

PHYSICS 208: GENERAL RELATIVITY

10 Jan

WEBSITE: faculty.ucr.edu/~flipt/P208-2016.htmlLECTURES: PHY 2104 TUE/THR 5:10 - 6:30Can we move
to reading or
conf. room?→ PROBLEM: missing 3 lectures
could do make up over exams?
or GOOGLE HANGOUTS?
or no make up, just work?GRADING / HW / EXAMS

- EVERYONE HERE IS A GRAD STUDENT TAKING THIS CLASS TO LEARN; I DON'T INTEND FOR ANYONE TO BE PENALIZED FOR THAT.

- NO EXAMS

- GRADING (A or B) BASED ON HW

WEEKLY,
DUE TUESUBJECT TO
CHANGE!

↑

↓

PLEASE DO MINIMUM work! (>50% HW)

→ WILL BE GRADED LOOSELY, BUT WILL BE IMPORTANT FOR ME TO CALIBRATE

- HOMEWORK WILL INCLUDE READING - PLEASE DO THIS, I WILL EXPECT YOU TO HAVE SEEN THE ASSIGNED MATERIAL AHEAD OF CLASS,
- AS YOU ALREADY KNOW: MOST OF YOUR LEARNING WILL COME FROM DOING PROBLEMS & DISCUSSING W/ EACH OTHER!!

LECTURE IS A WAY TO GUIDE THIS LEARNING,
BUT CANNOT REPLACE IT.

DISCUSSION/OH :

- no discussion see.
- office hrs: the 30 min before or after class
→ by appt.

(AFTER 5:30 PM ON MOST DAYS
I WILL BE AVAILABLE)

TEXTBOOK : official choice is HARTLE
see website for other
suggestions — I DON'T CARE
WHICH YOU USE.

TODAY: Assessment ("judgement free")
NO NEED FOR NAMES, FEEL FREE TO DISCUSS.

no gravity, no
curvatureREVIEW: SPECIAL RELATIVITY

↑ SPEED OF LIGHT IS CONSTANT

... & IMPLICATIONS OF THIS

RECALL THE USUAL IDEAS

- LENGTH CONTRACTION & TIME DILATION / $\vec{E} \leftrightarrow \vec{B}$
↑ "watching a train pass by"
- SIMULTANEITY IS NOT ABSOLUTE

→ gives all sorts of cute "paradoxes" that test our intuition

UNITS: I WILL ALMOST ALWAYS WORK IN UNITS WHERE $\boxed{c=1}$ (NATURAL UNITS)

YOU CAN USE DIMENSIONAL ANALYSIS'S TO UNIQUELY RESTORE FACTORS OF c TO GIVE "unnatural units" (eg MKS)

so I can write things like

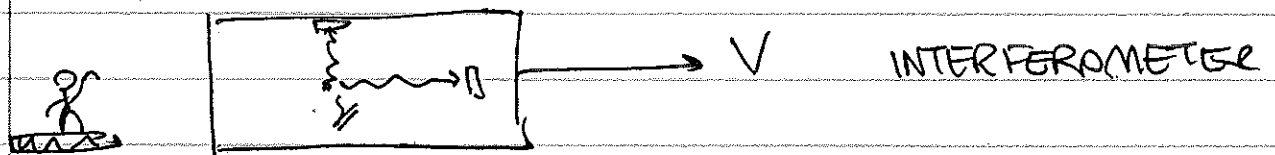
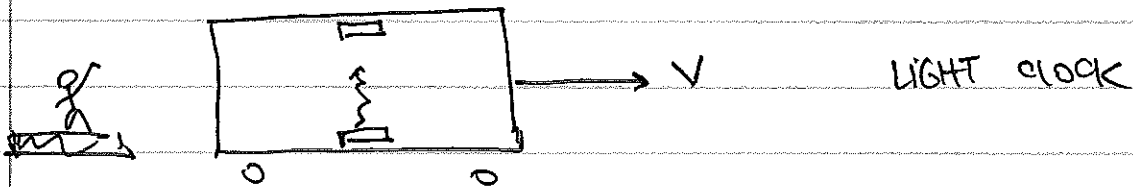
$$ds^2 = dt^2 - dx^2$$

