

Quantum Dynamics

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Goals:

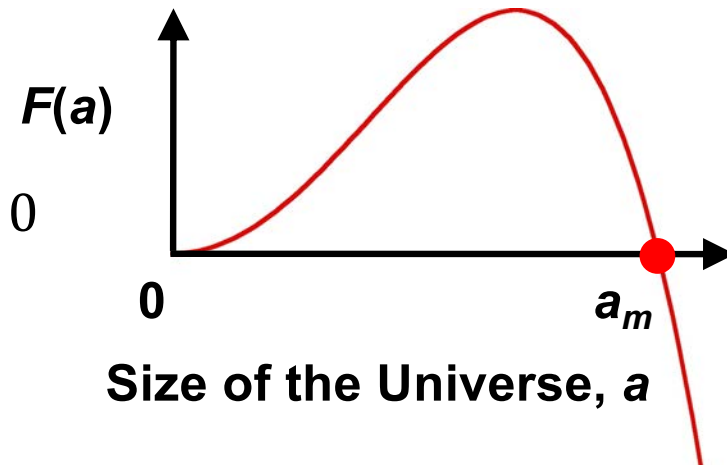
- 1. Partial differential equation**
- 2. Spectral method
(Fourier transform)**



Quantum Universe

- **Wheeler-deWitt equation**

$$\left[-\hbar^2 \frac{d^2}{da^2} + \underbrace{\left(\frac{3\pi c^3}{2G} \right)^2 \left(a^2 - \frac{a^4}{a_m^2} \right)}_{F(a)} \right] \psi(a) = 0$$



CREATION OF UNIVERSES FROM NOTHING

Alexander VILENKIN

Physics Department, Tufts University, Medford, MA 02155, USA

Received 11 June 1982

A cosmological model is proposed in which the universe is created by quantum tunneling from literally nothing into a de Sitter space. After the tunneling, the model evolves along the lines of the inflationary scenario. This model does not have a big-bang singularity and does not require any initial or boundary conditions.

Phys. Lett. **117B**, 25 ('82)

IS IT POSSIBLE TO CREATE A UNIVERSE IN THE LABORATORY BY QUANTUM TUNNELING?

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*Center for Theoretical Physics, Laboratory for Nuclear Science and Department of Physics,
Massachusetts Institute of Technology, Cambridge, MA 02139, USA*

Alan H. GUTH***

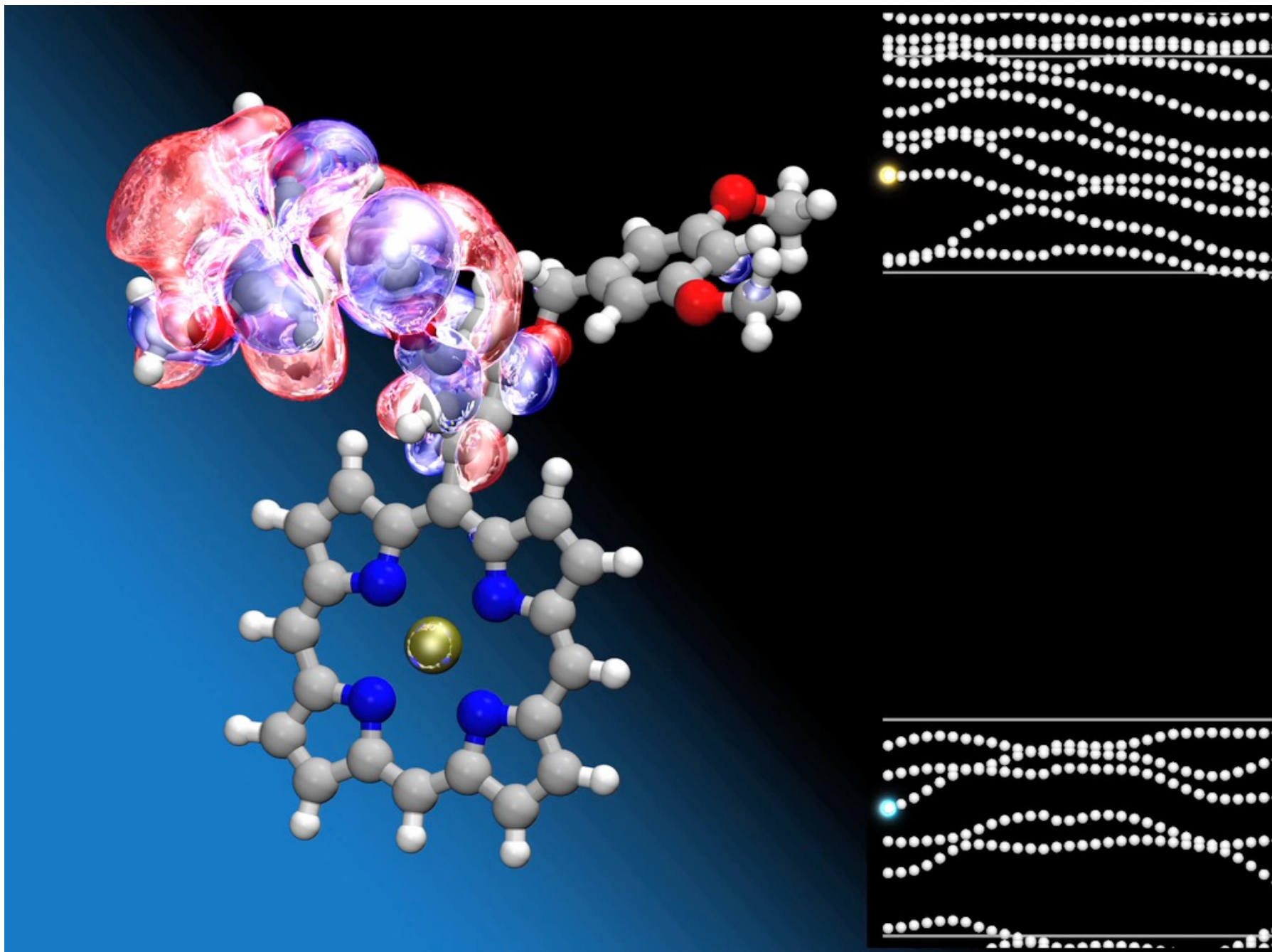
*Center for Theoretical Physics, Laboratory for Nuclear Science and Department of Physics,
Massachusetts Institute of Technology, Cambridge, MA 02139, USA
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Jemal GUVEN

