

# Ch10. Non linear constitutive equation

## 10.2 The theory of finite elastic deformation

Linear theory of elasticity is limited to cases the deformation gradient  $F$  are small.

Rubbers: behave elastically while undergoing large deformation.

$W = \rho_0 e$  : strain-energy function

$W = W(F_{iR}) = W(F) \leftarrow$  arbitrary dependence on the deformation.

Recall  $\frac{D}{Dt} F_{iR} = \frac{D}{Dt} \left( \frac{dx_i}{dX_R} \right) = \frac{dv_i}{dX_R} = \frac{dv_i}{dx_j} \frac{dx_j}{dX_R} = L_{ij} F_{jR}$

$$\rho \frac{De}{Dt} = T_{ij} \frac{dv_i}{dx_j} - \frac{d\rho}{dt} \frac{e}{\rho}$$

$$\Rightarrow T_{ij} \frac{dv_i}{dx_j} = \frac{\rho}{\rho_0} \frac{DW}{Dt} = \frac{\rho}{\rho_0} \frac{dW}{dF_{iR}} \frac{DF_{iR}}{Dt} = \frac{\rho}{\rho_0} \frac{dW}{dF_{iR}} \frac{dx_j}{dX_R} \frac{dv_i}{dx_j} \quad (\text{for all } \frac{dv_i}{dx_j})$$

$$\hookrightarrow \therefore T_{ij} = \frac{\rho}{\rho_0} \frac{dW}{dF_{iR}} \frac{dx_j}{dX_R} \quad \text{Constitutive equation for finite elasticity.}$$

$W(F_{iR})$  shouldn't change if a rigid-body rotation is superposed on the deformation

$$\bar{x} = M \cdot x \quad \text{rotation}$$

$$\bar{F}_{iR} = \frac{d\bar{x}_i}{dX_R} = M_{ij} \frac{dx_j}{dX_R} = M_{ij} F_{jR} \quad \text{or} \quad \bar{F} = M \cdot F \quad \leftarrow (\text{좌표계가 아닌, 실제 물체를 돌린 것})$$

$\hookrightarrow \underline{W(F) = W(M \cdot F)}$  for any  $M$ .

$W(F) = W(M \cdot R \cdot M)$  for any  $M$ , especially  $M = R^T$

$$= W(U) \rightarrow \boxed{W = W(C)} \quad \text{결론: rigid-body motion} \rightarrow W = W(C)$$

$\hookrightarrow 6 \text{ components}$

$$\begin{aligned} \frac{DW}{Dt} &= \frac{dW}{dC_{RS}} \frac{DC_{RS}}{Dt} = \frac{dW}{dC_{RS}} \frac{D}{Dt} \left( \frac{dx_i}{dX_R} \frac{dx_i}{dX_S} \right) = \frac{dW}{dC_{RS}} \left( \frac{dv_i}{dX_R} \frac{dx_i}{dX_S} + \frac{dx_i}{dX_R} \frac{dv_i}{dX_S} \right) \\ &= \left( \frac{dW}{dC_{RS}} + \frac{dW}{dC_{SR}} \right) \frac{dx_i}{dX_R} \frac{dv_i}{dX_S} = \left( \frac{dW}{dX_{RS}} + \frac{dW}{dX_{SR}} \right) \frac{dx_i}{dX_R} \frac{dx_j}{dX_S} \frac{dv_i}{dx_j} \end{aligned}$$

$$\text{Since } T_{ij} \frac{dv_i}{dx_j} = \frac{\rho}{\rho_0} \frac{DW}{Dt} \rightarrow \boxed{T_{ij} = \frac{\rho}{\rho_0} \frac{dx_i}{dX_R} \frac{dx_j}{dX_S} \left( \frac{dW}{dX_{RS}} + \frac{dW}{dX_{SR}} \right)}$$

Required general form of constitutive equation for finite elastic solid.

