AAE 538: Air-Breathing Propulsion

Lecture 5: Fundamentals of Compressible Flow (continued)

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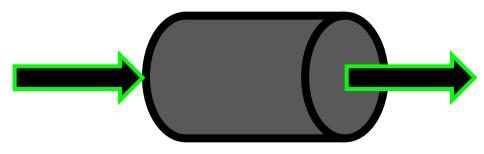


Extensions of 1-D Analysis

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Flow with Friction

- We've examined isentropic flows and shocks, which are iso-energetic discontinuities that cause a change in the flow entropy and a loss in total pressure. Now, we extend our analysis to compressible flows with friction.
- _____are a canonical case where the duct area is constant and the flow is adiabatic, but friction forces are present throughout.
 - Clearly many propulsion problems involve duct flow with friction
 - •
 - For now, we will develop the treatment of friction using these assumptions to highlight the effects on the thermo-physical properties of the flow.
- The addition of friction to the flow makes the process irreversible, and leads to a reduction in the stream thrust in the axial direction.





Beginning our analysis with the force-momentum balance

$$p_{o,1}A_1G_1 + F_x = p_{o,2}A_2G_2$$

Dividing the conservation of momentum expression by the

 From the energy equation, for an Adiabatic flow with no work transfer, we know that the ______.

$$T_{o,1} - \frac{\dot{W}}{\dot{m}c_p} + \frac{\dot{Q}}{\dot{m}c_p} = T_{o,2}$$

Therefore, assuming constant properties, the momentum equation reduces to:



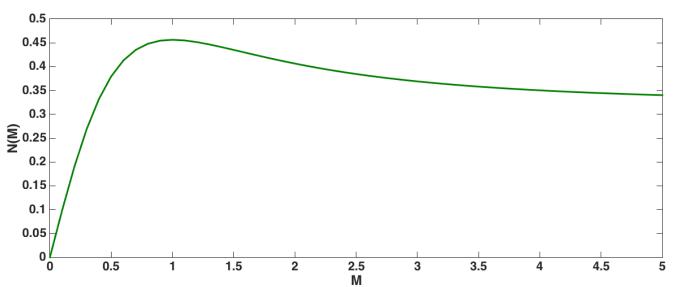
Pausing briefly to examine the implications of this:

$$\frac{1}{N(M_1)} + \frac{F_{\chi}}{p_{o,1}A_1D_1} = \frac{1}{N(M_2)}$$

• The ratio of the outlet (N_2) to inlet (N_1) thrust flow functions is equal to the ratio of the inlet (F_1) to exit (F_2) stream thrust.

In a Fanno Flow,

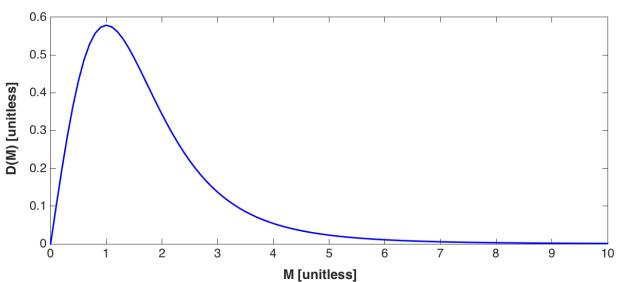




$$N = \frac{M\left(1 + \frac{\gamma - 1}{2}M^2\right)^{\frac{1}{2}}}{(1 + \gamma M^2)}$$

- Thus, friction raises a subsonic Mach number and lowers a supersonic Mach number.
 - 0
- This process is known as _______
 - O That is, for a given inlet flow condition, there is a maximum stream thrust loss (due to high friction of a long duct) before the flow reach M = 1.





The continuity equation for a Fanno flow can be simplified to show:

$$\dot{m} = \left(\frac{p_o A}{\sqrt{T_o}}\right) \sqrt{\frac{\gamma}{R}} D$$

Thus



- This can also be shown through the entropy relation
 - Defined with stagnation properties:

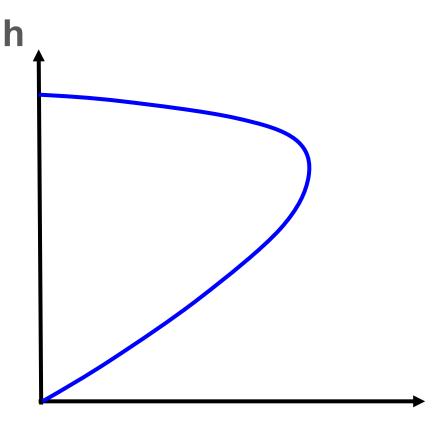
$$s_2 - s_1 = c_p \ln \left(\frac{T_{o,2}}{T_{o,1}} \right) - R \ln \left(\frac{p_{o,2}}{p_{o,1}} \right)$$

