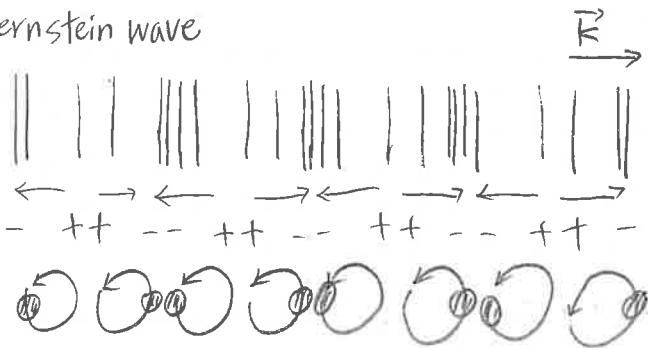


Fusion Plasma Theory 2

Lecture 15: Bernstein Waves and Mode Conversion

① Bernstein wave



collective rotation in phase of the electrons around their guiding centers sustain the longitudinal electron Bernstein wave (EBW)

② Electrostatic kinetic wave dispersion

$$\vec{k} \cdot \vec{E} \cdot \vec{k} = 0 \quad \text{w/ } \vec{k} = k_x \hat{x} + k_z \hat{z} = k_{\perp} \hat{x} + k_{\parallel} \hat{z}$$

$$\Rightarrow k_x^2 E_{xx} + 2k_x k_z E_{xz} + k_z^2 E_{zz} = 0 \quad (E_{xz} = E_{zx})$$

↓ some tedious analysis using Eqs (21)-(24) in Lecture 14.

$$1 + \sum_s \frac{w_{ps}}{k^2} \sum_{n=-\infty}^{\infty} \int \frac{J_n^2(z)}{w - k_{\parallel} v_{ts} - n\omega_{cs}} \left(\frac{n\omega_{cs}}{v_{ts}} \frac{df_{so}}{dv_{ts}} + k_{\parallel} \frac{df_{so}}{dv_{\parallel}} \right) dv = 0 \quad \text{where } z = \frac{k_{\perp} v_{ts}}{\omega_{cs}} \text{ and } k^2 = k_{\perp}^2 + k_{\parallel}^2$$

(Harris dispersion relation)

③ Electrostatic kinetic wave dispersion under Maxwellian

$$\begin{aligned} & k_x^2 + k_z^2 + \sum_s \frac{w_{ps}}{w^2} \gamma_{os} e^{-\lambda s} \sum_{n=-\infty}^{\infty} \left[k_x^2 T_{xx} + 2k_x k_z T_{xz} + k_z^2 T_{zz} \right] \\ & \quad \text{I part in } E_{zz}, E_{xx} \\ & = k_x^2 + k_z^2 + \sum_s \frac{w_{ps}}{w^2} \gamma_{os} e^{-\lambda s} \sum_{n=-\infty}^{\infty} \left[k_x^2 \frac{n^2}{\lambda s} J_n Z - k_x k_z \left(\frac{2}{\lambda s} \right)^{1/2} n Z' - k_z^2 \gamma_{ns} Z' \right] = 0 \\ & \quad (\lambda s = \frac{k_x v_{ts}}{2 w_{ps}}, \gamma_{ns} = \frac{w - n\omega_c}{k_z v_{ts}}, \text{ and } Z' = -2(1 + \gamma_{ns} Z)) \\ & 0 = k_x^2 + k_z^2 + \sum_s \frac{w_{ps}}{w^2} \gamma_{os} e^{-\lambda s} \sum_{n=-\infty}^{\infty} \text{In} \left[\frac{2w_c^2}{v_{ts}^2} n^2 Z - k_z^2 \gamma_{-ns} Z' \right] \\ & = k_x^2 + k_z^2 + \sum_s \frac{w_{ps}}{w^2} \gamma_{os} e^{-\lambda s} \sum_{n=-\infty}^{\infty} \text{In} \left[\frac{2w_c^2}{v_{ts}^2} n^2 Z + 2k_z^2 \gamma_{-ns} (1 + \gamma_{ns}) Z \right] \\ & = k_x^2 + k_z^2 + \sum_s \frac{w_{ps}}{w^2} \gamma_{os} e^{-\lambda s} \sum_{n=-\infty}^{\infty} \text{In} \left[\frac{2w_c^2}{v_{ts}^2} n^2 Z + 2k_z^2 \gamma_{os} + 2k_z^2 \left(\gamma_{os}^2 Z - \frac{n^2 w_c^2}{k_z^2 v_{ts}^2} Z' \right) \right] \\ & = k_x^2 + k_z^2 + \sum_s \frac{2w_{ps}}{w^2} k_z^2 \gamma_{os}^2 e^{-\lambda s} \sum_{n=-\infty}^{\infty} \text{In} (1 + \gamma_{os} Z) \quad \rightarrow \sum_{n=-\infty}^{\infty} e^{-\lambda s} \text{In} = 1 \\ & = k_x^2 + k_z^2 + \sum_s \frac{2w_{ps}}{w^2} k_z^2 \gamma_{os}^2 \left(1 + e^{-\lambda s} \sum_{n=-\infty}^{\infty} \text{In} Z \right) = 0 \end{aligned}$$

⊕ Bernstein dispersion relation

Using $k_z^2 \xi_{os}^2 = w^2/v_{ts}^2$, we obtain general dispersion for Bernstein Wave:

$$k_z^2 + \sum_s \frac{2w_{ps}^2}{v_{ts}^2} \left(1 + e^{-\lambda_s} \xi_{os} \sum_{n=-\infty}^{\infty} I_n(\lambda_s) Z(\xi_{ns}) \right) = 0$$

