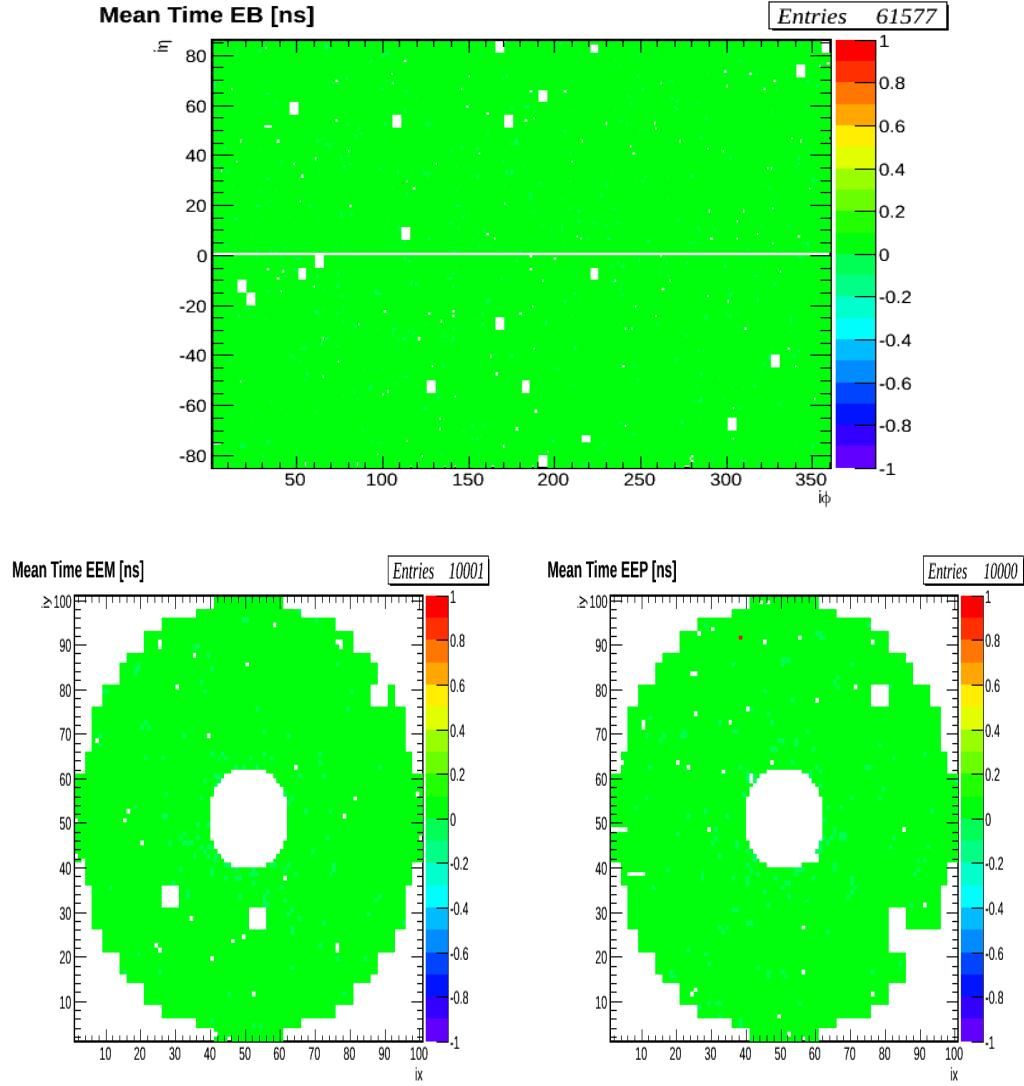


Figure 4.4: *Top 3*: Time calibration maps showing the distribution of mean time for each channel/PbWO₄ crystal in EB (top) and EE (below: EE-(left), EE+(right)) before calibration. *Bottom 3*: Time calibration maps showing the distribution of mean time after calibration. After calibration most crystals have an average time of zero(GREEN).

Hardware Time Calibration

Time offsets introduced during hardware interventions for ECAL front end electronics repairs are calibrated at the level of trigger towers. Most often the adjustment of the

hardware time latency happen during the process of CMS data recording. These time alignments can be done using either data from proton-proton collision or from laser.

Hardware Time Calibration With Collision

The approach for adjusting hardware time offsets during CMS data recording with ECAL for stable proton-proton collision beams involves interfering with the data recording process in order for the hardware time adjustments to be made. This interference during stable proton beams causes frequent data recording downtimes leading to data lost. Frequent down times is one of the causes for the difference between the luminosity recorded by CMS detector compared to luminosity delivered by the LHC. Even though we can argue that this procedure for adjusting the observed trigger tower time shifts using collision data is self efficient, the large CMS data recording downtime and loss of luminosity makes it less dependable especially during frequent machine interventions. To remedy the situation, a laser based time latency adjustment approach can be used.

Hardware Time Calibration With Laser

The ECAL laser system comprise of two lasers, a 440 nm wavelength (close to peak emission for PbWO₄ crystals) laser for monitoring crystal transparency losses and a 796 nm wavelength laser for monitoring readout electronics chain from photodetectors to the electronics (i.e. APDs to ADCs). Both lasers have a jitter of less than 4 ns and as a result, the time from the lasers is averaged over 600 event pulses. The time for each crystal from the laser is expected to be the same as the time from collision data and is represented as $T_{\text{MAX}}^{\text{APD}}$. The laser system is also equipped with a fast acquisition card called MATACQ. The time for each channel recorded using the Matacq is also averaged over 600 event pulse and denoted T_{MATACQ} . The difference, $T_{\text{MAX}}^{\text{APD}} - T_{\text{MATACQ}}$, of the two times, averaged over the 25 crystals of a Clock and Control Unit (CCU) is used as the time for each CCU, $t_{\text{CCU}} = \langle T_{\text{MAX}}^{\text{APD}} - T_{\text{MATACQ}} \rangle$. To obtain the time shift of 25 crystals belonging to the same Front End (FE) electronics, we monitor for change of this time average before (t_{CCU}^B) and after (t_{CCU}^A) hardware intervention during detector maintenance. The time difference, $\Delta t_{\text{CCU}} = t_{\text{CCU}}^A - t_{\text{CCU}}^B$, is averaged over all the 25 crystals, i.e. $\langle t_{\text{CCU}}^A - t_{\text{CCU}}^B \rangle$, is the time shift, and the time calibration constant for the CCU is of opposite sign so that after correcting, the average time

of the CCU is zero. This is done for all the 68 CCUs in a given supermodule (SM) or front end detector (FED). The global time shift of a given FED is caused by the non-homogeneous laser light distribution on all CCU or trigger towers. Each FED has 1,700 PbWO₄ crystals and we produced laser based time calibration constants for all its crystals. Using laser data, we are able to measure the time shift of each CCU to within 0.2 (0.5) ns EB (EE) in precision. Compared to calibration using collision data, there are no downtime in CMS data recording as we can adjust for any hardware time shift prior to stable proton beams. The full procedure including technical details for performing hardware latency adjustments online using collision data or laser data is well described in [40].

4.4.2 Time Bias

An additional source contributing to poor time measurement arise from energy related time bias. This bias originates from the time reconstruction algorithm. It is expected that the ratio approach for time reconstruction performs efficiently for all ADC counts. However, during data recording in LHC Run 1, it was observed that for high ADC counts produced by very energetic particles, an inherent bias in the time is introduced by the multi-gain pre-amplifier electronics for electromagnetic particles with energy above certain gain transition points. The full conversion of the energy of a particle recorded by a crystal in ADC counts to GeV is expressed as; $E_i = G \cdot S_i(t) \cdot C_i \cdot A_i$, where A is the amplitude from the pulse shape in ADC counts, G is the ADC-to-GeV conversion factor equal to 0.039(0.063) in EB(EE), C_i is the inter-calibration coefficients accounting for individual channel response to the electromagnetic shower and $S_i(t)$ is the correction term obtained from laser accounting for radiation-induced channel response. S_i changes over time. The first gain transition point (Gain-1) of the multi-gain pre-amplifier occurs at 4096 ADC counts corresponding to 159.744 GeV in EB and 258.048 GeV in EE. The subsequent Gain 6 and 12 transitions occur at energy values of TeV.

The ratio algorithm introduced time bias at gain transition points which need to be adjusted. These adjustments are made on an energy dependent basis during event reconstruction for a particular CMS event reconstruction software (CMSSW) release. Figure 4.5 shows the comparison between two CMS event reconstruction software releases, CMSSW44X, where these time bias corrections have not been made and CMSSW53X,

where the the corrections have been made. Evidently, the average time displayed in both CMSS44X and CMSS53X show the disappearance of the time bias for all particle energy in CMSSW53X after the time bias corrections were made.

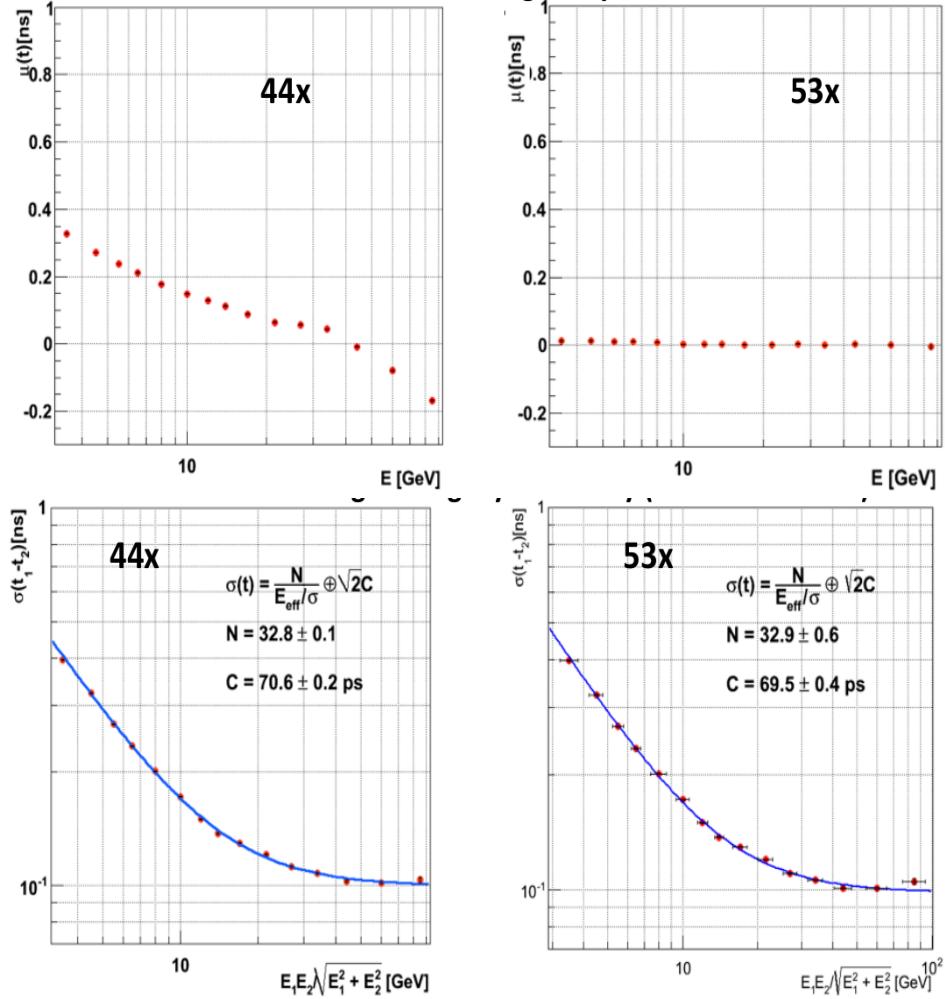


Figure 4.5: Distribution of mean time (μ , top row) and time deviation (σ , bottom row) as a function of crystal energy for EB prior (left) and after (right) time bias corrections depending on energy have been applied.

To investigate further for any additional time bias un related to the particle energy, we study how the crystal geometric position in ECAL covering from $\eta = 0$ to $\eta = 3.142$. The results shown in Figure 4.6 show no crystal position or η dependence. However, time bias of the order of 100 ps have been observed. There are efforts to understand

the source of these time bias with speculations that the time bias might be caused by electronics or loss in PbWO_4 crystal transparency due to radiation.

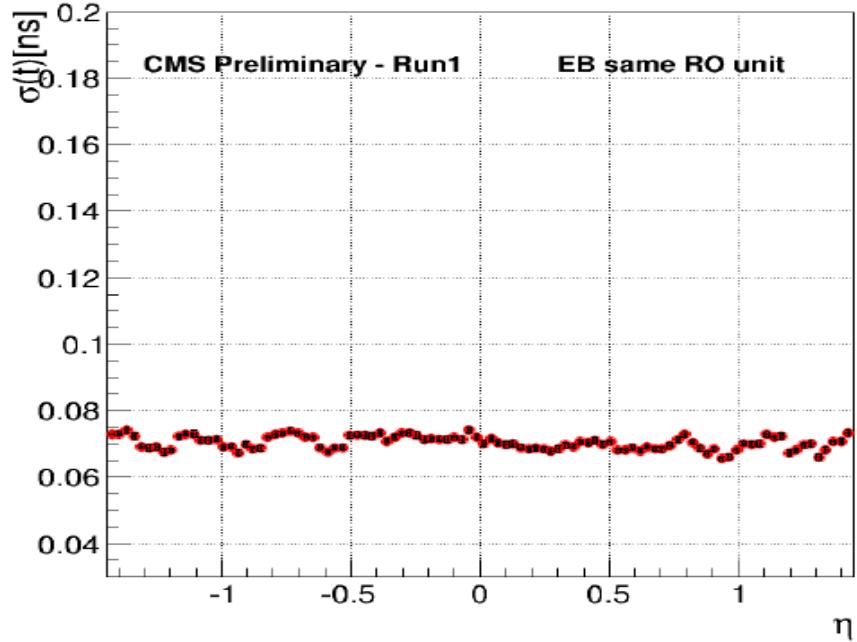


Figure 4.6: Distribution of the time standard deviation (σ) against crystal geometrical position, η , in the ECAL barrel. Almost flat distribution of σ with η .

4.4.3 ECAL Time Performance With Z Bosons

We evaluate the precision of ECAL time measurements during proton-proton collisions by studying the time measurement of a well understood physics process; the decay of the Z boson to an electron pair, i.e. $Z \rightarrow e^-e^+$. We use the standard deviation (time resolution), σ_{eff} , of the difference in arrival time of the two electrons to evaluate ECAL time performance. The standard deviation is obtained from the difference in the seed time, t_{seed} , of the electron electromagnetic shower of each electron after correcting for contributions from the bending of the electron travel path inside CMS magnetic field of 3.8 T. In Figure 4.7, we present the distribution of the time difference, $t_{electron1} - t_{electron2} = t_{seed1} - t_{seed2}$, of both electrons adjusted for time of flight corrections and in Figure 4.8, we show the time resolution or $\sigma_{eff}(t_1 - t_2)$, obtained from the time distribution of the seed crystal time without correcting for the bending of the electron's

flight path contributions. A time resolution of 232 ps in EB and 384 ps in EE is realized. However, if we remove the contributions from the spread in time, $\sigma(t_{\text{collision}})$, due to the finite time it takes for the two proton bunches of length 5.5 cm to collide; which is about $\sigma(t_{\text{collision}}) = \sigma(t_Z) = 183 \text{ ps}$, we obtain an improved time resolution of 142 ps in EB and 337 ps in EE. The selection for Z candidate events required that the electrons have a transverse energy bigger than 10 GeV and the reconstructed Z mass is within, $60 \text{ GeV} < m_{\text{inv}}(e_1, e_2) < 150 \text{ GeV}$ in order to ensure mostly good Z boson candidate events are used.

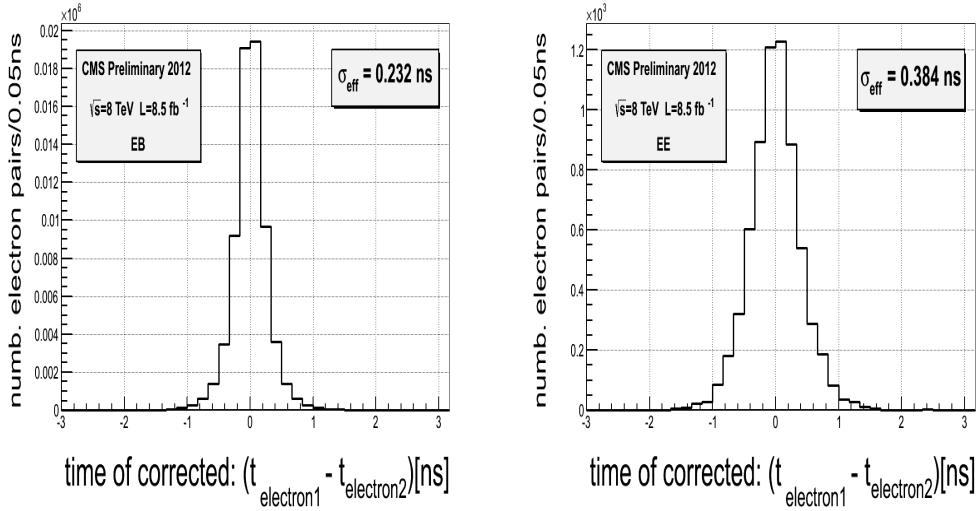


Figure 4.7: Ecal time difference between the two reconstructed electrons in $Z \rightarrow e^-e^+$ decay. The electron time is the seed (crystal with highest energy deposit) time with additional correction due to the time of flight of the electron in EB and EE

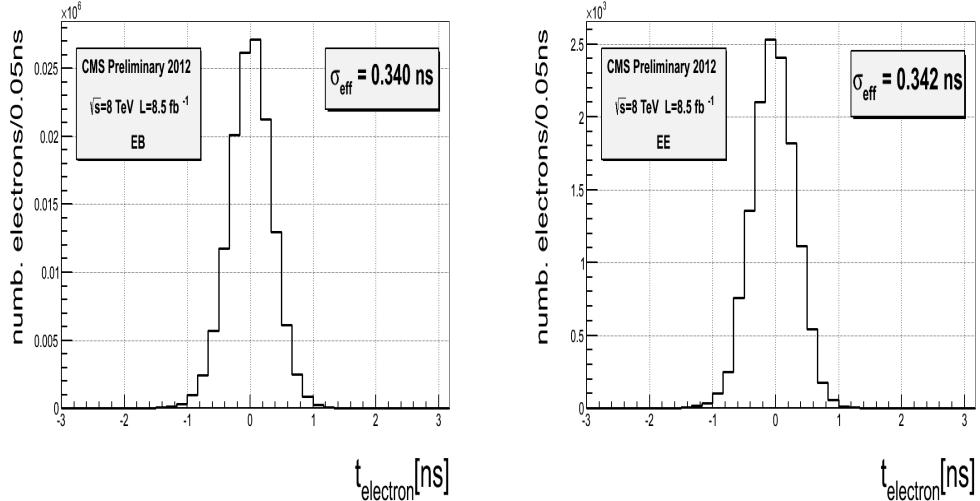


Figure 4.8: Ecal absolute time of a single reconstructed electron in $Z \rightarrow e^-e^+$ decay. The electron time is the seed (crystal with highest energy deposit) time of the electron in EB and EE

We also investigate the contributions attributed to the readout electronics on the time resolution using events with $Z \rightarrow e^-e^+$ decay. Figure 4.9(left), shows the time resolution obtained from measuring the electrons arrival time in the case where the seed crystal time of both electrons is read from the crystals belonging to same ReadOut (RO) electronics and compared to the other case where the seed crystal time is read from crystals belonging to two different readout electronics shown in Figure 4.9(right). The Constant term, C , for the same RO electronics is about 67 ps while that for different RO electronics is 130 ps indicating that electronic readout de-synchronization contributes to the worsening of the time resolution.

Despite, the observed 100 ps bias which seems to be related to readout electronics, we argue that, with a time resolution of $\sigma(t) \leq 400$ ps, the ECAL subdetector is very reliable for time measurements and can be used in the search for delayed electromagnetic particles produced in the decay of long-lived particles.

The ECAL time resolution for the entire LHC Run 1 of 2011 and 2012, comparing the absolute and single precision time measurements is summarized in table 4.1.

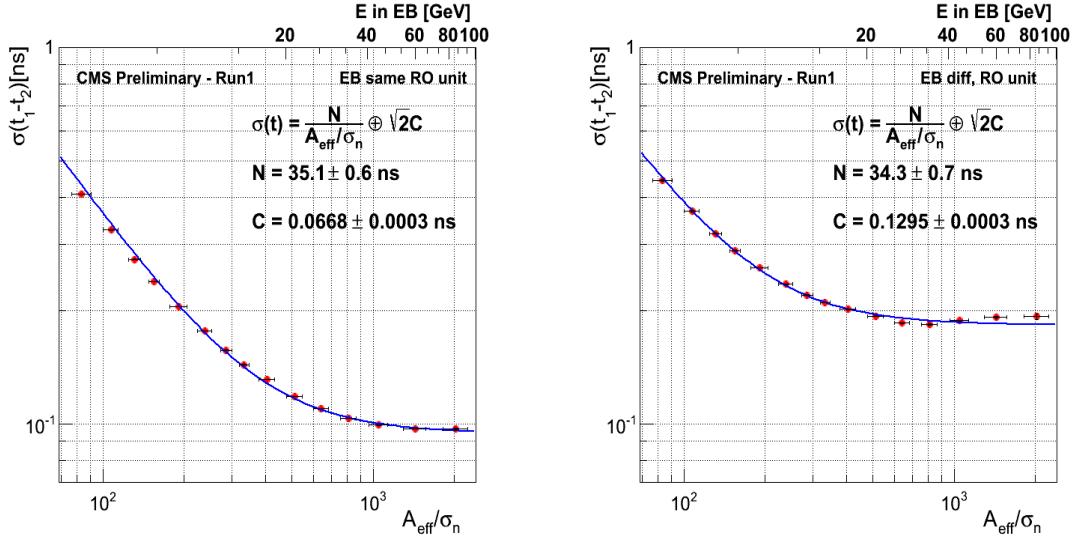


Figure 4.9: Timing resolution from: *left*: Two most energetic crystals in the same readout unit, *right*: Two most energetic crystals belonging to different readout units, as a function of effective amplitude($A_{\text{eff}} = A_1 A_2 / \sqrt{A_1^2 + A_2^2}$) normalized to noise in EB. Both crystals are from reconstructed electrons in $Z \rightarrow e^- e^+$ events.