0

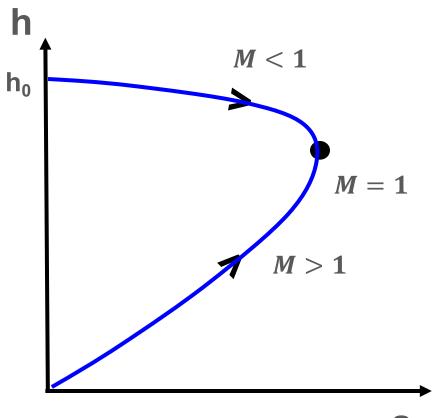
- Static pressure change depends on the starting Mach number of the flow
 - Subsonic:
 - Supersonic:



- By definition, Fanno flow is an iso-energetic with a constant area.
 - 0
 - 0
- Velocity change within the duct as a result friction is caused by a change in entropy
 - Friction can only act to entropy
 - Friction can only to push the flow toward choking.
 - Friction, alone cannot act to transition a flow between the sonic regimes







- The total (integral) amount of friction acting on the flow between 1 and 2 increases with _____.
- If *L* is large enough,
- If L is increased ______ for a given inlet stagnation enthalpy, then it is ______.
 How does it respond?
 - \circ

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- Since friction loss is a path-dependent process, we need to integrate the differential forms of the conservation equations to quantify property changes:
 - For a constant area channel, we find:

$$\frac{dM^{2}}{M^{2}} = \frac{\gamma M^{2} \left(1 + \frac{\gamma - 1}{2} M^{2}\right) f dx}{(1 - M^{2})} \frac{f dx}{D}$$

$$\frac{dp}{p} = \frac{-\gamma M^{2} (1 + (\gamma - 1) M^{2}) f dx}{2(1 - M^{2})} \frac{f dx}{D}$$

$$\frac{dT}{T} = \frac{dh}{h} = \frac{-\gamma (\gamma - 1) M^{4} f dx}{2(1 - M^{2})} \frac{f dx}{D}$$

$$\frac{d\rho}{\rho} = -\frac{du}{u} = \frac{-\gamma M^{2} f dx}{2(1 - M^{2})} \frac{f dx}{D}$$

$$\frac{ds}{R} = -\frac{dp_{o}}{p_{o}} = -\frac{\gamma M^{2} f dx}{2} \frac{f dx}{D}$$



 Integrating between states 1 and 2, we can relate the conditions between the initial and final state:

$$\frac{T_2}{T_1} = \left(\frac{1 + \frac{(\gamma - 1)}{2} M_1^2}{1 + \frac{(\gamma - 1)}{2} M_2^2}\right)$$

$$\frac{p_2}{p_1} = \frac{M_1}{M_2} \sqrt{\frac{T_2}{T_1}}$$

$$\frac{p_{o,2}}{p_{o,1}} = \frac{M_1}{M_2} \left(\frac{T_2}{T_1}\right)^{\frac{\gamma+1}{2(1-\gamma)}}$$

$$\frac{u_2}{u_1} = \frac{\rho_1}{\rho_2} = \frac{M_2}{M_1} \sqrt{\frac{T_2}{T_1}}$$

Property	M < 1	M > 1
Entropy, s		
Stagnation temperature, T _o		
Stagnation pressure, p_o		
Mach number, Ma		
Static Enthalpy, h		
Static Temperature, T		
Static Pressure, p		
Density, $ ho$		
Velocity, u		

Extensions of 1-D Analysis

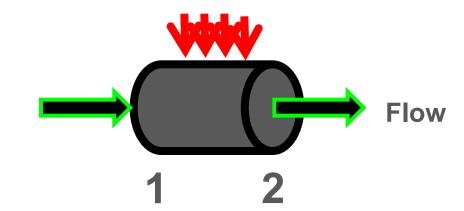
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Flow with Heat Transfer

- Another special case of non-isentropic flow that we need to examine is one where heat is added (or subtracted).
 - Specific applications in air-breathing engines:
 - •
 - •
 - •
 - Other (more far-reaching) applications
 - •
 - •
- corresponds to compressible flow in a frictionless, constant-area duct with heat transfer.



