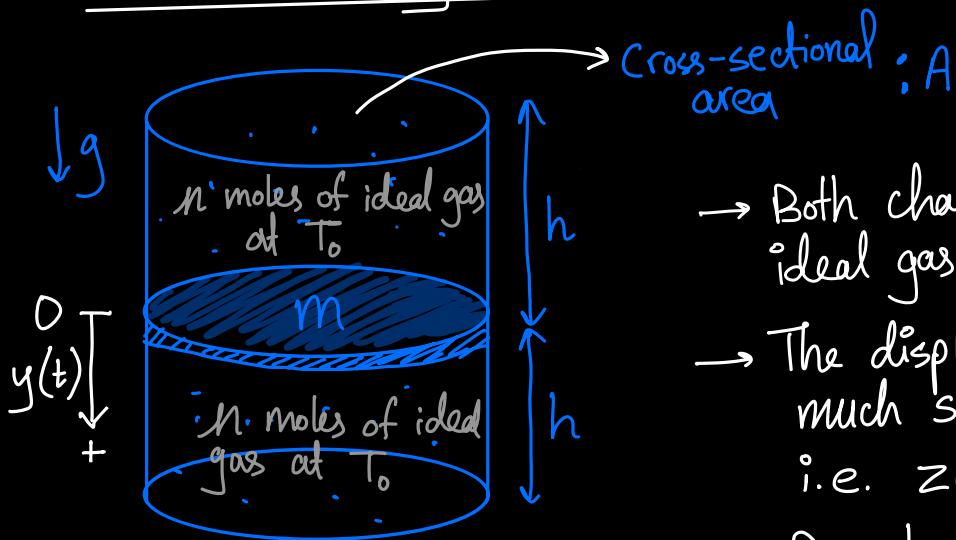


An interesting SHM!



→ Both chambers contain n moles of ideal gas at T_0 temperature.

→ The displacements of mass m are much smaller compared to l , i.e. $z(t) \ll l$.

→ Due to the previous fact along with the fact that there is no friction between the cylinder walls & m , we can consider the expansion and contractions of the gas chambers to be isothermal.

Question:

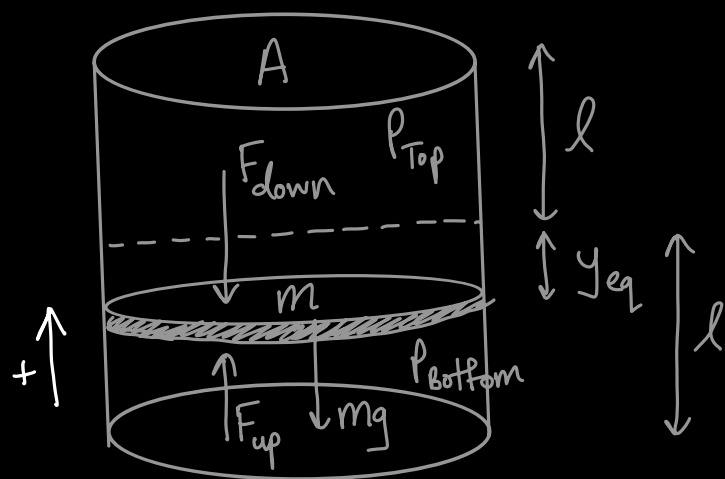
Find an expression for $y(t)$.

Solution:

As we can all guess, m will move in some kind of oscillatory motion. Initially, force due to pressures from both top and bottom chambers will be equal, as a result, m will accelerate downwards due to mg . However, this compresses the bottom chamber and increases the upwards force on m . Hence, after reaching the equilibrium position, m "overshoots" due to inertia and compresses the bottom chamber even more.

After reaching a turning point, it accelerates upwards towards the equilibrium position, only to "overshoot" upwards yet again.

It is always helpful to consider displacements of m with respect to the EQUILIBRIUM POSITION (acceleration = 0) in the cases of oscillatory motion.



The state equation of an ideal gas:

$$PV = nRT \Rightarrow P = \frac{nRT}{V}$$

In our case, all expansions and contractions are isothermal and thus temp. of both chambers are T_0 .

$$\begin{aligned} F_{\text{down}} &= p_{\text{top}} \cdot A \\ &= \frac{nRT_0}{A(h+y_{\text{eq}})} \cdot A \end{aligned}$$

$$\therefore F_{\text{down}} = \frac{nRT_0}{h+y_{\text{eq}}}$$

$$\begin{aligned} F_{\text{up}} &= p_{\text{bottom}} \cdot A \\ &= \frac{nRT_0}{A(h-y_{\text{eq}})} \cdot A \end{aligned}$$

$$\therefore F_{\text{up}} = \frac{nRT_0}{h-y_{\text{eq}}}$$

$$\therefore \sum F_y = F_{\text{up}} - F_{\text{down}} - mg \quad [\text{Upwards is positive}]$$

$$= \frac{nRT_0}{h-y_{\text{eq}}} - \frac{nRT_0}{h+y_{\text{eq}}} - mg$$

$$= nRT_0 \left(\frac{1}{h-y_{\text{eq}}} - \frac{1}{h+y_{\text{eq}}} \right) - mg$$

$$= nRT_0 \left(\frac{h + y_{eq} - (h - y_{eq})}{h^2 - y_{eq}^2} \right) - mg$$

$$= nRT_0 \left(\frac{2y_{eq}}{h^2 \left(1 - \frac{y_{eq}^2}{h^2} \right)} \right) - mg$$

Now, notice that since $y(t) \ll l$, the ratio $\frac{y_{eq}}{h} \approx 0$.