in natural units, these both have dim. of $(\text{length})^2$

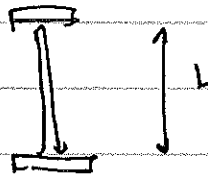
clocks & rulers on
moving train

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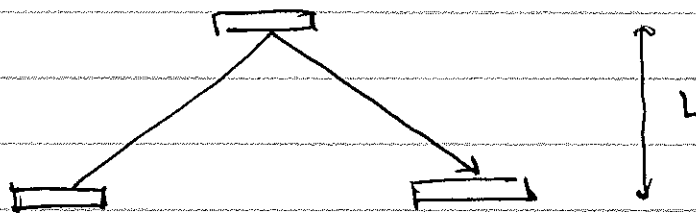
RESULT OF "TRAIN ANALYSIS": (see HW)



Key point: in train frame



in train station frame



BUT IN BOTH FRAMES, SPEED OF LIGHT
IS CONSTANT.

for v along \hat{x} direction

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← TIME DILATION

LEN. CONTR?? (GET TO THIS LATER)

$$t' = \gamma t - \gamma \beta x$$

$$x' = -\gamma \beta t + \gamma x$$

$$y' = y$$

$$z' = z$$

LORENTZ
BOOST
ALONG \hat{x}

$$\gamma = \frac{1}{\sqrt{1-\beta^2}}$$

$$\beta = v$$

We've set $c=1$, or else $t \rightarrow ct$

WE SIMPLIFY FURTHER } DEFINE

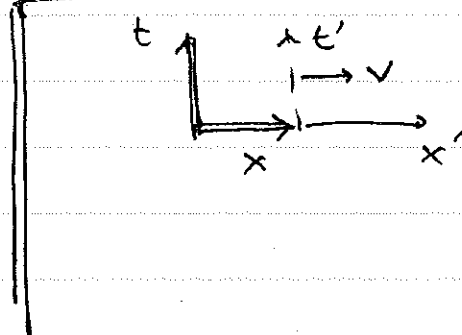
SIGN CONVENTION
IS FOR

$$x^0 = t$$

$$x^1 = x$$

$$x^2 = y$$

$$x^3 = z$$



to understand physics, sufficient
to consider 2D (x^0, x^1) plane.

$$\begin{pmatrix} x^0' \\ x^1' \end{pmatrix} = \underbrace{\begin{pmatrix} \gamma & -\gamma\beta \\ -\gamma\beta & \gamma \end{pmatrix}}_{\Lambda} \begin{pmatrix} x^0 \\ x^1 \end{pmatrix}$$

β can have
either sign.

SCALING: AS $v = \beta$ increases, $(0 \leq |\beta| \leq 1)$

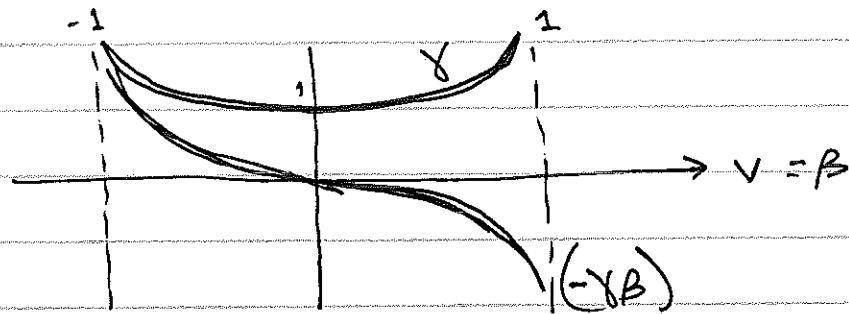
γ increases ; singular @ $\beta = 1$

$\gamma\beta$ increases ; zero @ $\beta = 0$

CHECK: For $\beta = 0$, $\Lambda = 11$. no transf.