Recall that  $\Pi = (\det F) F^{-1} \cdot T$

$$\text{From } T_{ij} = \frac{\rho}{\rho_0} \frac{dW}{dF_{iR}} F_{jR}, \quad T_{ji} = \frac{\rho}{\rho_0} F_{iR} \frac{dW}{dF_{jR}} \rightarrow [(\det F) F^{-1} \cdot T]_{Ri} = \frac{dW}{dF_{iR}} = \pi_{Ri}$$

$$\text{Also, } P = \Pi \cdot (F^{-1})^T = (\det F) F^{-1} \cdot T \cdot (F^{-1})^T$$

$$\text{From } T_{ij} = \frac{\rho}{\rho_0} \frac{dx_i}{dX_R} \left( \frac{dW}{dC_{RS}} + \frac{dW}{dC_{SR}} \right) \frac{dx_j}{dX_S}, \quad P_{RS} = \left( \frac{dW}{dC_{RS}} + \frac{dW}{dC_{SR}} \right)$$

Material symm: If rotational symmetry by  $\mathcal{Q}$

$$\mathbb{F} \rightarrow \mathcal{Q}^T \mathbb{F} \mathcal{Q} \Rightarrow \mathbb{C} = \mathbb{F}^T \mathbb{F} \rightarrow \mathcal{Q}^T \mathbb{C} \mathcal{Q}$$

$$\underline{W(\mathbb{C}) = W(\mathcal{Q}^T \mathbb{C} \mathcal{Q})} \quad (\text{if isotropic, this holds for all } \mathcal{Q})$$

$$\hookrightarrow W \text{ is } \underline{\text{invariant of } \mathbb{C}} \rightarrow \underline{W = W(I_1, I_2, I_3)} \quad \text{Similarly, } W(\mathbb{R}) = W(\mathcal{Q}^T \mathbb{R} \mathcal{Q})$$

$$\frac{dW}{dC_{RS}} = \frac{dW}{dI_1} \frac{dI_1}{dC_{RS}} + \frac{dW}{dI_2} \frac{dI_2}{dC_{RS}} + \frac{dW}{dI_3} \frac{dI_3}{dC_{RS}}$$

$$\left( \begin{array}{l} I_1 = \text{tr } \mathbb{C} = C_{RR} \\ I_2 = \frac{1}{2} \{ (\text{tr } \mathbb{C})^2 - \text{tr } \mathbb{C}^2 \} = \frac{1}{2} \{ C_{RR} C_{SS} - C_{RS} C_{RS} \} \\ I_3 = \det \mathbb{C} \end{array} \right)$$

$$\frac{dI_1}{dC_{RS}} = \frac{dC_{PP}}{dC_{RS}} = \delta_{PR} \delta_{PS} = \delta_{RS}$$

$$\begin{aligned} \frac{dI_2}{dC_{RS}} &= \frac{1}{2} \frac{d}{dC_{RS}} \{ C_{PP} C_{QQ} - C_{PQ} C_{PQ} \} = \frac{1}{2} \{ \delta_{PR} \delta_{PS} C_{QQ} + \delta_{QR} \delta_{QS} C_{PP} - 2 \delta_{PR} \delta_{QS} C_{PQ} \} \\ &= \delta_{RS} C_{PP} - C_{RS} = I_1 \delta_{RS} - C_{RS} \end{aligned}$$

$$I_3 = \frac{1}{3} \{ \text{tr } \mathbb{C}^3 - I_1 \text{tr } \mathbb{C}^2 + I_2 \text{tr } \mathbb{C} \} = \frac{1}{3} \{ C_{AB} C_{BC} C_{CA} - I_1 C_{AB} C_{BA} + I_2 C_{AA} \}$$

$$\begin{aligned} \frac{dI_3}{dC_{RS}} &= \frac{1}{3} \frac{d}{dC_{RS}} \{ C_{AB} C_{BC} C_{CA} - I_1 C_{AB} C_{BA} + I_2 C_{AA} \} \\ &= C_{RP} C_{SP} - I_1 C_{RS} + I_2 \delta_{RS} \end{aligned}$$

$$\boxed{T_{ij} = 2 \frac{P}{P_0} \frac{dx_i}{dx_R} \frac{dx_j}{dx_S} \left\{ \left( \frac{dW}{dI_1} + I_1 \frac{dW}{dI_2} + I_2 \frac{dW}{dI_3} \right) \delta_{RS} - \left( \frac{dW}{dI_2} + I_1 \frac{dW}{dI_3} \right) C_{RS} + \frac{dW}{dI_3} C_{RP} C_{PS} \right\}}$$

Constitutive equation for an isotropic finite elastic solid.