ECAL Timing Resolution

2011	
Absolute Time	Single Precision
$\sigma_{\text{eff}}(t_{\text{seed}})[\text{ps}]$	$\sigma_{\text{eff}}(t_{e1} - t_{e2})/\sqrt{2}[\text{ps}]$
EB 376	190
EE 356	282
2012	
Absolute Time	Single Precision
$\sigma_{\text{eff}}(t_{\text{seed}})[\text{ps}]$	$\sigma_{\text{eff}}(t_{e1} - t_{e2})/\sqrt{2}[\text{ps}]$
EB 340	164
EE 342	272

Table 4.1: ECAL timing resolution absolute time and single precision for 2011 and 2012 of LHC Run 1

Chapter 5

Event Reconstruction

5.1 Event Reconstruction Overview

An event is comprise of one or more particles produced in a particular proton-proton collision. Event reconstruction is the process of constructing particles and their four momenta using raw data read from the electronics of the different CMS subdetectors. The process begins in each subdetector using a local reconstruction algorithm and eventually extends to all subdetectors. In the calorimeter subdetectors, an energy clustering algorithm is used to produced energy clusters and in the tracker and muon chambers, a combination of tracking algorithms is used to construct particle tracks. By matching selected tracks to energy clusters, particles like photons, electrons, muons, jets and missing transverse energy (E_T^{miss}) are reconstructed.

5.2 Supercluster Reconstruction

The clustering algorithm produces a cluster of clusters called *supercluster*, of the energy of an electromagnetic shower. A supercluster is either a 3×3 or 5×5 crystal energy matrix. About 94% (97%) of the incident photon or electron energy is deposited in this 3×3 (5×5) matrix of crystals in the (η, ϕ) plane in barrel or (x, y) plane in endcap. The 3.8 T magnetic field and material in front of the calorimeter causes bremsstrahlung electrons and converted photons to deposit their energy in a cluster of crystals spread in ϕ . As a result, clustering algorithms start building superclusters with

a seed crystal(crystal with the maximum energy), continue within a narrow window in η by summing the crystal energies along the ϕ which is the direction of energy spread due to the magnetic field. Figure 5.1 is schematic picture showing the direction of the clustering process in the (η, ϕ) plane in barrel and the fraction of electromagnetic energy coverage in a typical supercluster.

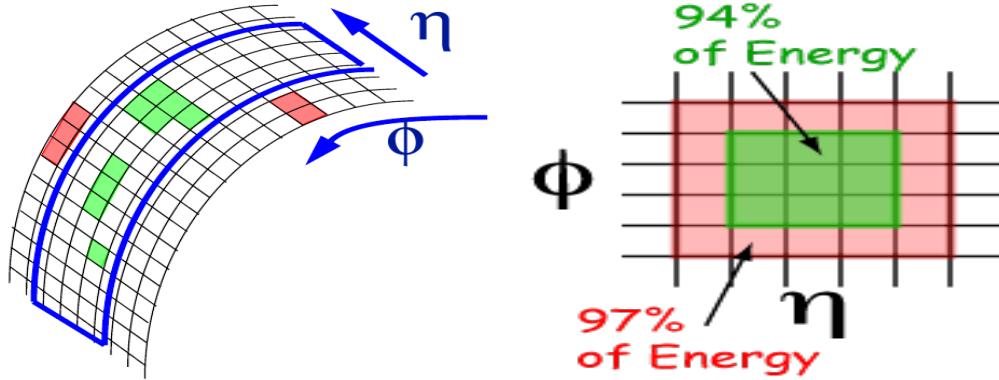


Figure 5.1: Superclustering algorithm direction in the (η, ϕ) plane in EB and fraction of electromagnetic shower energy coverage in a crystal energy matrix.

The two major clustering algorithms used in ECAL are the *hybrid* (EB) and *island* (EE) algorithms.

- **Hybrid Supercluster Algorithm:** This algorithm is used for making super clusters in the barrel (EB). It takes advantage of the $\eta - \phi$ geometry of barrel crystals by taking a fixed 3 or 5 crystals in η and dynamically search and sum separate crystals energy along ϕ . The Hybrid algorithm takes advantage of our knowledge of the lateral shower shape along the η direction. The supercluster consists of basic clusters, usually 3×3 crystals matrix.
- **Island Supercluster Algorithm:** This algorithm is used for making clusters in the endcap (EE). It begins by finding the seed crystal of the electromagnetic shower with maximum energy above a certain energy threshold. Using the seed crystal position, adjacent crystals are examined and added to a cluster until a rise in energy where a crystal belonging to another cluster or crystal that has no readout is reached. For each crystal to be added to the cluster, its energy read must be positive, it must not have been assigned to another cluster and the

previous crystal added in the same direction must have a higher energy. These non-overlapping clusters finally forms a “supercluster”.

5.3 Track and Vertex Reconstruction

Charge particle track reconstruction in CMS uses many algorithms. All these algorithms use the reconstruction positions (*hits*) during the passage of a charged particle in the silicon detectors, to determine the helix trajectories of charged tracks. These helix trajectory is used to measure the particle’s momentum and direction.

The main algorithm used for track reconstruction of proton-proton collisions is the Combinatorial Track Finder (CTF). The CTF proceeds in three stages: seeding, finding and fitting. During seeding, hits that are compatible with the interaction region above a lower p_T limit are used as possible candidates of the charge tracks. Pixel hits are the best track seeds. In the more forward region of the tracker detector, $2 < |\eta| < 2.5$, pixel and inner strips are used for better track finding. The track finding stage uses the Kalman Filter pattern recognition approach where, starting with the seeded parameters, the track trajectory is extrapolated to neighboring tracker layers and compatible hits are assigned to the track. During the fitting stage, the Kalman Filter algorithm is again applied where each candidate track is fitted using least-squares fitting in two stages. The first stage removes possible bias from the seeding stage while the next stage yields the best estimates of the track parameters at the original vertex. Other algorithms like the *iterative tracking algorithm* which is a general purpose tracking algorithms is used in association with customized CTF tracking algorithm to reconstruct the tracks of non-collisions events like cosmics and beam halo.

Similar to track reconstruction, vertex reconstruction involves two stages of vertex finding and fitting. During vertex finding, many tracks are grouped together into vertex candidates. The vertex finding algorithms used depends on whether it is finding primary or secondary vertex or the reconstruction of an exclusive particle decay. In vertex fitting, the best estimates of the vertex parameter like position, covariance matrix and track parameter constraint as well as the fit quality (chi-square, number of degrees of freedom, track weights) is used to distinguish among a given sets of tracks.

5.4 Photon and Electron Reconstruction

Photons are reconstructed from superclusters using the energy contained in a 5×5 crystals matrix surrounding the seed crystal. Since photons are neutral, photons are identified as superclusters in ECAL not associated to any tracks or reconstructed positions in the pixel tracker. To improve photon identification, several criteria are used to distinguish true photons from other particles misidentified as photons. The criteria to identify true photons consists of tracker isolation, ECAL isolation, hadron calorimeter isolation, hadronic to electromagnetic ratio and R_9 , which is the ratio of the electromagnetic energy contained in a 3×3 matrix to the supercluster energy. The R_9 variable is very useful in separating photons from the decay of π^0 with isolated photons since the photons from π^0 decay have a lower value of R_9 compared to isolated photons.

Similar to photons, electron candidates are found when a supercluster is associated to a track reconstructed in the silicon tracker detector and in particular, its inner most layers. In electron reconstruction, the seeding approach is either driven by ECAL or by the tracker. The ECAL driven seeding approach is very efficient for electrons with $p_T > 10 \text{ GeV}/c$. The track driven seeding approach uses a boosted decision tree to perform a pre-selection of the tracker clusters, in order to reduced fake electrons which are light hadrons. Isolated electrons with low p_T and non-isolated electrons (electrons embedded in jets) are reconstructed efficiently using the tracker driven seeded approach. When fitting the electron tracks, we must account for the different energy loss mechanisms of the electron compared to other charged particles. Since electrons energy loss is mainly through the non gaussian nature of its bremsstrahlung, the Gaussian Sum Filter algorithm is used to provide a good estimated the track momentum both at the ECAL surface and at the interaction point.

During proton-proton collisions, in addition to events with true photons and electrons, are events produced from low energy proton-proton collisions called *minimum biased events*, events produced from the radiation of the quarks and gluons called *underlying events*, and events from multiple proton-proton interactions called *Pile Up* (PU) events, contributing to the reconstructed photon or electron energy. Additional contributions to this energy comes from poor detector calibration, poor supercluster or track

reconstruction, faulty electronics and crystal transparency loss due to radiation. The true photon or electron energy must be adjusted for these contributions during reconstruction.

An estimate of the energy deposited by an electromagnetic particle in the ECAL, $E_{e/\gamma}$, can be approximated using Equation 5.1;

$$E_{e/\gamma} = F_{e/\gamma} \cdot [G \cdot \sum_i S_i(t) \cdot C_i \cdot A_i], \quad (5.1)$$

where A_i is the signal amplitude in ADC counts, C_i is the inter-calibration coefficient, $S_i(t)$ is the time-dependent corrections for response variable, usually obtain from laser, G is the global scale calibration allowing to go from energy in ADC counts to GeV and $F_{e/\gamma}$ is the particle energy corrections for geometric, clustering and other effects. The sum is over all the crystals belonging to the photon or electron supercluster. To obtain the true electron or photon energy, energy adjustments which depend on η through $F_{e/\gamma}$ are applied during supercluster reconstruction to account for detector energy mis-measurements caused by cracks between crystals and electronic noise. In Figure 5.2, we show comparisons between cases where no energy adjustments were made to those where energy adjustments (in the form of crystal inter-calibration and laser monitoring corrections for crystal transparency loss) have been made, through measuring the mass of the Z boson. We see improvements on measuring the Z mass, 91 GeV/cc, after the inter-calibration (IC) and laser monitoring (LM) corrections were made. Figure 5.3, shows the case where adjustments are made during supercluster reconstruction. Once again the Z mass is well reconstructed after introducing these corrections.

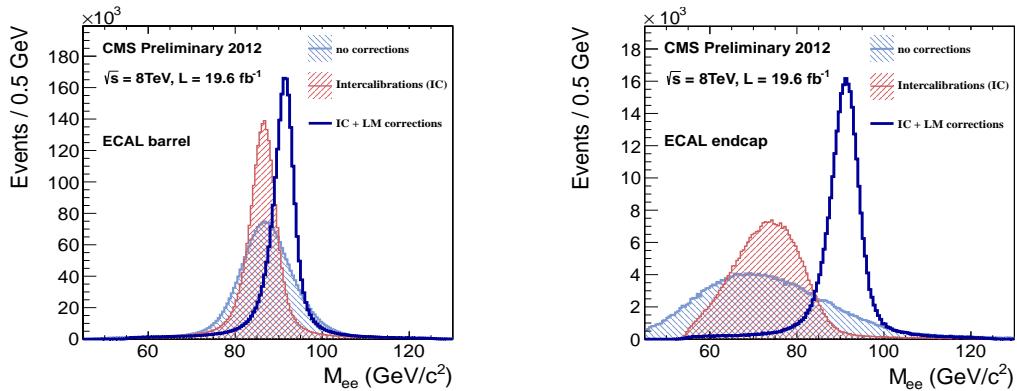


Figure 5.2: Z mass distribution from $Z \rightarrow e^+e^-$ decay showing improvement in the measurement of the Z mass after performing energy adjustments to account for intrinsic spread in crystal, photo-detector response and time-dependent corrections to compensate for channel response loss for EB (right) and EE (left)

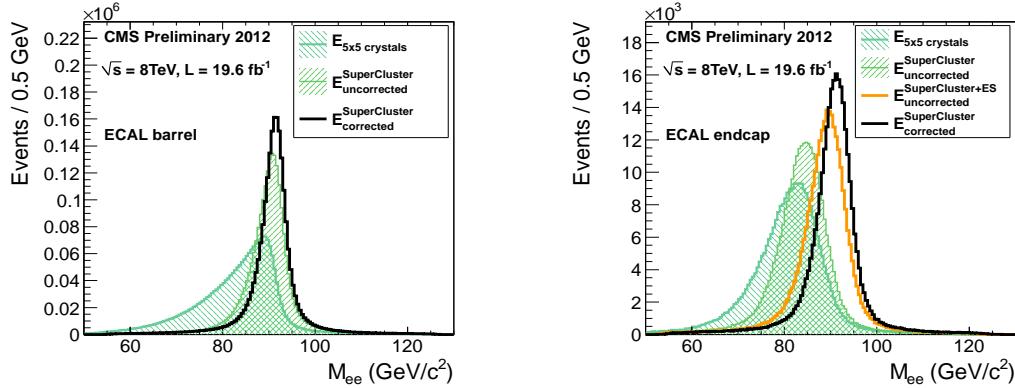


Figure 5.3: Z mass reconstructed using electron superclusters shows improvement in Z mass measurement after applying energy adjustment at superclusters for EB (right) and EE (left).

In Table 5.1, we present a summary of variables constructed using information of the spread of the electromagnetic shower in η and ϕ , the ratio of the energy deposited in HCAL to ECAL, the track p_T and ECAL E_T , isolation in ECAL, HCAL and tracker, the ratio of the energy sums over 3×3 and 5×5 matrices centered on the highest energy crystal of the seed cluster; $R_9 = E_{3 \times 3}/E_{5 \times 5}$ or $R_9 = \sum E_9 / \sum E_{\text{Supercluster}}$ and impact parameter, $d0$, which is the minimum separation of the electron track computed with respect to the reconstruction vertex and track transverse momentum, p_T . Photon and electron selection using these variables, applied during electron and photon identification, have been shown to produce very good quality electron and photon candidates with efficiency above 70%.

Simple Cut Based Electron Photon Identification

ID Variable	Electron	Photon
H/E	0.05(EB), 0.10(EE)	0.05
$ \Delta\eta_{in} $	0.005(EB), 0.007(EE)	0.015(EB)
$ \Delta\phi_{in} $	0.09(EB), 0.09(EE)	N/A
$\sigma_{i\eta i\eta}$	0.01(EB), 0.03(EE)	0.011(EB), 0.03(EE)
Pixel Veto	No	Yes
$ d0 (vertex)$	0.02(EB), 0.02(EE)	Veto
$ dZ (vertex)$	0.1(EB), 0.1(EE)	0.02 (cm)(Veto)
$ 1/E - 1/p $	0.05(EB), 0.05(EE)	N/A
PF isolation / p_T (cone dR=0.3)	0.15(EB), 0.10(EE)	N/A
ECAL Isolation	same	$4.2 + 0.006 * E_T^\gamma +$ $0.183 * \rho(\text{EB})$
HCAL Isolation	same	$2.2 + 0.0025 * E_T^\gamma +$ $0.062 * \rho$
TRACK Isolation	same	$2.0 + 0.001 * E_T^\gamma +$ $0.0167 * \rho$
Rho corrected PF photon isolation	N/A	$1.3 + 0.005 * p_T^\gamma(\text{EB})$

Table 5.1: Simple cut-based selection criteria for electron and photon identification.

5.5 Muon Reconstruction

Muon tracks are reconstructed using the all-silicon inner tracker (tracker tracks) and the muon system (standalone tracks). The standalone tracks are reconstructed using reconstructed positions (hits) in the muon system consisting of the Drift Tubes (DT) in the barrel ($|\eta| < 0.9$), Cathode Strip Chambers (CSC) in the endcaps ($1.2 < |\eta| < 2.4$) and Resistive Plate Chambers (RPC) in the overlap region ($0.9 < |\eta| < 1.2$). There are two independent muon reconstruction approaches: *Global muon reconstruction (Outside-in)* and *Tracker muon reconstruction (Inside-out)*. For Global muon reconstruction, each standalone-muon track is matched to a tracker track by comparing the parameters of

the two tracks propagated to a common surface. The global muon track is fitted combining hits from the tracker track and standalone-muon track using the Kalman-filter algorithm. For the tracker muon reconstruction, all tracks with $p_T > 0.5 \text{ GeV}/c$ and total momentum $p > 2.5 \text{ GeV}/c$ are considered as possible muon candidates and are extrapolated to the muon system taking into consideration the magnetic fields, the average expected energy loss in the calorimeters and multiple Coulomb scattering in the detector material to locally reconstruct segments in the muon system. A combination of different muon algorithms depending on the muon p_T , provides a robust and efficient muon identification.

Using the beam spot as constraint of the muon's vertex, it is possible to distinguish between muons produced from proton-proton collisions from those produced from cosmic *cosmic muons* and proton beams interacting with the gas in beam pipe and proton dump muon splash events produced 150 m upstream the CMS detector during proton beam dump. The proton beam related muons are referred to as *beam Halo muons*. Figure 5.4 show an illustration of trajectories of different muon sources interacting with the CMS detector.

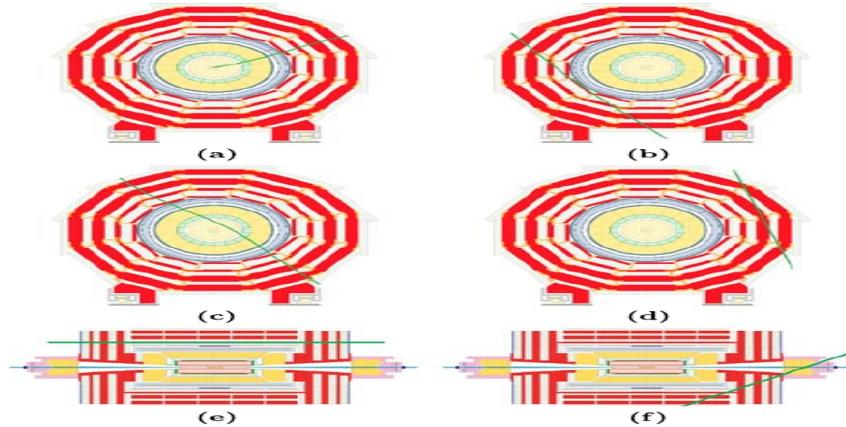


Figure 5.4: Illustration of muons from proton-proton collision, cosmic rays and beam halo. (a) Muons from collision propagating from the center and moving outwards, (b) Cosmic muons traveling through the detector leaving signals in opposite hemispheres of the muon system, (c) Cosmic muons leaving signals in the tracker and opposite hemispheres, (d) cosmic muons entering and leaving the detector without passing through the muon detector layers, (e) beam halo muons penetrating the detector and leaving signals in the endcaps and (f) Cosmic muons entering the detector through the endcap and leaving through the barrel and which can happen in a *vice-versa* manner.

Particle Flow Algorithm

The *Particle Flow* (PF) algorithm is an algorithm for reconstructing particles combining information from the tracker, ECAL, HCAL and muon chambers of the CMS detector. It uses a combination of different algorithms comprising of calorimeter clustering, tracking and extrapolation to calorimeters, muon identification, electron pre-identification and linking of topological elements, for reconstructing a list of particles which include photons, charge hadrons, neutral hadrons, muons and electrons. The same list of particles is subsequently used to reconstruct composite “particles” like jets, E_T^{miss} and taus. The versatility of the PF algorithm is the reason why it was introduced for reconstructing Jets, missing transverse energy (E_T^{miss}) where complete information of the event content from every subdetector is needed. The PF algorithm uses tracks, electron energy seeds, 4-momentum, super cluster energy calibration, bremsstrahlung tracks for electron and photon reconstruction making it extremely efficient at minimizing electron and photon misidentification. For E_T^{miss} reconstruction where full reconstruction of all the particles belonging to an event is necessary, the PF algorithm is very reliable.

5.6 Jet Reconstruction

A jet is a spray of particles arising from the hadronization of colored particles. Because jets are made of many particles like electrons, photons, hadrons, charged and neutral electromagnetic and hadronic showers, they are best reconstructed using the particle flow algorithm. Jets reconstructed using the PF algorithm are called *Particle Flow Jets*. Using calorimeter towers as input, jets can also be reconstructed using the Anti- k_T clustering algorithm which combines four vectors according to their relative transverse momentum within a standard cone size of $R = 0.5$ in the (η, ϕ) plane. The quality of a reconstructed jet depends on the jet identification selection criteria known as the *JetID*. High quality jets are required to have an electromagnetic energy fraction (EMF) > 0.01 , within the ECAL fiducial region of $|\eta| < 2.6$, the number of calorimeter cells containing more than 90% of jet energy (N_{jet}^{90}) must be > 1 , the fraction of jet energy in the hottest Hybrid photodetector (HPD) unit in HCAL readout within a jet (f_{HPD}) must be > 0.98 , the charge hadron fraction (CHF) > 0.0 if within $|\eta| < 2.4$, the neutral hadron fraction (NHF) < 1.0 , the charge electromagnetic fraction (CEF) < 1.0 , and

neutral electromagnetic fraction $z(NEF) < 1.0$. These jetId selection requirements have been shown to remove mis-reconstructed jets arising from spurious energy deposition in subdetectors with good efficiency. The jet energy is often mis-measured due to cracks in the detector, non-linear responses in the calorimeters, electronic noise, poor detector calibration and additional energy from PU. Therefore, adjusting the jet energy through *jet energy corrections* (JEC) is often performed guaranteeing a reliable measurement of the jet energy during reconstruction.

5.7 Missing Transverse Energy Reconstruction

It is very challenging to detect particles which do not interact with the detector. Such particles include very weakly interacting particles like neutrinos (ν), neutralinos ($\tilde{\chi}_1^0$) and gravitino (\tilde{G}). In CMS, the presence of this kind of particles in an event is inferred through a measurable quantity called *missing transverse energy*. Measurable quantities like transverse momentum (p_T) and cone size, $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2}$, are used to identify and distinguish particles traveling in the transverse plane, (x, y) or (η, ϕ) , from the longitudinal plane, along the z direction. Employing the conservation of momentum in the transverse plane, we can use the imbalance in transverse momentum required to ensure conservation of total transverse momentum of an event to infer the presence of an undetectable particle traveling in the transverse plane. The amount or quantity of this transverse momentum imbalance or missing transverse momentum is called missing transverse energy (MET) or E_T^{miss} . The presence of reconstructed E_T^{miss} in an event indicates the presence of very weakly interacting particle(s) which escape the detector undetected.

We define E_T^{miss} as the magnitude of the negative vector sum of the uncorrected transverse energy deposit of all the particles in an event.

$$E_T^{\text{miss}} = \left| - \sum_n (E_n \sin \theta_n \cos \theta_n \hat{\mathbf{i}} + E_n \sin \theta_n \sin \theta_n \hat{\mathbf{j}}) \right| = |\not{E}_T^x \hat{\mathbf{i}} + \not{E}_T^y \hat{\mathbf{j}}| \quad (5.2)$$

Where, n is the sum over all calorimeter energy deposits including energy deposits in towers, reconstructed energies (hits) or generator level particle energies. In order to measure E_T^{miss} accurately, a particle detector should be as near hemispherical as possible to

allow for complete measurement of the transverse momentum of all the particles belonging to an event. The CMS detector, provides a near hemicity with its nearly 4π solid angle coverage. This near 4π solid angle is achieved through the hadronic forward (HF) detector with little space allowing for the passage of the proton beams. Since according to conservation of momentum, the total transverse momentum before and after collision should be the same, by measuring the p_T of every detectable particle belonging to an event after proton-proton collision, the imbalance in p_T required to ensure the conservation of transverse momentum in the event, is ascribed to the momentum carried away by undetected particle(s). This idea is used to infer the presence of neutrinos W boson decay to an electron and an anti-electron neutrino; $W \rightarrow e + \bar{\nu}$, where the missing transverse momentum is used to infer the presence of the undetectable neutrino.