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C.U. A. Postal 70-543, 04510 Mexico D.F. Mexico*

Nucl. Phys. **B339**, 417 ('90)

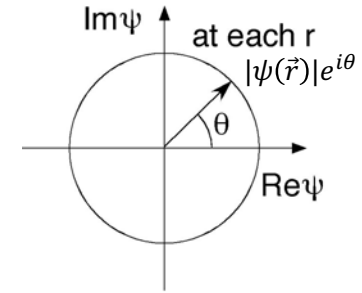
Photoexcited Electron Dynamics



Wave Equation

- Complex wave function**

$$\psi(\vec{r}, t) = \text{Re}\psi(\vec{r}, t) + i\text{Im}\psi(\vec{r}, t) \in \mathbb{C} \quad (i = \sqrt{-1})$$



- Probability**

$$P(\vec{r}, t) = \psi^*(\vec{r}, t)\psi(\vec{r}, t) = |\psi(\vec{r}, t)|^2 = |\text{Re}\psi(\vec{r}, t)|^2 + |\text{Im}\psi(\vec{r}, t)|^2$$

$$\psi^*\psi = (\psi_0 - i\psi_1)(\psi_0 + i\psi_1) = \psi_0^2 + \psi_1^2 \geq 0$$

- Normalization**

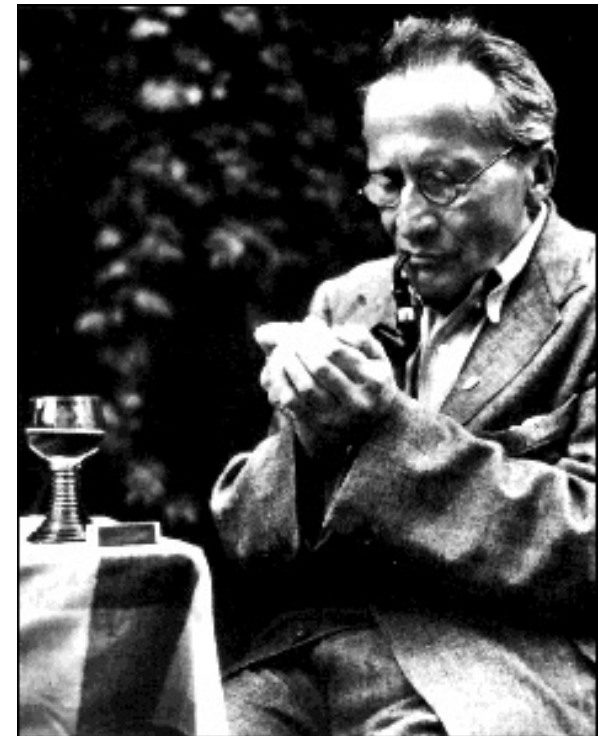
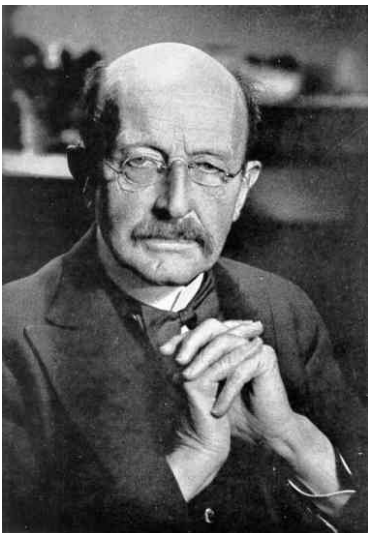
$$\int dx \int dy \int dz |\psi(\vec{r}, t)|^2 = 1$$