"continuously connected to the identity"

SO PURELY FROM ASYMPTOTICS, WE HAVE



SO A BOOST IN \hat{x} DIRECTION LOOKS LIKE

$$\Lambda = \begin{pmatrix} \gamma & -\gamma\beta \\ -\gamma\beta & \gamma \end{pmatrix} \quad \begin{matrix} 1 \leq \gamma < \infty & -\infty < \gamma\beta < \infty \end{matrix}$$

infinitesimal boost: ($\beta \ll 1$)

$$\Lambda - 1 \approx \begin{pmatrix} \mathcal{O}(\beta^2) & -\beta \\ -\beta & \mathcal{O}(\beta^2) \end{pmatrix}$$

↑
this infinitesimal transf.
is called a generator
of the boost

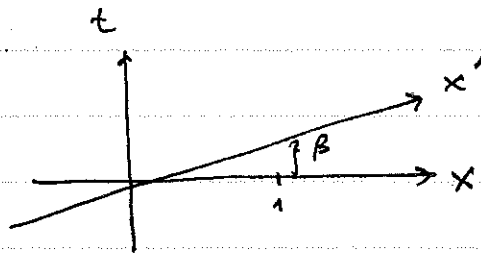
GIVES GROUP STRUCTURE OF
LORENTZ SYMMETRY

WHAT DO (x', t') COORDINATES LOOK LIKE ON (x, t) PLANE?

PLOT (x', t') AXES

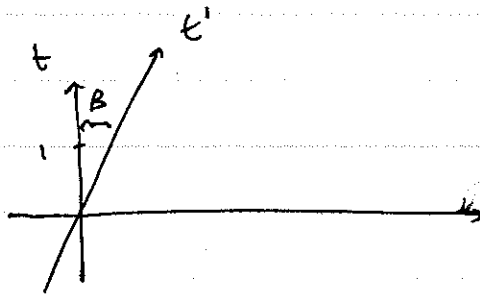
$$x' \text{ AXIS} \leftrightarrow t' \equiv 0 = \gamma t - \gamma \beta x$$

$$t = \beta x$$



$$t' \text{ AXIS} \leftrightarrow x' \equiv 0 = \gamma x - \gamma \beta t$$

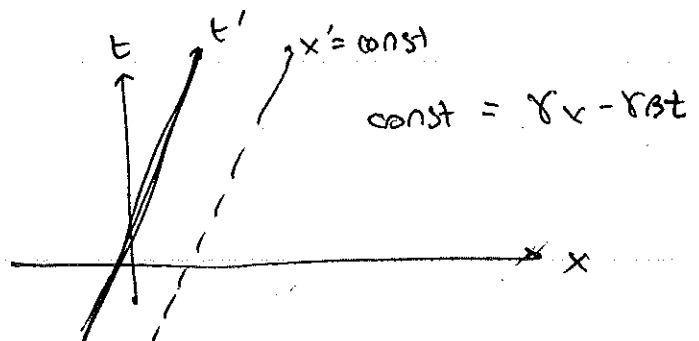
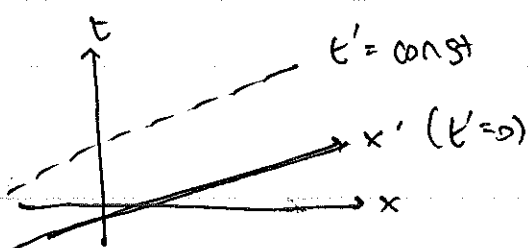
$$x = \beta t \quad (\text{or } t = \frac{1}{\beta} x)$$



