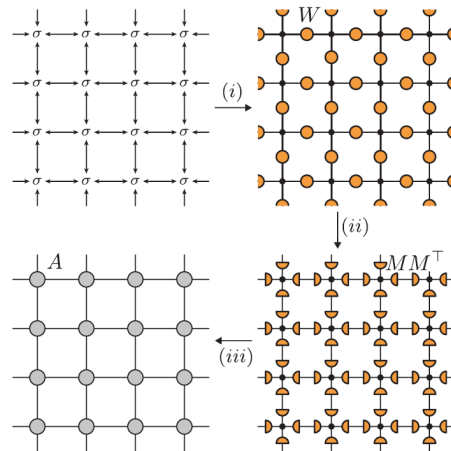


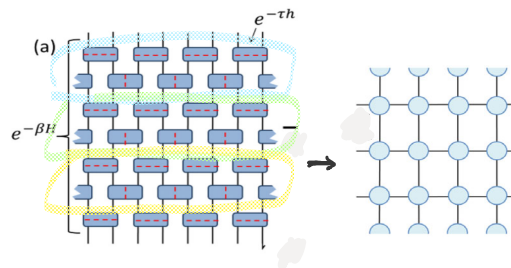
Goal: Compute 2D contractions by coarse-graining RG schemes (instead of transfer matrix schemes)

Applications:

Partition functions of
2D classical models:



Imaginary time evolution of
1D quantum models:



[Levin2007] Levin, Nave: proposed original idea for TRG for classical lattice models.
Local approach: truncation error is minimized only locally.

[Jiang2008] Jiang, Weng, Xiang: adapted Levin-Nave idea to 2D quantum ground state projection via imaginary time evolution. Local approach: truncation is done via 'simple update'. TRG is used to compute expectation values.

[Xie2009] Jiang, Chen, Weng, Xiang; and [Zhao2010] Zhao, Xie, Chen, Wei, Cai, Xiang: Propose 'second renormalization' (SRG), a global approach taking account renormalization of environmental tensor ('full update'). Reduced truncation error significantly.

[Xie2012] Xie, Qin, Zhu, Yang, Xiang: different coarse-graining scheme, using higher-order SVD, employing both local and global optimization schemes.

[Zhao2016] Zhao, Xie, Xiang, Imada: coarse-graining on finite lattices.

[Evenbly2019] Lan, Evenbly: propose core tensor renormalization group (CTRG), which rescales lattice size linearly (not exponentially), but at much lower cost, $\mathcal{O}(\chi^4)$ (rather than $\mathcal{O}(\chi^6)$).

Spin Hamiltonian:

$$H(\{\sigma\}) = \sum_{\langle i,j \rangle} h(\sigma_i, \sigma_j)$$

$$\sigma_i = \{\uparrow, \downarrow\} = \{+1, -1\}$$

Classical partition function:

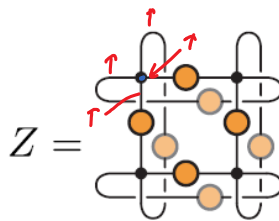
$$Z = \sum_{\{\sigma\}} e^{-\beta H(\{\sigma\})} = \sum_{\{\sigma\}} \bigotimes_{\langle i,j \rangle} W_{\sigma_i \sigma_j}$$

$$h(\sigma_i, \sigma_j) = -\sigma_i \sigma_j J_{ij}$$

Bond weights:

$$W_{\sigma_i \sigma_j} = e^{-\beta h(\sigma_i, \sigma_j)} = \begin{pmatrix} W_{\uparrow\uparrow} & W_{\uparrow\downarrow} \\ W_{\downarrow\uparrow} & W_{\downarrow\downarrow} \end{pmatrix} = \begin{pmatrix} e^\beta & e^{-\beta} \\ e^{-\beta} & e^\beta \end{pmatrix} =: \begin{array}{c} W \\ \sigma_i \quad \sigma_j \end{array}$$

For 2x2 lattice (with periodic conditions):



$$\text{with } \delta_{abcd} = \begin{array}{c} a \quad b \\ | \quad | \\ d \quad c \end{array}$$