- Schrödinger (partial differential) equation**

$$i\hbar \frac{\partial}{\partial t} \psi(\vec{r}, t) = \left[-\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \psi(\vec{r}, t)$$

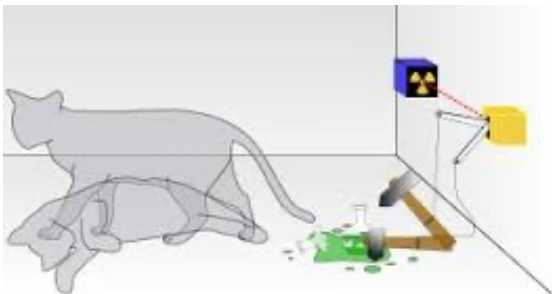
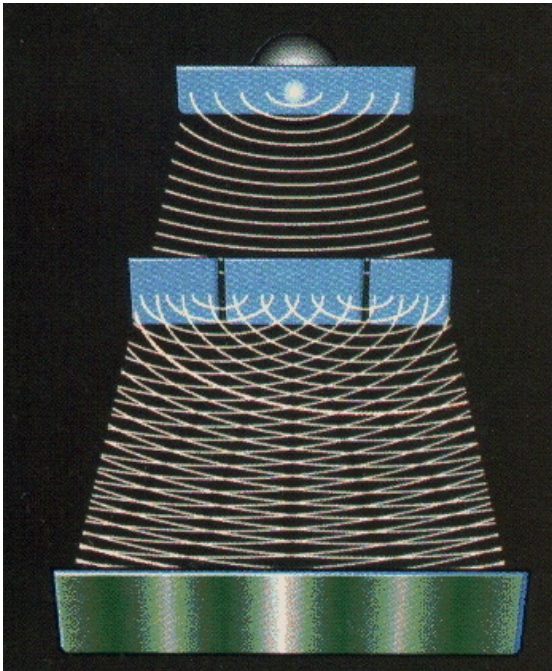
Laplacian: $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$

Planck constant: $\hbar = 1.05457 \times 10^{-27} \text{ g}\cdot\text{cm}^2/\text{s}$



Single-Electron Double-Slit Experiment

What wave?



<http://rdg.ext.hitachi.co.jp/rd/moviee/doubleslite-n.mpeg>

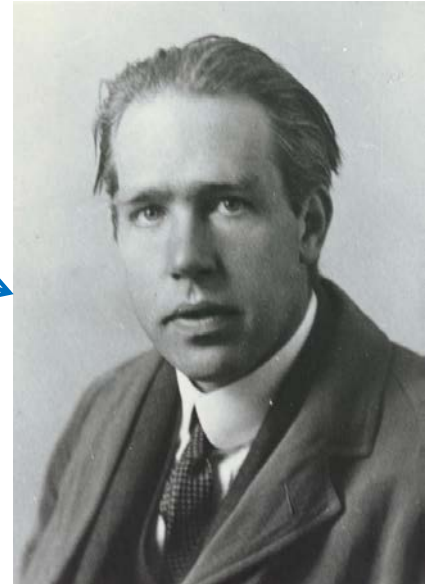
Akira Tonomura (Hitachi, Ltd.)