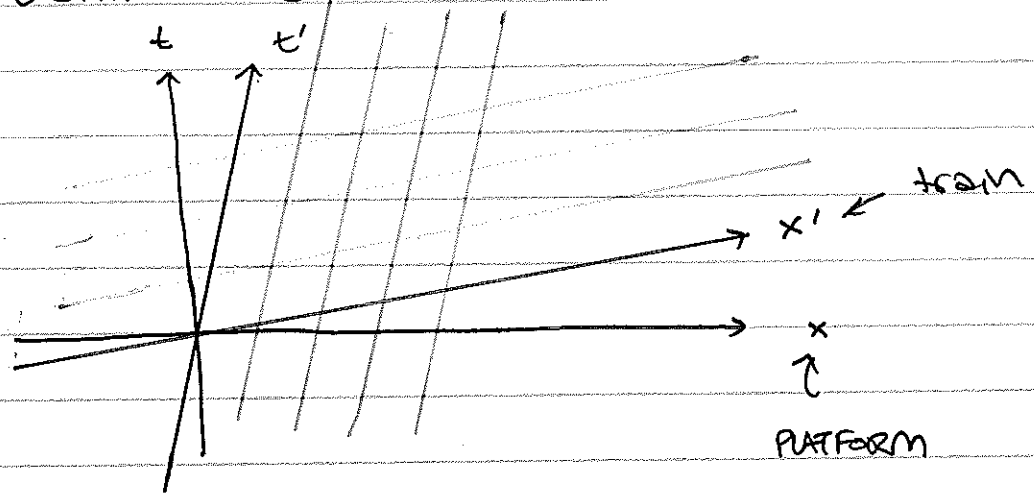
shifts of axes
↓

INES OF CONSTANT x' OR t' ARE SIMILAR

$$\text{const} = \gamma t - \gamma \beta x$$



AND SO, FOR A BOOST IN THE $+\hat{x}$ DIRECTION,
A CARTESIAN GRID IS TRANSFORMED TO :

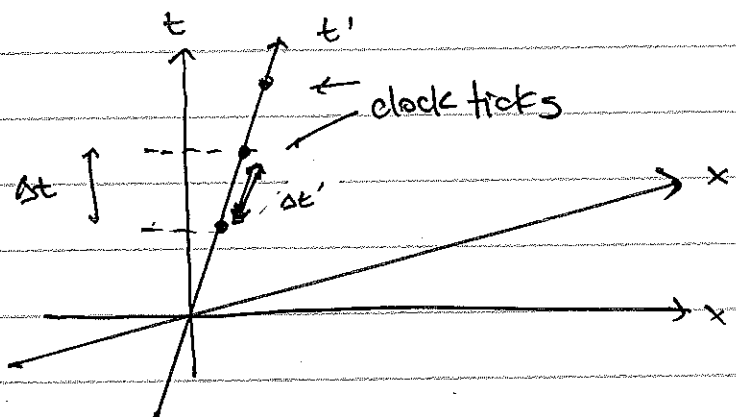


THIS CONTAINS ALL THE INFORMATION OF THE BOOST.

TIME DILATION

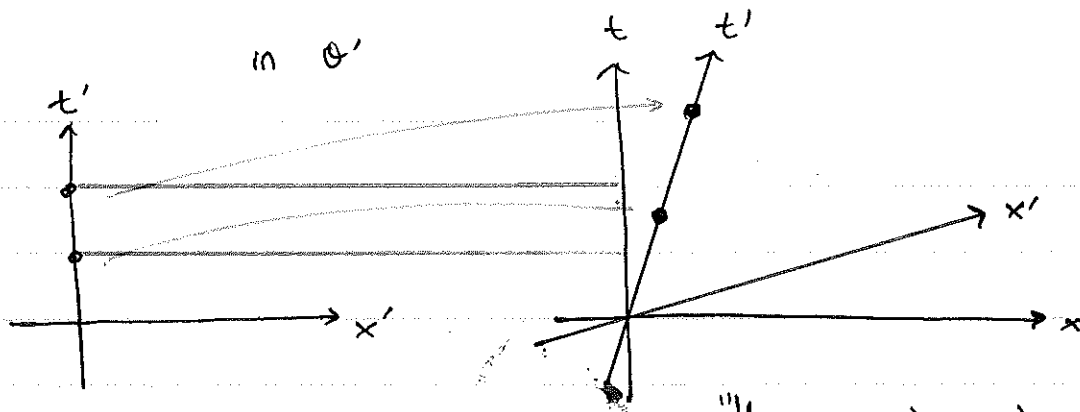
LET \mathcal{O}' BE FRAME OF CLOCK (on train)

