

Fluctuations of the rapidity distribution

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In the coarse-grained Lieb-Liniger model, expectation values w.r.t an ensemble are given by

$$\langle \hat{A} \rangle_w = \frac{\int \mathcal{D}\rho e^{L(S[\rho]-W[\rho])} A[\rho]}{\int \mathcal{D}\rho e^{L(S[\rho]-W[\rho])}} \quad (1)$$

where $A[\rho]$ is the expectation value of \hat{A} in an eigenstate with rapidity distribution ρ , $\mathcal{D}\rho$ denotes the functional integral over all possible rapidity distributions, L is the length of the system, $S[\rho]$ is the Yang-Yang entropy, which is of the form

$$S[\rho] = \int s(\nu(\theta)) \rho_s(\theta) d\theta, \quad (2)$$

and the weight functional that defines the ensemble is

$$W[\rho] = \int w(\theta) \rho(\theta) d\theta. \quad (3)$$

Since the system size L is assumed to be large, the expectation value (5) can be evaluated by saddle point, leading to the Yang-Yang equation,

$$w(\theta) = \frac{\delta S[\rho]}{\delta \rho(\theta)}. \quad (4)$$

This fixes all one-point functions. If we are interested in two-point functions, $\langle \hat{A}\hat{B} \rangle^{\text{conn.}} = \langle \hat{A}\hat{B} \rangle - \langle \hat{A} \rangle \langle \hat{B} \rangle$, then

$$\langle \hat{A}\hat{B} \rangle_w^{\text{conn.}} = \int d\theta \int d\theta' \frac{\delta A}{\delta \rho(\theta)} \frac{\delta B}{\delta \rho(\theta')} \langle \delta \rho(\theta) \delta \rho(\theta') \rangle \quad (5)$$

where

$$\begin{aligned} \langle \delta \rho(\theta) \delta \rho(\theta') \rangle &= \frac{\int \mathcal{D}\delta \rho e^{\frac{L}{2} \int d\lambda \int d\lambda' \frac{\delta^2 S_{YY}}{\delta \rho(\lambda) \delta \rho(\lambda')} \delta \rho(\lambda) \delta \rho(\lambda')} \delta \rho(\theta) \delta \rho(\theta')}{\int \mathcal{D}\delta \rho e^{\frac{L}{2} \int d\lambda \int d\lambda' \frac{\delta^2 S_{YY}}{\delta \rho(\lambda) \delta \rho(\lambda')} \delta \rho(\lambda) \delta \rho(\lambda')}} \\ &= \frac{1}{L} \left[-\frac{\delta^2 S_{YY}}{\delta \rho \delta \rho} \right]^{-1} (\theta, \theta'). \end{aligned} \quad (6)$$

After a calculation (see below), I arrive at the following formula,

$$\langle \delta\rho(\theta)\delta\rho(\theta') \rangle = \frac{1}{L} \int \frac{[\delta(\theta - \lambda) + \nu(\theta)C(\theta, \lambda)][\delta(\theta' - \lambda) + \nu(\theta')C(\theta', \lambda)]}{-s''(\nu(\lambda))} \rho_s(\lambda) d\lambda, \quad (7)$$

with

$$C(\theta, \lambda) = [\frac{1}{2\pi} \Delta(\lambda - \cdot)]^{\text{dr}}(\theta). \quad (8)$$

This result is compatible with formula (1.1) in arXiv:1705.08141 by Herbert and Benjamin, who say that it is ‘an immediate consequence of the (generalized) thermodynamic Bethe ansatz formalism’. To check the equivalence with my formula, one specifies $s(\nu) = -\nu \log \nu - (1 - \nu) \log(1 - \nu)$ so that $\frac{1}{-s''(\nu)} = \nu(1 - \nu)$, and then one uses the fact that

$$h^{\text{dr}}(\lambda) = \int d\theta h(\theta) [\delta(\theta - \lambda) + \nu(\theta)C(\theta, \lambda)].$$

Derivation. To first order we have

$$\begin{aligned} \delta(S - W) &= \int [s'(\nu(\theta))(\delta\rho(\theta) - \nu(\theta)\delta\rho_s(\theta)) + s(\nu(\theta))\delta\rho_s(\theta) - w(\theta)\delta\rho(\theta)] d\theta \\ &= \int [(s'(\nu) - w)\delta\rho + (s(\nu) - \nu s'(\nu))\delta\rho_s] d\theta \end{aligned}$$

and, remarkably, to second order

$$\delta(S - W) = \int s''(\nu) (\delta\nu)^2 \rho_s(\theta) d\theta.$$

It is a miracle that the result is diagonal in $\delta\nu(\theta)$. There must be a deeper meaning to this, but for now I don’t know what it is. Since $\delta^2 S$ is diagonal, it can be inverted straightforwardly:

$$\begin{aligned} \langle \delta\nu(\theta)\delta\nu(\theta') \rangle &= \frac{1}{L} \left[-\frac{\delta^2(S - W)}{\delta\nu\delta\nu} \right]^{-1}(\theta, \theta') \\ &= \frac{1}{L} \frac{1}{-s''(\nu(\theta)) \rho_s(\theta)} \delta(\theta - \theta') \end{aligned} \quad (9)$$

where the last δ is the Dirac delta function, not a differential. We also have

$$\delta\rho(\theta) = \int [\delta(\theta - \lambda) + \nu(\theta)C(\theta, \lambda)] \rho_s(\lambda) \delta\nu(\lambda) d\lambda \quad (10)$$

where $C(\theta, \lambda)$ is defined in (8). This follows as usual from the definition of the dressing. Putting everything together, one gets formula (7).