Atomic Unit

Length, energy & time in atomic unit

$$\begin{array}{l} \frac{e^2}{r} = \frac{\hbar^2}{mr^2} \\ E = \frac{e^2}{r} \\ E = \frac{\hbar}{t} \end{array} \left\{ \begin{array}{ll} \vec{r} = \frac{\hbar^2}{me^2} \vec{r}' & \frac{\hbar^2}{me^2} = 0.529177 \text{ \AA} \quad \text{Bohr} \\ V = \frac{me^4}{\hbar^2} V' & \frac{me^4}{\hbar^2} = 27.2116 \text{ eV} \quad \text{Hartree} \\ t = \frac{\hbar^3}{me^4} t' & \frac{\hbar^3}{me^4} = 0.0241889 \text{ fs} \end{array} \right.$$

CGS
Gaussian unit



Time-dependent Schrödinger equation in atomic unit

$$i\hbar \frac{\partial}{\partial t} \psi(\vec{r}, t) = \left[-\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \psi(\vec{r}, t)$$



$$i \frac{\partial}{\partial t'} \psi(\vec{r}', t') = \left[-\frac{\nabla'^2}{2} + V(\vec{r}') \right] \psi(\vec{r}', t')$$

Two-Dimensional Electron

- Schrödinger equation (in atomic unit)

$$i \frac{\partial}{\partial t} \psi(x, y, t) = H \psi(x, y, t)$$

- Hamiltonian operator

$$H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \frac{1}{2} \frac{\partial^2}{\partial y^2} + V(x, y)$$

$$= T_x + T_y + V$$



The Nobel Prize in Physics 1985

"for the discovery of the quantized Hall effect"



Klaus von Klitzing

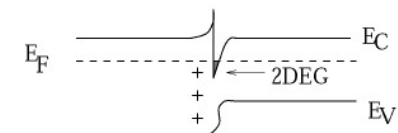
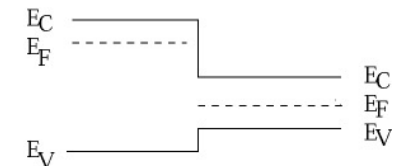
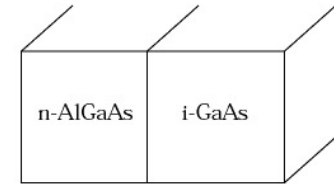
Federal Republic of Germany

Max-Planck-Institut für Festkörperforschung
Stuttgart, Federal Republic of Germany

b. 1943

- Periodic boundary condition

$$\begin{cases} \psi(x + L_x, y) = \psi(x, y) \\ \psi(x, y + L_y) = \psi(x, y) \end{cases}$$



The Nobel Prize in Physics 1998

"for their discovery of a new form of quantum fluid with fractionally charged excitations"



Robert B. Laughlin

1/3 of the prize
USA

Stanford University
Stanford, CA, USA
b. 1950



Horst L. Störmer

1/3 of the prize
Federal Republic of Germany

Columbia University
New York, NY, USA
b. 1949



Daniel C. Tsui

1/3 of the prize
USA

Princeton University
Princeton, NJ, USA
b. 1939
(in Henan, China)

Layered Materials Genome

- Atomically-thin layered materials will dominate materials science in this century

Geim & Grigorieva, *Nature* **499**, 419 ('13)