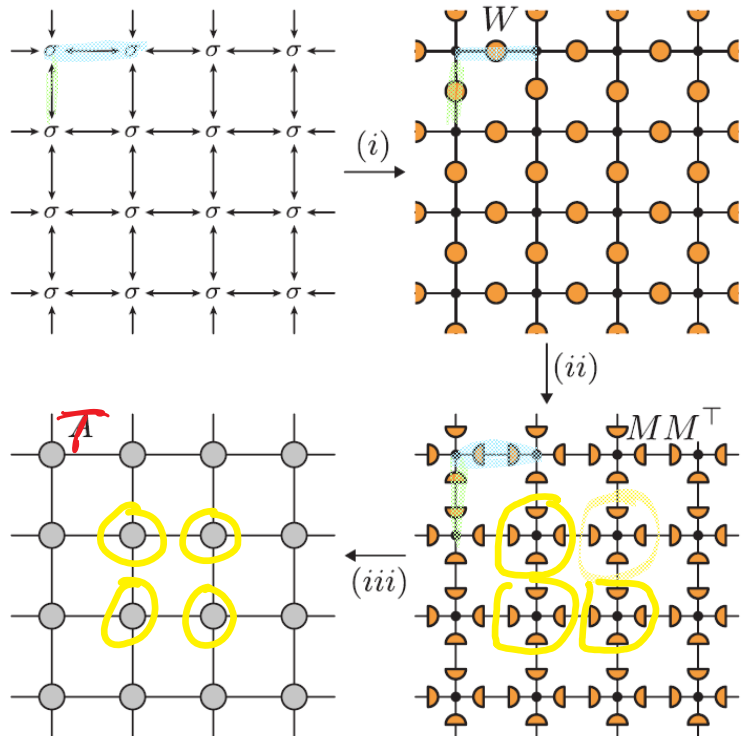


For infinite 2D lattice, we obtain a 2D tensor network:

$$W_{\sigma_i \sigma_j} = \begin{array}{c} W \\ \sigma_i \quad \sigma_j \end{array} = \begin{array}{c} M M^\top \\ \sigma_i \quad \sigma_j \end{array}$$

$$M = \begin{pmatrix} \sqrt{\cosh \beta} & \sqrt{\sinh \beta} \\ \sqrt{\cosh \beta} & -\sqrt{\sinh \beta} \end{pmatrix}$$

$$\begin{aligned} \mathcal{T}_{lurd} &= \begin{array}{c} l \quad u \\ | \quad | \\ d \quad r \end{array} = \begin{array}{c} l \quad u \\ | \quad | \\ \sigma_i \quad \sigma_j \\ | \quad | \\ d \quad r \end{array} \\ &= \sum_{ijkl \in \{\pm, -\}} \delta_{ijkl} M_{il} M_{ju} M_{kr} M_{ld} \\ &= \sum_{\sigma} M_{\sigma l} M_{\sigma u} M_{\sigma r} M_{\sigma d} \end{aligned}$$



Technical challenge: contract this infinite tensor network!

Do SVD on \mathcal{T} in two different ways:

$$\mathcal{T} = \text{diagram} \approx \text{diagram} = U \Sigma V^\dagger$$

$$\mathcal{T} = \text{diagram} \approx \text{diagram} = U' \Sigma' V'^\dagger$$

(ignore red shading)

Iterate until $\mathcal{T}^{(s)}$ converges as $s \rightarrow \infty$
(reaches fixed point)

$$Z = \text{Tr} \mathcal{T}^{(\infty)} = \text{diagram}$$

figure from [Hauru2018]

Structure of \mathcal{T}^∞ can be used to characterize different phases [Gu2009].

Proxy for thermal density matrix:

$$\Gamma = \frac{\text{diagram}}{\text{diagram}} \Rightarrow \text{eigenvalues } \lambda_\alpha$$

von Neumann entropy:

$$S = -\sum_\alpha |\lambda_\alpha| \log(|\lambda_\alpha|)$$

Degeneracy counter:

$$X = \frac{(\text{diagram})^2}{\text{diagram}}$$

has different values in trivial or non-trivial phases

TRG has issues:	does not fully remove local loop correlations (see [Hauru2018])
	computing 'environment' of given site involve tracking all layers of the iteration scheme

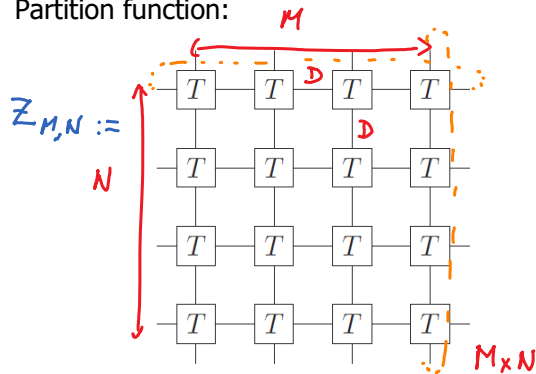
2. 2D contractions via Variational Uniform Matrix Product States (VUMPS)