0



Rayleigh Flow



Beginning our analysis with the energy equation, we see

$$T_{o,1} - \frac{\dot{W}}{\dot{m}c_p} + \frac{\dot{Q}}{\dot{m}c_p} = T_{o,2}$$

The force-momentum balance, in this case is simplified

$$p_{o,1}A_1G_1 + F_x = p_{o,2}A_2G_2$$

Dividing the conservation of momentum expression by the continuity equation,

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Rayleigh Flow

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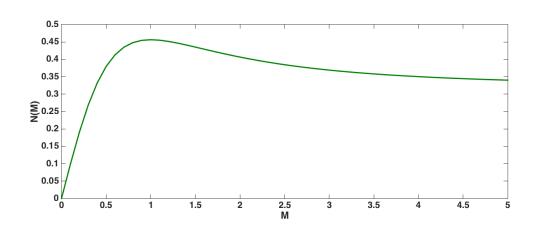
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• In a Rayleigh flow,

where

- When heat is transferred
 the flow, the Mach
 number is driven toward
- This process is known as



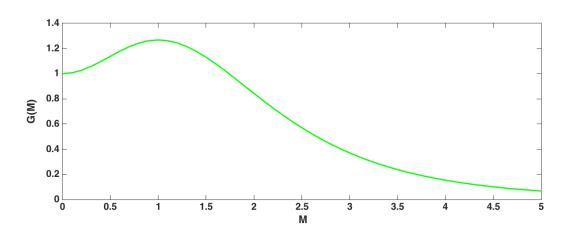
- That is, heat addition to the fluid acts to ______ the Mach number of a _____ flow and _____ the Mach number of a _____ flow towards M=1 (choking).
- While heat addition acts to raise the total temperature of the flow, it also lowers the
 ______. To show this, we again consider the momentum equation for a flow without friction:

$$p_{o,1}A_1G_1 = p_{o,2}A_2G_2$$

Rayleigh Flow



- Examining the behavior of the G(M) thrust flow function:
 - \circ The maximum at M=1 requires that, for a uniquely subsonic or supersonic flow,
 - o Therefore,



 We can again correlate this to entropy change through our understanding of the stagnation pressure as our 'entropy' variable.

$$s_2 - s_1 = c_p \ln \left(\frac{T_{o,2}}{T_{o,1}} \right) - R \ln \left(\frac{p_{o,2}}{p_{o,1}} \right)$$

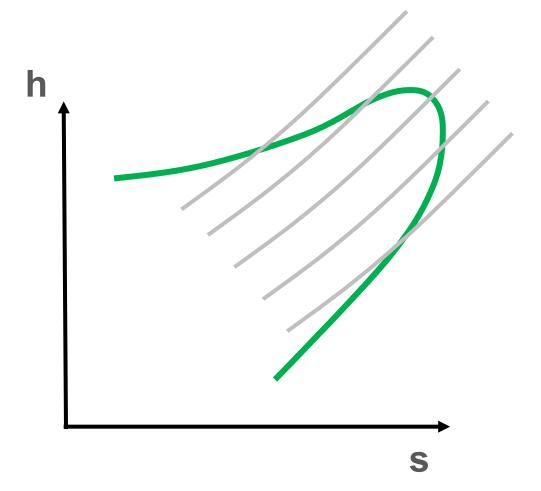
Property Changes in Rayleigh Flow



 The relationship between the flow (static) enthalpy and entropy can be shown by the Rayleigh line, where:

$$\frac{ds}{c_p} = \frac{M^2 - 1}{\gamma M^2 - 1} \frac{dh}{h}$$

Note that:



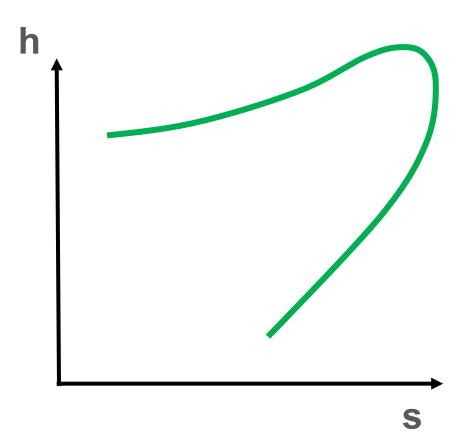
Property Changes in Rayleigh Flow



Heat addition corresponds to an _____ in the entropy while cooling causes a _____ in entropy

0

- For the case of heat subtraction $(dT_o < 0)$, changes in Mach number must drive the flow
 - deceleration of a subsonic flow and acceleration of a supersonic flow.
- Note that, under these conditions, we have the ability to decelerate a supersonic flow to a subsonic flow:



Property Changes in Rayleigh Flow



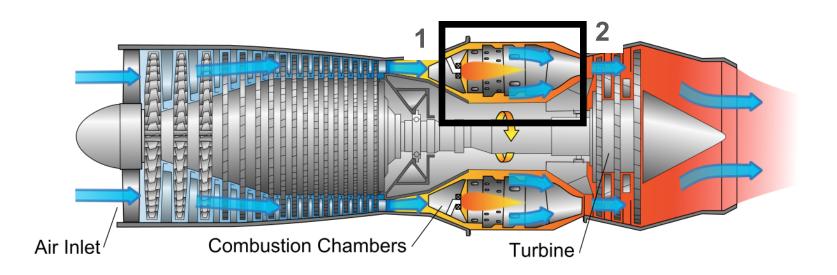
Property	Heating		Cooling	
	M < 1	M > 1	M < 1	M > 1
Entropy, s				
Stagnation temperature, T_o				
Static Temperature, T				
Mach number, Ma				
Stagnation Pressure, p_o				
Static Pressure, p				
Velocity, u				
Density, $ ho$				

- Heat addition $(\frac{T_{0,2}}{T_{0,1}} > 1)$ always decreases P_o .
 - The opposite is also true.
- The sensitivity of P_o to changes in T_o is the lowest when M is low.