$$\underline{\underline{\mathbb{T} = 2(I_3)^{-\frac{1}{2}} \mathbb{F} \cdot \{ (W_1 + I_1 W_2 + I_2 W_3) \mathbb{I} - (W_2 + I_1 W_3) \mathbb{C} + W_3 \mathbb{C}^2 \} \cdot \mathbb{F}^T}}$$

$$(\text{Using } \mathbb{B} = \mathbb{F} \cdot \mathbb{F}^T \text{ \& } \mathbb{C} = \mathbb{F}^T \mathbb{F})$$

$$\underline{\underline{\mathbb{T} = 2(I_3)^{-\frac{1}{2}} \{ (W_1 + I_1 W_2 + I_2 W_3) \mathbb{B} - (W_2 + I_1 W_3) \mathbb{B}^2 + W_3 \mathbb{B}^3 \}}}$$

$$(\text{using } \mathbb{B}^3 - I_1 \mathbb{B}^2 + I_2 \mathbb{B} - I_3 \mathbb{I} = 0, \text{tr } \mathbb{B}^n = \text{tr } \mathbb{C}^n)$$

$$\mathbb{T} = 2(I_3)^{-\frac{1}{2}} \{ I_3 W_3 \mathbb{I} + (W_1 + I_1 W_2) \mathbb{B} - W_2 \mathbb{B}^2 \}$$

### 10.3 A non linear viscous fluids

Recall : Newtonian viscous fluid

$$T_{ij} = -p(\rho, \theta) \delta_{ij} + B_{ijkl}(\rho, \theta) D_{kl}$$

If fluid is at rest,  $T_{ij} = -p(\rho, \theta) \delta_{ij}$

If motion is rotated,  $F = \left( \frac{dx_i}{dx_R} \right) \Rightarrow \underline{F \rightarrow Q^T F Q}$



Likewise,  $\underline{\frac{D}{Dt} F \rightarrow Q^T \frac{D}{Dt} F Q}$

Also,  $\underline{F^{-1} \rightarrow Q^T F^{-1} Q}$

$$D_{kk} = \frac{1}{2} \left( \frac{dv_k}{dx_k} + \frac{dv_k}{dx_k} \right) = \frac{1}{2} \left( \frac{dv_k}{dx_R} \frac{dx_R}{dx_k} + \frac{dv_k}{dx_R} \frac{dx_R}{dx_k} \right) = \frac{1}{2} \left( \frac{dF}{dt} F^{-1} + (F^{-1})^T \frac{dF^T}{dt} \right)_{kk}$$

$$\therefore \underline{D \rightarrow Q^T D Q}$$

If isotropic,  $\underline{T \rightarrow Q^T T Q}$

$$\therefore Q_{ia} T_{ab} Q_{bj} = -p \delta_{ij} + B_{ijkl} Q_{kc} D_{cd} Q_{de}$$

$$Q_{ei} (Q_{ia} T_{ab} Q_{bj}) Q_{jf} \Rightarrow T_{ef} = -p \delta_{ef} + \underbrace{B_{ijkl} Q_{ei} Q_{aj} Q_{dk} Q_{de}}_{B_{efcd}} D_{cd}$$

$\therefore B_{ijkl}$  : isotropic ,

$$T_{ij} = \{-p + \lambda D_{kk}\} \delta_{ij} + 2\mu D_{ij} \quad , \text{or,} \quad \underline{T = (-p + \lambda \text{tr} D) \mathbb{I} + 2\mu D.}$$

\* Non-newtonian fluid

$$T_{ij} = T_{ij} \left( \frac{dv_p}{dx_k}, \rho, \theta \right) \quad , \text{ or,} \quad \underline{T = \Pi(\mathbb{L}, \rho, \theta)}$$

$$\text{Since } \mathbb{L} = \mathbb{D} + \mathbb{W} \quad , \quad \underline{T = \Pi(\mathbb{D}, \mathbb{W}, \rho, \theta)}$$

Original motion :  $\underline{x} = \underline{x}(\underline{X}, t)$  ,  $\underline{v} = \underline{v}(\underline{x}, t)$

New motion :  $\underline{\bar{x}} = \underline{M}(t) \underline{x}(\underline{X}, t) \leftrightarrow \underline{x} = \underline{M}^T \underline{\bar{x}}$

