1-1000 =

## ELECTROWEAK LEPTONIC SECTOR

- · spacetime symmetry & translation
- UN Hypercharge > GAUGE SYMMETRIES
  SU(2) Weak + these are special

(furdomental)

remark: GAUCE SYM () FORCE

further. the spin-1 is in the ADMINT ROP.

229/32 pm 1 - NIGE EI particle that talks to things charged under the gauge symmetry

F B a net itself charged

80: U(1): Bu ~~~ BM

SU(2): Will Will SU(2) -> triplet

remart: so before we've made any choices about PARTICLE CONTENT, WE AUTOMATICALLY GOT & PARTICES (really 4) FOR FREE!

particles

LEPTON DOUBLET (Lt) & >- Lag L. (Lais = er

RH (AMT) - ELECTRON (E) (E) Y

HPBOS Doublet (H+) 1 - ->- H = (H+)

[and: What are the Feynman tules?

## Drawing the rules

I all of them

Hyperonage boson talks to hyperonarged particles

- THIS MIL LOOKS LIVE QUE
- WA ARE MORE INTERESTING. TALKS TO BUCE) CHAPGED OBJECTS: L & H

$$\sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 2 & -2 \end{pmatrix}, \quad \sum_{i=1}^{N} \frac{1}{2} \begin{pmatrix} 1 & 1$$

observe? recall (HARMONIC OSCILLATOR) ON DEPONE:

$$T = \frac{1}{2}\sigma^{-2} = \frac{1}{2}(\sigma^{1} - i\sigma^{2}) = (10)$$

Why this should look terrilier:

T+(?) = ( b) C"1+>" " (11)"

RAISING OFFRATER

 $T - \binom{1}{6} = \binom{6}{6}$ 

WHERING OP.

 $T^3\begin{pmatrix}a\\b\end{pmatrix}=\begin{pmatrix}a\\-b\end{pmatrix}$ 

a front, "H abanta. 42k2: MAYL 12 ADAS CHARGE UNDER SHEETS

unsurprisingly, these are related by commutation relations

... in turn, related to how W's talk to each other

the raising Nowering structure in the TA matrices is an intimate part of what it means to be non-Abelian.

SU(8) is more complicated. other groups are more complicated stil ...

but the structure of the symmetry is encoded in now the T's commute will each other.

m more [mternal] directrons lowering

In the same way, one may write

 $\mathcal{W}^{\pm} = (\mathcal{W}' \pm i\mathcal{W}^2) / \mathcal{I}_{\mathcal{Z}}$ C womalize

> \$ ; con have charge

Why would you to this? we'll see ... but e this level we're just picking a weird basis.

WHAT BOES THE # CHARGE REFER TO?			
$1 \text{ w}^2$ $+ i \text{ w}^2$ Wike	$ W^1\rangle \pm i W^2\rangle$		
	mixed state		
neither of these have I charge.	<b>e</b>		
BUT WE KNOW WIT ~ EA	BC ( PM WAS + )		
WE WE T	details here don't matter, anly		
totally ontisym.	that we can form an invariant		
80: A31 8=2 C=8	- W. M.		
	ore homes M3		
AHL IT wors like W' is OHA remember this week's hom			
(T3)9 6 ~ (1-1)			
acts on doublets like a c each component			
e e = (H') = (	e10 H2 )		
	1		
in fo	act: ABECIAN!		
observe: there is a U(1) sym	imetry inside su(2)		

## THIS FAR:

- nothing has. mass
- · fermions (spin-1/2) don't have mass because their chinal quantum this are respected

( d' à maires are separate.

- . Hrags (spin-o) could have mass. we're not touching that yet.
- · We have chosen a work basis for W's

6 W1,2,3 - W±, W3

breaks the abrievs sym. but a this point, sym is still there.

HOW THE W INTERACTS WI DOURLETS:

$$(T'+)^{9} = (0)$$

 $W^{\dagger} = W^{\dagger} + 1 \text{ otherwise}$   $\begin{cases} W^{\dagger} & W^{\dagger} \\ W^{\dagger} & W^{\dagger} \end{cases}$ 

 $H_{t}$ 

for LEPTON: W+	W-	<b>υ</b> .	
vit e. e.			
now we're getting comeaniere! like a theory that we've !	seen be	torc.	
QUESTION: What are allowed this	185 - Jeri	Mi'eV	
LORENTZ: either other or	84	80%	
h resultes some 4-momentum (as gai know 4-m ths) H	Eas	£ **	
X Pt Sh			
las den knom trow An	3 SU(2) 01)		
BUT: HYPER CHARGE?			
Ya-13	does not Y not	work!	ed ! !
Lt Y:1 FJ			

so this vertex is not allowed.

6ther Lorentz structure contract arral spinors (H+) a Ha LY=42 (UPWACD) HL. (=- } 7=1 LBB 回く (百十) 沒 ENB Sis Esi Sa ope; duone on shinos hit each other! LET'S UNPACK THE DOUBLETS (not shown: 14-c.) (升)<sup>S</sup> (H+)1 2 1 call this ept : E wend interaction  $\mathcal{M}$ very interesting connects very hamself "electron" to 124. We howen't yet Shown that's

these are sur femonte particles 121

HERE'S WHERE IT RET'S GRAZY - The BIG TWIST ATHE COASS

new rule: the Higgs line can end.

racuum expectatras

SPECIFICALLY: H - (a+ib)

this thing "ends'

IMPLICATION S

(H+) =0

e. 5-e

E

ez + x > ez

in some frames
it's LH, in
strers it is
pett... sot a
chirality not a
good growtum #!

("chiral sym breaking")

SMIXES EL = EÉ (MIXING CHIPALITIES) SPIN 1/2 LH => PA+ ... HUS IS (MSS).