1) parallel propagating wave. $k_z \rightarrow 0, \lambda_s \rightarrow 0$

$$1 + \sum_s \frac{2w_{ps}^2}{k_{||}^2 v_{ts}^2} \left(1 + \xi_{os} Z(\xi_{os}) \right) = 1 - \sum_s \frac{w_{ps}^2}{k_{||}^2 v_{ts}^2} Z'(\xi_{os}) = 0$$

2) Perpendicular propagating wave $k_{||} \rightarrow 0, \xi_{ns} \rightarrow \infty$

$$1 + \sum_s \frac{2w_{ps}^2}{k_z^2 v_{ts}^2} \left(1 - e^{-\lambda_s} \xi_{os} \sum_{n=-\infty}^{\infty} \frac{I_n(\lambda_s)}{\xi_{ns}} \right) = 0$$

(*)

$$(*) = 1 - e^{-\lambda_s} \xi_{os} \left(\frac{I_0}{\xi_{os}} + \sum_{n=1}^{\infty} \left(\frac{1}{\xi_{ns}} + \frac{1}{\xi_{-ns}} \right) I_n \right)$$

$$= 1 - e^{-\lambda_s} I_0 - e^{-\lambda_s} \xi_{so} \sum_{n=1}^{\infty} \frac{2W_{||} k_{||} v_{ts}}{w^2 - n^2 w_c^2} I_n \quad \xi_{so} = \frac{w}{k_{||} v_{ts}}$$

$$= 1 - e^{-\lambda_s} I_0 - e^{-\lambda_s} \sum_{n=1}^{\infty} \frac{2w^2 - 2n^2 w_c^2 + 2n^2 w_c^2}{w^2 - n^2 w_c^2} I_n$$

$$= 1 - e^{-\lambda_s} \underbrace{\left(I_0 + 2 \sum_{n=1}^{\infty} I_n \right)}_{=1} - e^{-\lambda_s} \sum_{n=1}^{\infty} \frac{2n^2}{(w/w_c)^2 - n^2} I_n$$

(* note $\lambda_s = \frac{1}{2} k_z^2 \xi_{os}^2$)

$$\Rightarrow 1 = \sum_s \frac{w_{ps}^2}{w_c s^2} \frac{2}{\lambda_s} \sum_{n=1}^{\infty} e^{-\lambda_s} I_n(\lambda_s) \frac{n^2}{(w/w_{cs})^2 - n^2} \quad \dots (13)$$

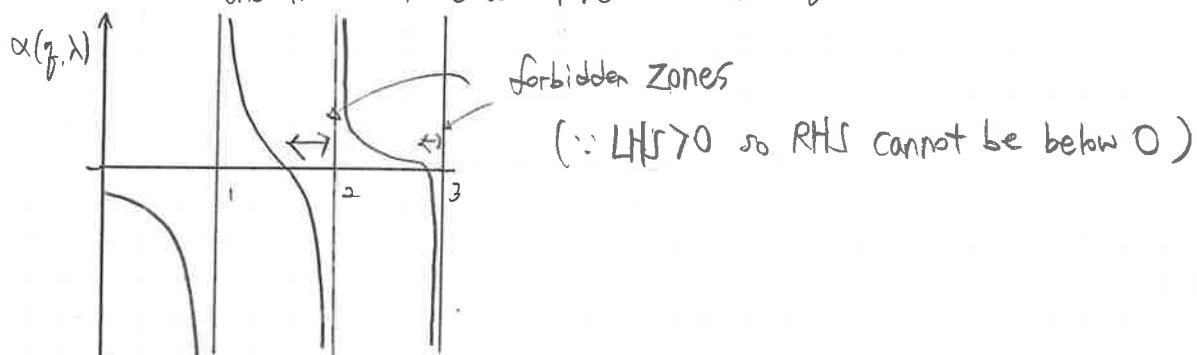
⑤ Frequency characteristics of Electron Bernstein wave

Consider the high frequency mode where the ion contribution can be ignored due to $\omega/\omega_i \gg 1$. Define $q_f = \omega/\omega_{ce}$ and drop 'e' subscript.

$$\Rightarrow \left[\frac{k_\perp^2 v_{te}^2}{2\omega_{pe}^2} = 2 \sum_{n=1}^{\infty} e^{-\lambda} I_n(\lambda) \frac{n^2}{q_f^2 - n^2} \right] \equiv \alpha(q_f, \lambda)$$

Given a $\lambda > 0$, α as a function of q_f is shown below.

↔ indicates that "its frequencies are just above each cyclotron harmonic frequencies" and that "there are forbidden zones just below them".



⑥ Dispersion characteristics of Electron Bernstein Wave

The fluid limit of EBW can be obtained by $\lambda \rightarrow 0$ and $I_n(\lambda) \sim \frac{1}{n!} \left(\frac{\lambda}{2}\right)^n$

From eq (13), only $n=1$ remains finite:

$$1 = \frac{\omega_{pe}^2}{\omega_{ce}^2} \cancel{\neq} \cancel{\lambda} \frac{1}{(\omega/\omega_{ce})^2 - 1} = \frac{\omega_{pe}^2}{\omega^2 - \omega_{ce}^2} \rightarrow \begin{array}{l} \text{Upper hybrid resonance Condition} \\ (\text{x-wave resonance condition}) \\ (S=0) \end{array}$$