Heat Addition in a Constant Area Duct

In a section of a duct between points 1 and 2, heat is added in a volumetric sense; i.e. not through the walls and into the boundary layer, but uniformly across the complete passage in such a manner that the temperature field remains constant across the duct at every axial location. Given the pressure, temperature, and velocity at location 1, compute the pressure, temperature, and velocity at location 2.





As always, we start with the three basic relations for 1-D Compressible Flow

$$\dot{m} = \left(\frac{p_{o,1}A_1}{\sqrt{T_{o,1}}}\right)\sqrt{\frac{\gamma}{R}}D_1 = \left(\frac{p_{o,2}A_2}{\sqrt{T_{o,2}}}\right)\sqrt{\frac{\gamma}{R}}D_2$$

$$p_{o,1}A_1G_1 + F_x = p_{o,2}A_2G_2$$

$$T_{o,1} - \frac{\dot{W}}{\dot{m}c_p} + \frac{\dot{Q}}{\dot{m}c_p} = T_{o,2}$$

Making the following assumptions/approximations:

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We reduce these expressions to the form:



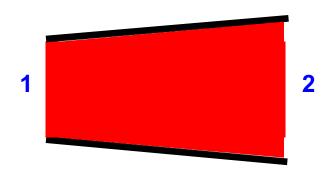




Constant Pressure Heat Addition

Compressible flow textbooks will often directly state that heat addition will cause Mach numbers of subsonic flows to increase and will cause a decrease in Mach number for supersonic flows. The fact is, that for typical conditions encountered in propulsive engines, the exact opposite can also be true. The difference is a matter of boundary conditions.

Another simple, one-dimensional problem concerns flow through a duct with heat addition occurring at constant pressure; i.e. the static pressure between locations 1 and 2 remains constant.



Again, we start from the conservation equations and simplify through the following assumptions/approximations:



Mass:

$$\dot{m} = \rho_1 u_1 A_1 = \rho_2 u_2 A_2$$

or, in terms of stagnation quantities

$$\dot{m} = \left(\frac{p_{o,1}A_1}{\sqrt{T_{o,1}}}\right)\sqrt{\frac{\gamma}{R}}D_1 = \left(\frac{p_{o,2}A_2}{\sqrt{T_{o,2}}}\right)\sqrt{\frac{\gamma}{R}}D_2$$

Momentum:

$$\rho_1 u_1^2 A_1 + p_1 A_1 + F_x = \rho_2 u_2^2 A_2 + p_2 A_2$$

or, in terms of stagnation quantities

$$p_{o,1}A_1G_1 + F_x = p_{o,2}A_2G_2$$

Energy:

$$T_{o,1} - \frac{\dot{W}}{\dot{m}c_p} + \frac{\dot{Q}}{\dot{m}c_p} = T_{o,2}$$















Solution for

$$\frac{T_{o,2}}{T_{o,1}} = 2.0$$

M_1	M_2	A_2/A_1	T_2/T_1	u_2/u_1	$ ho_2/ ho_1$	p_{2}/p_{1}	$p_{o,2}/p_{o,1}$
0.100	0.071	2.002	2.002	1.000	0.500	1.000	0.997
0.200	0.141	2.008	2.008	1.000	0.498	1.000	0.986
0.400	0.281	2.032	2.032	1.000	0.492	1.000	0.946
0.600	0.417	2.072	2.072	1.000	0.483	1.000	0.884
0.800	0.548	2.128	2.128	1.000	0.470	1.000	0.805
1.000	0.674	2.200	2.200	1.000	0.455	1.000	0.716
1.200	0.793	2.288	2.288	1.000	0.437	1.000	0.624
1.400	0.905	2.392	2.392	1.000	0.418	1.000	0.534
1.600	1.010	2.512	2.512	1.000	0.398	1.000	0.450
1.800	1.106	2.648	2.648	1.000	0.378	1.000	0.374
2.000	1.195	2.800	2.800	1.000	0.357	1.000	0.308
2.500	1.387	3.250	3.250	1.000	0.308	1.000	0.183
3.000	1.539	3.800	3.800	1.000	0.263	1.000	0.106
4.000	1.754	5.200	5.200	1.000	0.192	1.000	0.035