$\Delta x' = 0$ \leftarrow clock is stationary in train



hey- that doesn't look like
time dilation?!

$$\begin{array}{c} \beta \\ 1 \end{array} \quad \sqrt{1+\beta^2} > 1$$



"this is not just a rotation"

SANITY CHECK

t' AXIS

$$x' = 0 = -\gamma \beta t + \gamma x \Rightarrow \gamma x = \gamma \beta t$$

$$\text{then: } t' = \gamma t - \gamma \beta x$$

$$= \gamma t - \gamma \beta^2 t$$

$$= \frac{1 - \beta^2}{\sqrt{1 - \beta^2}} t$$

$$= \frac{1}{\gamma} t$$

$$\leftarrow \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

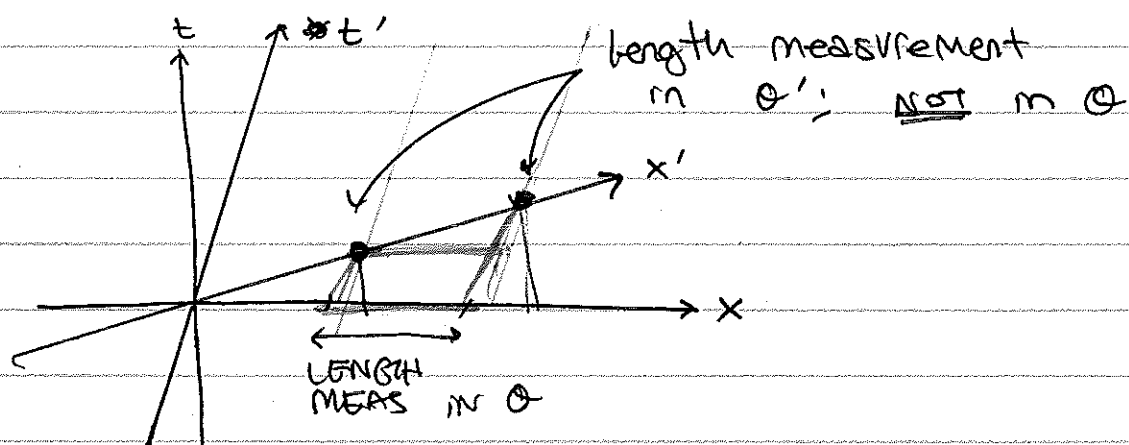
$$\boxed{t = \gamma t'} > t'$$

[see mut-hyperbolae]

LENGTH CONTRACTION : more subtle

↓ "events" in spacetime

LENGTH MEASUREMENT IS 2 POINTS
 (t_1, x_1) & (t_2, x_2) s.t. $t_2 - t_1 = 0$



OBSERVE ASYMMETRY

TIME DILATION : $\Delta x' = 0$

not $\Delta t' = 0$

LENGTH CONTRACTION : $\Delta t = 0$

[maybe I should expand this section]

SIMULTANEITY → out the window

↳ this is why we must work
 in spacetime vs. space & time

WANT TO GET BACK TO THIS FROM GEOMETRIC POV -
FIRST, REVIEW "TRIVIAL" GEOMETRY

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GEOMETRY: REVIEW OF 2D EUCLIDEAN SPACE

↑ "ordinary flat space"

VECTORS: WE'LL USE INDICES (mathematicians hate this)

$$\underline{V} = V^i \underline{e}_{(i)} \quad \leftarrow \text{SUMMATION NOTATION}$$