The Nobel Prize in Physics
2010

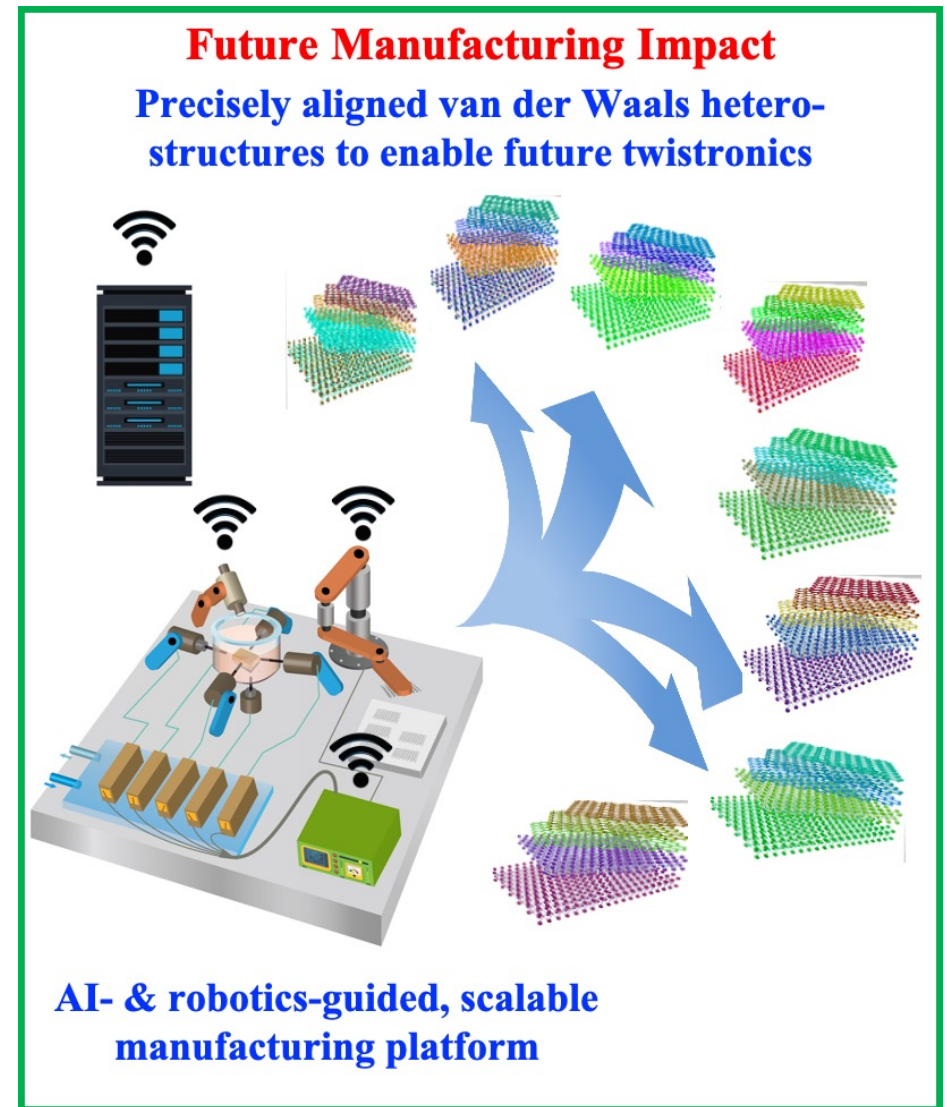


© The Nobel Foundation.
Photo: U. Montan
Andre Geim
Prize share: 1/2



© The Nobel Foundation.
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Konstantin
Novoselov
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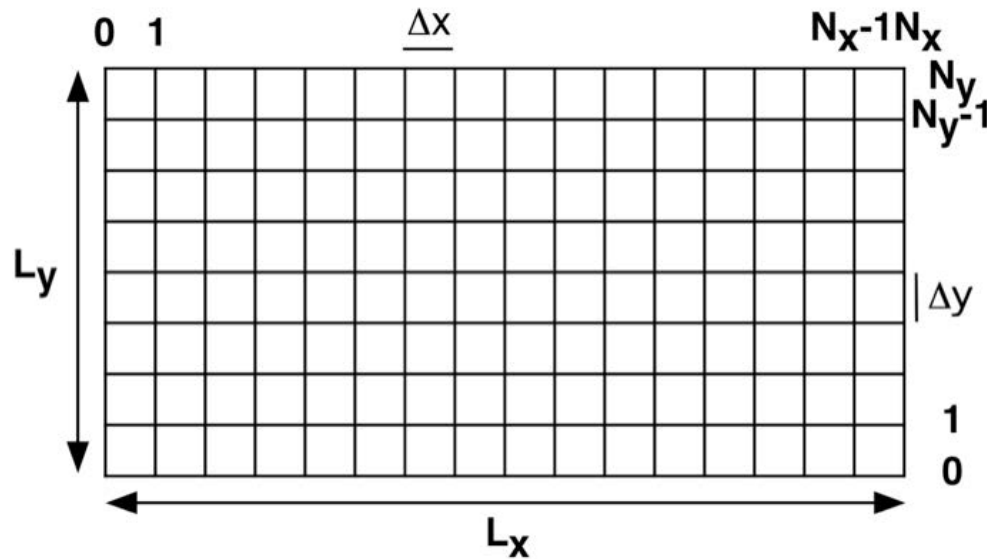
- Tuning material properties in desired ways by building heterostructures composed of unlimited combinations of atomically thin layers in a way similar to genomics



<https://aiqma.netlify.app>

Spatial Discretization

- **Regular 2D mesh:** $\psi_{jk} = \psi(j\Delta x, k\Delta y)$ ($\Delta x = L_x/N_x$ & $\Delta y = L_y/N_y$)



- **Finite differencing**

$$\begin{cases} (T_x \psi)_{j,k} = -\frac{1}{2} \frac{\psi_{j-1,k} - 2\psi_{j,k} + \psi_{j+1,k}}{(\Delta x)^2} \\ (T_y \psi)_{j,k} = -\frac{1}{2} \frac{\psi_{j,k-1} - 2\psi_{j,k} + \psi_{j,k+1}}{(\Delta y)^2} \\ (V\psi)_{j,k} = V_{j,k} \psi_{j,k} \end{cases}$$

$$\begin{aligned} \frac{\partial^2}{\partial x^2} \psi(x, y, t) &= \frac{\partial \psi(x+\Delta/2, y, t) / \partial x - \partial \psi(x-\Delta/2, y, t) / \partial x}{\Delta} \\ &= \frac{\frac{\psi(x+\Delta, y, t) - \psi(x, y, t)}{\Delta} - \frac{\psi(x, y, t) - \psi(x-\Delta, y, t)}{\Delta}}{\Delta} \\ &= \frac{\psi(x+\Delta, y, t) - 2\psi(x, y, t) + \psi(x-\Delta, y, t)}{\Delta^2} \end{aligned}$$