Since particle interactions involving supersymmetry and other beyond the standard model processes involves large E_T^{miss} due to the production of undetectable particle(s) during proton collisions, E_T^{miss} is the desirable quantity to use in many particle searches involving the possible production of new very weakly interacting particles. Standard Model particle interactions produces small E_T^{miss} . In any particle search which involves the production of new weakly interacting particle, it is imperative to accurately measure E_T^{miss} , as some standard model interaction processes like machine induced backgrounds from halo muons, mis-measured energy and mis-reconstructed particles and anomalous signals like spike can lead to large E_T^{miss} mimicking the signal for new particle(s).

5.8 Anomalous Signals

Anomalous signals are signals read from the electronics called *spikes* with abnormally large energy deposits in the avalanche photodiodes without any scintillation in the PbWO₄ crystals. They are produced when neutrons or charged hadrons like protons, strike directly ionizing the silicon of the photodiode producing an electronic signal in the absence of any crystal scintillation.

Because of no scintillation, they appear much earlier (negative) in ECAL measured arrival time and often populate the earlier time of the reconstructed hit (*rechit*) time distribution. Their energy deposit range from a few GeV to ECAL saturation energy of about 1.7 TeV. Since they do not electromagnetically shower in PbWO₄ crystals,

their electromagnetic energy shower shape is very isolated, with only one or two crystals contributing to the energy cluster. Spikes may also have positive time appearing late or delayed in their arrival time at ECAL populating the tails of the rechit time distribution. The late arrival time is due to the slow propagation of neutrons through the CMS detector.

Numerous test beam, collision data and simulation studies,[45], have been carried out towards understanding the properties of events with spikes and how they can be tagged and removed. These studies, reveal that spikes can be identified and removed using a topological energy variable constructed as $1 - \frac{E_4}{E_1}$ known as the “*Swiss-Cross*” (SX). E_1 is the energy deposit of the central (highest energy) crystal and E_4 is the sum total of the energy of the neighboring four crystals in an (η, ϕ) plane. A selection cut in $\text{SX} > 0.95$ rejects more than 99% of isolated spikes with transverse energy greater than 10 GeV with very little impact on the efficiency of selecting electromagnetic EM showers. Other topological variables like $1 - \frac{E_2}{E_6}$ and $1 - \frac{E_2}{E_9}$, where E_2 is the sum of the energy of two crystals sharing the energy deposited and $E_6(E_9)$ is the sum of the neighboring 6(pairs-of)(9) crystals in the (η, ϕ) plane are used. The $1 - \frac{E_2}{E_6}$ variable is used to identify isolated spikes whose energy deposit spread in two adjacent crystals while the $1 - \frac{E_2}{E_9}$ is used to identify non-isolated spikes i.e. spikes which are found embedded in a supercluster. It has also been shown that applying selection cuts on the rechit time of ± 3 ns leads to more than 90% efficiency for rejecting spikes. However, in this thesis, we do not require such selection cuts on the rechit time as these rechits include rechits of possible delayed electromagnetic particles with arrival time beyond 3 ns produced during proton-proton collisions.

Chapter 6

Search Analysis for Long-Lived Particles

6.1 Analysis Strategy

This analysis is about the search for events with at least a single late arrival time photon at the Electromagnetic Calorimeter (ECAL) and large missing transverse energy (E_T^{miss}). It uses a counting method where an excess in the number of events with photon time above a defined ECAL timing threshold, to the expected number of events from background processes, indicates the presence of a new physics phenomena. Such a phenomena is the decay of a massive long-lived neutral particle into a late photon and large E_T^{miss} which is not common with standard model interactions.

We expect most of our background events of this search to arise from non-collision rather than proton-proton collision events.

6.1.1 Signal and Background Events

We expect a typical signal event of the decay of a massive neutral long-lived particle, to be produced in a proton-proton (pp) collision at the LHC and detected with the CMS detector, according to the prediction of a benchmark Gauge Mediating Supersymmetry Breaking (GMSB) model described as the “Snowmass Points and Slope 8” (SPS8).

In this model, the massive long-lived neutral particle is the Next-to-Lightest Supersymmetric Particle (NLSP) which is the lightest neutralino ($\tilde{\chi}_1^0$). The $\tilde{\chi}_1^0$ decay into a photon (γ) and the Lightest Supersymmetric Particle (LSP) called the gravitino (\tilde{G}), $\tilde{\chi}_1^0 \rightarrow \gamma + \tilde{G}$. The \tilde{G} is undetected because it does not interact with the CMS detector material. Its presence is indirectly inferred by measuring the missing transverse momentum (E_T^{miss}). The photon, because of the long-lived nature of the $\tilde{\chi}_1^0$, arrives late at the ECAL compared to photons produced from nominal pp collisions.

In a LHC collider, the $\tilde{\chi}_1^0$ can be singly or pair produced in the direct interaction of the quarks/gluons (partons) inside the proton, $q\bar{q} \rightarrow \tilde{\chi}^\pm \tilde{\chi}_1^0$ or $q\bar{q} \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$, respectively. However, in this analysis, we are interested in cases where the $\tilde{\chi}_1^0$ is produced indirectly from the cascade decay of heavier supersymmetric particles like gluino (\tilde{g}) and squark (\tilde{q}), $gg \rightarrow \tilde{g}\tilde{g} \rightarrow q\bar{q}q\bar{q}\tilde{\chi}_1^0\tilde{\chi}_1^0$ and $q\bar{q} \rightarrow \tilde{q}\tilde{q} \rightarrow \bar{q}q\tilde{\chi}_1^0\tilde{\chi}_1^0$, respectively. This is because the gluino/squark production cross-section is largest at a pp collider like the LHC, and in association with the neutralino decay, the event also contains multiple high- p_T jets which provide an added signal requirement useful for suppressing background events with no jets. Thus, our typical signal event should comprise of at least a single delayed high- p_T photon, large E_T^{miss} and multiple high- p_T jets. Finding such an event with the CMS detector will be evidence of a new physics phenomena.

Our background events could be produced from pp collisions which we call *collision background events* and from non- pp collisions which we call *non-collision background events*, like the so-called beam halo, cosmic muons and spikes. Most of our background events are from non-collision background events with late photons and large E_T^{miss} .

Collision background events can mimic the $\tilde{\chi}_1^0$ decay signal producing late photons and large E_T^{miss} in cases where the ECAL time or E_T^{miss} is mismeasured or the collision event has real E_T^{miss} . For example, multijets and QCD events, inclusive Z+jets/W+jets events, inclusive top-anti-top ($t\bar{t}$)+jets events and inclusive ZZ/WW/WZ+jets events can either have true E_T^{miss} ; like events with $Z \rightarrow \nu\bar{\nu}$ and $W \rightarrow e\bar{\nu}_e$ decays, where the undetected neutrino (ν) leads to true E_T^{miss} , or fake E_T^{miss} (instrumental E_T^{miss}), where there is no undetected particle pertaining to the event but rather E_T^{miss} arising because of poor reconstruction of the energy of particles, as it is with the case for QCD events. The late photon arises when one of the jets or an electron is misidentified as a photon and its ECAL time is misreconstructed. The other jets in the event satisfy the high- p_T

multijets requirement.

6.1.2 Signal Modeling

We produced simulated high-energy-physics signal events according to the SPS8 benchmark model using general-purpose event generators based on Monte Carlo (MC) methods of numerical computations. The process of event generation starts with the production of *SUSY Les Houches Accord* (SLHA) files using a SUSY software package called **ISASUSY**[54]. These SLHA files contain the masses, interaction couplings, decay widths and all possible decay channels and their Branching Ratio (*BR*) of every supersymmetric particles. Each file also contain the value of the fundamental GMSB model parameters;

$$\left\{ \Lambda, M_{\text{mess}}, N_5, \tan(\beta), sgn(\mu), C_{\text{grav}} \right\} \quad (6.1)$$

spanning all the possible, in terms of particle kinematics, supersymmetric particle production and decay configurations or phase space, as defined by the model, in the production of events. In the SPS8 model, the choice of parameters is such that

$$sgn(\mu) = 1, \tan(\beta) = 15, N_5 = 1, M_{\text{mess}} = 2\Lambda, \quad (6.2)$$

where C_{grav} and Λ are not fixed . By varying, respectively, C_{grav} and Λ , we can exploit different decay scenarios where the $\tilde{\chi}_1^0$ has a different lifetime and mass. For a specific choice of C_{grav} and Λ , we produce a signal sample where the events have a mean lifetime ($c\tau$) and mass of the $\tilde{\chi}_1^0$. A special software package called **HDECAY**, is used to handle the decay of all supersymmetric particles including the $\tilde{\chi}_1^0$.

The events are generated according to the physics process, which we described in detail in section 2.2.1, where the $\tilde{\chi}_1^0$ production and decay to photon and gravitino is through the cascade decay of the gluino/squarks;

$$p + p \rightarrow \tilde{g}\tilde{g}, \tilde{q}\tilde{q} \rightarrow [1 \text{ or } 2 \text{ cascade decays}] \rightarrow 2\tilde{\chi}_1^0 + \text{jets} \rightarrow 2\gamma + 2\tilde{G} + \text{jets}. \quad (6.3)$$

We use **PYTHIA-6** [55] as the MC event generator. It takes as input an SLHA file and generates events of supersymmetric particles, produced at the LHC pp collider with center of mass energy, $\sqrt{S} = 8$ TeV. The interaction of these supersymmetric particles with the

CMS detector is simulated using the GEANT4 [56], particle detector simulation software. A full physics event is reconstructed from energy deposits (hits) in the CMS detector using the CMS event reconstruction Software (CMSSW). To minimize any disagreement between MC and data, we use the same CMSSW release version (**CMSSW_5_3_29**) in the MC event reconstruction with the same detector conditions as the recorded data.

6.1.3 Datasets

The data sample used in this analysis was produced during LHC Run 1 in 2012 with proton-proton (pp) collisions at the center of mass energy, $\sqrt{S} = 8 \text{ TeV}$. We used data equivalent to total integrated luminosity of 19.1 fb^{-1} recorded by the CMS detector.

Data Samples

Our collision data samples comprise of triggered single photon events from luminosity sections certified as good. In Table 6.1, we present these data samples showing their corresponding integrated luminosity. The *json* file with the list of certified good luminosity sections is **Cert_8TeVPromptReco_Collisions12_JASON.txt**.

Data Sample	Recorded Luminosity [fb $^{-1}$]
/Run2012B/SinglePhoton/	5.1
EXODisplacedPhoton-PromptSkim-v3	
/Run2012C/SinglePhoton/	6.9
EXODisplacedPhoton-PromptSkim-v3	
/Run2012D/SinglePhoton/	7.1
EXODisplacedPhoton-PromptSkim-v3	
/SingleElectron/Run2012A-22Jan2013-v1/AOD	5.2
/DoubleElectron/Run2012C-22Jan2013-v1/AOD	4.8

Table 6.1: Data samples and their corresponding integrated luminosity totaling 19.1 fb^{-1} used in the our delayed photon search analysis

Monte Carlo Samples

The MC samples were produced with *Summer 2012* prescription of the calibration and alignment status of the CMS detector and pile up conditions at 8 TeV.

The GMSB SPS8 signal samples had 50,000 events for each mean lifetime ($c\tau$) and SUSY breaking scale (Λ_m) or mass of $\tilde{\chi}_1^0$ ($m_{\tilde{\chi}_1^0}$). The sample are produced for different $c\tau$ ranging from 50 cm to 1000 cm for each Λ or $m_{\tilde{\chi}_1^0}$ point. We vary Λ from 100 TeV to 220 TeV which is equivalent to $m_{\tilde{\chi}_1^0}$ ranging from $139 \text{ GeV}/c^2$ to $314 \text{ GeV}/c^2$. Table 6.2 shows a summary of our signal MC samples used in this analysis with the number of events, cross-section and branching ratio for $\tilde{\chi}_1^0$ production and decay to γ and \tilde{G} for each SUSY breaking scale.

The γ +jet MC samples were generated for different momentum of the colliding partons and normalized to the 19.1 fb^{-1} integrated luminosity. The cross-sections, p_T of the photon (\hat{p}_T) radiated by the colliding parton and the number of events in each sample is summarized in Table 6.3. The \hat{p}_T range from $50 \text{ GeV}/c$ to $800 \text{ GeV}/c$ was covered.

Λ [TeV]	$c\tau$ (mm)	σ_{LO} (pb)	Number of Events	Branching Ratio
100	500-10,000	0.368	50,000	0.9444
120	500-10,000	0.133	50,000	0.9042
140	500-10,000	0.0574	50,000	0.8711
160	500-10,000	0.0277	50,000	0.8464
180	500-10,000	0.0145	50,000	0.8282
220	500-10,000	0.0044	50,000	0.8282

Table 6.2: Signal GMSB SPS8 Monte Carlo samples for different Λ with $50 \text{ cm} < c\tau < 1000 \text{ cm}$ and Branching Ratios (BR) studied in this analysis

$\hat{p_T}$	σ_{LO} (pb)	Number of Events
$50 \sim 80$	3322.3	1995062
$80 \sim 120$	558.3	1992627
$120 \sim 170$	108.0	2000043
$170 \sim 300$	30.1	2000069
$300 \sim 470$	2.1	2000130
$470 \sim 800$	0.212	1975231

Table 6.3: The $\gamma+$ jets samples studied in this analysis

We checked that the GMSB signal samples were correctly generated by PYTHIA-6, by looking at the number of photons, E_T^{miss} and number of jets in each signal event. We also measured the mean lifetime of the $\tilde{\chi}_1^0$ by performing a fit analysis to the distribution of the $\tilde{\chi}_1^0$ mean lifetime computed from its transverse distance traveled before decay. This distance is computed using its production and decay vertices. By comparing the computed $c\tau$ to its theoretical value supposedly used in the event generation, we are able to validate that each signal MC sample.

We also observed that most of the events had at least a single photon (left plot) and at least 2 jets (right plot) shown in the top plots of Figure 6.5. Comparing different signal samples with $c\tau = 2000 \text{ mm}, 4000 \text{ mm}, 6000 \text{ mm}$ of $\Lambda = 180 \text{ TeV}$ and a $\gamma+$ jet with $120 < \hat{p_T} < 170 \text{ GeV}/c$ sample, we observed that the E_T^{miss} (shown in the bottom plot of the same figure) from signal events was larger than the E_T^{miss} of events from the $\gamma+$ jet sample agreeing with our expectation that the $\gamma+$ jet sample has mostly fake E_T^{miss} while signal samples have large E_T^{miss} due to the $\tilde{\chi}_1^0$ decay to gravitino.

These observations gave us confidence that the signal samples were properly generated and indeed most of the signal events had at least one $\tilde{\chi}_1^0$ which decayed to a photon and gravitino.

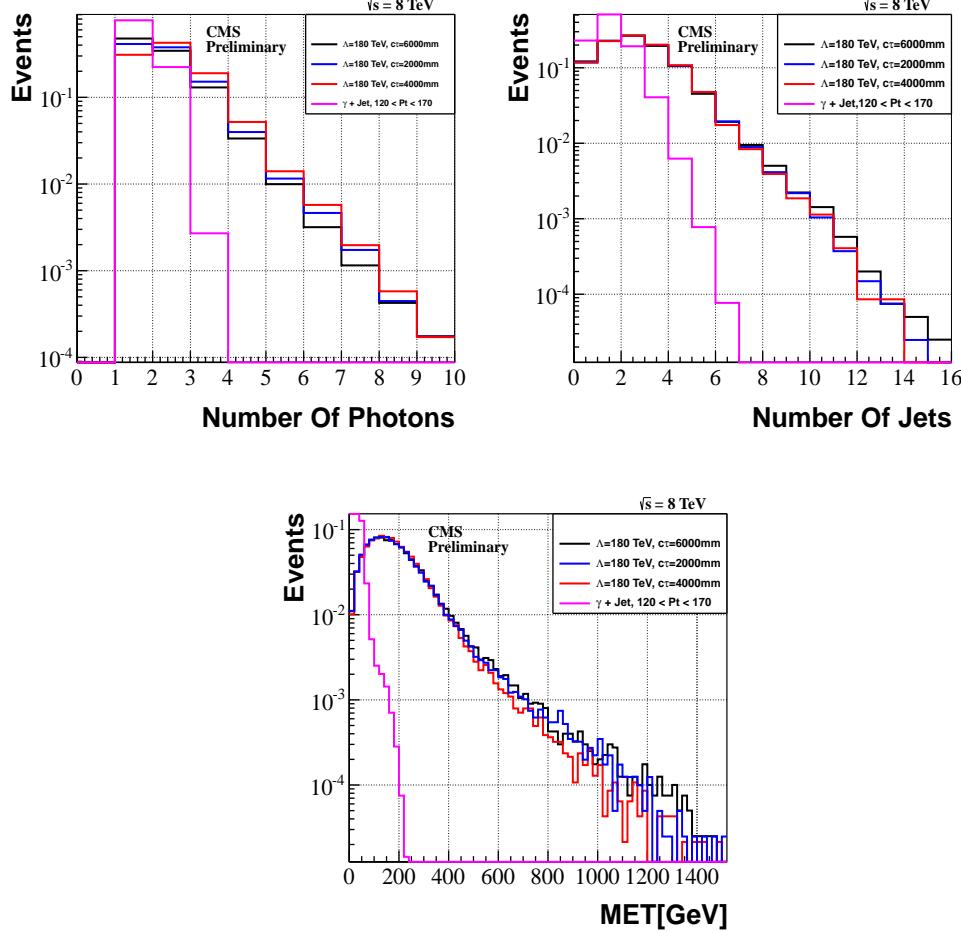


Figure 6.1: Number of photons (top left), Number of jets (top right) and E_T^{miss} (bottom) for events with $\tilde{\chi}_1^0$ decay to γ and \tilde{G} for different $c\tau = 2000 \text{ mm}, 4000 \text{ mm}, 6000 \text{ mm}$ points of $\Lambda = 180 \text{ TeV}$ of the SPS8 model. A $\gamma+\text{jet}$ with $120 < p_T < 170 \text{ GeV}/c$ sample shown for comparison with the signal samples.

6.2 ECAL Timing

In this section, we describe how the photon arrival time is measured, how we adjust the MC time so that it agrees with data. We also described the neutralino lifetime and discuss the possible reasons why the photon is delayed and finally we discuss the satellite bunches and argue that they could be a possible background source to late photons.

6.2.1 Photon Time Measurement

The presence of spikes, noisy crystals and pile-up events, demand a robust method for measuring the photon arrival time at ECAL and since ECAL time is our main observable for distinguishing background from signal events, such a method must be capable of reducing timing bias which may arise from such anomalous events. As a result, we studied different methods for measuring the photon arrival time at ECAL.

The electromagnetic shower of an electromagnetic particle spreads across several crystals (energy and time measurements from several channels) belonging to the particle's supercluster containing all of the particle's energy. Using the supercluster, the arrival time of an electromagnetic particle can be defined using either the reconstructed time (t_{reco}) of a single crystal, the *seed crystal* (crystal with the highest energy deposit), or a weighted average time calculated using the reconstructed time and its uncertainty of each crystal of the supercluster. We write, t_{seed} , for the seed time and, t_{Ave} , for the average time defined as

$$t_{Ave} = \frac{\sum_{i=1}^N \frac{t_{reco}^i}{\sigma_i^2}}{\sum_{i=1}^N \frac{1}{\sigma_i^2}}, \quad (6.4)$$

where N is the total number of crystals of the supercluster, $t_{reco,i}$ and σ_i are the time and uncertainty on the reconstructed time of each channel, respectively. Figure 6.2 shows a comparison of a photon time measured as the seed time, t_{seed} , and as the average time, t_{Ave} . Both distributions have been normalized to total number of events.

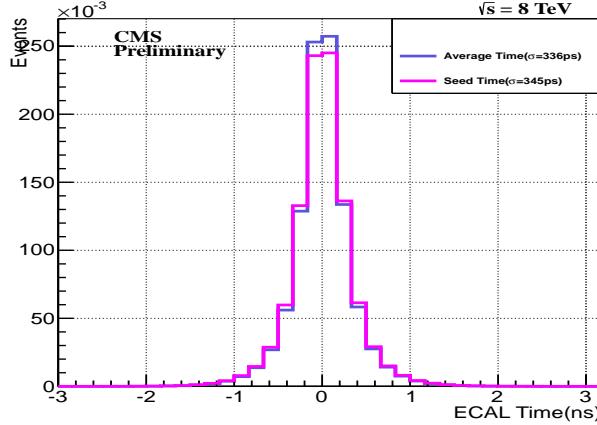


Figure 6.2: Measuring the photon time using either seed time (black) or average time (blue). σ of the Gaussian fit from seed time slightly better average time which is computationally intensive.

The width (σ) of both Gaussian distributions are similar. We observed $\sigma = 345$ ps for the seed time compared to $\sigma = 336$ ps for the average time distributions.

In this analysis, we used the seed time for the photon ECAL time. However we also used the χ^2 on the photon ECAL time computed using the average time for identifying electromagnetic particles with bad timing. For example, when one or more of the crystals in a supercluster is poorly time calibrated or embedded with a spike, the photon time from the weighted average time measurement can be very biased which is a good reason to reject such a supercluster.

The χ^2 on the photon ECAL time calculated as

$$\chi^2 = \sum_{i=0}^N \frac{(t_{reco}^i - t_{Ave})^2}{\sigma_i^2} \quad (6.5)$$

where, N is the number of crystals in the photon supercluster, t_{reco}^i and σ_i are the time and uncertainty on the reconstructed time for each crystal and t_{Ave} is the mean time defined in Equation 6.4, is a reliable quantity to legitimized the photon time and to distinguish fake photons (jets misidentified as photons) and spikes from true photons. Figure 6.3 shows the profile of the pulse shape (left) of an identified spike with that of a

true photon. A distribution of the normalized χ^2 against the ECAL time is also shown in the right plot of the same figure. Photons with a signal pulse shape profile as that of spikes are associated with large values of χ^2 and usually have negative ECAL time.

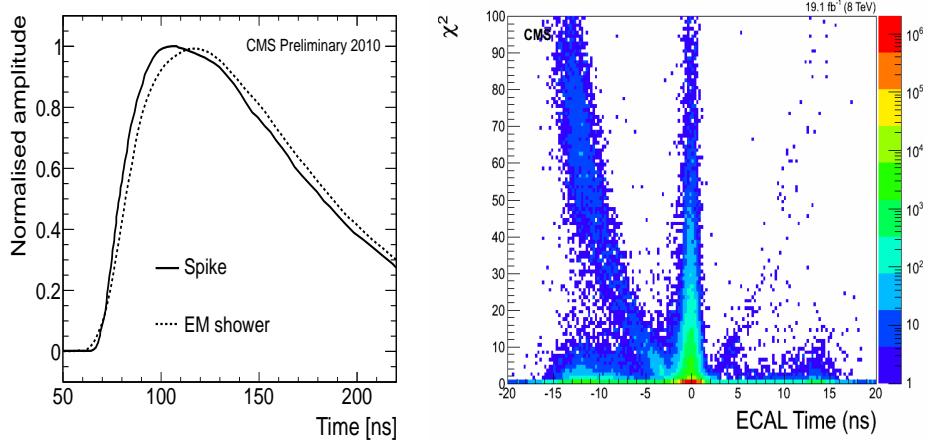


Figure 6.3: Pulse shape profile(left) showing a spike (solid line) and a true EM shower (dashed line) from data. The χ^2 against ECAL time distribution(right) shows spikes misidentified as photons with very large χ^2 and negative ECAL time. The region of $\chi^2 > 4$ is mostly dominated by spikes events.