↑ ABSTRACT OBJECT
REPRESENTING
THE (i) UNIT VECTOR

$$V^i \underline{e}_{(i)} \equiv V^1 \underline{e}_{(1)} + V^2 \underline{e}_{(2)} + \dots$$

eg if $\underline{V} = \begin{pmatrix} 1 \\ 3 \end{pmatrix}$, then this is:

$$\underline{V} = \underset{\substack{\nearrow \\ V^1}}{1} \underbrace{\begin{pmatrix} 1 \\ 0 \end{pmatrix}}_{\underline{e}_{(1)}} + \underset{\substack{\nearrow \\ V^2}}{3} \underbrace{\begin{pmatrix} 0 \\ 1 \end{pmatrix}}_{\underline{e}_{(2)}}$$

the components of a vector have upper indices.
sometimes we'll be sloppy & identify
 V^i with the vector itself.

WE CAN TAKE 2 VECTORS & MAKE A NUMBER
that is: there is a linear map

$$g: V \times V \rightarrow \mathbb{R}$$

vector space (\mathbb{R}^2 in this case)

EUCLID.
FLAT SPACE
METRIC

$$g(\underline{v}, \underline{w}) \equiv v_i w^i = v_1 w^1 + v_2 w^2$$

LOWER INDEX OBJECT: dual vector
row vector
etc.

can think of lower index object
as $g(\underline{v}, \cdot) \leftrightarrow v_i$

function: $V \rightarrow \mathbb{R}$

so v_i is a linear function of vectors
 $v_i(w^j) = v_i w^j$

gives the usual dot product.

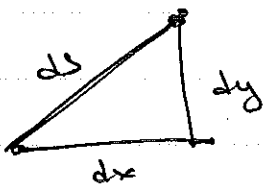
$$\underline{v} \cdot \underline{w} = \underline{v}^T \underline{w} = v_i w^i = v^1 w^1 + v^2 w^2$$

"matrix"
mult

stop thinking like this

LET US WRITE THIS IN A PECULIAR WAY

$$ds^2 = dx^2 + dy^2$$



infinitesimal length
between 2 points

$$= g_{ij} dx^i dx^j$$

$$\uparrow g_{ij} = \begin{pmatrix} 1 & \\ & 1 \end{pmatrix}$$

METRIC

(related to measuring
distances)

Then: $V_i = g_{ij} V^j$ METRIC LOWERS INDICES

$$g_{ij} V^i W^j = V^1 W^1 + V^2 W^2$$

METRIC AS A MAP: $V \times V \rightarrow \mathbb{R}$

$$g^{ij} g_{jk} = \delta^i_k \quad \text{metric} \rightarrow \text{inverse}$$

$$g^{ij} V_j = V^i \quad \text{inverse metric raises indices}$$

VERY LOOSELY can think of lower index object as a ROW VECTOR.

$$V_i W^i = \underbrace{(V_1 \ V_2)}_{\vec{V}^T} \begin{pmatrix} W^1 \\ W^2 \end{pmatrix} = V_1 W^1 + V_2 W^2$$

(this is a crutch; BETTER: ROW VECTOR IS AN EXAMPLE OF A LOWER INDEX OBJECT ON EUCLIDEAN SPACE)

GIVEN THE METRIC OF EUCLIDEAN SPACE, THERE ARE SPECIAL TRANSFORMATIONS THAT LEAVE 'DOT PRODUCTS' UNCHANGED:

$$(V_1 \ V_2) \begin{pmatrix} C_\theta & -S_\theta \\ S_\theta & C_\theta \end{pmatrix} \begin{pmatrix} C_\theta & S_\theta \\ -S_\theta & C_\theta \end{pmatrix} \begin{pmatrix} W^1 \\ W^2 \end{pmatrix}$$

$$R^i_j V_j \rightarrow (R_\theta \underline{V})^T \quad (R_\theta \underline{W}) \leftarrow R^i_j W_j$$

ROTATION, R_θ , CAN BE DEFINED AS A TRANSFORMATION THAT LEAVES $V_i W^i$ INV.