$\hookrightarrow$  time dependent rigid rotation

$$\Rightarrow \underline{\bar{v}} = \underline{\frac{D}{Dt} \bar{x}} = \dot{M} \underline{x} + M \underline{v}$$

$$\underline{\bar{L}}_{ij} = \frac{d\bar{v}_i}{d\bar{x}_j} = \frac{d\bar{v}_i}{dx_k} \frac{dx_k}{d\bar{x}_j} = \left( M_{ip} \delta_{pk} + M_{ip} \frac{dv_p}{dx_k} \right) M_{kj}^T \quad , \text{ or,} \quad \underline{\bar{L} = (\dot{M} + M \cdot \mathbb{L}) M^T}$$

$$\begin{cases} \underline{\bar{D}} = \frac{1}{2} (\underline{\bar{L}} + \underline{\bar{L}}^T) = \frac{1}{2} (\dot{M} \cdot M^T + M \cdot \dot{M}^T) + \frac{1}{2} M (\mathbb{L} + \mathbb{L}^T) M^T \\ \underline{\bar{W}} = \frac{1}{2} (\underline{\bar{L}} - \underline{\bar{L}}^T) = \frac{1}{2} (\dot{M} \cdot M^T - M \cdot \dot{M}^T) + \frac{1}{2} M (\mathbb{L} - \mathbb{L}^T) M^T \end{cases}$$

From  $M \cdot M^T = I \Rightarrow \bar{D} = M D M^T$   
 $\rightarrow \dot{M} \cdot M^T + M \cdot \dot{M}^T = 0 \Rightarrow \underline{\bar{W} = M (M^T \cdot \dot{M} + \dot{M}^T \cdot M) \cdot M^T}$

We need  $\bar{T} = M \cdot T \cdot M^T$ ,  $\bar{T} = T(\bar{D}, \bar{W}, \rho, \theta)$

$\therefore T\{M D M^T, M (M^T \cdot \dot{M} + \dot{M}^T \cdot M) \cdot M^T, \rho, \theta\} = M \cdot T(D, W, \rho, \theta) \cdot M^T$  for any  $M$ .

Consider the case  $M = I$ ,  $\dot{M} \neq 0$

$T\{D, \dot{M} + \dot{M}^T, \rho, \theta\} = T(D, W, \rho, \theta) \rightarrow \therefore T$  should be indep. of  $W$ .

$\hookrightarrow T(M \cdot D \cdot M^T, \rho, \theta) = M T(D, \rho, \theta) M^T$

In general,  $T = -pI + \alpha D + \beta D^2$

where  $p, \alpha, \beta$  are function of  $\rho, \theta$ , and invariants of  $D$ .

$(\text{tr} D, \frac{1}{2}\{(\text{tr} D)^2 - \text{tr} D^2\}, \det D)$   
 $(\text{tr} D, \text{tr} D^2, \text{tr} D^3)$

\* Appendix : Representation theorem for an isotropic tensor function of a tensor

Theorem :  $T$  is an isotropic tensor function of  $D$  iff

$T = \alpha I + \beta D + \gamma D^2$  where  $\alpha, \beta, \gamma$  are scalar functions of  $\text{tr} D, \text{tr} D^2, \text{tr} D^3$ .

(a) : Sufficiency (denote  $\bar{T} = M T M^T$ ,  $\bar{D} = M D M^T$ )

$T(M \cdot D \cdot M^T) = M \cdot T(D) \cdot M^T$   
for all  $M$ .

$M \cdot T(D) \cdot M^T = M (\alpha I + \beta D + \gamma D^2) M^T$

$= \alpha I + \beta M D M^T + \gamma M D^2 M^T$

$= \alpha I + \beta \bar{D} + \gamma \bar{D}^2 = T(\bar{D})$

(b) : necessity

Assume  $T(M \cdot D \cdot M^T) = M T(D) M^T$  for all  $M$ .

Consider the case  $D$  is diagonal.  $D = \begin{pmatrix} D_1 & & \\ & D_2 & \\ & & D_3 \end{pmatrix} \rightarrow T_{ij} = T_{ij}(D_1, D_2, D_3)$

(i) Choose  $M = \begin{pmatrix} 1 & & \\ & -1 & \\ & & -1 \end{pmatrix}$ , then  $\bar{D} = M D M^T = D$ , and  $\bar{T} = \begin{pmatrix} T_{11} & -T_{12} & -T_{13} \\ -T_{12} & T_{22} & T_{23} \\ -T_{13} & T_{23} & T_{33} \end{pmatrix} = T(\bar{D}) = T(D)$

$\therefore T_{12} = T_{13} = 0$ , Likewise  $T_{23} = 0$  (by other choice of  $M$ )

→ If  $D$  is diagonal, so is  $\Pi$ . →  $D$  and  $\Pi$  have the same principal axes.