From the χ^2 Vs ECAL time plot (right plot of Figure 6.3), we observe that most of the photons have ECAL time around zero but also many of the photons with large ECAL time have large normalized χ^2 values which is expected where the time measurements by non-seeded crystals is inconsistent with the time measurement of the seed crystal. We also observed that, a cut in $\chi^2 < 4$, can reduce with 99.2% efficiency, photon contributions from spikes and misidentified jets where a neutron embedded in the jet hitting directly the photo-detectors produces a spike.

Monte Carlo Time

It is challenging to properly simulate ECAL time for MC events so that it captures the conditions of the ECAL detector during data recording. As a result, the time from MC events does not entirely represent the status of ECAL time in data, in terms of mean time and time resolution. We extract the adjustments for MC event time needed to match the mean time and timing resolution of data using a MC γ +jet samples of selected 1

or 2-jets events. Only events with isolated in-time, $t_\gamma < 2$ ns, with $p_T > 80$ GeV are selected. We perform the adjustment by shifting the mean time and smearing the timing resolution by an additional Gaussian convolution on the photon time of MC events so as to match the mean time and time resolution of data. The photon ECAL time from data and MC γ +jet samples both shown in Figure 6.4; where we compare before (left plot) and after (right plot) the adjustment on MC ECAL time was made, agree quite well.

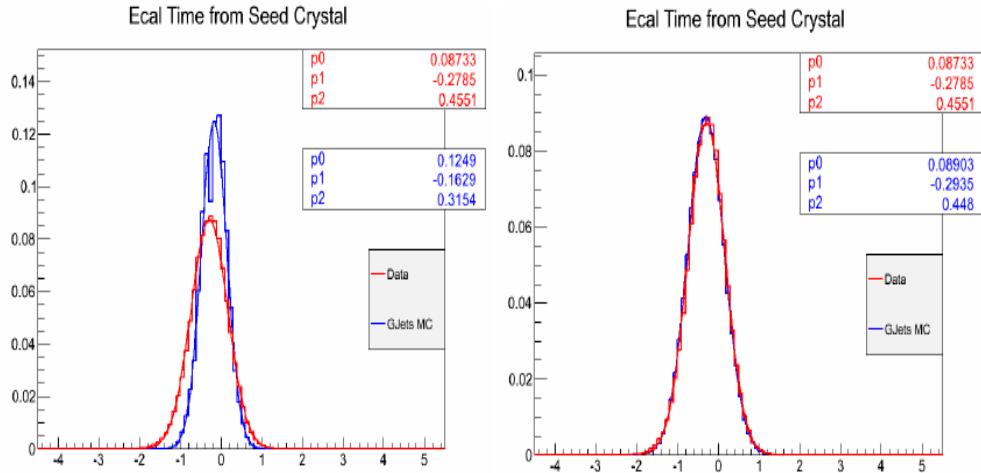


Figure 6.4: ECAL time distributions of in-time photons from MC γ + jets (blue) and data (red) samples before (left) and after (right) we adjusted the photon time from MC.

6.2.2 Lightest Neutralino's Lifetime

The distance traveled, in the CMS detector, by the $\tilde{\chi}_1^0$ before it decays, mentioned previously section 2.3.1, is given as

$$L = \left(c\tau_{\tilde{\chi}_1^0} \right) \cdot (\gamma\beta) = \left(c\tau_{\tilde{\chi}_1^0} \right) \cdot \left(\frac{p}{m_{\tilde{\chi}_1^0}} \right). \quad (6.6)$$

This distance depends on the momentum (p) and mean lifetime ($c\tau_{\tilde{\chi}_1^0}$) of the $\tilde{\chi}_1^0$. Large values of $c\tau_{\tilde{\chi}_1^0}$ means, this distance is large and if larger than the ECAL radius, can be interpreted as the $\tilde{\chi}_1^0$ traveled outside the ECAL volume before it decayed, making detecting the delayed photon no longer possible. On the other hand, small values of

$c\tau_{\tilde{\chi}_1^0} < 10$ cm means the $\tilde{\chi}_1^0$ decayed early and the photon is not late enough such that its detection using ECAL time measurements alone is not reliable. High- p values means the $\tilde{\chi}_1^0$ is boosted and could travel out of the ECAL before it decay, making the photon again undetectable. On the other hand, low- p values means the $\tilde{\chi}_1^0$ is less boosted and traveling slow enough for its decay to happen inside the ECAL volume and the photon delayed arriving at ECAL. In Figure 6.5, we show a distribution of the momentum of the $\tilde{\chi}_1^0$ in the transverse ($x - y$) plane (transverse momentum ($p_T^{\tilde{\chi}_1^0}$)), the $\tilde{\chi}_1^0$ transverse distance traveled before it decay, the transverse momentum of the photon (p_T^γ) and photon's estimated time (T_γ) using only generator level information. These distributions are for different Λ and $c\tau_{\tilde{\chi}_1^0}$ points of the SPS8 model. We observed that, the $p_T^{\tilde{\chi}_1^0}$ increases with increase values of Λ , from $\Lambda = 100$ TeV to 220 TeV, which agrees with our expectation. As Λ increases along with the masses of the gluino/squark and $\tilde{\chi}_1^0$, the $p_T^{\tilde{\chi}_1^0}$ also increases since a massive gluino/squark decay into the $\tilde{\chi}_1^0$. In the same way, increasing the mass of $\tilde{\chi}_1^0$ ($m_{\tilde{\chi}_1^0}$) through increasing Λ , leads to increase in the photon p_T . For a given value of $\Lambda = 180$ TeV, which means $p_T^{\tilde{\chi}_1^0}$ is fixed, the transverse distance traveled by the $\tilde{\chi}_1^0$ before decay (shown in the top right plot of Figure 6.5) and photon time (shown in the bottom right plot of the same figure) increased with increasing value of $\tilde{\chi}_1^0$ mean lifetime, $c\tau = 50$ cm to 600 cm in the same way, qualitatively. These observations support our expectation that the photon is delayed primarily because of the long lifetime of the $\tilde{\chi}_1^0$. However, one can argue that this is not entirely true and we will exploit this further in a detailed study of the source of delayed photons in the next section.

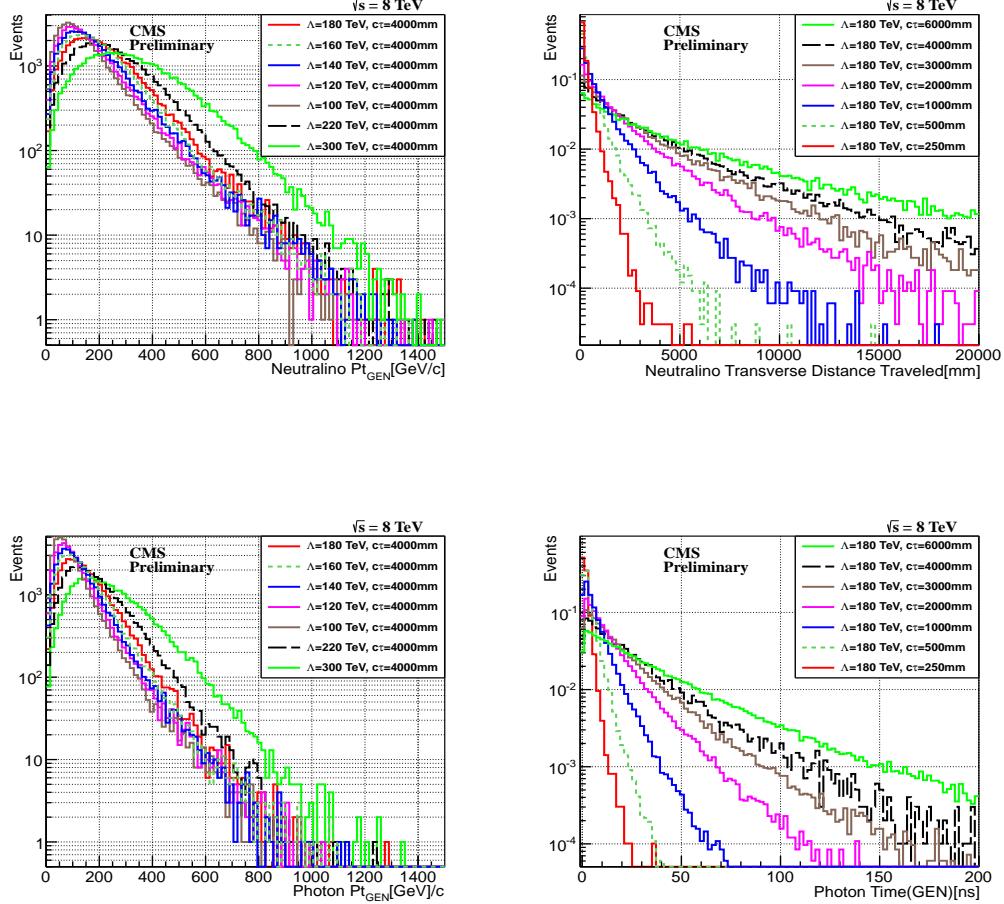


Figure 6.5: Neutralino transverse momentum ($p_T^{\tilde{\chi}_1^0}$) distribution (top left) and transverse distance traveled (top right). Transverse momentum (bottom left) and time (bottom right) of photon from neutralino decay for different Λ and $c\tau$ points in GMSB SPS8 model.

The photon from the decay of $\tilde{\chi}_1^0$ can arrive late at ECAL for either one of the following reasons: first, because the $\tilde{\chi}_1^0$ is traveling slow i.e. with boost, $\beta = \frac{p_{\tilde{\chi}_1^0}}{m_{\tilde{\chi}_1^0}} << 1$, and second, because the $\tilde{\chi}_1^0$ was produced with significant boost such that the photon traveled to ECAL through a non-direct flight path from the nominal pp interaction point. We distinguish between these two sources of delayed photons by estimating the photon arrival time at ECAL in each scenario using the distance traveled by $\tilde{\chi}_1^0$ before decay

and the distance traveled by the photon from the decay point to the point of detection at ECAL. Figure 6.6 (left) is a schematic representation showing how we estimate the photon arrival time at ECAL in each of the possible travel flight path representing the different source of delayed photons. The estimated photon arrival ECAL time in each scenario is given as follows:

- For slow moving neutralinos: $\Delta t_1 = (L1/c\beta) - (L1/c)$
- For non-direct traveled flight path: $\Delta t_2 = (L1 + L2 - L3)/c$
- ECAL measured time = $\Delta t_1 + \Delta t_2$

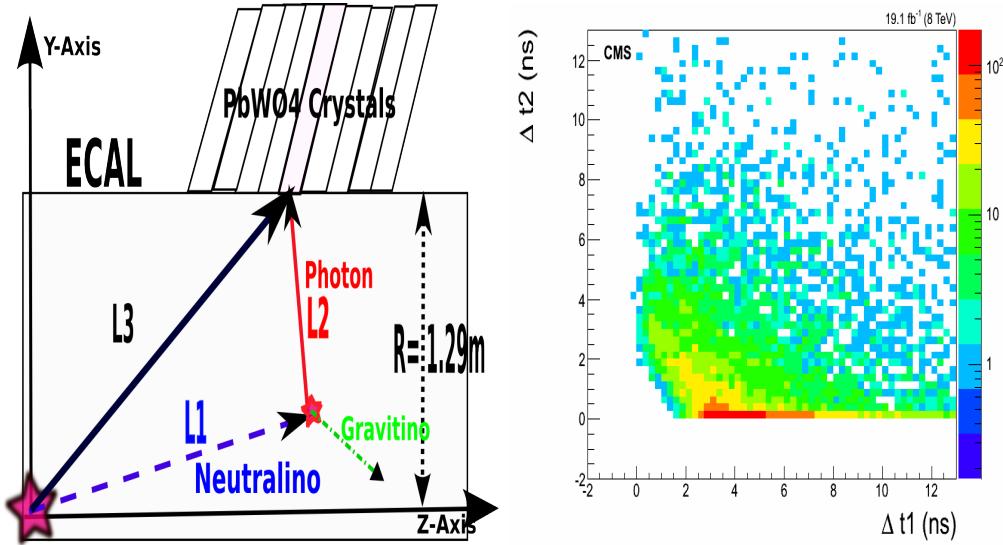


Figure 6.6: Schematic diagram showing $\tilde{\chi}_1^0 \rightarrow \gamma + \tilde{G}$ decay topology within the ECAL volume of the CMS detector (Left). The estimated photon arrival time at ECAL from the decay of $\tilde{\chi}_1^0$ in the SPS8 model with $m_{\tilde{\chi}_1^0} = 256 \text{ GeV}/c^2$ and $c\tau = 600 \text{ cm}$ (Right).

The $\tilde{\chi}_1^0$ is traveling with velocity, $v = c\beta$, where c is the speed of light in vacuum. The distribution of the estimated photon ECAL arrival times Δt_1 and Δt_2 , is plotted shown in Figure 6.6(right), where the color intensity represents the photon population. We observed that most of the late arrival time photons are from the decay of slow moving $\tilde{\chi}_1^0$ compared to those from non-direct flight path to ECAL. This proves that a good number of $\tilde{\chi}_1^0$ produced with low momentum such that the ratio $\frac{p_{\tilde{\chi}_1^0}}{m_{\tilde{\chi}_1^0}} \ll 1$ produced

most of our detectable delayed photons, using ECAL time measurements. While $\tilde{\chi}_1^0$ with very long lifetimes ($c\tau$) produced with high momentum will very likely escape the ECAL without detection unless their decay happens within the ECAL volume such that the delayed photon arrives at the ECAL in a non-direct flight path.

6.2.3 Satellite Bunches

We observed (see Figure 6.7) a 2.5 ns discrete pattern in the photon ECAL time of events, produced in pp collisions, with photon $p_T > 50 \text{ GeV}/c$ and belonging to the barrel and endcaps.

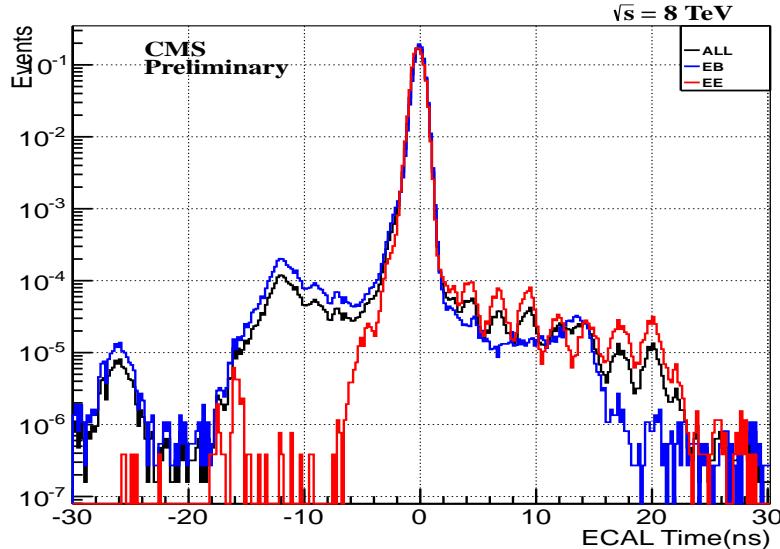


Figure 6.7: ECAL timing distribution of photons in barrel (EB), endcap (EE) and all of ECAL (ALL ECAL) with $p_T > 50 \text{ GeV}$ from data. A 2.5 ns delay timing structure is observed in endcaps.

Most of the photons with this signature belong to the endcap ($1.47 < \eta < 3.0$) with a few in the barrel ($|\eta| < 1.47$). These out-of-time photons are produced from beam halo muons from protons in *Ghost/Satellite* bunches or RF cavity, described in section 3.1.4. We consider these events as background to our delayed photon signal. In addition to developing methods to tag and veto these events, a simple estimate of their contribution can be made using the ratio of the number of protons in the ghost

RF cavities or bunches to the main proton bunches taken from the proton population in the RF cavity of the LHC proton filling profile. From this ratio, we estimate that one in every 100,000 photons observed in ECAL particularly the endcaps, is from beam halo muons produced by ghosts bunches. Their presence is not clearly visible in the barrel as in the endcaps as they will be highly smeared in the barrel because their time is dependent on η which is not the case in the endcaps.

Our observation of this phenomenon together with the sub-par ECAL time resolution in the endcaps compared to the barrel, means, this analysis only uses photons belonging to the barrel.

We estimate, the contribution of these events to our delayed photon background in the signal using a control sample of events with electromagnetic particles of moderate p_T range in the barrel like events with $Z \rightarrow e^+e^-$ decay. The details of this study and the results obtained is discussed in the collision background estimation section of this analysis.

6.2.4 E_T^{miss} Adjustments

During energy clustering in the formation of superclusters used for event reconstruction by the particle flow (PF) algorithm, “out-of-time” energy deposits in ECAL are excluded. The reason is because the algorithm avoids energy deposits from particles not produced from the main proton-proton bunch collisions like cosmic muons and machine induced backgrounds which are normally out-of-time. As a result of the exclusion, the out-of-time photon E_T is not included in the calculation of missing transverse energy or \cancel{E}_T (from now on, we will be using for convenience, \cancel{E}_T instead of E_T^{miss} as the symbol for missing transverse energy) of the event. This exclusion introduces some differences in the calculation of \cancel{E}_T for in-time ($|t_\gamma| < 3.0$ ns) and out-of-time photon events.

Because energy deposits from events with out-of-time photons could be from a possible signal event, and we are searching for events with large arrival time photons, we correct for this exclusion by adding back the out-of-time photon’s E_T vector to the particle flow reconstructed \cancel{E}_T (PF-MET) vector for events with out of time photons. We introduced an additional missing transverse energy variable defined as $\vec{E}_T^\gamma = \cancel{E}_T + \vec{E}_T$, and use it in our final event selection to avoid any bias in our event selection, particularly for events with out-of-time photons.

6.3 Event Selection

Our event selection happens in two stages. The first stage is online at the L1 and HLT trigger levels where only triggered single photon events are selected while the second stage happens offline where our signal-like event selection requirement is applied.

The offline event selection criteria is designed to select signal-like events whose event topology comprise of at least a single photon, multiple jets and large \cancel{E}_T . The multiple jets arise from the cascade decay of gluino or squark to other quarks/gluons in addition to the neutralino. We require multiple jets in our event selection as this helps suppress non-collision background events like cosmic and beam halo muons, and pp collision background events like the ElectroWeak (EWK) background events of $W \rightarrow e + \nu$ and $t\bar{t}$ decays, where the top (t) decays to a b quark and a W boson 100% of the time. These events with W boson can produce fake photons (misidentified electron as photon) and real \cancel{E}_T arising from the presence of the neutrino (ν).

Other collision background events producing fake photons (jet misidentified as photon) like multijets and QCD events can also be suppressed by selecting purely hadronic jets i.e. jets containing a greater portion of hadronic energy fraction from pions and kaons. The single late photon and large \cancel{E}_T is from $\tilde{\chi}_1^0 \rightarrow \gamma + \tilde{G}$ decay. We select only photons in the barrel ($|\eta_\gamma| < 1.479$). This suppresses out-of-time photons from the so-called halo muons from ghost/satellite proton bunches populating mostly the endcaps ($1.5 < \eta_\gamma < 3.0$). We also require the photons to be of high- p_T and this helps suppress misidentified out-of-time photon contribution from EWK events.

The large \cancel{E}_T requirement is useful for suppressing $\gamma +$ jets and QCD events with apparent \cancel{E}_T or *fake* \cancel{E}_T which arise from energy misreconstruction and cracks in the detector. Out-of-time photon contribution from spikes can be reduced by applying electromagnetic shower shape selection on S_{Minor} and S_{Major} at the HLT trigger level.

6.3.1 Trigger Selection

In our online event selection, we select events passing our online higher level trigger (HLT): `HLT_DisplacedPhoton65_CaloIdVL_IsoL_PFMET25`, which is seeded by a level 1 trigger: `HLT_L1SingleEG12`. The HLT trigger was developed primarily to trigger on events with delayed photons with the requirement that the accepted event contains

at least one photon with p_T more than $65\text{ GeV}/c$ and \cancel{E}_T (without any out-of-time energy deposit bias) above 25 GeV . The minor axis of the photon electromagnetic shower must not be spread across many crystals in any direction. This is implemented as $0.1 < S_{Minor} < 0.4$ of the photon. We studied the HLT efficiency and turn-on (efficiency becomes nearly 100%) curve in the photon p_T and event \cancel{E}_T and separately for p_T and \cancel{E}_T using another trigger HLT_Photon50_CaloIdVL_IsoL as the denominator. The HLT_Photon50_CaloIdVL_IsoL selects photon candidates which are very loose isolated with $p_T > 65\text{ GeV}/c$ and must also satisfy our offline photon selection requirement summarized in Table 6.4. The HLT event selection efficiency in p_T is defined as the fraction of offline reconstructed photons to photons triggered by the HLT_Photon50_CaloIdVL_IsoL trigger within $\Delta R < 0.5$ while the efficiency in \cancel{E}_T uses the HLT_Photon50_CaloIdVL_IsoL trigger which has no \cancel{E}_T and jet multiplicity requirement for the denominator and the HLT_DisplacedPhoton65_CaloIdVL_IsoL_PFMET25 for the numerator. The results of the trigger efficiency measurements in photon p_T and event \cancel{E}_T shown in Figure 6.8 indicate that the event selection efficiency is 100% for events with photon $p_T > 80\text{ GeV}/c$ and $\cancel{E}_T > 60\text{ GeV}$.

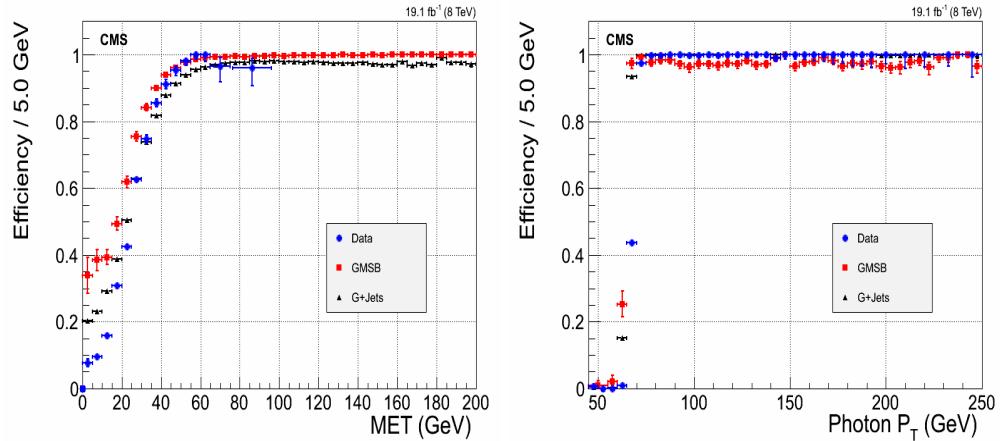


Figure 6.8: Our HLT trigger efficiency turn-on curves in photon p_T (right) and event \cancel{E}_T (left). The $\gamma+\text{jets}$ samples require photon $p_T > 170\text{ GeV}/c$.

