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Final states dipoles with initial spectators

These have to satisfy the following conditions:

- 1) they must subtract soft singularity;
- 2) they have to match on the soft limit;
- 3) we should write them using Ellis-Campbell mapping;

To start, we define the collinear limit.

$$M^2 \rightarrow M_\mu P_{gg}^{\mu\nu} M_\nu, \text{ where } P_{gg}^{\mu\nu} \text{ is}$$

$$P_{g_i g_j}^{\mu\nu} \sim \left[-g^{\mu\nu} \left(\frac{\xi}{1-\xi} + \frac{1-\xi}{\xi} \right) - 2(1-\xi)\xi \frac{k_\perp^\mu k_\perp^\nu}{k_\perp^2} \right]$$

where ξ & k_\perp^μ are defined as:

$$p_i^\mu = (1-\xi)p^\mu + k_\perp^\mu - \frac{k_\perp^2 n^\mu}{(1-\xi)(2pn)} \quad \& \quad p^\mu, p^2=0 \text{ is the collinear direction.}$$

$$p_j^\mu = \xi p^\mu - k_\perp^\mu - \frac{k_\perp^2 n^\mu}{\xi 2pn}$$

We now use the fact that

$$\pi_{ij}^\mu \equiv p_i^\mu - (1-\xi)p_j^\mu \Rightarrow k_\perp^\mu \quad \text{and} \quad 2p_i \cdot p_j = -\frac{k_\perp^2}{\xi(1-\xi)}$$

$$\Rightarrow P_{g_i g_j}^{\mu\nu} \sim \left[-g^{\mu\nu} \left(\frac{\xi}{1-\xi} + \frac{1-\xi}{\xi} \right) + (1-\xi) \frac{\pi_{ij}^\mu \pi_{ij}^\nu}{p_i p_j} \right].$$

To relate ξ with Campbell-ELLIS

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note that
$$\begin{aligned} p_i \cdot t &= (1-\xi) p^t \\ p_j \cdot t &= \xi p^t \end{aligned} \Rightarrow \frac{p_i \cdot t}{p_j \cdot t} = \frac{1-\xi}{\xi}$$

inter
limit
 $p_i p_j \rightarrow 0$

We have
$$p_i \cdot t = \frac{m_t^2}{2} (1-r^2)^2 (1-z)$$

$$p_j \cdot t = (t - p_i - p_j)^2 = m_t^2 r^2 \Rightarrow$$

$$0 = m_t^2 - m_t^2 r^2 - 2t p_i - 2t p_j + 2p_i p_j \Rightarrow$$

$$m_t^2 (1-r^2) - m_t^2 (1-r^2)^2 (1-z) = \cancel{m_t^2} 2t p_j + 2p_i p_j$$

$$\Rightarrow 2t p_j = m_t^2 (1-r^2) z + O(2p_i p_j) \Rightarrow$$

$$\frac{p_i \cdot t}{p_j \cdot t} = \frac{1-z}{z} \Rightarrow \underline{z \leftrightarrow \xi}$$

Now, we can also modify π_{ij}^M by taking

$$\pi_{ij}^M = \xi p_i^M - (1-\xi) p_j^M \rightarrow z p_i^M - (1-z) p_j^M$$

$$\rightarrow \frac{2t p_j}{m_t^2 (1-r^2)} p_i^M - \frac{2t p_i}{m_t^2 (1-r^2)} p_j^M \equiv$$

$$\boxed{\tilde{\pi}_{ij}^M \equiv \frac{2}{m_t^2 (1-r^2)} (t p_j p_i^M - t p_i p_j^M)} \Rightarrow \tilde{\pi}_{ij}^M = (t p_j) p_i^M - (t p_i) p_j^M$$

$$D_{g_i g_j} = \frac{-1}{2 p_i p_j} \left\{ -g^{\mu\nu} \frac{z}{1-z} + (1-\xi) \frac{2}{m_t^4 (1-r^2)^2} \frac{\tilde{\pi}_{ij}^M}{p_i p_j} \right\}$$