[Fishman2018]

TRG-I.2

Goal: contract $M \times N$ tensor network (for given T); ultimate take $N \times M \rightarrow \infty \times \infty$

Partition function:



$$=: (\mathcal{K}^M)^N \quad (1)$$

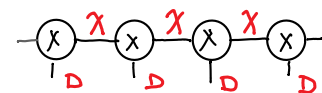
\mathcal{K} = partition function per site

each row contributes a factor \mathcal{K}^M

$$Z_{M,N+1} = \left[\begin{array}{c} \text{grid of } T \text{ tensors} \\ \text{with } N+1 \text{ columns} \end{array} \right] \approx Z_{M,N} \cdot \mathcal{K}^M = \left[\begin{array}{c} \text{grid of } T \text{ tensors} \\ \text{with } N \text{ columns} \end{array} \right] \cdot \mathcal{K}^M \quad (2)$$

$\left[\begin{array}{c} \approx \\ \approx \text{ becomes } = \\ \text{for } M, N \rightarrow \infty \end{array} \right]$

In limit $N \rightarrow \infty$, represent $Z_{M,N}$ by an 'upper boundary MPS':



Then:

$$\left[\begin{array}{c} \text{upper boundary MPS} \\ \text{with } M \text{ rows} \end{array} \right] \approx \left[\begin{array}{c} \text{upper boundary MPS} \\ \text{with } 1 \text{ row} \end{array} \right] \cdot \mathcal{K}^M \quad (3)$$

'fixed-point condition'

'row-to-row transfer matrix'

In limit, $M \rightarrow \infty$, $\dots \text{X} \text{X} \text{X} \text{X} \dots$ is translationally invariant. Express it in canonical form:

$$= \dots \text{A} \text{A} \text{C} \text{B} \dots = \dots \text{A} \text{A} \text{A} \text{B} \text{B} \dots = \dots \text{A} \text{C} \text{B} \text{B} \dots \quad (4)$$

with

$$\left[\begin{array}{c} \text{A} \\ \text{A}^* \end{array} \right] = [\quad , \quad \left[\begin{array}{c} \text{A} \\ \text{B} \end{array} \right] = 1, \quad \left[\begin{array}{c} \text{B} \\ \text{B}^* \end{array} \right] =] \quad (5)$$

left-normalization overall normalization right-normalization

while C, Λ satisfy the 'gauge conditions':
which must hold on all sites.

$$\left[\begin{array}{c} \text{C} \\ \text{A} \end{array} \right] = \left[\begin{array}{c} \text{A} \\ \text{B} \end{array} \right] = \left[\begin{array}{c} \text{A} \\ \text{A} \end{array} \right] \quad (6)$$

[illegible]

Diagram illustrating the approximation of a quantum circuit. The left side shows a sequence of gates T (represented by boxes) connected by a horizontal line, with a blue arrow indicating a transformation. The right side shows a simplified representation of the same sequence, with a blue arrow indicating a transformation and a red arrow indicating a measurement or output.

Given T , $(6,7,8)$ are to be solved for $\frac{A}{V}$, $\frac{C}{Q}$, $\frac{A}{O}$, $\frac{B}{V}$

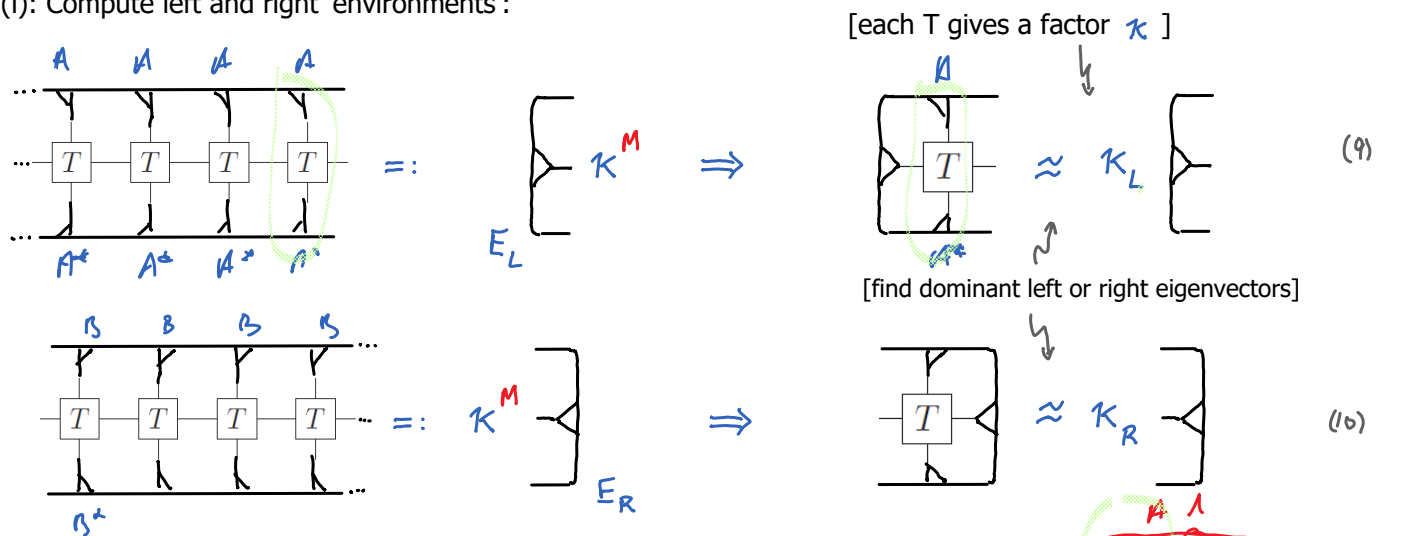
So contraction of infinite tensor network has been reduced to self-consistent solution of four equations!