6.3.2 Offline Selection

Our offline event selection, applied on HLT triggered single photon events, require that the leading photon have $p_T > 80\text{ GeV}/c$ and if the event has more than one photon, the

sub-leading photon have $p_T > 45 \text{ GeV}/c$. Electromagnetic showers initiated by charged hadrons are rejected by requiring $E^{\text{HCAL}}/E^\gamma < 0.05$, where E^{HCAL} is the sum of the energy in the HCAL towers directly behind the ECAL photon supercluster within a $\Delta R < 0.15$, with $\Delta R = \sqrt{(\eta - \eta^\gamma)^2 + (\phi - \phi^\gamma)^2}$, and E^γ the photon energy in ECAL. Electrons are rejected by requiring the absence of hits in the first two layers of the pixel detector that is consistent with an electron track matching the observed location and energy of the photon candidate (this is known as pixel veto requirement).

The photon candidates must satisfy three *isolation* requirements that reject photons produced in hadronic decays: (1) $Iso_{\text{TRK}} < 0.2 \text{ GeV}$, where Iso_{TRK} is the sum of the p_T of tracks compatible with the primary event vertex in an annulus $0.015 < \Delta R < 0.40$, excluding a strip half η width of 0.015 and the additional inner cone of size 0.04, optimized to exclude reconstructed tracks from $Z \rightarrow e^+e^-$ events, centered around the line of vertex pointing to the photon supercluster. The exclusion is to remove the photon's own energy if it converts into an e^+e^- pair; (2) $Iso_{\text{ECAL}} < 4.5 \text{ GeV}$, where Iso_{ECAL} is the transverse energy deposited in ECAL in an annulus $0.045 < \Delta R < 0.4$ centered around the photon ECAL supercluster, excluding the strip half η width of 0.02 and the additional inner cone of size 0.045 centered around the ECAL supercluster position; (3) $Iso_{\text{HCAL}} < 4.0 \text{ GeV}$, where Iso_{HCAL} is the transverse energy of HCAL tower in an annulus of $0.15 < \Delta R < 0.40$ centered about the ECAL supercluster position. If the photon is very close to a track within a range in $\Delta R(\gamma, \text{track}) < 0.6$ it is rejected. This is to avoid misidentifying charge particles as photons and to prevent double counting jets with high electromagnetic energy component as photons. The photon must also be isolated from any other particle in a cone size of $\Delta R(\gamma, \text{particle}) < 0.4$. The size of the photon electromagnetic shower along the minor axis (S_{Minor}) must be within 0.12 and 0.38 to suppress photons embedded in hadronic jets.

Only photons belonging to the barrel (EB) region i.e. $|\eta_\gamma| < 1.479$, are accepted to avoid many out-of-time photon candidates from ghost/satellite beam halo which belong to the endcap (EE) shown in Figure 6.7.

Topological selection cuts, $1 - E_6/E_2 < 0.98$ and $1 - E_4/E_1 < 0.98$ on the photon energy deposit helps suppress direct interactions in the ECAL APDs by charged particles and neutrons producing anomalous photons known as spikes. Our full photon selection criteria is presented in Table 6.4.

For jets, we select jets with $\eta_{jet} < 2.4$, require that the leading jet in the event has a $p_T > 35 \text{ GeV}/c$ and the number of jets in an event must be at least 1. This helps suppress non-collision background events without jets. The jets are reconstructed using the PF algorithm and identified based on the identification selection criteria summarized in Table 6.5, where the charge electromagnetic fraction (CHF) and the neutral electromagnetic fractions (NEF) must make up a greater portion of the jet sub-structure ($> 99\%$), the neutral energy fraction (NEF) must be less than 99% in order that the jet is not easily misidentified as a photon. A jet near a photon object within a cone of 0.3 is rejected. From our HLT trigger event selection efficiency plot in Missing Transverse Energy (MET) shown in Figure 6.8, we observed that a missing transverse energy of at least 60 GeV for \cancel{E}_T and \cancel{E}_T^γ is enough to suppress $\gamma + \text{jets}$ and QCD events with apparent missing transverse energy.

Photon Selection Criteria	
Criteria	Requirement
Event leading photon must have $p_T(\gamma^1)$	$> 80 \text{ GeV}$
Other photons in event must have $p_T(\gamma^{2,3,\dots})$	$> 45 \text{ GeV}$
$ \eta_\gamma , (\text{Barrel Only}),$	$< 3.0 (< 1.5)$
S_{Minor}	$0.12 \leq S_{Minor} \leq 0.38$
E^{HCAL}/E^γ	< 0.05
$\Delta R(\gamma, track)$	> 0.6
$I_{\text{HCAL}}^{\text{Iso}}, I_{\text{ECAL}}^{\text{Iso}}, I_{\text{TRK}}^{\text{Iso}}$	$< 4.0 \text{ GeV}, < 4.5 \text{ GeV}, < 0.2 \text{ GeV}$
Photon Isolation cone size $\Delta R(\gamma, \text{particle})$	< 0.4
Topological Spike cuts	$1 - E_6/E_2 < 0.98, 1 - E_4/E_1 < 0.98$

Table 6.4: The photon identification and selection criteria used in this analysis

Jet PF identification selection criteria	
Criteria	Requirement
Jet p_T	$> 35 \text{ GeV}$
Number of Jet constituents	> 1
Charge EM energy fraction (CEF)	> 0.99
Neutral Hadron energy fraction (NHF)	< 0.99
Neutral EM energy fraction (NEF)	< 0.99
If $ \eta $ of jet is > 2.4 , Charge Hadron energy fraction (CHF)	> 0
If $ \eta $ of jet is > 2.4 , Charge multiplicity (NCH)	> 0
$\Delta R(\gamma, \text{jet}) = \sqrt{(\phi_\gamma - \phi_{\text{jet}})^2 + (\eta_\gamma - \eta_{\text{jet}})^2}$	> 0.3
E_T , \not{E}_T^γ	$> 60 \text{ GeV}$

Table 6.5: The Jet ID and MET selection used in this analysis

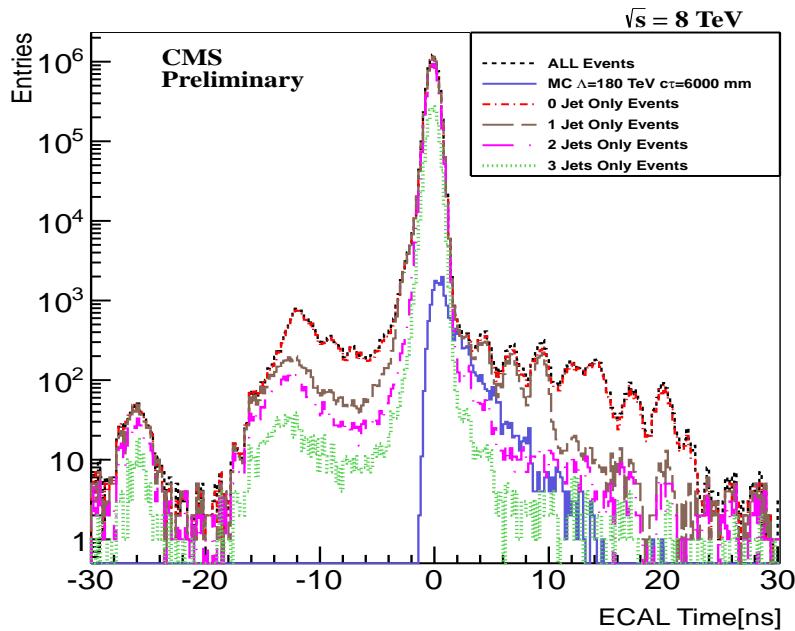


Figure 6.9: Comparing ECAL time distribution of events with different jet multiplicity from small sample of data and a single GMSB $\Lambda = 180 \text{ TeV}$ and $c\tau = 6000 \text{ mm}$ sample. Accepted photons must have $p_T > 60 \text{ GeV}$.

In summary, our signal events are events whose topology comprise of $\geq 1 \gamma + \geq 2 \text{ jets} + \cancel{E}_T > 60 \text{ GeV} + \cancel{E}_T^{\gamma} > 60 \text{ GeV}$. In Figure 6.9, we observed that most of our background events, especially non-collision background events like beam halo, are mostly zero and one jet events. Thus, we use these zero and one jet events as a control sample to do our background study.

6.4 Background Estimation

Most of our background events with out-of-time photons are non-collision events produced from many different sources. In order to qualify and quantify the different sources, we compare in-time ($|t_\gamma| < 2 \text{ ns}$) photon candidates to out-of-time ($t_\gamma < -3 \text{ ns}$ and $t_\gamma > 3 \text{ ns}$) photon candidates. By also comparing photon candidates of events with different number of jets, we were able to uncover the different background sources and better quantify their contribution. In Figure 6.10, we present scatter plots of the photon's ECAL time against η (left) and against ϕ (right) for events with photon $p_T < 60 \text{ GeV}$, belonging to both the barrel and endcap regions and event $\cancel{E}_T > 25 \text{ GeV}$.

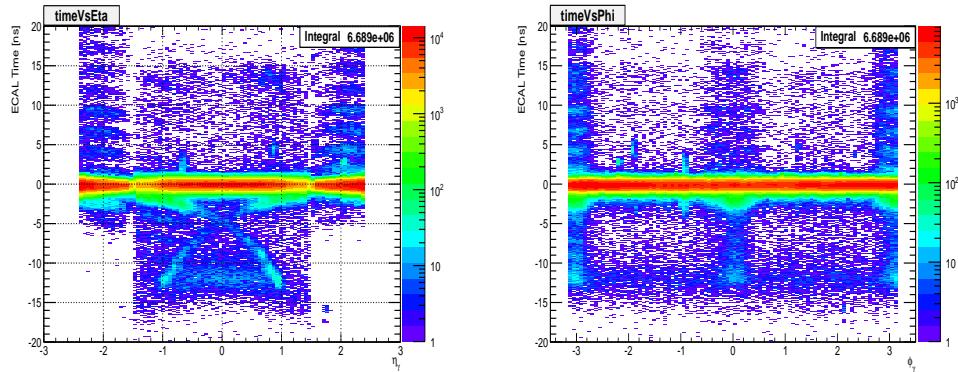


Figure 6.10: ECAL time against η (left) and ECAL time against ϕ (right) for photons with $p_T > 60 \text{ GeV}$ from data.

From these plots, we observe that, although most of the photons arrive with ECAL time about zero, a good number of photons are out-of-time and from very different sources, for example, the cross-like feature seen in the plot on the left, is particular to photons with earlier ECAL time while the step-like pattern in ECAL time extends to

all out-of-time photons but is more prominent to photons which belong to the endcap regions ($1.479 < \eta_\gamma < 3$) and have late ECAL time. The high concentration of out-of-time photons at $\phi = 0 \pm \pi$ on the right plot indicate a unique source of non-collision events which is not random. We conclude that the possible sources of background photon candidates with out-of-time can be split into photons from beam halo muons, because of the cross-like feature and high population at $\pi = 0 \pm \pi$, cosmic muons, due to the random distribution of out-of-time photons throughout ECAL and spikes, due to the high concentration of photons with ECAL time of about -12.5 ns. The step-like feature can be attributed to photons arising from ghost/satellite bunch collisions while a small contribution of out-of-time photons is from QCD and electroweak events. To reduce the contamination of possible signal out-of-time photons from spikes which contribute the most to events with out-of-time photons particularly with ECAL time around $t_\gamma \approx -12$ ns, we restrict our event selection to photons with ECAL time $2.0 < t_\gamma < 13.0$ ns and $-10.0 < t_\gamma < -3$ ns, making our search restricted to signal events with photon ECAL time between $2.0 < t_\gamma < 13.0$ ns.

For our background estimation study, we split the background events into two categories: Collision and Non-Collision events. We first identify and reject photon candidates from beam halo, cosmic muons and spikes and then estimate the residual non-collision and collision background photon candidates using the **ABCD** background estimation technique.

6.4.1 Collision Background

QCD Photons

Events from satellite/ghost proton bunches described in section 3.1.4, produce out-of-time photons which can be present in the barrel. Because these photons are produced from collisions just like QCD events, we refer to these photons as QCD photons. It is challenging to define a strategy for rejecting this background events. Our approach, after rejecting non-collsion events, is to estimate their contribution to signal using the **ABCD** background estimation method. We also perform a separate background estimation using a control sample of Z events, and show that background contributions from events with QCD photon candidates from satellite/ghost is almost negligible.

6.4.2 Non-Collision Background

Halo Photons

Protons in the main and sometimes ghost/satellite bunches can, through inelastic scattering with residual gas molecules like H₂ and CO₂ in beam pipe, produce pions which later decay into muons called *Beam Halo* muons, traveling with energy of about a few TeV. These energetic muons radiate energetic photons called *Halo photons* in the calorimeter through a process called *bremssstrahlung*. It is also possible that beam halo muons can be produced when protons scatter off Tertiary Collimators (TCT), due to proton beam deflection by the magnets, just before entering the CMS detector. The halo muons can travel nearly parallel to the main proton bunch but often stir away from the nominal orbit in the transverse direction spreading mostly in the horizontal plane, with respect to the CMS detector coordinates, due to betatron oscillations. Despite beam cleaning, a sizable population of beam halo muons remains eventually producing energetic photons in the calorimeters. A scatter plot of the photon ECAL time against ϕ shown earlier in the right plot of Figure 6.10, show that most of these beam halo muons enter the ECAL in the horizontal plane at $\phi = 0, \pm\pi$. The rate of halo photons in the general proton *Beam-Induced Background* (BIB) events depend on the beam current and the operational conditions of the LHC like the machine optics, collimator settings, residual gas densities and RF cavity filling scheme.

The halo muons before entry into ECAL, produce hits which can be reconstructed into muon tracks using segments in the Cathode Strip Chambers (CSC) Endcap muon detectors. The reconstructed hits in the CSC segments can be associated with a halo photon supercluster in ECAL within some narrow opening angle in ϕ . The geometry of ECAL allows most of the halo muons in the endcaps but also enough in the barrel. The resulting halo photons are usually out-of-time compared to photons produced directly from nominal pp collisions and most of them have earlier arrival time in cases where the beam halo muons arrive at the ECAL crystals before pp collisions occur at the interaction point. This arrival time can be estimated from the unique flight path of the beam halo muons with respect to the arrival time of photons from pp collisions as

$$t_{\text{ECAL}}^{\text{expected}} = -1/c \left(\pm Z_{\text{cluster}} + \sqrt{Z_{\text{cluster}}^2 + R_{\text{cluster}}^2} \right), \quad (6.7)$$

where Z_{cluster} is the point where the halo muon hits ECAL or longitudinal distance along z -axis of the halo photon supercluster position from the nominal interaction point, R is the radial distance of the supercluster from the beam line which is the radius of ECAL equal to 1.29 m and c is the speed of light in vacuum. The estimated halo muon ECAL arrival time can be re-arrange to become

$$t_{\text{ECAL}}^{\text{expected}} = -\frac{R}{2c} \exp^{(-\eta)} \quad (6.8)$$

showing the explicit dependence on η , for the beam halo muon entry point in ECAL. Comparing this estimated time shown in Figure 6.11 as the two red lines, with observation from data for the 2-dimensional plot of ECAL time *vs* η of photon candidates, show a nice agreement especially for earlier arrival time photon candidates.

By matching halo muon hit positions in CSC segments to photon supercluster positions in the ECAL calorimeter in ϕ , since halo muons spread mostly in the horizontal plane or azimuthal (ϕ) direction, we are able to match halo muons to their corresponding halo photons. We use the quantity, $\Delta\phi(\text{CSC Seg}, \gamma)$, which is the difference in ϕ between the CSC segment and the photon supercluster position in ECAL, to express this matching. A plot of $\Delta\phi(\text{CSC Seg}, \gamma)$ for in-time and out-of-time photons is shown in the left plot of Figure 6.11. We see that out-of-time photons have small $\Delta\phi(\text{CSC Seg}, \gamma)$ confirming that some out-of-time photons are produced by beam halo muons.

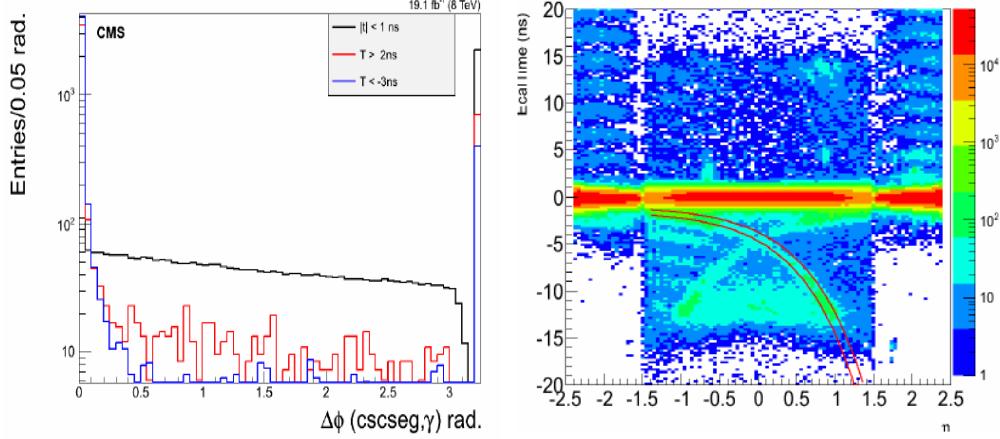


Figure 6.11: (Left) ECAL time $V.s$ $\Delta\phi(\text{CSC Seg}, \gamma)$ for in time(black) and out-of-time(red and blue) photons. (Right)Photon ECAL time $V.s$ η , expected halo photon time is shown as two red lines.

To estimate the performance of using $\Delta\phi(\text{CSC Seg}, \gamma)$ for tagging events with halo photons, we use a halo photon candidate sample of selected events with photons in the endcaps where we expect mostly halo photon candidates and with ϕ_γ around $\phi = 0, \pm\pi$. We were able to tag with $\Delta\phi(\text{CSC Seg}, \gamma) < 0.05$, a good number of photon candidates as halo photons in the endcaps shown in the Figure 6.12 comparing halo photons in the endcaps to tagged halo photons in the endcaps.

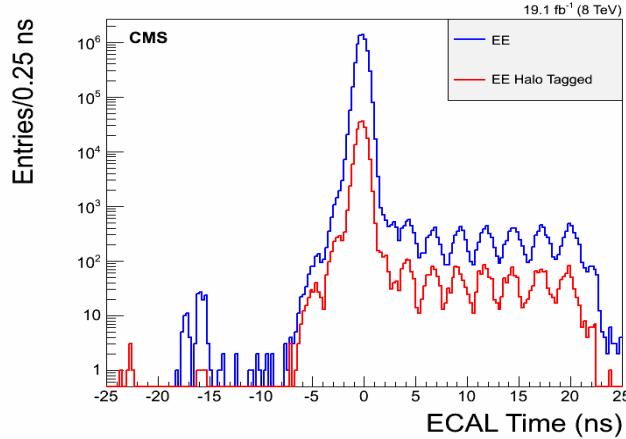


Figure 6.12: Using $\Delta\phi(\text{CSC Seg}, \gamma) < 0.05$ to tag photons with $\phi_\gamma = 0, \pm\pi$ in the endcaps. A good portion of endcap photon candidates are tagged.

Cosmic Photons

Cosmic muons produced in cosmic rays, like beam halo muons, with sufficient energy will radiate (bremsstrahlung) photons in ECAL. We refer to these photons as *Cosmic Photons*. Unlike halo muons, cosmic muons can arrive at ECAL from any direction at any time. We expect cosmic photons in the barrel produced from cosmic muons, to have hits in the Drift Tubes (DT) segments found in the muon barrel behind the calorimeters. Using DT segments and photon supercluster $\eta - \phi$ position in ECAL, we can match cosmic muon hits in DT segments to ECAL photon superclusters within a narrow window in $\Delta\eta$ and $\Delta\phi$. The DT position used in the calculation of $\Delta\eta$ and $\Delta\phi$ is a projection of the muon trajectory using the direction of the DT segment to the outer surface of ECAL because of the large space between the muon barrel and the ECAL. The two dimensional distribution for $\Delta\eta(\text{DT Seg}, \gamma)$ and $\Delta\phi(\text{DT Seg}, \gamma)$ of this matching, for events with out-of-time photons; $t_\gamma > 2$ ns and $t_\gamma < -3$ ns, is shown in the right plot of Figure 6.13. We compared to $\Delta\eta(\text{DT Seg}, \gamma)$ and $\Delta\phi(\text{DT Seg}, \gamma)$ distribution for in-time ($|t_\gamma| < 1$ ns), shown in the left plot, still in the same figure, we observed that most out-of-time photons have a small $\Delta\eta$ and $\Delta\phi$. Comparing the $\Delta\eta$ and $\Delta\phi$ 2-dimensional distributions of these out-of-time photons to photons from a pure cosmic muons sample (data taken when no proton-proton collisions is happening), shown in Figure 6.14, we found a similar small $\Delta\eta$ and $\Delta\phi$ occupancy for the true cosmic muons sample. We conclude that small $\Delta\eta(\text{DT Seg}, \gamma)$ and $\Delta\phi(\text{DT Seg}, \gamma)$ can be used to tag and reject events with cosmic photons.

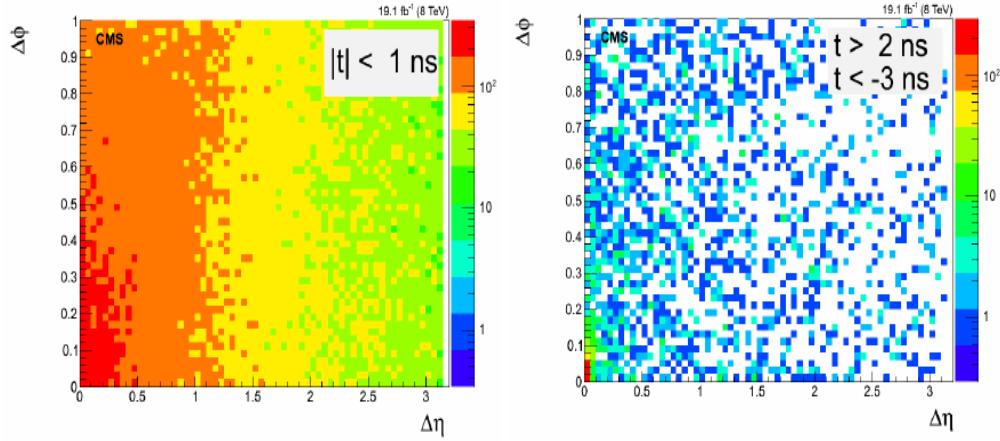


Figure 6.13: Scatter plot showing $\Delta\eta(\text{DT Seg}, \gamma)$ against $\Delta\phi(\text{DT Seg}, \gamma)$ for out-of-time ($t_\gamma > 2$ ns and $t_\gamma < -3$ ns) photons (right) compared to in-time ($|t_\gamma| < 1$ ns) photons (left). Cosmic photon candidates have small $\Delta\eta$ and $\Delta\phi$.