USEFUL: OBJECTS w/ INDICES TRANSFORM

UPPER INDEX i : R^i_j ~~V^i_j~~

LOWER INDEX i : R^j_i ~~V_j~~

"
 ~~V_j~~ R^j_i ;

explicitly $R^i_j = \begin{pmatrix} R^1_1 = C_0 & R^1_2 = S_0 \\ R^2_1 = -S_0 & R^2_2 = C_0 \end{pmatrix}$

"COLUMN VEC" : $V^i \rightarrow R^i_j V^j = R^1_1 V^1 + R^1_2 V^2$
 $= \begin{pmatrix} R^1_1 V^1 + R^1_2 V^2 \\ R^2_1 V^1 + R^2_2 V^2 \end{pmatrix}$
 $= \begin{pmatrix} C_0 V^1 + S_0 V^2 \\ -S_0 V^1 + C_0 V^2 \end{pmatrix}$

"ROW VEC" : $V_i \rightarrow V_j R^j_i = V_1 R^1_i + V_2 R^2_i$
 $= \begin{pmatrix} V_1 R^1_1 + V_2 R^2_1, \\ V_1 R^1_2 + V_2 R^2_2 \end{pmatrix}$
 $= \begin{pmatrix} C_0 V_1 - S_0 V_2, \\ S_0 V_1 + C_0 V_2 \end{pmatrix}$

FOR A MORE GENERAL OBJECT

(p, q) tensor: $T^{i_1 \dots i_p}_{j_1 \dots j_q}$

multi-linear map
in the sense that

$$T: \underbrace{V \times \dots \times V}_q \times \underbrace{V^* \times \dots \times V^*}_p \rightarrow \mathbb{R}$$

q vectors p dual vec

$$T^{i_1 \dots i_p}_{j_1 \dots j_q} V_{(1)}^{j_1} \dots V_{(q)}^{j_q} W_{(1)}^{i_1} \dots W_{(p)}^{i_p}$$

HOW DOES IT TRANSFORM?

R^i_j for EACH UPPER INDEX $T^{i_1 \dots i_p}_{j_1 \dots j_q}$

R^j_i for EACH LOWER INDEX $T^{i_1 \dots i_p}_{j_1 \dots j_q}$

↑
GENERALIZE TO ROTATIONS IN, say, \mathbb{R}^3

IMPORTANT: THINGS w/ NO FREE INDICES DO NOT TRANSFORM (not frame-dep.)

$$V \cdot W = \underbrace{V^i W_i}_{\text{no free indices}} \text{ is invariant}$$

↑
same under rotations

GEOMETRY & RELATIVITY: MINKOWSKI SPACE

LORENTZ TRANSFORM

$$\Lambda = \begin{pmatrix} \gamma & -\gamma\beta \\ -\gamma\beta & \gamma \end{pmatrix} \leftarrow \text{does not look like a rotation!}$$

WHAT TYPE OF METRIC (WHAT KIND OF SPACE) IS INVARIANT UNDER Λ ?

PROPOSE: $ds^2 = dt^2 - dx^2$

$$\rightarrow (\gamma dt' - \gamma\beta dx')^2 - (-\gamma\beta dt' + \gamma dx')^2$$

$$= \gamma^2 (dt')^2 - 2\gamma^2\beta dt' dx' + \gamma^2\beta^2 (dx')^2 - \gamma^2\beta^2 (dt')^2 + \underline{\quad} - \gamma^2 (dx')^2$$

\uparrow
note: $\gamma^2(1-\beta^2) = 1$

$$= (dt')^2 - (dx')^2 \quad \checkmark$$

MINKOWSKI SPACE (FLAT SPACETIME)

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REDERIVE LORENTZ TRANSFORMATION FROM GEOMETRY

$$ds^2 = \eta_{\mu\nu} dx^\mu dx^\nu$$
$$\equiv dt^2 - dx^2$$

$$x^\mu = \begin{pmatrix} x^0 \\ x^1 \end{pmatrix} = \begin{pmatrix} t \\ x \end{pmatrix}$$

when this was a + sign in euclidean space, we recognized that $ds^2 = \text{const}$ gives a circle.

