

This document introduces the implementation of the von Mises based model that employs a Voce (1955) type isotropic hardening rule, a Armstrong–Frederick (Armstrong and Frederick, 1966) type kinematic hardening rule that supports multiplicative back stresses as proposed by Chaboche and Rousselier (1983) and a Peric (1993) type viscosity rule.

References:

1. <https://doi.org/10.1017/S0368393100118759>
2. <https://doi.org/10.1179/096034007X207589>
3. [https://doi.org/10.1016/0749-6419\(89\)90015-6](https://doi.org/10.1016/0749-6419(89)90015-6)
4. <https://doi.org/10.1002/nme.1620360807>

1 Yield Function

A von Mises yielding function is used.

$$F = \sqrt{\frac{3}{2}}|\eta| - k = q - k, \quad (1)$$

in which $\eta = s - \beta$ is the shifted stress, s is the stress deviator, β is the back stress and k is the isotropic hardening stress.

2 Flow Rule

The associated plasticity flow is adopted. The plastic strain rate is then

$$d\epsilon^p = \gamma \frac{\partial F}{\partial \sigma} = \sqrt{\frac{3}{2}}\gamma n, \quad (2)$$

where $n = \frac{\eta}{|\eta|}$. The corresponding accumulated plastic strain rate is

$$dp = \sqrt{\frac{2}{3}}d\epsilon^p : d\epsilon^p = \gamma. \quad (3)$$

3 Plastic Multiplier

The rate of plastic multiplier is defined as

$$\frac{\gamma}{\Delta t} = \dot{\gamma} = \frac{1}{\mu} \left(\left(\frac{q}{k} \right)^{\frac{1}{\epsilon}} - 1 \right), \quad (4)$$

in which μ and ϵ are two material constants. Equivalently, it is

$$q \left(\frac{\Delta t}{\Delta t + \mu\gamma} \right)^{\epsilon} - k = 0. \quad (5)$$

4 Hardening Rules

An exponential function with a linear component is used for isotropic hardening stress.

$$k = \sigma_y + k_l p + k_s - k_s e^{-mp}. \quad (6)$$

The corresponding derivative is

$$\frac{dk}{d\gamma} = k_l + k_s m e^{-mp}. \quad (7)$$

The rate form of back stress $\beta = \sum \beta^i$ is defined as

$$d\beta^i = \sqrt{\frac{2}{3}} a^i d\varepsilon^p - b^i \beta^i dp.$$

In terms of γ , it is $d\beta^i = a^i \gamma n - b^i \gamma \beta^i$. The incremental form is thus

$$\beta^i = \beta_n^i + a^i \gamma n - b^i \gamma \beta^i, \quad \beta^i = \frac{\beta_n^i + a^i \gamma n}{1 + b^i \gamma}. \quad (8)$$

5 Incremental Form

The shifted stress can be computed as

$$\eta = s - \beta = 2GI_d \left(\varepsilon^{tr} - \varepsilon_n^p - \sqrt{\frac{3}{2}} \gamma n \right) - \beta = s^{tr} - \sqrt{6} G \gamma n - \sum \frac{\beta_n^i + a^i \gamma n}{1 + b^i \gamma} \quad (9)$$

with $s^{tr} = 2GI_d (\varepsilon^{tr} - \varepsilon_n^p)$. Hence,

$$|\eta| n + \sqrt{6} G \gamma n + \sum \frac{a^i \gamma}{1 + b^i \gamma} n = s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}, \quad |\eta| + \sqrt{6} G \gamma + \sum \frac{a^i \gamma}{1 + b^i \gamma} = \left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right|.$$

Eventually,

$$|\eta| = \left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right| - \sqrt{6} G \gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma}. \quad (10)$$

Then η can be expressed as,

$$\eta = \frac{\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right| - \sqrt{6} G \gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma}}{\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right|} \left(s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right).$$

It is equivalent to

$$\eta = \left(\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right| - \sqrt{6} G \gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma} \right) u.$$

It is easy to see that $n = u$ with $u = \frac{s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma}}{\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right|}$. The derivatives of u are

$$\frac{\partial u}{\partial \gamma} = \frac{1}{\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right|} \sum \frac{b^i}{(1 + b^i \gamma)^2} \left(\beta_n^i - u : \beta_n^i u \right), \quad \frac{\partial u}{\partial \varepsilon^{tr}} = \frac{1}{\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right|} 2G (I_d - u \otimes u). \quad (11)$$

Furthermore,

$$\frac{\partial q}{\partial \gamma} = \sqrt{\frac{3}{2}} \sum \frac{b^i u : \beta_n^i - a^i}{(1 + b^i \gamma)^2} - 3G. \quad (12)$$

6 Scalar Equation Iteration

For viscoplasticity, the yield function F is not necessarily zero. The rate form of plastic multiplier is used here for iteration.

$$R = q \left(\frac{\Delta t}{\Delta t + \mu \gamma} \right)^\epsilon - k = 0.$$

The corresponding derivatives are then

$$\frac{\partial R}{\partial \gamma} = \left(\frac{\Delta t}{\Delta t + \mu \gamma} \right)^\epsilon \left(\frac{\partial q}{\partial \gamma} - \frac{q \epsilon \mu}{\Delta t + \mu \gamma} \right) - \frac{dk}{d\gamma}, \quad (13)$$

$$\frac{\partial R}{\partial \epsilon^{tr}} = \left(\frac{\Delta t}{\Delta t + \mu \gamma} \right)^\epsilon \sqrt{6} G u : I_d = \left(\frac{\Delta t}{\Delta t + \mu \gamma} \right)^\epsilon \sqrt{6} G u, \quad (14)$$

with

$$k = \sigma_y + k_l (p_n + \gamma) + k_s \left(1 - e^{-m(p_n + \gamma)} \right), \quad (15)$$

$$q = \sqrt{\frac{3}{2}} \left(\left| s^{tr} - \sum \frac{\beta_n^i}{1 + b^i \gamma} \right| - \sqrt{6} G \gamma - \sum \frac{a^i \gamma}{1 + b^i \gamma} \right). \quad (16)$$

7 Consistent Tangent Stiffness

For stiffness, ϵ^{tr} is now varying, then

$$\frac{\partial R}{\partial \epsilon^{tr}} + \frac{\partial R}{\partial \gamma} \frac{d\gamma}{d\epsilon^{tr}} = 0, \quad \frac{d\gamma}{d\epsilon^{tr}} = - \left(\frac{\partial R}{\partial \gamma} \right)^{-1} \frac{\partial R}{\partial \epsilon^{tr}}. \quad (17)$$

Since the stress can be written as

$$\sigma = E(\epsilon^{tr} - \epsilon^p) = E(\epsilon^{tr} - \epsilon_n^p - \Delta \epsilon^p) = E(\epsilon^{tr} - \epsilon_n^p) - \sqrt{6} G \gamma u. \quad (18)$$

The derivative is

$$\begin{aligned} \frac{d\sigma}{d\epsilon^{tr}} &= E - \sqrt{6} G \left(\gamma \frac{\partial u}{\partial \epsilon^{tr}} + \left(u + \gamma \frac{\partial u}{\partial \gamma} \right) \frac{d\gamma}{d\epsilon^{tr}} \right) \\ &= E + \sqrt{6} G \left(\left(u + \gamma \frac{\partial u}{\partial \gamma} \right) \left(\frac{\partial R}{\partial \gamma} \right)^{-1} \frac{\partial R}{\partial \epsilon^{tr}} - \gamma \frac{\partial u}{\partial \epsilon^{tr}} \right). \end{aligned} \quad (19)$$