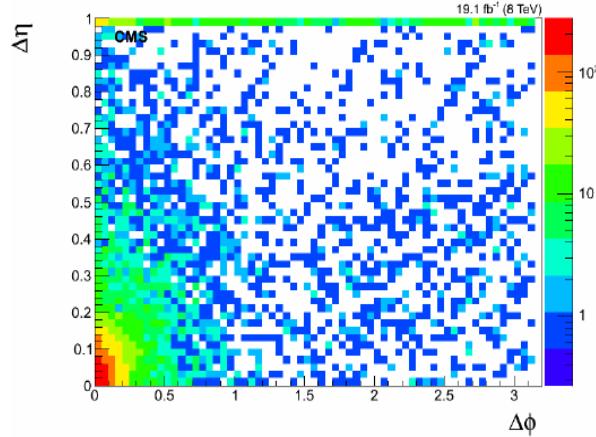


Figure 6.14: Scatter plot of $\Delta\eta(\text{DT Seg}, \gamma)$ against $\Delta\phi(\text{DT Seg}, \gamma)$ for photons from pure cosmic muon data. Small $\Delta\eta$ and $\Delta\phi$ are cosmic photons.

Anomalous Photons: Spikes

Neutrons and charge hadrons can at times deposit their energy directly to the APDs instead of going through the crystal scintillation process. Such signals read from the APDs are *anomalous* and referred to as *spikes*. Spikes produced from pp collisions and satisfying photon selection requirement can easily be identified as photons. A spike supercluster consists of very few crystals; most often one or two crystals. Spikes can also overlap with real photons and found embedded in a real photon supercluster or

buried inside of jets. Such embedded spikes are challenging to identify. By looking at the signal pulse height, the pulse shape for a spike signal appear different from a signal form an electromagnetic particle. Because of this difference in pulse shape profile, the reconstructed ECAL time for spikes usually have large calculated χ^2 values.

The arrival time of spikes is much earlier (negative), usually about $t \approx -12.0$ ns than for photons produced in nominal pp collisions depositing their energy in the APD through crystal scintillation. This is because of the absence of the scintillation process which takes on average 10 ns. The few crystals containing the spike cluster energy deposits make it possible for spikes to be identified using a energy topological selection quantity, $1 - \frac{E_4}{E_1}$, also know as *swiss-cross* variable. By looking at a distribution shown in right plot of Figure 6.15, of the swiss-cross for in-time and for events from a spike sample (events with photon with time $t = -12$ ns), we observed that most spikes have about 98% of their energy deposited in a single crystal.

Using the number of crystals in a photon supercluster and comparing the number of crystals in in-time photons to halo photons and spike candidate photons (photons with $1 - \frac{E_4}{E_1} > 0.98$) shown in the left plot of Figure 6.15, we observed that most spikes including spikes embedded in electromagnetic candidates have less than 7 crystals in their supercluster. A combination of the swiss-cross, number of crystals in supercluster, calculated χ^2 from ECAL time, S_{Major} and S_{Minor} (both S_{Major} and S_{Minor} describe the shape of the electromagnetic shower) is useful for identifying and rejecting events with spikes.

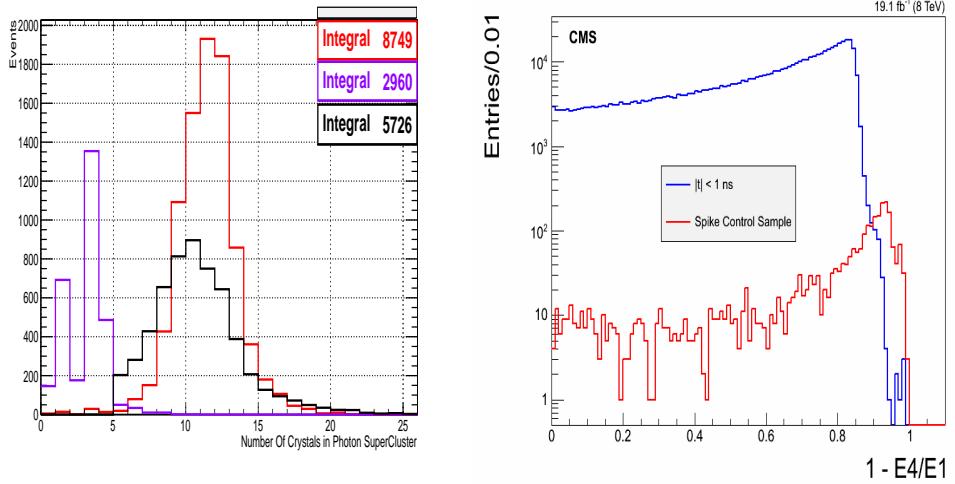


Figure 6.15: *Number of crystals* in photon supercluster plot (Left) for in-time photon candidates (black), spike photon candidates (magenta) and halo photon candidates (red). The spike photon candidates are selected with an energy swiss-cross variable ($1 - E_4/E_1$) (right) shown comparing in-time photons ($|t_\gamma| < 1.0$) to spike candidate sample selected with S_{Minor} and S_{Major} variables.

6.4.3 Event Vetoing, Performance and Fake Rates

Using the quantities studied for tagging events with halo, cosmic and spike photons, we tag and veto non-collision background events as follows:

- An event with a halo photon is tagged and vetoed if a CSC segment for $|\eta| > 1.6$ is found within 0.05 radian of the azimuthal angle with the photon supercluster, i.e. a photon found within $\Delta\phi(\text{CSC Seg}, \gamma) < 0.05$ is vetoed. We found that we are able to veto events with a halo photon with 91% efficiency and 3% mis-tag rate using this photon selection requirement.
- An event with cosmic photon is tagged and vetoed with 75.5% efficiency and 1.4% mis-tag rate if the photon can be matched to a DT segment within $\Delta\eta(\text{DT Seg}, \gamma) < 0.1$ and $\Delta\phi(\text{DT Seg}, \gamma) < 0.1$.
- An event with spike is vetoed if the photon has ECAL time $\chi^2 > 4$, Number of crystals < 7 , $1 - E_4/E_1 > 0.98$, $S_{\text{Major}} < 0.6$ and $S_{\text{Minor}} < 0.17$ with only 0.4% mis-tag rate.

Presented in Table 6.6, is a summary for the mis-tag rates of the different non-collision background sources.

Background Source	Fake Rate(%)
<i>Halo Photons</i>	≈ 3
<i>Cosmic Photons</i>	≈ 1.4
<i>Spikes</i>	≈ 0.4

Table 6.6: Fake rates for different non-collision vetoing.

The result of the event tagging is shown in Figure 6.16. We observe that the majority of non-collision background events are events with a halo photon. Very few late arrival time photons are produced from spikes. There is also some significant contribution from cosmic photons. The most interesting observation is the residual out-of-time background (in red) which could not be tagged.

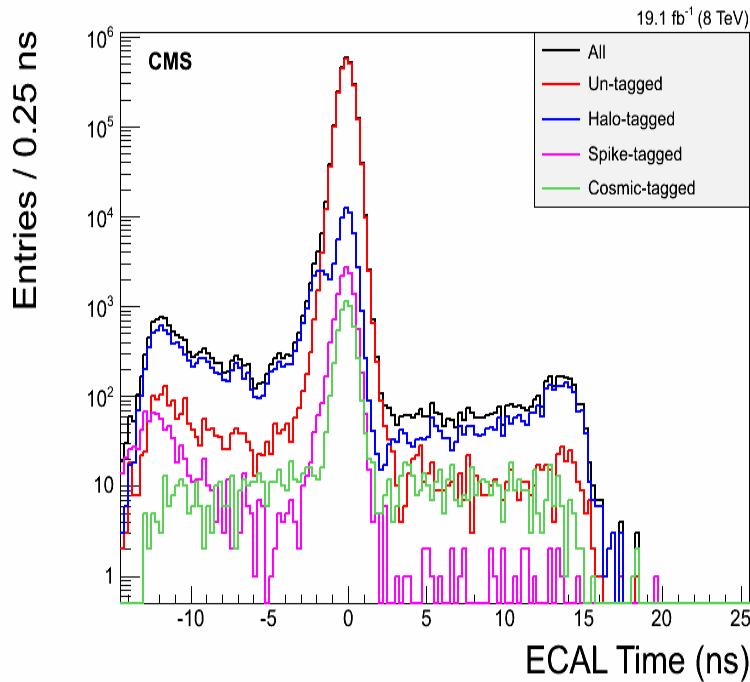


Figure 6.16: ECAL photon time for a data sample of 0 and 1-jet events showing our tagging performance for non-collision background events.

These residual untagged events comprise of some non-collision events, since the tagging and vetoing are not 100% efficient, events with QCD photons from ghost/satellite

collisions and events with $W \rightarrow e\bar{\nu}_e$ decay with the fake photon passing the photon selection requirement. In order to quantify our estimates of these different sources contributing to the untagged background events, we use control samples defined with \cancel{E}_T^γ and \cancel{E}_T .

We expect signal events, because of the undetected gravitino from the neutralino decay, to have large \cancel{E}_T^γ and \cancel{E}_T , by large, we mean above 60 GeV, and events with $W \rightarrow e\bar{\nu}_e$ decay to equally have large \cancel{E}_T^γ and \cancel{E}_T because of the undetected neutrino. The non-collision (cosmic, halo, spike) and collision (ghost/sattelite or QCD) background events can be categorized into high- p_T and low- p_T events. For high- p_T non-collision events, we expect these events to have large \cancel{E}_T , due to the exclusion of the energy deposits from the out-of-time photons in the missing transverse energy reconstruction, and small \cancel{E}_T^γ when the large transverse energy contribution from the energy deposits of the out-of-time photons is re-introduced, while for low- p_T non-collision events, we expect these events to have both \cancel{E}_T^γ and \cancel{E}_T small, since the out-of-time photon transverse energy, which was excluded, was in the first place small. For high- p_T out-of-time collision events which usually have small missing transverse energy, we expect their \cancel{E}_T to be small and after re-introducing the energy deposits from out-of-time high- p_T photons, we expect their \cancel{E}_T^γ to be large, while, using the same argument made for low- p_T non-collision events, we expect low- p_T collision events to have both small \cancel{E}_T^γ and \cancel{E}_T . A summary of our expectations for \cancel{E}_T^γ and \cancel{E}_T for the different expected background events sources with possible contribution to the residual untagged events with out-of-time photon is presented in Table 6.7.

Event Sample	\cancel{E}_T	\cancel{E}_T^γ
Signal Events	Large	Large
$W \rightarrow e\bar{\nu}_e$ Events	Large	Large
High- p_T Non-Collision(Mostly Beam Halo) Events	Large	Small
Low- p_T Non-Collision Events	Small	Small
High- p_T Collision(QCD/Ghost) Events	Small	Large
Low- p_T Collision Events	Small	Small

Table 6.7: Summary of missing transverse expectation for events with out-of-time photons.

Using control samples defined using \cancel{E}_T and \cancel{E}_T^γ for events with in-time ($|t_\gamma| < 2.0$ ns) photons, events with photon time, $t_\gamma > 3.0$ ns and $t_\gamma < -3.0$ ns, where each control sample is defined purposely to enhance the contribution of either collision or non-collision background events to that control sample while simultaneously suppressing contributions from the other, we perform a background estimation in the signal control sample ($t_\gamma > 3.0$ ns, $\cancel{E}_T^\gamma > 60$ GeV and $\cancel{E}_T > 60$ GeV) using the so-called ABCD background estimation method. We verify that the background estimation method is performing as expected using a data sample of zero and one jet events.

6.4.4 Background Estimation with ABDC Method

Since we expect most of our background events to come from non-collision events, we estimate the number of events we expect from non-collision events in the signal control sample using the ABCD technique. We also estimate the possible contamination, we expect from collision events, to the control samples used in the ABCD method.

1. Non-Collision Background Estimation:

To estimate the number of background events from non-collision, we define Control Samples (CS), labeled as ABCD, in photon ECAL time and \cancel{E}_T^γ for selected events with $\cancel{E}_T > 60$ GeV. Events with $\cancel{E}_T > 60$ GeV define a CS where the contribution from collision (QCD) background events is suppressed, since most collision background events have small \cancel{E}_T , as we saw in Table 6.7. The control samples; A and B, shown in Table 6.8, are events with, respectively, $\cancel{E}_T^\gamma < 60$ GeV and $\cancel{E}_T^\gamma > 60$ GeV and photon ECAL time, $-10.0 < t_\gamma < -3.0$ ns, and CS C have events with $\cancel{E}_T^\gamma < 60$ GeV and photon ECAL time, $3.0 < t_\gamma < 13.0$ ns.

Non-Collision	$\cancel{E}_T^\gamma < 60$ GeV	$\cancel{E}_T^\gamma > 60$ GeV
$3.0 < t_\gamma < 13.0$ ns.	C	D
$-10.0 < t_\gamma < -3.0$ ns	A	B

Table 6.8: ABCD Control Samples (CSs) definitions used for estimating non-collision background events in the signal CS D. Events must satisfy the $\cancel{E}_T > 60$ GeV selection requirement.

Using these control samples for the ABCD method, and with the assumption that

$\frac{N_D}{N_C} = \frac{N_B}{N_A}$, the number of non-collision background events with out-of-time photons expected in the signal CS D is given as

$$N_D^{non-col} = \left(\frac{N_B}{N_A} \right) \cdot N_C, \quad (6.9)$$

where N_B , N_A and N_C are the number of events observed in B, A and C control samples in that order and $N_D^{non-col}$ is the number of non-collision background events we expect in the signal CS D.

2. Collision Background Estimation:

We estimate the number of collision background events contaminating CSs defined in Table 6.8, by once again using the ABCD technique with the Control Samples (CSs) used defined in photon ECAL time and \cancel{E}_T for selected events with $\cancel{E}_T^Y > 60 \text{ GeV}$. We are interested in the contamination from collision background events to CS B and D of Table 6.8, since these are the CSs where we expect events with large missing transverse energy (see Table 6.7). Most collisions events have photons with in-time ($|t_\gamma| < 2 \text{ ns}$), therefore, we use CSs defined with in-time photons to estimate the number of possible collision events with out-of-time photons contributing to CSs B and D. The CSs used for the new ABCD method to perform the collision background estimation are labeled as A' B' C' D' and defined as shown in Table 6.9.

Collision	$\cancel{E}_T < 60 \text{ GeV}$	$\cancel{E}_T > 60 \text{ GeV}$
$3.0 < t_\gamma < 13.0 \text{ ns.}$	C'	D'
$-2.0 < t_\gamma < 2.0 \text{ ns}$	I'	I
$-10.0 < t_\gamma < -3.0 \text{ ns}$	A'	B'

Table 6.9: A' B' C' D' and I I' Control Samples (CSs) used for estimating collision background events contamination to CSs B and D. Events here must satisfy $\cancel{E}_T^Y > 60 \text{ GeV}$ selection requirements.

The number of collisions events contributing to the CS B, N_B^{col} , is estimated as

$$N_B^{col} = N_{B'} = \left(\frac{I}{I'} \right) \cdot N_{A'}, \quad (6.10)$$

while the number of events contributing to the CS D, N_D^{col} , is estimated as

$$N_D^{col} = N_{D'} = \left(\frac{I}{I'} \right) \cdot N_{C'}, \quad (6.11)$$

where the general assumption is that $\frac{N_{B'}}{N_{A'}} = \frac{N_I}{N_{I'}}$ and $\frac{N_{D'}}{N_{C'}} = \frac{N_I}{N_{I'}}$, with each N_i being the number of events in each CS $i = A', B', C', D', I, I'$.

3. Combined Background Estimation:

Now that we have estimates for both collision and non-collision event contributions, we can estimate the total number of background events expected in the signal CS D (Events with $E_T > 60$ GeV, $E_T^{\gamma} > 60$ GeV and $3.0 < t_{\gamma} < 13.0$ ns), as

$$N_D^{Total} = \left(\frac{N_B - N_B^{col}}{N_A} \right) \cdot N_C + N_D^{col} = N_D^{non-col} + N_D^{col} \quad (6.12)$$

where $N_D^{Total} = N_D^{non-col} + N_D^{col}$ is the total background events estimated in our signal CS D coming from non-collision and collision background events.

Background Estimation Technique Validation

We performed a validity test of our background estimation method using a data sample of 0 and 1-jet events, where we do not expect any signal events. A statistical agreement in the number of expected background events obtained using our background estimation method and the number of events observed in our signal CS D, will indicate provide a positive affirmation that the method. The accepted 0 and 1-jet events used must pass the same event selection requirements as potential signal events described in Tables 6.4 and 6.5. The event yields in each control sample including the tagging of beam-halo, cosmic and spike events is shown in Table 6.10.

Control Sample	Yield	Beam Halo	Cosmic	Spike
A	852	5075	237	65
B	39	300	17	1
C	359	1508	368	9
D	10	22	30	0
<i>A'</i>	8	1	1	0
<i>C'</i>	2	0	0	0
<i>I</i>	35464	-	-	-
<i>I'</i>	1446522	-	-	-
<i>B'</i>	-	-	-	-
<i>D'</i>	-	-	-	-

Table 6.10: Event yields used in validity test for the ABCD background estimation method using 0 and 1-jet events sample. Beam halo/cosmic/spikes yields are obtained from tagged events. Events must pass photon, jet and E_T^{miss} selection requirements.

Using the event yields in Table 6.10 for each CS, and Equations 6.10, 6.11 and 6.12, we obtain the following estimates for the expected number of events in signal CS D as

$$\begin{aligned} N_B^{col} &= \frac{35464}{1446522} \times 8 = 0.20^{+0.08}_{-0.06} \\ N_D^{col} &= \frac{35464}{1446522} \times 2 = 0.05^{+0.05}_{-0.02} \\ N_D^{Total} &= \left(\frac{39 - 0.20}{852} \times 359 \right) + 0.05 = 16.40^{+3.04}_{-2.63}. \end{aligned}$$

The uncertainty are statistical uncertainties based on the event statistics in each CS. Our expected number of background events in signal CS D is $16.40^{+3.04}_{-2.63}$ which is, within the statistical uncertainties, agreeable with the 10 events we observed satisfying all our event selection requirements and belonging to the CS D. This give us confidence that our method is working and should be used in our background estimation method for potential signal events from data sample.

6.4.5 Background Estimation Cross Check

Another method for estimating background events with out-of-time photons from collision, which we did, is using $Z \rightarrow e^+e^-$ events. Since we expect most of the events with electron candidates from Z decay to be in-time due to the prompt nature of the decay of Z bosons, we use events with electron candidates from a **SingleElectron** and **DoubleElectron** data samples of 2012, where we expect the contributions from non-collision events to be very small, to do the estimation. These events are selected such that out-of-time energy deposits in ECAL from electron candidates are accepted. The background events with out-of-time energy deposits are mostly events produced from the Drell-Yan process, where electron candidates can be matched randomly to produce a di-electron mass close to the mass of Z and poorly reconstructed out-of-time energy deposits in ECAL. To further reduce any out-of-time contribution from beam halo and cosmic events, which can occur simultaneously with true pp collision events, we require that events with Z candidates must satisfy the following event selection requirements: the two electron candidates of the Z boson must each have a $p_T > 30\text{ GeV}/c$, di-electron mass, $|m_{e^+e^-} - 91| > 61\text{ GeV}/cc$, and both electrons must be in the barrel, i.e. $|\eta_{e^-}| < 1.479$ and $|\eta_{e^+}| < 1.479$.

From these selected Z candidate events, we define a signal event sample where the Z candidates have a well defined di-electron mass between $76 < |m_{e^+e^-}| < 100\text{ GeV}/c^2$ and a background or sideband event sample where the di-electron mass is either between $50\text{ GeV}/c^2 < m_{e^+e^-} < 76\text{ GeV}/c^2$ or $100\text{ GeV}/c^2 < m_{e^+e^-} < 130\text{ GeV}/c^2$.

The electron ECAL arrival time is taken to be the seed crystal time adjusted to account for the electron time of flight. The seed crystal must satisfy the recommended crystal (reconstructed hit) cleaning criteria which requires that the seed crystal is not a spike, is not noisy and has been properly time calibrated. A scatter plot of the arrival ECAL time against η (top plots) and ϕ (bottom plots) of the electron candidate of Z is shown in Figure 6.17. A clear difference in the scatter distributions of events is seen comparing events from the **Single/DoubleElectron** data sample (plots on the right) which do not have the familiar beam halo features (the X-shape and increase event population around $\phi = 0, \pm\pi$) to events from **SinglePhoton** data sample (plots on the left). This confirms our expectation that the candidate Z event sample is free from contamination from non-collision events and thus a good sample for estimating

out-of-time event contributions from collisions.

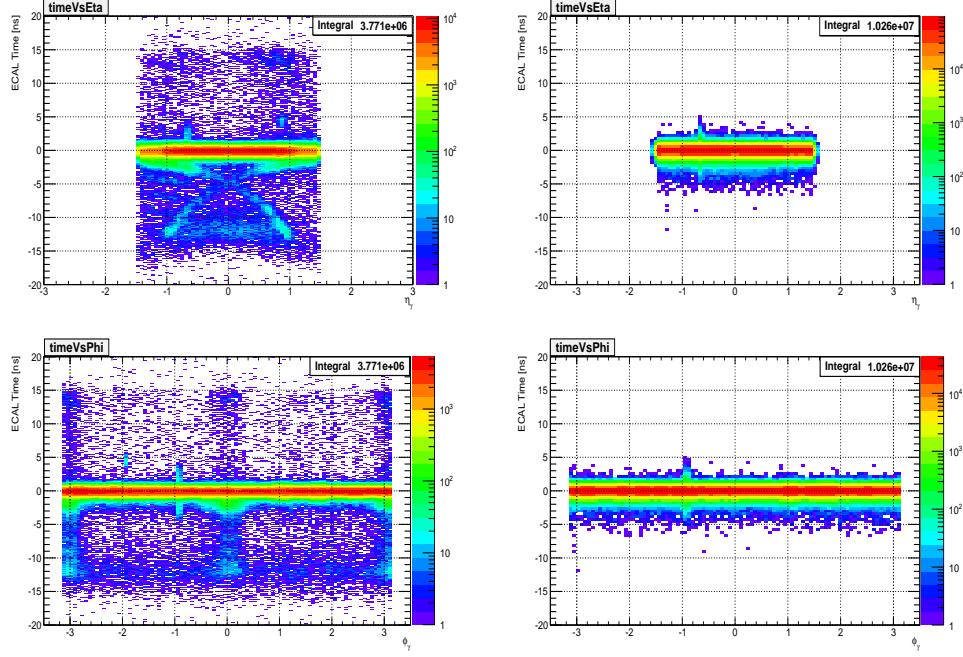


Figure 6.17: ECAL time $Vs \eta$ (top plots) and $Vs \phi$ (bottom plots) for photon candidates from SinglePhoton dataset (left) compared to electron candidates from the DoubleElectron dataset (right). All photons or electron candidates are in barrel subdetector. Most of the photons with $\phi = 0, \pm\pi$ are halo photons which are not observed in the Z boson candidate sample.