for MINKOWSKI SPACE,

ds^2 is a hyperbola.

so USE HYPERBOLIC TRIG FUNCTIONS

$$\cosh x = \frac{1}{2}(e^x + e^{-x})$$

$$\sinh x = \frac{1}{2}(e^x - e^{-x})$$

$$\boxed{\cosh^2 x - \sinh^2 x = 1}$$

INDEED

$$\begin{pmatrix} t \\ x \end{pmatrix} \rightarrow \begin{pmatrix} \cosh \eta & \sinh \eta \\ \sinh \eta & \cosh \eta \end{pmatrix} \begin{pmatrix} t \\ x \end{pmatrix}$$

is an invariance

~~is~~

η is transformation parameter.

check: let $g = \begin{pmatrix} t \\ x \end{pmatrix} \xrightarrow{\Lambda} \begin{pmatrix} ct + sx \\ st + cx \end{pmatrix}$ 1.9

where $c = \cosh R$
 $s = \sinh R$

the length is $\|g\|^2 = g \cdot g = t^2 - x^2$

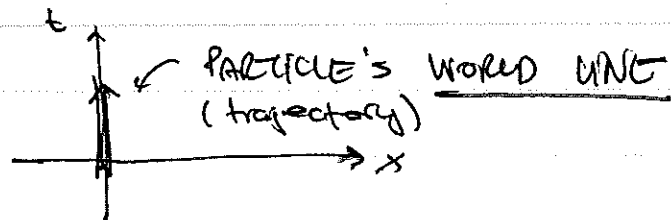
$$\begin{aligned}
 &= (ct + sx)^2 - (st + cx)^2 \\
 &= c^2 t^2 + 2cs tx + s^2 x^2 \\
 &\quad - s^2 t^2 - 2cs tx + c^2 x^2 \\
 &= t^2 - x^2 \quad \checkmark
 \end{aligned}$$

BUT WHAT TO MAKE OF R ?

CONNECT TO PHYSICS

IN THE REST FRAME OF A PARTICLE,

$$\Delta x / \Delta t = 0$$



NOW DO LORENTZ TRANSFORM:

$$\begin{pmatrix} t' \\ x' \end{pmatrix} = \begin{pmatrix} t \cosh R + x \sinh R \\ t \sinh R + x \cosh R \end{pmatrix}$$

ORIGIN IS NOT TRANSFORMED.

BUT FOR $\Delta t \neq 0$

$$\Delta x = 0$$

$$\begin{pmatrix} \Delta t' \\ \Delta x' \end{pmatrix} = \begin{pmatrix} \Delta t \cosh R \\ \Delta t \sinh R \end{pmatrix}$$

then $\boxed{\frac{\Delta x'}{\Delta t'} = \tanh R} \rightarrow R \text{ is } \underline{\underline{\text{RAPIDITY}}}$

VELOCITY, $v = \beta$

next use :

$$\cosh^2 R - \sinh^2 R = 1$$

$$1 - \frac{\tanh^2 R}{\beta^2} = \frac{1}{\cosh^2 R} \quad \leftarrow \begin{array}{l} \text{top left} \\ \text{elem of} \\ \Lambda. \end{array}$$

$$\Rightarrow \cosh^2 R = \frac{1}{1 - \beta^2}$$

$$\Lambda = \begin{pmatrix} \cosh R = \frac{1}{\sqrt{1 - \beta^2}} & 0 \\ 0 & \frac{1}{\sqrt{1 - \beta^2}} \end{pmatrix}$$

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then use: $\sinh R = \tanh R \cosh R$

\uparrow \uparrow
 $\pm \beta$ γ

$$\Lambda = \begin{pmatrix} \gamma & \pm \gamma \beta \\ \pm \gamma \beta & \gamma \end{pmatrix}$$

\uparrow sign ambiguity corresponds
to which way we
boost.

(Analogous to which
direction we rotate θ)