As such, the ratio $\frac{y_{eq}^2}{h^2} = \left(\frac{y_{eq}}{h} \right)^2$ is even smaller.

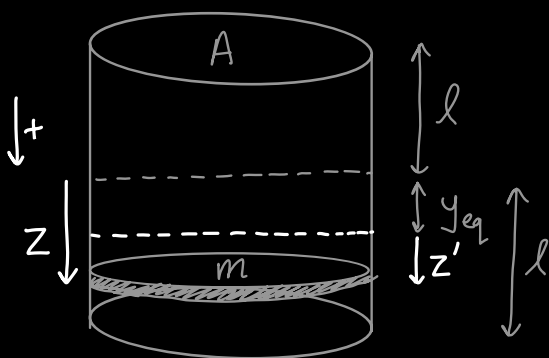
$$\therefore \Sigma F_y = nRT_0 \left(\frac{2y_{eq}}{h^2} \right) - mg \rightarrow \textcircled{i} \left[\text{Note that this is true for all displacement } y(t) \right]$$

For y_{eq} , $\Sigma F_y = 0$.

$$\therefore nRT_0 \left(\frac{2y_{eq}}{h^2} \right) = mg$$

$$\Rightarrow \boxed{y_{eq} = \frac{mgh^2}{2nRT_0}}$$

Now that we have obtained an expression for y_{eq} , we will consider all displacements of m relative to this.



At a displacement of z' from equilibrium, m is at a displacement of:

$$z = y_{eq} + z'$$

from the original (middle) position.

Hence, by (i):

$$\begin{aligned}\underbrace{\sum F_y}_{\substack{\text{upwards} \\ \text{is} +}} &= nRT_0 \left(\frac{2(y_{eq} + z')}{h^2} \right) - mg \\ &= \frac{2nRT_0}{h^2} (y_{eq} + z') - mg \\ &= \frac{2nRT_0}{h^2} \left(\frac{mgh^2}{2nRT_0} + z' \right) - mg \\ &= \cancel{mg} + \frac{2nRT_0 z'}{h^2} = \cancel{mg}\end{aligned}$$

For z and z' : Downwards is +

$$\therefore F_{\text{net up}} = \left(\frac{2nRT_0}{h^2} \right) z'_{\text{down}}$$

Assuming downwards as ultimate + (as in question):

$$\begin{aligned}F_{\text{net}} &= - \left(\frac{2nRT_0}{h^2} \right) z' \\ \therefore ma &= - \left(\frac{2nRT_0}{h^2} \right) z' \end{aligned} \quad \left| \quad \begin{aligned}\therefore a &= - \left(\frac{2nRT_0}{mh^2} \right) z' \\ \text{This is clearly reminiscent} \\ &\text{of the equation of SHM!}\end{aligned}\right.$$

SHM: $a = -\omega^2 x$ \rightarrow displacement from equilibrium

$$\therefore \omega^2 = \frac{2nRT_0}{mh^2}$$

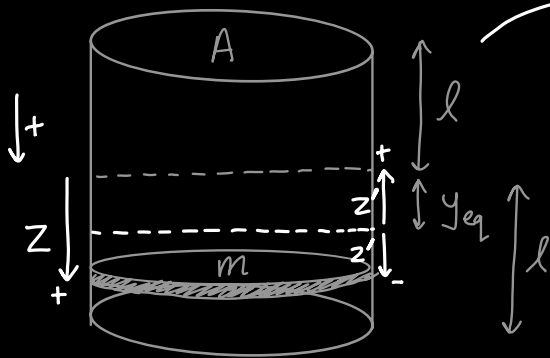
$$\therefore \boxed{\omega = \frac{1}{h} \sqrt{\frac{2nRT_0}{m}}}$$

Now, observe that since this motion has been proven to be simple harmonic, the amplitude of the SHM must be y_{eq} .

Hence:
$$z'(t) = A \cdot \cos(\omega t)$$

$$= \frac{mgh^2}{2nRT_0} \cdot \cos\left(\frac{t}{h} \sqrt{\frac{2nRT_0}{m}}\right)$$

Now, before we arrive at an expression for $z(t)$, observe that at $z'(t)$ is positive ($= +\frac{mgh}{2nRT_0} = +A$) at $t=0$. However, at $t=0$, $z(t) = 0$. As time passes, $z'(t)$ decreases to 0 and then becomes negative below y_{eq} .



Hence, we should negate $z'(t)$ before adding it to y_{eq} to get our required value of $z(t)$.

$$\therefore z(t) = y_{eq} - z'(t)$$

$$= y_{eq} - y_{eq} \cdot \cos(\omega t)$$

$$= y_{eq} (1 - \cos(\omega t))$$

Hence, we have:

$$\therefore z(t) = \frac{mgh^2}{2nRT_0} \cdot \left(1 - \cos\left(\frac{t}{h} \sqrt{\frac{2nRT_0}{m}}\right)\right)$$

Fajaz
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