Now, we discuss integration of the -3-

tensor $\frac{\hat{\pi}_{ij}^{\mu} \hat{\pi}_{ij}^{\nu}}{m_t^4 (1-r^2)^2 (p_i p_j)^2}$ We have

$$t^{\mu} \hat{\pi}_{\mu, ij} = t p_j t p_i - t p_i t p_j = 0 \Rightarrow$$

$$\left\langle \frac{\hat{\pi}_{ij}^{\mu} \hat{\pi}_{ij}^{\nu}}{m_t^4 (1-r^2)^2 (p_i p_j)^2} \right\rangle \equiv T^{\mu\nu}(t, \hat{b}) \text{ \& } T^{\mu\nu} t_{\mu} = 0$$

We write $T^{\mu\nu}(t, \hat{b}) = A \left(g^{\mu\nu} - \frac{t^{\mu} t^{\nu}}{t^2} \right) + B \hat{b}_{\perp}^{\mu} \hat{b}_{\perp}^{\nu}$,

where $\hat{b}_{\perp}^{\mu} = \left(\hat{b}^{\mu} - \frac{(\hat{b} t) t^{\mu}}{t^2} \right) \Rightarrow$

Now, to find results for A & B, we do:

$$g^{\mu\nu} T_{\mu\nu} = (d-1)A + B \hat{b}_{\perp}^2$$

$$\hat{b}_{\perp}^2 = \left(\hat{b}^{\mu} - \frac{(\hat{b} t) t^{\mu}}{t^2} \right)^2 = \frac{(\hat{b} t)^2}{t^2} - 2 \frac{\hat{b} t \cdot \hat{b} t}{t^2} = - \frac{(\hat{b} t)^2}{t^2}$$

Now $(t - \hat{b})^2 = m_t^2 r^2 \Rightarrow m_t^2 (1-r^2) = 2 t \hat{b} \Rightarrow$

$$t \hat{b} = \frac{m_t^2 (1-r^2)}{2} \Rightarrow$$

$$\boxed{\hat{b}_{\perp}^2 = - \frac{m_t^2}{4} (1-r^2)^2} \Rightarrow$$

$$\boxed{g^{\mu\nu} T_{\mu\nu} = (d-1)A = \frac{B m_t^2}{4} (1-r^2)^2}$$

Also, contract with $\hat{b}^\mu \Rightarrow$ -4-

$$\hat{b}^\mu \hat{b}^\nu T_{\mu\nu} = -A \frac{(t\hat{b})^2}{t^2} + B \frac{(\hat{b}t)^4}{t^4} \equiv \frac{(\hat{b}t)^2}{t^2} \left(B \frac{(\hat{b}t)^2}{t^2} - A \right)$$

$$\boxed{\hat{b}^\mu \hat{b}^\nu T_{\mu\nu} \equiv \frac{m_t^2 (1-r^2)^2}{4} \cdot \left(B \frac{m_t^2 (1-r^2)^2}{4} - A \right)}$$

To an extent that I did this correctly, we have

$$\begin{aligned} \left\langle \frac{\hat{\pi}_{ij}^\mu \hat{\pi}_{ij}^\nu}{(p_i p_j)^2} \right\rangle &\equiv m_t^2 \left[\frac{(1+r)^2}{12\epsilon} \mp \frac{1}{6} \frac{r^2 (r^4 - 3r^2 + 4) \ln(r)}{(1-r^2)(1-r)^2} \right. \\ &\quad \left. - \frac{1}{3} (1+r)^2 \ln(1+r) + \frac{(3r^2-5)(3r^2-1)}{24(1-r)^2} \right] \left(g_{\mu\nu} \frac{t_\mu t_\nu}{m_t^2} \right) \\ + \left(\frac{1}{3(1-r)^2 \epsilon} \mp \frac{2}{3} \frac{r^4 (r^2-3) \ln r}{(1-r^2)^3 (1-r)^2} - \frac{4}{3} \frac{\ln(1+r)}{(1-r)^2} \right. \\ &\quad \left. + \frac{(9r^4 - 14r^2 + 9)}{6(1-r^2)^2 (1-r)^2} \right) \left(\hat{b}^\mu - \frac{(\hat{b}t)}{t^2} \frac{t^\mu}{t} \right) \left(\hat{b}^\nu - \frac{(\hat{b}t)}{t^2} \frac{t^\nu}{t} \right) \end{aligned}$$

Now, $\langle \rangle$ here stands for

$$\int_0^1 dz \int_0^{y_{\max}} dy \, y^{-\epsilon} (y_{\max} - y)^{-\epsilon} \times \text{but no other} \\ \times (z + r^2(1-z))^{-\epsilon}$$