$(6,7,8)$ have the same structure as when finding ground state of infinite uniform system.


So, solution strategy developed for 'variational uniform matrix product states' (VUMPS) applies:

Repeat following three steps until convergence [with A, C, λ, B from previous iteration as input]:



(i): Compute left and right 'environments':



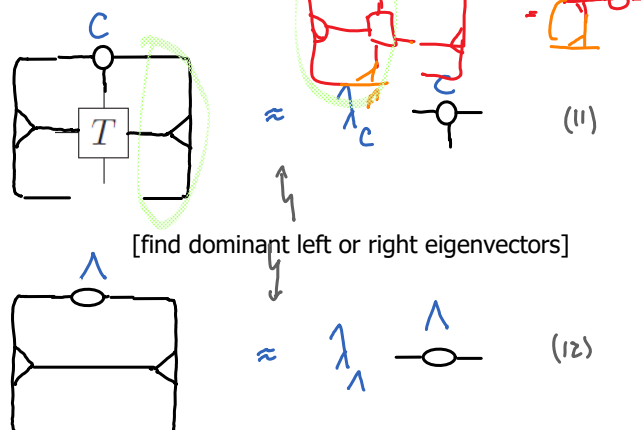
(ii) Solve for central tensor and bond tensor:

(7) contracted with ...  ...

and expressed through environmental tensors, implies:

(8), contracted with...  ...  ...

and expressed through environmental tensors, implies:



At or near fixed point: $\lambda_c \approx \lambda_1 K_L \approx \lambda_1 K_R$ [this follows by contracting (11) with $\frac{1}{\lambda}$ or $\frac{1}{\lambda^2}$]

(iii) From $\overset{\checkmark}{C}$, $\overset{\checkmark}{\Lambda}$ found in (ii), find new $\overset{A}{\nabla}$, $\overset{B}{\nabla}$ that best satisfy (6), $\|C - A\Lambda\|^2$

$$\|C - \Lambda B\|^2 \quad \text{and} \quad \overset{C}{\bigcirc} = \overset{\Lambda}{\bigcirc} \overset{B}{\nabla} \quad (13)$$

$\|C\|^2 + \|\Lambda B\|^2 - 2\|C \cdot \Lambda B\|$
i.e. that maximize $\|C \cdot \Lambda B\|$

and $\overset{C}{\bigcirc} = \overset{A}{\nabla} \overset{\Lambda}{\bigcirc}$ (14)

To that end, do SVDs:

and choose new

$$J = D$$

Unk this note, $= \sum_i S_i$

Repeat (i), (ii), (iii) until convergence, measured, e.g., by change in singular values of Λ .

There may be alternative schemes for finding optimal isometries $\overset{A}{\nabla}$ and $\overset{B}{\nabla}$ that satisfy (13), see 'Riemannian optimization', see [Hauru2021], [Li2023]. Those papers discuss how to optimize a cost function w.r.t. a tensor satisfying an isometry condition. Here, the cost functions would be

$$\|\overset{C}{\bigcirc} - \overset{\Lambda}{\bigcirc} \overset{B}{\nabla}\|^2 \quad \text{and} \quad \|\overset{C}{\bigcirc} - \overset{A}{\nabla} \overset{\Lambda}{\bigcirc}\|^2 \quad (17)$$

and the isometry conditions are Eqs. (5).