

Preliminaries

"PERIODIC TABLE"

↳ 1st gen matter
4 fund. forces
HIGGS

+ ELECTRIC CHARGES

RULES for QED + Z + γ

HW1b: #1 ERIC
#2 NICK

↑ e-less \rightarrow e#
 μ -less \rightarrow μ #

NEW RULES:

P^+ : uud
 $\approx 1 \text{ GeV}$

<div style="border: 1px solid black; padding: 5px; display: inline-block;"> $U \quad +2/3$ $d \quad -1/3$ </div>		$\times 3$
		$\approx 5 \text{ MeV}$ $\approx 5 \text{ MeV}$
$\nu_e \quad 0$		$\approx 1 \text{ eV?}$
$e \quad -1$		$\approx \frac{1}{2} \text{ MeV}$

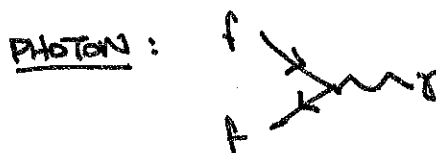
SPIN $1/2$

γ	0
Z	$\approx 100 \text{ GeV}$
W^\pm	$\approx 100 \text{ GeV}$
g	0

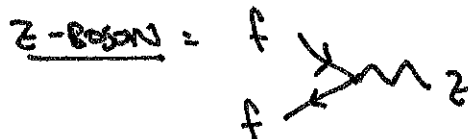
SPIN-1

h
 $\approx 100 \text{ GeV}$

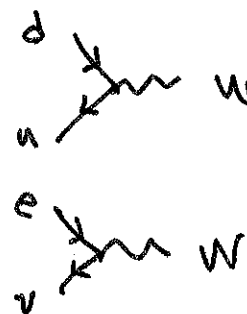
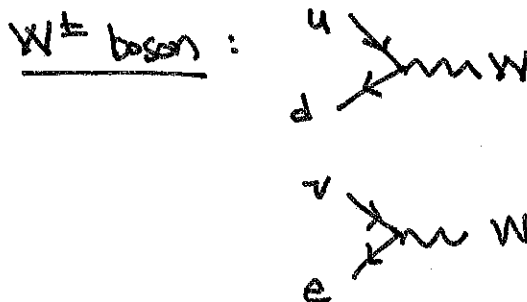
SPIN-0



for: $f = u, d, e$ (anything w/ electric charge)



for: $f = u, d, \nu, e$
↳ is there "Z charge"?



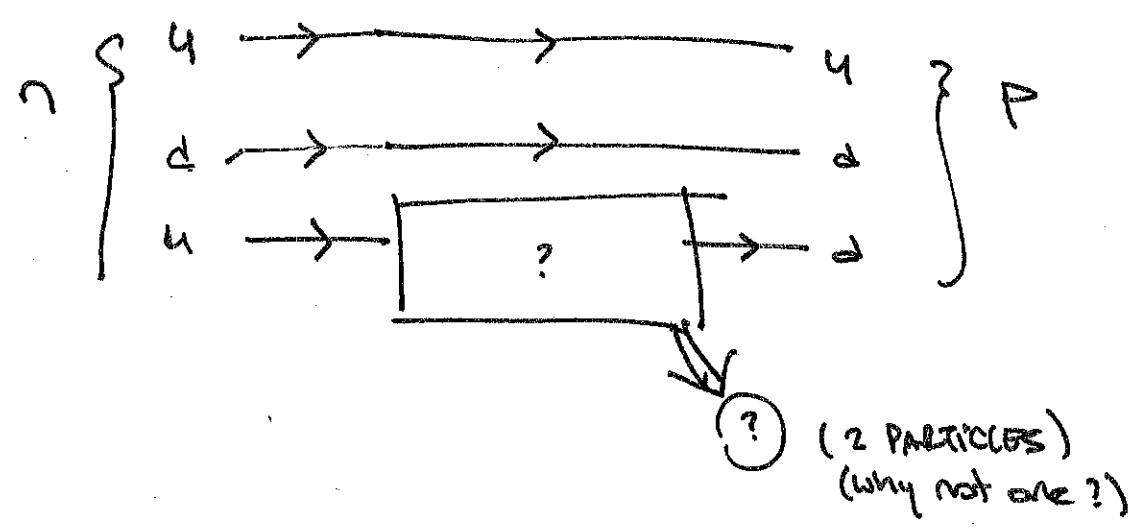
CHECK: CHARGES. (W^- GOING IN = W^+ OUT)

ANNOUNCE:

KUNTAI on THU.
CROSS SEC.