Temporal Propagation

- Formal solution to the Schrödinger equation:** $\frac{\partial}{\partial t}\psi(t) = -iH\psi(t)$

$$\psi(t + \Delta t) = \exp(-iH\Delta t)\psi(t) \quad \exp(-i\hat{H}t) = \sum_{n=0}^{\infty} \frac{(-it)^n}{n!} \hat{H}^n$$

- Split-operator method (Trotter-expansion): unitary!**

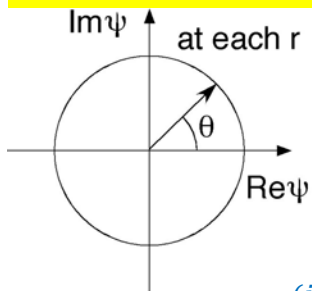
$$\begin{aligned} \psi(t + \Delta t) &= \exp(-i(T_x + T_y + V)\Delta t)\psi(t) \stackrel{T_x T_y = T_y T_x}{\approx} \exp(-i(T_x + T_y)\Delta t) \exp(-iV\Delta t) \psi(t) \\ &= \exp(-iV\Delta t/2) \exp(-iT_x\Delta t) \exp(-iT_y\Delta t) \exp(-iV\Delta t/2) \psi(t) + O([\Delta t]^3) \end{aligned}$$

Split in a way each operator is easily exponentiated

- Potential propagator (mesh point-by-point complex-number multiplications)**

$$(u_0 + iu_1)(\psi_0 + i\psi_1) = (u_0\psi_0 - u_1\psi_1) + i(u_1\psi_0 + u_0\psi_1)$$

$$(\exp(-iV\Delta t/2)\psi)_{jk} = \exp(-iV_{jk}\Delta t/2)\psi_{jk}$$



Rotation

$$\begin{aligned} &= [\cos(V_{jk}\Delta t/2) - i \sin(V_{jk}\Delta t/2)] [\text{Re}\psi_{jk} + i \text{Im}\psi_{jk}] \\ &= [\cos(V_{jk}\Delta t/2) \text{Re}\psi_{jk} + \sin(V_{jk}\Delta t/2) \text{Im}\psi_{jk}] \\ &\quad + i[-\sin(V_{jk}\Delta t/2) \text{Re}\psi_{jk} + \cos(V_{jk}\Delta t/2) \text{Im}\psi_{jk}] \end{aligned}$$

$$\begin{aligned} \exp(ia) &= \sum_{n=0}^{\infty} \frac{(ia)^n}{n!} = \left(1 - \frac{a^2}{2!} + \frac{a^4}{4!} - \dots\right) + i\left(a - \frac{a^3}{3!} + \frac{a^5}{5!} - \dots\right) \\ &= \cos(a) + i\sin(a) \end{aligned}$$

$$\begin{bmatrix} \cos(V_{jk}\Delta t/2) & \sin(V_{jk}\Delta t/2) \\ -\sin(V_{jk}\Delta t/2) & \cos(V_{jk}\Delta t/2) \end{bmatrix} \begin{bmatrix} \text{Re}\psi_{jk} \\ \text{Im}\psi_{jk} \end{bmatrix}$$

Kinetic Propagator

- Mesh-point coupling

$$T_x \psi_{j,k} = b\psi_{j-1,k} + 2a\psi_{j,k} + b\psi_{j+1,k}$$

- Tridiagonal matrix representation

$$T_x = \begin{bmatrix} 2a & b & & & & & b \\ b & 2a & b & & & & \\ & b & 2a & b & & & \\ & & \ddots & \ddots & \ddots & & \\ & & & b & 2a & b & \\ & & & & b & 2a & b \\ & & & & & b & 2a \end{bmatrix}$$



Note the periodic boundary condition

$$\begin{cases} a = 1/2(\Delta x)^2 \\ b = -1/2(\Delta x)^2 \end{cases}$$

Space Splitting Method (SSM)