The di-electron mass (left plots) and electron arrival time (right plots) of both electron candidates of the Z boson for signal (blue) and sideband (red) events is shown in Figure 6.18. The out-of-time background contribution to the true Z events sample is estimated using a *sideband subtraction method* described as follows:

- We fit the di-electron candidate mass distribution of the sideband control sample with a polynomial function (see top plots of Figure 6.19) to extract the background template.
- Using the background template, we obtain a scale factor as

$$\text{Scale Factor} = \frac{N}{M_1 + M_2}, \quad (6.13)$$

shown in the bottom plot of Figure 6.19, which we used to determine the correct contribution of the background events to the signal event sample.

- After scaling the electrons arrival ECAL time distribution (red distribution on the bottom right plot in Figure 6.18) of the sideband control sample using the extracted scale factor, we subtract the scaled time distribution from the time distribution (blue distribution of right plot in Figure 6.18) of Z electron candidates from the signal sample. The resulting ECAL time distribution, shown in Figure 6.20 is the time distribution from true Z events. This ECAL time distribution is used to estimate the number of events with out-of-time electromagnetic particles from collision.
- The ratio of the total number of events with the electron ECAL time, $t > 3$ ns, to those with electron ECAL time, $|t| < 2$ ns, i.e. $N_{t>3\text{ ns}}/N_{|t|<2.0\text{ ns}}$, gives an estimate of the fraction of out-of-time background events from collision.

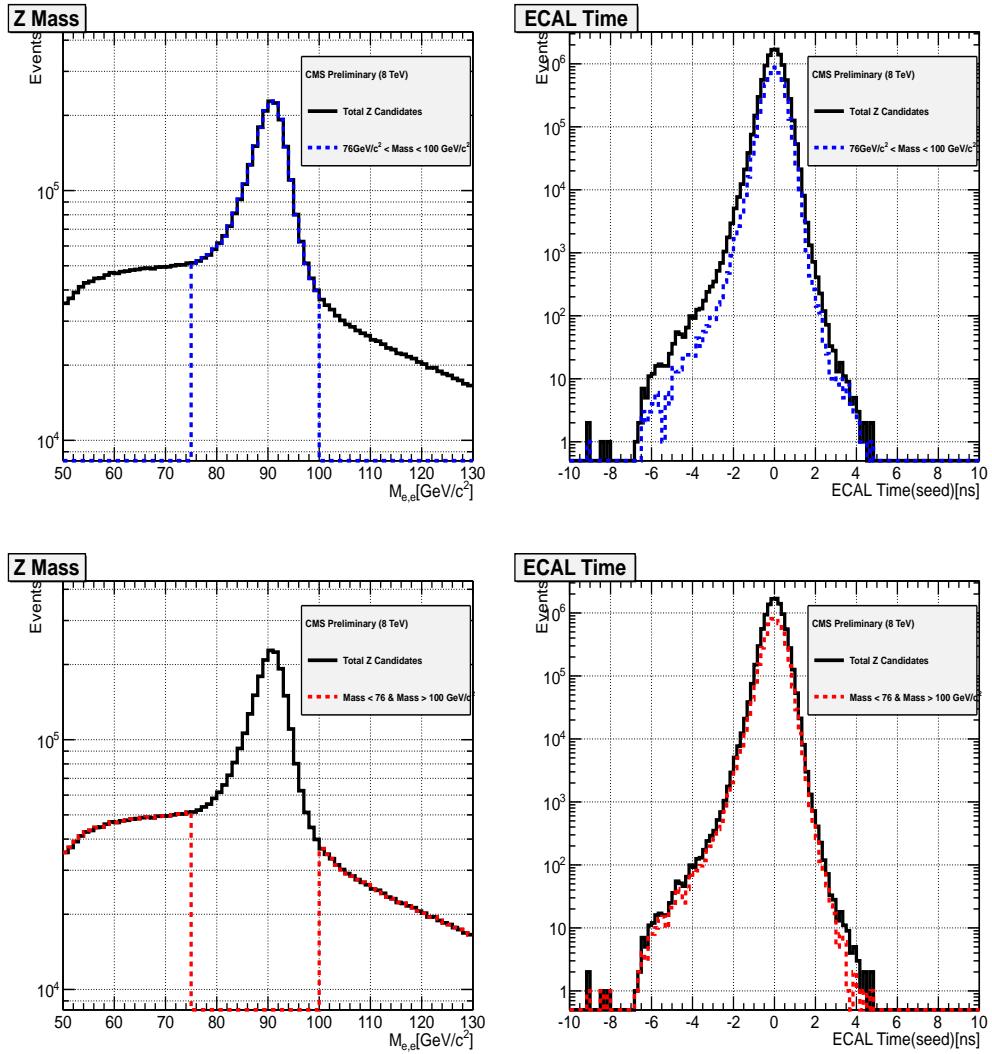


Figure 6.18: Di-electron mass distribution (left) and the time (right) of the two electron candidates for the signal $76 < m_{e^+e^-} < 100 \text{ GeV}/c^2$ Z boson sample and for sideband sample ($50 < m_{e^+e^-} < 76 \text{ GeV}/c^2$ and $100 < m_{e^+e^-} < 130 \text{ GeV}/c^2$). Events are from the Single/DoubleElectron data sample.

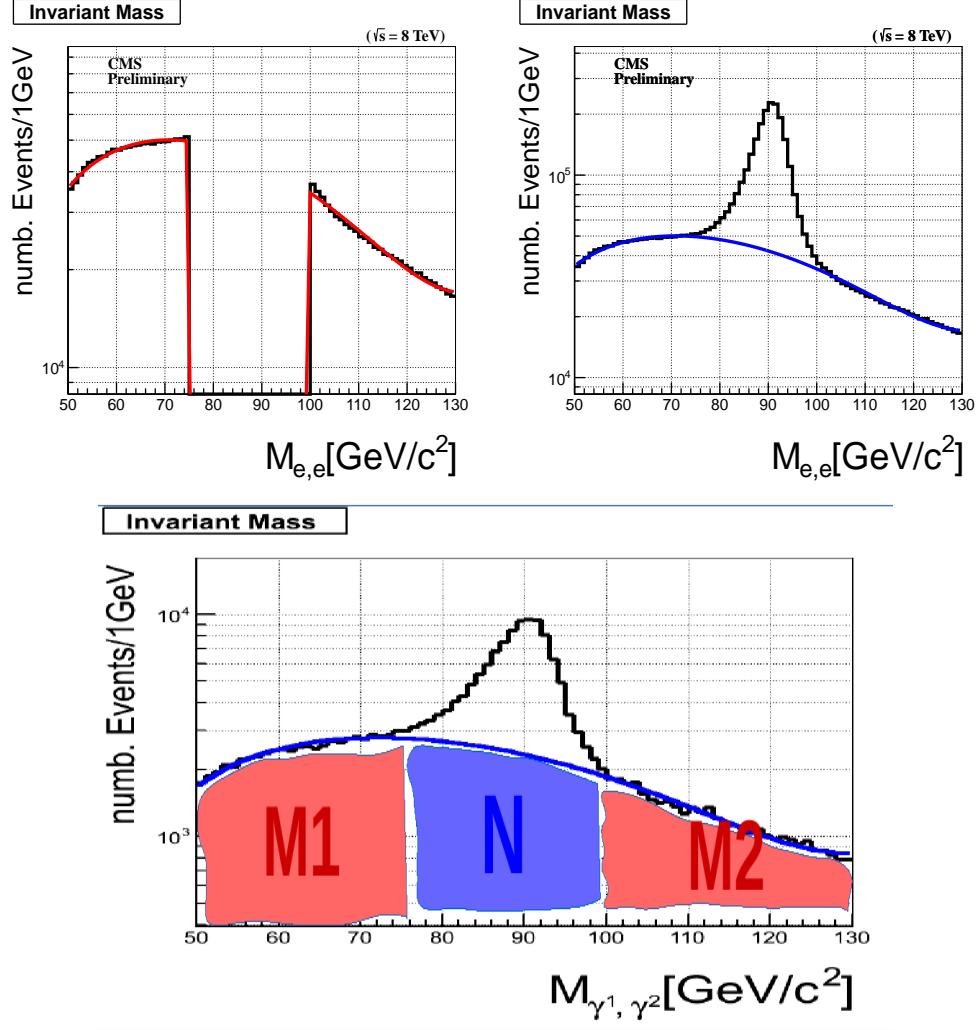


Figure 6.19: *Top:* Polynomial fit (red & blue lines) of the di-electron invariant mass of candidate Z events for the sideband control sample (left) only and all Z events (right). *Bottom:* Fitting the total Z event sample with template extracted from fitting the sideband CS to extract the scale factor use for estimating the background event contribution in true Z signal events sample.

The estimated ratio for out-of-time events to in-time collision events for the final time distribution of signal electron candidates after the normalized sideband subtraction, shown in Figure 6.20 is about, $N_{t>3 \text{ ns}}/N_{|t|<2.0 \text{ ns}} = 1.29^{+0.374}_{-0.325} \times 10^{-5}$.

To obtain an estimate of the number of background events from collisions in the signal control sample D, N_D^{col} , of the $\tilde{\chi}_1^0$ decay, we multiply this estimated ratio to the

total number of in-time events (CS *I*), shown in Table 6.12, passing our final signal event selections requirements, which is 28282. This gives $N_D^{col} = 0.366_{-0.092}^{+0.106}$ events. Comparing this to the value obtained with the ABCD method, which is $\frac{28283}{1446522} = 0.0391_{-0.047}^{+0.039}$ events, we find that the two methods of estimating the number of background events from collision are not exactly equal.

We cannot interpret this inequality as a disagreement between both methods for estimating the number of background events from collision. On the contrary, this non-agreement can be explained. We have not applied any E_T^{miss} selection requirements for the Z candidate events selection. A simple cut in missing transverse energy, $E_T^{\text{miss}} > 60 \text{ GeV}$, could further reduce the estimated ratio, $N_{|t|>3 \text{ ns}}/N_{|t|<2.0 \text{ ns}}$, to a very smaller number than what we got. In fact, the smallness of this ratio confirms our speculation that indeed the contribution from collision events with out-of-time electromagnetic particles to our background events is almost negligible (less than a single event) such that the uncertainties on the ratio of out-of-time to in-time events used in our ABCD collision background estimation is irrelevant. This means, the number $N_D^{col} = 0.366_{-0.092}^{+0.106}$, can be used for our estimated number of background events from collision and our final results will not change by a lot.

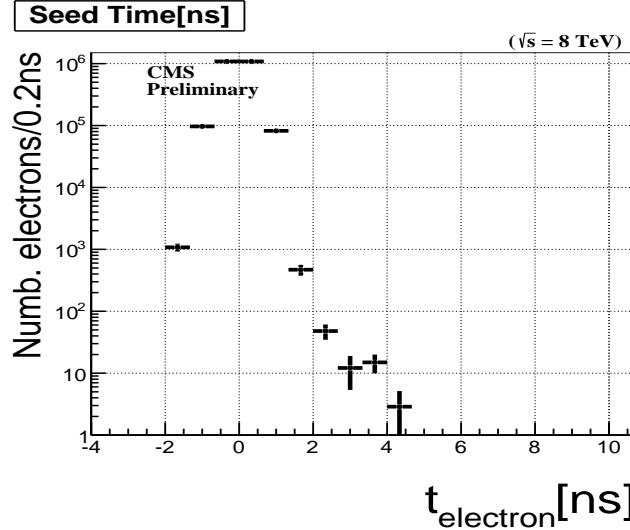


Figure 6.20: ECAL time distribution of genuine Z bosons after background contribution is subtracted.

6.5 Systematic Studies

Our event selection comprise of the following selection cuts: photon p_T greater than $80\text{ GeV}/c$, jet p_T greater than $35\text{ GeV}/c$ and missing energy greater than 60 GeV . The same selection requirements applied to our MC signal event sample should guarantee a good event selection efficiency. Any difference in the photon p_T , jet p_T and missing transverse energy between MC and data, will be a source of systematic uncertainties on the efficiency of selecting signal events. These uncertainties can come from quantities like jet energy scale (JES), jet energy resolution (JER), electron-photon energy scale, instrumentation related and missed or energy deposits not clustered during missing transverse energy reconstruction, photon ECAL arrival time bias and ECAL time resolution. The different sources of uncertainties considered in this analysis is presented in Table 6.11. We obtain these uncertainties by varying the nominal values of each quantity, while keeping the others fixed, by 1σ deviation and counting the number of events passing our event selection requirements. Our largest uncertainty is from ECAL timing bias on the absolute reference time (zero ns) of the ECAL timing used for measuring the photon arrival time. This has the largest impact on our analysis since our analysis is based on counting the number of events with photon ECAL time above 3 ns. The next largest uncertainties are from energy deposits not clustered. This affects the photon energy scale, missing energy scale, jet energy scale and resolution.

The uncertainty on the photon energy scale in the barrel was estimated to be 4.0% based on measuring the photon energy of events with $Z \rightarrow \mu\mu\gamma$ decay, where the muon radiates a photon in a process known as the final-state radiation (FSR)[58].

The uncertainty on the E_T^{miss} resolution uses a conservative estimate from E_T^{miss} measurements in QCD events [59].

The uncertainty on the ECAL time resolution was obtained by comparing the mean time of photons of events from $\gamma +$ jet MC sample to events from data with photon ECAL time, $|t_\gamma| < 2\text{ ns}$. The difference is found to be of the order of 200 ps per event. The systematic uncertainty on luminosity measurement has the recommended value of 2.2% provided by CMS and LHC luminosity measurements while the uncertainty from parton density functions (PDF) is evaluated using the re-weighting technique which uses the Master Equation of CTEQ65 model set described in [60].

Another important source of systematic is in the selection of the background control samples and our background estimation and individual systematic in the tagging and mistagging of non-collision background and estimating background contributions from collisions. We combined these separate uncertainty contributions into a single background estimation uncertainty referred here as our background estimation uncertain. However, since our background estimation is data-driven, most of the uncertainties arising due to systematic can be neglected (less than 1% impact) as they cancel out. This is our largest source of uncertainty in our analysis which we estimate to vary it upward by 223% and downward by 51% as a statistical uncertainty in the ABCD background estimation method. The large background statistical uncertainty is due to very low event statistics. However, despite the large background estimation uncertainty, the signal selection uncertainties are the most significant uncertainties which impact our final results. These uncertainties are used as nuisance parameters in the calculation of the upper limit on observed signal cross-section (σ_{UL}).

Source	Uncertainty(%)
ECAL absolute time (0.0 ns)	< 10.0%
ECAL time resolution (0.5 ns)	< 5.0%
Unclustered energy deposits	< 9.0%
Photon energy scale	< 4.0%
Jet energy scale (JES)	< 9.0%
Jet energy resolution (JER)	< 9.0%
E_T^{miss} resolution	< 2.8%
PDF uncertainty	< 1.70%
Background estimation uncertainty	51.0% to 223%
Luminosity (4.5%)	< 2.2%

Table 6.11: Summary of systematic uncertainties for signal efficiency and background estimation in this analysis and applied to our final results.

6.6 Results

After running our analysis on the SinglePhoton data samples requiring events with at least 2 jets, at least one photon with ECAL time $3.0 < t_\gamma < 13.0$ ns, $\cancel{E}_T > 60$ GeV, and $\cancel{E}_T > 60$ GeV, we observed a single event passing all our event selection requirements as a signal event. This event has one photon and two jets. The photon has a transverse momentum of 224 GeV/c and an ECAL time of 12.17 ns.

Our final expected number of background events estimated is $0.093^{+0.301}_{-0.047}$. This number was computed as

$$\begin{aligned} N_B^{col} &= \frac{28283}{605496} \times 3 = 0.140^{+0.108}_{-0.061} \\ N_D^{col} &= \frac{28283}{605496} \times 2 = 0.093^{+0.093}_{-0.047} \\ N_D^{Total} &= \left(\frac{1 - 0.14}{3} \times 0 \right) + 0.093 = 0.093^{+0.301}_{-0.047}. \end{aligned}$$

using the event yields and background events vetoed presented in Table 6.12.

Control Sample	Yield	Beam Halo	Cosmic	Spike
D	1	0	0	0
C	0	2	1	0
B	1	1	1	1
A	3	6	0	0
<i>A'</i>	2	0	0	0
<i>C'</i>	4	0	0	0
<i>I</i>	28283	-	-	-
<i>I'</i>	669013	-	-	-

Table 6.12: Event yields used in our final background estimation for signal candidate events with at least a single photon and at least 2 jets, $\cancel{E}_T > 60$ GeV, and $\cancel{E}_T > 60$ GeV.

With no significant excess over expected number of background events, we find upper limits on the number of event counts at 95% confidence limit using the statistical CLs limit finding method.

Chapter 7

Statistical Analysis

7.1 Limit Computation

The upper limit calculation procedure used in this analysis is the CLs technique. We fed carefully estimated amounts of background and signal with systematic to obtain the limit. The variable for which the 95% upper limit is set unlike previous experiments is based entirely on the neutralino proper decay length, $c\tau_{\tilde{\chi}_1^0}$.

The method we used in our upper limit calculation is by first performing a Hypothesis test and then use the result of this test to derived our confidence intervals. We do the following:

- We define a NULL hypothesis (H_0) and the Alternate hypothesis (H_1). If we had several other hypothesis, we will defined them also.
- Select a Test statistics ($t(x)$), where x is the data.
- Select a corresponding test statistics calculator.
- Use the result of the hypothesis test to compute the interval by inverting the result of the hypothesis test.

First, we describe the acceptable technique in experimental high energy physics for computing *p-values* used in any search and discovery experiment.

7.1.1 CLs Technique

The CL_s technique [61] is attributed as the standard technique or framework for computing the confidence or exclusion intervals in a search and discovery experiment. It has been shown to work during the search for the Higgs boson at LEP and recently in the discovery of the scalar boson in 2012, by both CMS and ATLAS experiments with the mass of this boson being: $m_H = 125.36 \pm 0.37(\text{stat.Unc}) \pm 0.18(\text{syst.Unc})$.

This method has been implemented in a unique statistical software package called *HiggsCombine* with the goal of providing direct access to a variety of robust statistical methods with optimised performance for computing limits or confidence intervals. *HiggsCombine* [63] is the official standard tool recommended by the CMS statistical committee and CMS Higgs group for calculating limits in any CMS search and discovery analysis. It takes as input estimates on the number or distribution of signal and background and the observed number or distribution from data and produces an upper limit in the production cross section of a given physics process for a given value of a parameter of interest (POI). *Higgscombine* tool has the advantage in that, it allows for the possibility to use several different statistical methods of calculating the upper limit. This way, one can make comparison and simple checks for any inconsistency. In this analysis, we used an Asymptotic [64] and HybridNew (a hybrid of Frequentist and Bayesian methods),[63], to calculate our observed upper limits. The purpose of the using the CL_s method is to compute reliable upper limits in a search scenario when the observed signal is very small compared to the background. In the CL_s technique, one uses not the p-value (CL_{s+b}) but rather divide this by CL_b (which is 1 minus the p-value for background only hypothesis). The reason for this is to define a conditional probability conditioned to the scenario of observing only background or background only hypothesis. The CL_s is formally defined as:

$$CL_s = \frac{CL_{s+b}}{CL_b} = \frac{p_{s+b}}{1 - p_b} \quad (7.1)$$

where $s + b$ means signal and background.

7.1.2 Statistical Test Formalism

The Neyman-Pearson Theorem states that the likelihood ratio gives the most powerful hypothesis test. Therefore, we construct our test statistics t_μ as a function of the observed data, as a likelihood ratio. In a search analysis, one defines the null hypothesis H_0 describing only known processes, or the background only which is to be tested against an alternate hypothesis H_1 defined as a background and signal. However in the computation of upper limits:

- H_0 being the NULL hypothesis includes the background and signal ($s + b$) while
- H_1 being the ALTERNATE hypothesis includes only the background (b).

Using these, two hypothesis we quantify the level of agreement between our observed data with either of the hypothesis by computing a p -value (p -value if the probability under the assumption of a given hypothesis, of finding data of equal or greater incompatibility with the predictions of the given hypothesis). A given hypothesis is then regarded as being excluded if its p -value is observed below a given threshold. In particle physics, this threshold value for the p -value is 0.05 corresponding to 95% of confidence level (CL). The CMS accepted method of computing upper upper limit is based on mix of frequentist-hybrid significance test using the profilelikelihood ratio as a test statistics (HybridNew method). The parameter of interests in in our case the rate (cross section) of signal process as well as *nuissance parameters* as systematics for the background and signal models. This parametrized systematics effects results, as is always the case, to loss in sensitivity.

In this search experiment, for each event in the signal, we measured the timing of the photon as our observable. We use this value to construct a histogram $\mathbf{n} = (n_1, \dots, n_N)$. The expectation value for each value of n_i can be written as:

$$E[n_i] = \mu s_i + b_i \quad (7.2)$$

where μ is the parameter which determines the signal strength, when $\mu = 0$ means background-only and when $\mu = 1$ then we have the signal and background hypothesis.

The the mean number of entries in the i th bin from signal and background are given as:

$$s_i = s_{tot} \int_{bin,i} f_s(t; \theta_s) \quad b_i = b_{tot} \int_{bin,i} f_b(t; \theta_b) \quad (7.3)$$

with the functions $f_s(t; \theta_s)$ and $f_b(t; \theta_b)$ being the probability density functions (Pdfs) of the variable t for the signal and background events and θ_s and θ_b representing the parameters which characterise the shapes of the pdfs. s_{tot} and b_{tot} represents the total mean numbers of signals and backgrounds while the integrals represent the probabilities for an event to be found in bin i . $\theta = (\theta_s, \theta_b, b_{tot})$ denote all nuisance parameters (systematic uncertainties) while s_{tot} is the signal normalization is fixed to the value predicted by the nominal signal model.

The likelihood function is the product of the Poisson probabilities for all bins:

$$\mathcal{L}(\mu, \theta) = \prod_{r=1}^N \frac{(\mu s_r + b_r)^{n_r}}{n_r!} e^{-(\mu s_r + b_r)} \cdot \mathcal{G}(\theta) \quad (7.4)$$

where $\mathcal{G}(\theta)$ is a discrete (Poisson) distribution of the nuisance parameters. This distribution can be different for different nuisance parameter.

Using the likelihood function, the profilelikelihood ratio is then defined as:

$$\lambda(\mu) = \frac{\mathcal{L}(\mu, \hat{\theta})}{\mathcal{L}(\hat{\mu}, \hat{\theta})} \quad (7.5)$$

Here $\hat{\theta}$ is the the value of θ that maximizes \mathcal{L} for a specified μ , thus, it is referred to as the *conditional maximum-likelihood estimator* (CMLE) of θ (given as a function of μ). While $\mathcal{L}(\hat{\mu}, \hat{\theta})$ is the maximized (unconditional) likelihood function with $\hat{\mu}$ and $\hat{\theta}$ being its *maximum likelihood* (ML) estimators. The nuisance parameters broadens the profilelikelihood as a function of μ relative to what is expected if their values where fixed and this reflects in the loss of sensitivity or information about μ due to the systematic uncertainties.