Note card :

draw the diagram for neutron decay



Q: why not $P \rightarrow n + (?)^*$

LAST TIME vectors & tensors

1. DEFINE A SYMMETRY eg 2D ROTATIONS ($SO(2)$)
2. there are UPPER ^{COL VECTOR} and LOWER ^{ROW VECTOR} INDICES

nb: continuous symmetries only ("Lie Groups")

- 3 there are ROTATIONS & "anti-ROTATIONS"

$$R^i_j$$

$$\bar{R} = R^T$$

1st upper,
2nd lower

$$\text{for 2D ROT: } \bar{R}^i_j = R^j_i$$

4. given an object w/ indices, each
upper index transforms w/ R
lower index ————— \bar{R}

reminder : REPEATED UPPER & LOWER INDICES
MEAN A SUM.

eg. ROTATING A VECTOR

$$R^i_j = \begin{pmatrix} c_\theta & s_\theta \\ -s_\theta & c_\theta \end{pmatrix}$$

eg this is R^1_2

$$V^k = \begin{pmatrix} V^1 \\ V^2 \end{pmatrix} \quad \leftarrow \text{UPPER INDEX}$$

$$V^k \rightarrow (RV)^k = R^k_j V^j = R^k_1 V^1 + R^k_2 V^2$$

↑
DUMMY INDEX

this just reproduces matrix mult.

eg. ROW VECTORS TRANSFORM "OPPOSITELY"

$$W_k \rightarrow (W\bar{R})_k = \bar{R}^j_k W_j = \bar{R}^1_k W_1 + \bar{R}^2_k W_2$$

↑
in "matrix" notation,
ORDER MATRICES.

↑
no order DOESN'T MATTER,
these are "just #'s"

$$\bar{R} = \begin{pmatrix} c_\theta & -s_\theta \\ s_\theta & c_\theta \end{pmatrix}$$

$$(W\bar{R})_1 = \bar{R}^1_1 W_1 + \bar{R}^2_1 W_2$$

$$= c_\theta W_1 + s_\theta W_2$$

$$(W\bar{R})_2 = -s_\theta W_1 + c_\theta W_2$$

5. AN OBJECT W/ NO INDICES AFTER CONTRACTIONS
IS INVARIANT W/RT THE SYMMETRY.

↑ REPEATED
UPPER/LOWER

eg. $W_k V^k = W_1 V^1 + W_2 V^2$

$$(W_1, W_2) \begin{pmatrix} V^1 \\ V^2 \end{pmatrix} = \underline{W}^T \underline{V}$$

if WE TRANSFORM:

$$(W\bar{R})_k (R V)^k = (cW_1 + sW_2, -sW_1 + cW_2) \begin{pmatrix} cV^1 + sV^2 \\ -sV^1 + cV^2 \end{pmatrix}$$

$$= c^2 W_1 V^1 + c s W_1 V^2 + c s W_2 V^1 + s^2 W_2 V^2$$

$$+ s^2 W_1 V^1 - c s W_1 V^2 - c s W_2 V^1 + c^2 W_2 V^2$$

$$= W_1 V^1 + W_2 V^2 \quad \checkmark$$

alternatively: $W_j \bar{R}^j_k R^k_l V^l = W_j V^j \quad \checkmark$

BUT $\bar{R} = R^T = R^{-1}$

UNITARY
(ORTHOGONAL)

$$\bar{R}^j_k R^k_l = \delta^j_l$$

$$\begin{pmatrix} c & -s \\ s & c \end{pmatrix} \begin{pmatrix} c & s \\ -s & c \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

object w/ more indices:

$$(\ddot{})^{ik} \rightarrow R^i_m R^k_n (\ddot{})^{mn} \\ = R^1_1 R^1_1 (\ddot{})^{11} + R^1_2 R^1_1 (\ddot{})^{21} + \dots$$

NEW OBJECT: METRIC (\dagger inverse metric)

g_{ij} : takes upper index, converts to lower index

$$\boxed{g_{ij} V^j \equiv V_i}$$

similarly: inverse metric

$$\boxed{g^{ij} V_j \equiv V^i} \quad \dagger \quad g^{ij} g_{jk} \equiv \delta^i_k$$

for \mathbb{R}^2 : $g_{ij} = \delta_{ij}$ $g^{ij} = \delta^{ij}$

these are different!
different still from δ^i_j !

from $\mathbb{R}^2 \rightarrow$ Minkowski SPACE
 SPACETIME

$\mathbb{R}^2 \rightarrow \mathbb{R}^3$: R^i_j is more complicated.
 \bar{R}^i_j is still $\bar{R} = R^T$
 g_{ij} is still δ_{ij} (if so forth)
 ... but $i, j = 1, 2, 3$

from $\mathbb{R}^1 \rightarrow \mathbb{R}^{1,1}$: \nearrow one space, one time

$R^i_j \rightarrow \Lambda^\mu_\nu \leftarrow$ BOOST/ROT.

$$\Lambda^\mu_\nu = \begin{pmatrix} \gamma & -\gamma\beta \\ -\gamma\beta & \gamma \end{pmatrix} \quad \gamma = \frac{1}{\sqrt{1-\beta^2}}$$

$\bar{\Lambda}$ is now Λ^{-1} A LITTLE DIFFERENT
 (most of the time its DUAL)
 SPACETIME SYM.