- 2×2 block-diagonal decomposition & split-operator exponentiation

$$T_x = \begin{bmatrix} 2a & b & & & & & b \\ b & 2a & b & & & & \\ & b & 2a & b & & & \\ & & \ddots & \ddots & \ddots & & \\ & & & b & 2a & b & \\ & & & & b & 2a & b \\ b & & & & & b & 2a \end{bmatrix} \quad \begin{cases} \varepsilon_n^+ = \frac{1}{2} \left[\exp\left(-\frac{i\Delta t}{n}(a+b)\right) + \exp\left(-\frac{i\Delta t}{n}(a-b)\right) \right] \\ \varepsilon_n^- = \frac{1}{2} \left[\exp\left(-\frac{i\Delta t}{n}(a+b)\right) - \exp\left(-\frac{i\Delta t}{n}(a-b)\right) \right] \end{cases}$$

$$= \frac{1}{2} \begin{bmatrix} a & b & & & & & \\ b & a & & & & & \\ & & a & b & & & \\ & & b & a & & & \\ & & & \ddots & \ddots & & \\ & & & & a & b & \\ & & & & b & a & \end{bmatrix} + \begin{bmatrix} a & & & & & & b \\ & a & b & & & & \\ & b & a & & & & \\ & & & \ddots & & & \\ & & & & a & b & \\ & & & & b & a & \\ & & & & & & a \end{bmatrix} + \frac{1}{2} \begin{bmatrix} a & b & & & & & \\ b & a & & & & & \\ & & a & b & & & \\ & & b & a & & & \\ & & & \ddots & \ddots & & \\ & & & & a & b & \\ & & & & b & a & \end{bmatrix}$$

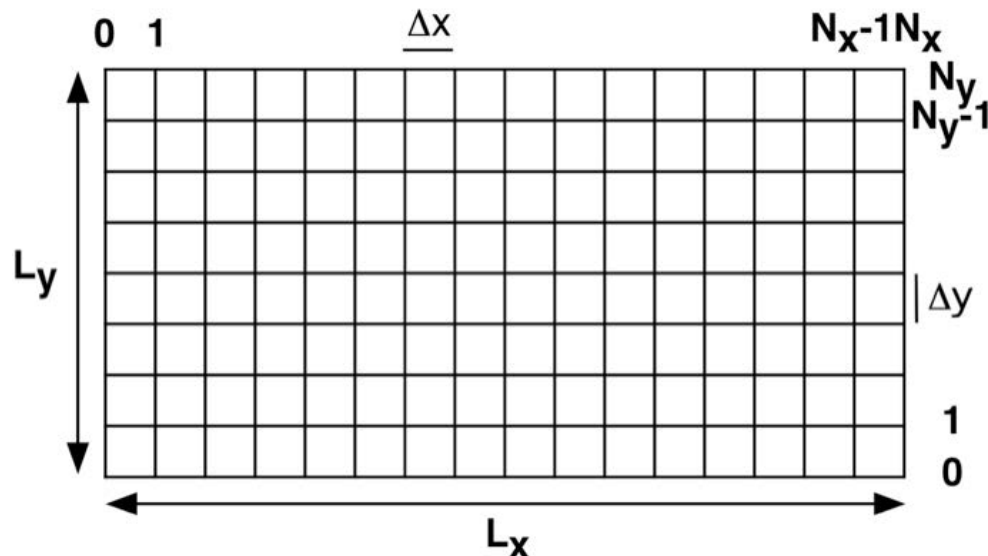
$$\exp(-i\Delta t T_x) = U_x^{(\text{half})} U_x^{(\text{full})} U_x^{(\text{half})} + O([\Delta t]^3) =$$

$$\begin{bmatrix} \varepsilon_2^+ & \varepsilon_2^- & & & & & \\ \varepsilon_2^- & \varepsilon_2^+ & & & & & \\ & & \varepsilon_2^+ & \varepsilon_2^- & & & \\ & & \varepsilon_2^- & \varepsilon_2^+ & & & \\ & & & \ddots & \ddots & & \\ & & & & \varepsilon_1^+ & \varepsilon_1^- & \\ & & & & \varepsilon_1^- & \varepsilon_1^+ & \\ & & & & & \ddots & \ddots \\ & & & & & & \varepsilon_1^+ & \varepsilon_1^- \\ & & & & & & \varepsilon_1^- & \varepsilon_1^+ \end{bmatrix} \begin{bmatrix} \varepsilon_1^+ \\ \varepsilon_1^- \\ \varepsilon_2^+ \\ \varepsilon_2^- \\ \vdots \\ \varepsilon_1^+ \\ \varepsilon_1^- \end{bmatrix} \begin{bmatrix} \varepsilon_1^- & \varepsilon_1^+ \\ \varepsilon_2^- & \varepsilon_2^+ \\ & \varepsilon_2^- & \varepsilon_2^+ \\ & & \varepsilon_2^+ & \varepsilon_2^- \\ & & & \ddots & \ddots \\ & & & & \varepsilon_2^+ & \varepsilon_2^- \\ & & & & \varepsilon_2^- & \varepsilon_2^+ \end{bmatrix}$$