$$\rightarrow T_{11} = T_1 = F(D_1, D_2, D_3), \quad T_{22} = T_2 = F_2(D_1, D_2, D_3), \quad T_{33} = T_3 = F_3(D_1, D_2, D_3)$$

(ii) Choose  $M = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$

$$\bar{D} = M D M^T = \begin{pmatrix} D_2 & & \\ & D_3 & \\ & & D_1 \end{pmatrix}, \quad \bar{\Pi} = M \Pi M^T = \begin{pmatrix} T_2 & & \\ & T_3 & \\ & & T_1 \end{pmatrix}$$

$$\therefore \Pi(\bar{D}) = \begin{pmatrix} F(D_2, D_3, D_1) \\ F_2(D_2, D_3, D_1) \\ F_3(D_2, D_3, D_1) \end{pmatrix} = \bar{\Pi} = \begin{pmatrix} T_2 & & \\ & T_3 & \\ & & T_1 \end{pmatrix} = \begin{pmatrix} F_2(D_1, D_2, D_3) \\ F_3(D_1, D_2, D_3) \\ F(D_1, D_2, D_3) \end{pmatrix}$$

$$\hookrightarrow \begin{cases} F(D_2, D_3, D_1) = F_2(D_1, D_2, D_3) & T_1 = F(D_1, D_2, D_3) \\ F_2(D_2, D_3, D_1) = F_3(D_1, D_2, D_3) & \Rightarrow T_2 = F(D_2, D_3, D_1) \\ F_3(D_2, D_3, D_1) = F(D_1, D_2, D_3) & T_3 = F(D_3, D_1, D_2) \end{cases}$$

$$\alpha(D_1, D_2, D_3) + \beta(D_1, D_2, D_3) D_1 + \gamma(D_1, D_2, D_3) D_1^2 = F(D_1, D_2, D_3) = T_1$$

$$\alpha(D_1, D_2, D_3) + \beta(D_1, D_2, D_3) D_2 + \gamma(D_1, D_2, D_3) D_2^2 = F(D_2, D_3, D_1) = T_2$$

$$\alpha(D_1, D_2, D_3) + \beta(D_1, D_2, D_3) D_3 + \gamma(D_1, D_2, D_3) D_3^2 = F(D_3, D_1, D_2) = T_3$$

(iii) Choose  $M = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$

$$\bar{D} = M D M^T = \begin{pmatrix} D_2 & & \\ & D_1 & \\ & & D_3 \end{pmatrix}, \quad \bar{\Pi} = M \Pi M^T = \begin{pmatrix} T_2 & & \\ & T_1 & \\ & & T_3 \end{pmatrix}$$

$$\Pi(\bar{D}) = \begin{pmatrix} F(D_2, D_1, D_3) \\ F_2(D_2, D_1, D_3) \\ F_3(D_2, D_1, D_3) \end{pmatrix} = \bar{\Pi} = \begin{pmatrix} F_2(D_1, D_2, D_3) \\ F(D_1, D_2, D_3) \\ F_3(D_1, D_2, D_3) \end{pmatrix}$$

(Recall)

$$F_2(D_1, D_2, D_3) = F(D_2, D_3, D_1) \Rightarrow \underline{F(D_2, D_3, D_1) = F(D_2, D_1, D_3)}$$

$$F_3(D_1, D_2, D_3) = F(D_3, D_1, D_2)$$

결론:  $F$ 는 맨 앞의 Argument  $D_1$ 에 의존  
(뒤의  $D_j, D_k$ 는 순서 바뀌어도 괜찮)

(iv) we can find  $\alpha(D_1, D_2, D_3), \beta(D_1, D_2, D_3), \gamma(D_1, D_2, D_3)$  such that

$$\alpha + \beta D_1 + \gamma D_1^2 = F(D_1, D_2, D_3) = T_1$$

$$\alpha + \beta D_2 + \gamma D_2^2 = F(D_2, D_3, D_1) = T_2$$

$$\alpha + \beta D_3 + \gamma D_3^2 = F(D_3, D_1, D_2) = T_3$$

$$\begin{pmatrix} T_1 \\ T_2 \\ T_3 \end{pmatrix} = \alpha \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} + \beta \begin{pmatrix} D_1 \\ D_2 \\ D_3 \end{pmatrix} + \gamma \begin{pmatrix} D_1^2 \\ D_2^2 \\ D_3^2 \end{pmatrix}$$

$$(T_3) \quad (1) \quad (p_3) \quad (D_3^2)$$