$$\bar{\Lambda}^\mu_\nu = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix}$$

$$g_{ij} \rightarrow \eta_{\mu\nu} = \begin{pmatrix} 1 & \\ & -1 \end{pmatrix}$$

eg $P^\mu = \begin{pmatrix} E \\ \mathbf{p} \end{pmatrix}$

$$P_\mu = \eta_{\mu\nu} P^\nu = (E, -\mathbf{p})$$

$$P^\mu \rightarrow \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} E \\ \mathbf{p} \end{pmatrix} \\ = \Lambda^\mu{}_\nu P^\nu$$

$$P^2 \equiv P_\mu P^\mu = P_0 P^0 + \mathbf{p} \cdot \mathbf{p}$$

$$= E^2 - (-\mathbf{p}) \cdot \mathbf{p}$$

$$= E^2 - p^2 \quad \leftarrow \text{invariant!}$$

transf.



$$\tilde{\Lambda}^\nu{}_\mu P_\nu \Lambda^\mu{}_\sigma P^\sigma = P_\nu \underbrace{\tilde{\Lambda}^\nu{}_\mu \Lambda^\mu{}_\sigma}_{\equiv \delta^\nu{}_\sigma} P^\sigma$$

check:

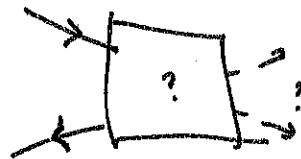
$$\begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} \gamma & -\gamma\beta \\ -\gamma\beta & \gamma \end{pmatrix} = \begin{pmatrix} \gamma^2 - \gamma^2\beta^2 & 0 \\ 0 & \gamma^2 - \gamma^2\beta^2 \end{pmatrix}$$

$$= \begin{pmatrix} \frac{1-\beta^2}{1-\beta^2} & 0 \\ 0 & 1 \end{pmatrix}$$

eg.



← m Feynman Diag.



What is a quantity that
is invariant & conserved?

$$S \equiv (P_{e^-} + P_{e^+})^2 \quad \leftarrow \text{TOTAL 4-momentum, squared}$$

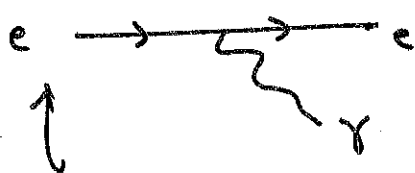
↑ stupid name, but that's what it is

IN CM FRAME: $P_{e^-} = (E^{cm}, p)$

$$P_{e^+} = (E^{cm}, -p)$$

$$S = (2E^{cm}, 0)^2 \\ = 4(E^{cm})^2$$

can an electron spontaneously emit a γ ?



$$P = (E, p)$$

$$S = P^2 = M_e^2$$

$$K_e = (E', k) \quad \leftarrow K_e^2 = M_e^2 \text{ (on shell)}$$

$$K_\gamma = (E_\gamma, E_\gamma \hat{n})$$

↑
IN CM FRAME,
 $P = (M_e, 0)$

$$E_\gamma \hat{n} = -k$$

$$S = (K_e + K_\gamma)^2 = k^2 + K_\gamma^2 + 2k \cdot K_\gamma \\ = M_e^2 + 0 + 2\vec{k} \cdot (-E_\gamma \hat{n}) \\ = M_e^2 - 2E_k E_\gamma < M_e^2$$

INTERNAL SYMMETRIES : generaliz. of ROTATIONS

for us: $SU(N)$ symmetries

$N \times N$ UNITARY MATRICES w/ $\det = 1$

$$\uparrow$$

$$U^\dagger U = \mathbb{1} \quad \text{s.t.} \quad \bar{U} = U^{-1} = U^\dagger$$

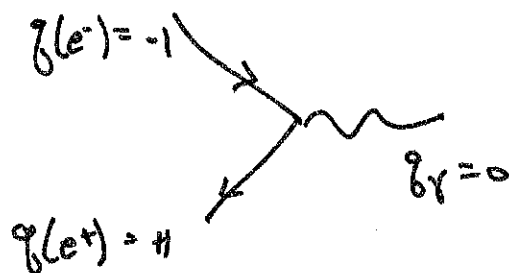
UPPER INDEX: U TRANSF

LOWER INDEX: U^\dagger TRANSF

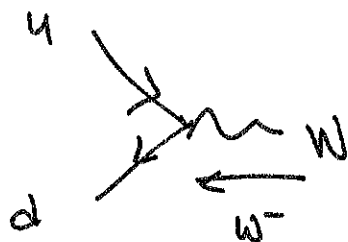
simplest case: $U = e^{i\theta} \leftarrow N=1$

RULE: a particle w/ charge q going into a vertex
picks up a phase $e^{i\theta q}$.

RULE: if a symmetry is good, then
the phases going into a vertex
must be zero.



$$e^{i\theta \cdot (-1)} e^{i\theta \cdot (+1)} e^{i\theta \cdot 0} = 1$$



$$q_u + q_d + q_{W^-} = 0 \quad \checkmark$$