7.1.3 Test Statistics and p -values

The above expression for $\lambda(\mu)$ shows that $0 \leq \lambda \leq 1$, where λ close to 1 indicates a very good agreement between the data and hypothesis value of μ . The test statistics to be used for our statistical test is defined as:

$$t_\mu = -2 \ln \lambda(\mu) \quad (7.6)$$

It is important to note that the test statistics approach to any statistical test is favourable because just by looking at the values of the test statistics, higher values corresponds to increasing incompatibility between the data and the value of μ which is from the signal hypothesis. This incompatibility or disagreement between the data and a given hypothesis is quantified by calculating the probability or p -value as:

$$CL_{s+b} = p_u = \int_{t_{\mu,obs}}^{\infty} f(t_\mu|\mu) dt_\mu \quad (7.7)$$

where, $t_{\mu,obs}$ is the value of the test statistics t_μ obtained from the data and $f(t_\mu|\mu)$ is a pdf constructed from t_μ depending on the signal strength μ . The set of values for μ that are rejected because their p -value is below a specified threshold value α lying on either sides of those not rejected gives a two sided confidence interval of μ and if just on one side of the ones not rejected gives an upper limit on the rejected values of μ .

In the background only scenario i.e $\mu = 0$, the test statistics is defined as:

$$q_\mu = \begin{cases} -2 \ln \lambda(0), & \hat{\mu} \geq 0 \\ 0, & \hat{\mu} \leq 0 \end{cases}$$

where $\lambda(0)$ is the profilelikelihood ratio for $\mu = 0$ defined in 7.5. and again to quantify the disagreement between the background-only hypothesis ($\mu = 0$) and the data is given by the p -value as:

$$CL_b = p_0 = \int_{q_{0,obs}}^{\infty} f(q_0|0) dq_0 \quad (7.8)$$

where $f(q_0|0)$ denotes the pdf if the test statistics q_0 under the background-only ($\mu = 0$) hypothesis. Figure 7.1 shows a sampling distributions of the test statistics and how the

p-values can be extracted from these distributions.

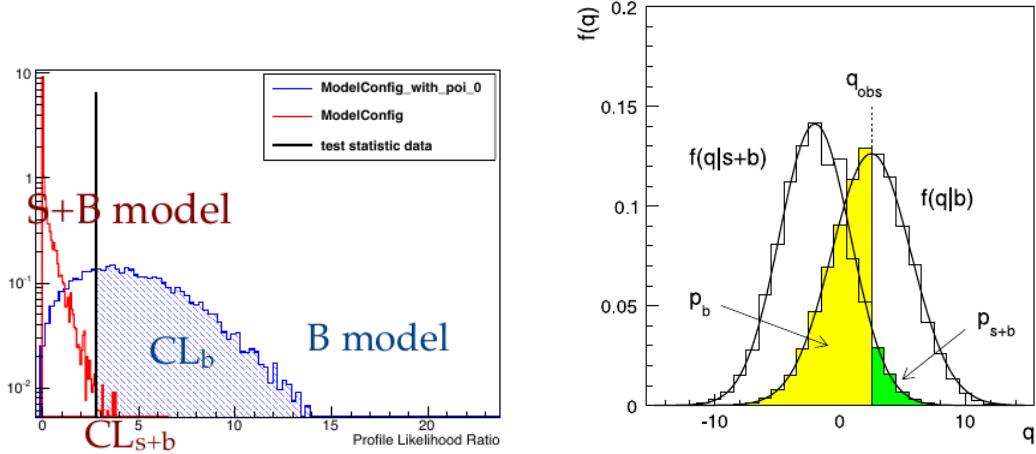


Figure 7.1: Sampling distributions for $f(t_\mu|\mu)$ showing how one extracts the p -values. left: is the using a analytic of the Asymptotic method and right: is from the HybridNew method.

In addition to the p -value, for expressing the disagreement between the data and a given hypothesis, the Higgscombine tool also provides a quantity known as the *significance* (\mathcal{Z}). \mathcal{Z} and the p -value have a very non-linear relation. Once can defined that relation using a two-sided fluctuation if a Gaussian variable σ , with 5σ significance corresponding to a p -value of $p = 5.7 \times 10^{-7}$ to denote a discovery. Since, we have not observed any significant excess of events over our standard model background, we will not mention a lot about significance in this thesis, but rather talk about p -values as they are indispensable in computing limits.

The important question is always, how does one obtain an expression or a distribution of the test statistics and $f(t_\mu|\mu)$ from the likelihood function? To answer this question, the HiggsCombine tool was developed which consist of various ways of both analytically (e.g the Asymptotic statistical method [64]) or through numerical integration or Monte Carlo computation (e.g the HybridNew statistical method) obtain the test statistics and $f(t_\mu|\mu)$. We have shown the limit computation results of both methods as used in this analysis. As an example, the pdf $f(q_\mu|\mu)$ of the test statistics (q_μ) obtained though the **Asymptotic** statistical method as given in [64] is:

$$f(t_\mu|\mu') = \Phi\left(\frac{\mu - \mu'}{\sigma}\right)\delta(t_\mu) + \frac{1}{2}\frac{1}{\sqrt{2\pi}}\frac{1}{t_\mu} \exp\left[-\frac{1}{2}\left(\sqrt{t_\mu} - \frac{\mu - \mu'}{\sigma}\right)^2\right] \quad (7.9)$$

where result to a half-chi-square distribution when $\mu = \mu'$.

In subtle point worth mentioning is that in the HybridNew approach, systematics uncertainties are taken into account through the Bayesian prior density $\pi(\theta)$, and the distribution of the test statistics is computed under the assumption if the Bayesian model of average given as:

$$f(q) = \int f(q|\theta)\pi(\theta)d\theta$$

and the prior pdf $\pi(\theta)$ is obtained from some measurements characterised by a given likelihood function $\mathcal{L}_\theta(\theta)$ which is then used to find the prior using Bayes' Theorem. Unlike other cases where systematic uncertainties are taking as being part of the data and incorporated directly through $\mathcal{G}(\theta)$ as shown in equation 7.4. Nevertheless, they arrive at the same result.

In summary, the hypothesis test is performed using a given statistical method on each value of a chosen parameter of interest (POI)(usually denoted μ). The p -value if obtained from the sampling distribution of the test statistics being used. Can either obtain this test statistics analytically or through Monte Carlo computation and numerical integration. By plotting the p-value as a function of the POI, we obtain the p-value curve (in this case the $CL_s = \frac{CL_{s+b}}{CL_b}$). The value of μ which has a p-value α (e.g 0.05) is the upper limit (for 1-dimensional limits,2-dimensional limits gives lower and upper limits) of $1 - \alpha$ confidence interval (e.g 95%).

combined result

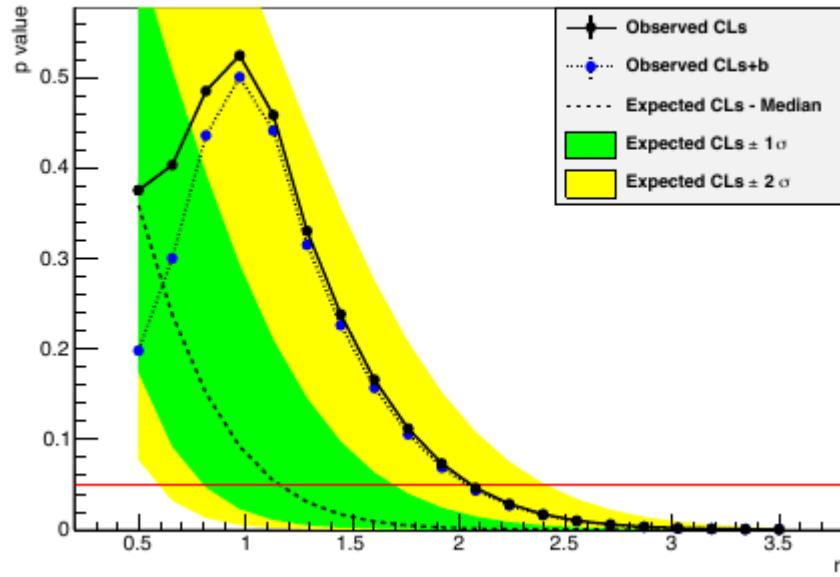


Figure 7.2: Distribution of p -values showing how upper limit on μ is extracted for a given threshold probability.

Chapter 8

Limit Interpretation

Using the CL_s technique, the HiggsCombine tool produces an upper limit along with the expected limit at different quantiles as the signal strength computed which is a ratio of Number of Signal events over the Number of Expected signal events i.e

$$r = \frac{N^{Obs}}{N_{expect}} \quad (8.1)$$

and using the equation as given in chapter 3 on the cross-section $\sigma = \frac{N}{\varepsilon \cdot A \cdot \mathcal{L}}$ and hence the observed cross-section upper limit is given as:

$$\sigma_{UL}^{Obs} = \frac{r \cdot N^{expect}}{\varepsilon \cdot A \cdot \mathcal{L}} \quad (8.2)$$

where \mathcal{L} is the integrated luminosity (19 fb^{-1}) and ε and A are the signal selection efficiency and Acceptance respectively. In addition to the observed limits (Solid black line), the uncertainties on the expected limits at 68%/16% ($\pm 1\sigma$) and at 98%/2.5% ($\pm 2\sigma$) provide the **GREEN** and **YELLOW** respectively, the error from the median (50%) expected limits (dashed red line) shown in figure 8.2.

8.1 Signal Efficiency and Acceptance

After our final selection on MC signal sample, the number of signal events passing our selection requirements for different neutralino lifetime and same mass, $\Lambda = 180 \text{ TeV}$, is

given in table 8.1. We use these numbers and those for other neutralino lifetimes and masses to compute our exclusion limits.

SPS8 GMSB Signal	Number of Events
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 250$ mm)	0.2096
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 500$ mm)	4.5423
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 1000$ mm)	6.3646
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 2000$ mm)	6.3968
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 4000$ mm)	6.1442
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 6000$ mm)	4.6498
GMSB(SPS8) ($\Lambda = 180$ TeV, $c\tau = 12000$ mm)	2.918

Table 8.1: Final number for $\Lambda = 180$ TeV GMSB SPS8 MC signal events events passing our selection cuts.

The efficiency times acceptance ($\varepsilon \times A$) combined as one is seen the figure 8.1.

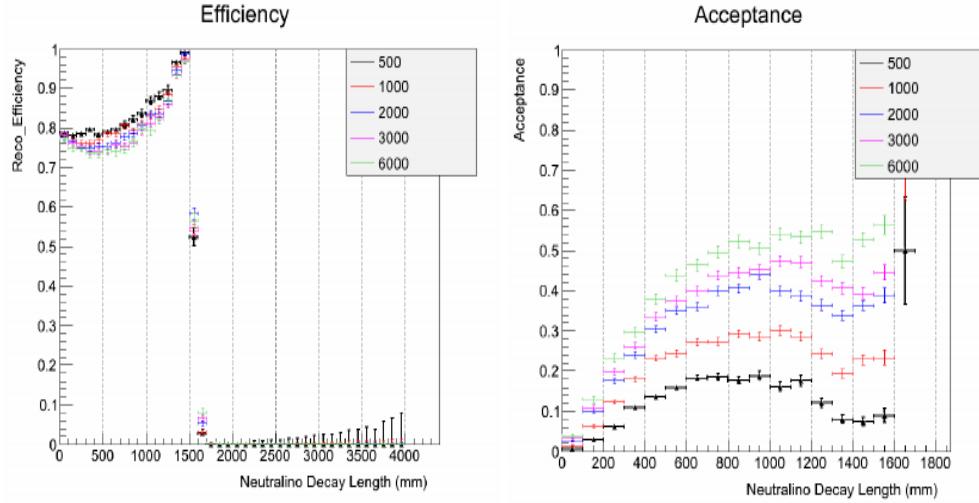


Figure 8.1: The reconstruction and selection efficiency (left) \times Acceptance ($t > 3$ ns) (right) against transverse decay length in laboratory frame for different $c\tau$ points.

The $\frac{N_{expect}}{\mathcal{L}}$ defines the expected signal cross section which is obtained from a given signal model. In our scenario, our choice of signal model we want to produce exclusion limits on the possible production and decay of a long-lived particle described by this signal model is GMSB. Thus the interpretation of our search analysis is given within the context of any GMSB model with a long-live neutral particle decaying to a photon and gravitino. Such a model is the minimal GMSB or the SPS8 model and the general GMSB model. However, the results provided are based on interpretation within the context of the SPS8 model. In GMSB, the neutralino $\tilde{\chi}_1^0$ is the NLSP and decays to the gravitino \tilde{G} the LSP (as a result of R-parity conservation) in association with a very energetic photon γ . Because of the smallness in mass difference between the $\tilde{\chi}_1^0$ and the \tilde{G} as well as the coupling, the $\tilde{\chi}_1^0$ decay to \tilde{G} is delayed and as a result, the photon emitted can arrive late in the calorimeter crystals. Measuring the arrival time of the photon on ECAL crystals, we can extract important parameters of theory of GMSB.

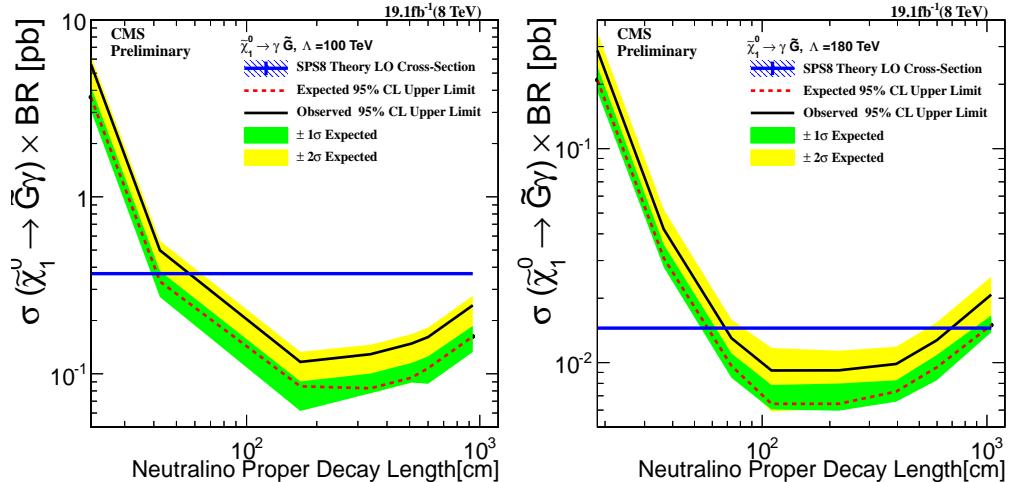


Figure 8.2: Neutralino production cross section against proper delay length upper limit interpretation in SPS8 model. (Left) $\Lambda = 100$ TeV, (Right)(left) $\Lambda = 180$ TeV

In the SPS8 model, the parameter space for long-live neutralinos is governed by $\Lambda_m - c\tau$ 2-dimensional parameter space. For each Λ_m point, we have a fixed neutralino mass with different proper lifetimes $c\tau$. We have obtained limits for Λ_m ranging from 100 TeV

to 180 TeV corresponding to lightest neutralino mass $m_{\tilde{\chi}_1^0}$ between 90 to 255 GeV/c^2 and proper lifetime $c\tau$ ranging from 250 to 12000 mm corresponding to $\tau_{\tilde{\chi}_1^0}$ from 0.8 ns to 40 ns.

For a given value of $\Lambda_m = 180$ TeV, we have a lightest neutralino production cross section times branching ratio plot shown in figure 8.2, showing that the ECAL detector is sensitive to lightest neutralinos of mass $m_{\tilde{\chi}_1^0} = 255$ GeV/c^2 and life time upto 30 ns and we are 95% confident that we have not missed any neutralino whose mass is $m_{\tilde{\chi}_1^0} = 255$ GeV/c^2 and lifetime is $\tau \leq 30$ ns.

For a given lifetime of $\tau = 20$ ns, we can also obtain upper limits on the production cross section times branching ratio when compared against their theoretically expected values for a lightest neutralino with mass ranging from $m_{\tilde{\chi}_1^0} = 90$ GeV/c^2 to $m_{\tilde{\chi}_1^0} = 255$ GeV/c^2 . The observed upper limit on this cross section is $\sigma_{\tilde{\chi}_1^0}^{UP} \geq XX$ pb with proper lifetime of $\tau = 30$ ns.

Using both the mass and proper lifetime of the lightest neutralino, we present possible 2-dimensional limits simultaneously on $m_{\tilde{\chi}_1^0}$ or Λ_m and $c\tau$ or τ in the SPS8 model, comparing this with the result of previous experiments. This is shown in Figure 8.3.

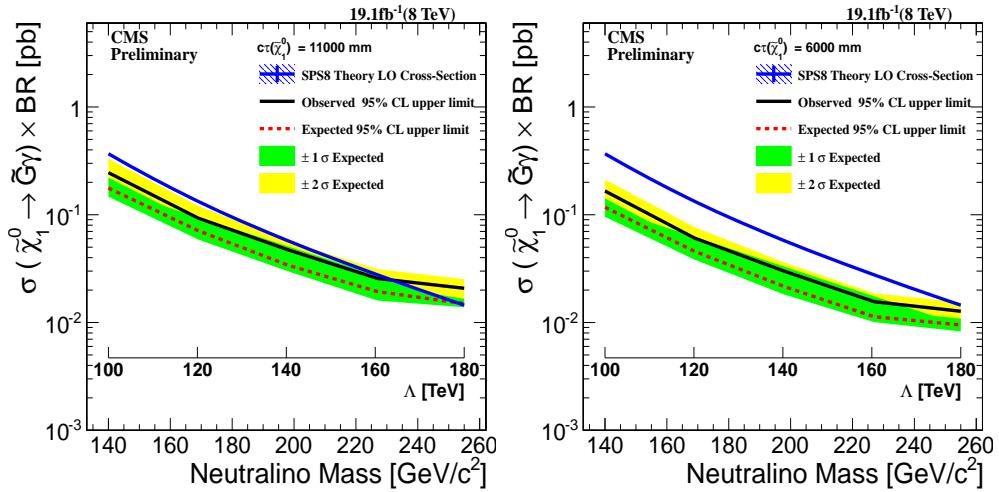


Figure 8.3: Neutralino production cross section against neutralino mass upper limit at 95% confidence levels interpretation in SPS8 model.(Left) $C\tau = 11000$ mm, (Right) $C\tau = 6000$ mm

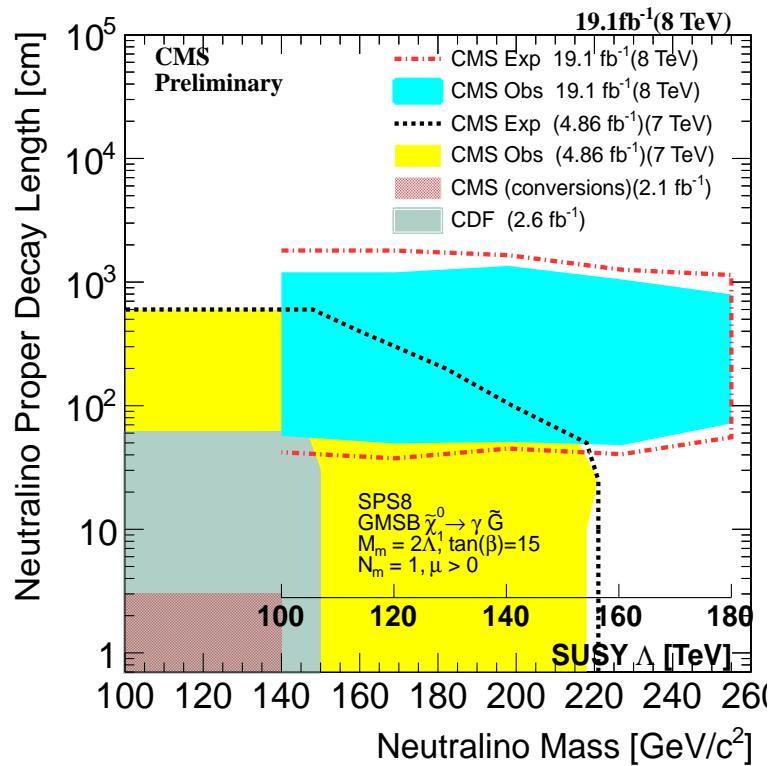


Figure 8.4: Neutralino two dimensional exclusion limit of neutralino mass (Λ) against proper delay length upper limit interpretation in SPS8 model in the decay $\tilde{\chi}_1^0 \rightarrow \gamma + \tilde{G}$ with limits from previous experiments shown.

Chapter 9

Conclusion

We have performed a search analysis for NMLLP decaying to photons using the time of arrival of the photon as measured by the ECAL sub detector of the CMS detector. Haven fail to find any significant signal of delayed photons over the standard model background, we interpreted our results in SUSY models with NMLLP like SPS8 of minimal GBSM or general GMSB models. We showed that, neutralinos whose production and decay mechanism is described in the SPS8 mGMSB model, with $m_{\tilde{\chi}_1^0} \leq 235 \text{ GeV}/c^2$ and $\tau_{\tilde{\chi}_1^0} \leq 35 \text{ ns}$ are ruled out of existence at 95% confidence level using the 2012 8 TeV LHC dataset. This corresponds to an upper limit of $\sigma_{\tilde{\chi}_1^0}^{UP} \geq 0.02 \text{ pb}$ on the production cross section times branching ratio in a hadron collider. In addition, we mention some of the limitations in this particular analysis from a detector point of view and how in future studies can be improved. We hope that in the future, with increase in center of mass energy of the LHC collider as well as luminosity and an improve in timing resolution beyond what is already very reliable, we will surely find a new fundamental particle whose dynamics cannot be described by the already very successful standard model of particle physics.

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Appendix A

Glossary and Acronyms

Care has been taken in this thesis to minimize the use of jargon and acronyms, but this cannot always be achieved. This appendix defines jargon terms in a glossary, and contains a table of acronyms and their meaning.

A.1 Glossary

- **Cosmic-Ray Muon (CR μ)** – A muon coming from the abundant energetic particles originating outside of the Earth’s atmosphere.
- **SUSY** – A theoretical model based on a fundamental symmetry called supersymmetry in which the fermions and bosons can exchange their spin, extending the standard model to account for the stability in the observed Higgs boson mass and to also predicting the existence of many extra new particles which could be candidates of dark matter.

A.2 Acronyms

Table A.1: Acronyms

NMLLP	Neutral Massive Long-Lived Particles.
DM	Dark Matter.
DE	Dark Energy.
GMSB	Gauge Mediated Supersymmetry Breaking
LHC	Large Hadron Collider
CMS	Compact Muon Solenoid
CR μ	Cosmic-Ray Muon