Data Structures in Program qd.c

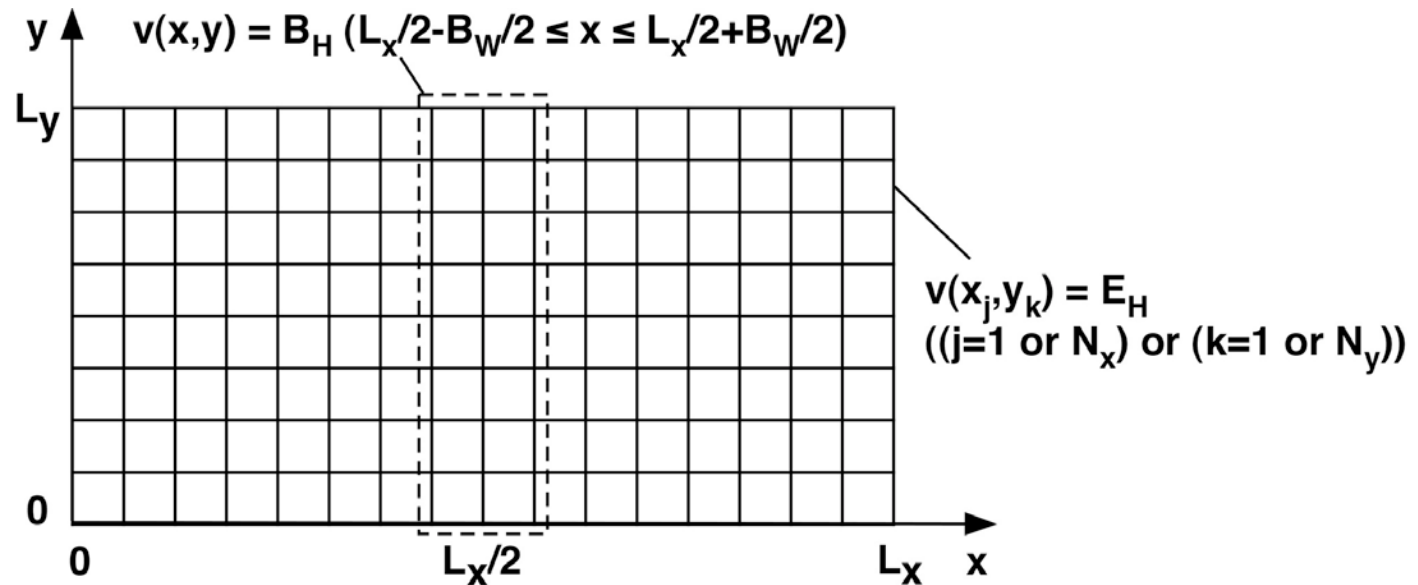
- **Wave function:** `psi[NX+2][NY+2][2]`
- **Periodic boundary condition by auxiliary elements**

```
for (sy=1; sy<=NY; sy++)  
  for (s=0; s<=1; s++) {  
    psi[0][sy][s] = psi[NX][sy][s];  
    psi[NX+1][sy][s] = psi[1][sy][s];}  
for (sx=1; sx<=NX; sx++)  
  for (s=0; s<=1; s++) {  
    psi[sx][0][s] = psi[sx][NY][s];  
    psi[sx][NY+1][s] = psi[sx][1][s];}
```



Potential Propagator in qd.c

- **Potential barrier:** $v[NX+2][NY+2]$



- **Potential propagator:** $\exp(-iV\Delta t/2)$, $u[NX+2][NY+2][2]$

- **Potential propagation:** $\psi \leftarrow \exp(-iV\Delta t/2) \psi$

```
for (sx=1; sx<=NX; sx++)  
    for (sy=1; sy<=NY; sy++) {  
        wr=u[sx][sy][0]*psi[sx][sy][0]-u[sx][sy][1]*psi[sx][sy][1];  
        wi=u[sx][sy][0]*psi[sx][sy][1]+u[sx][sy][1]*psi[sx][sy][0];  
        psi[sx][sy][0]=wr;  
        psi[sx][sy][1]=wi;};
```

Kinetic Propagator in qd.c

$$\begin{aligned}\left(U_x^{(\text{half})}\psi\right)_{i,j} &= \varepsilon_2^- \delta_{\text{mod}(i,2),0} \psi_{i-1,j} + \varepsilon_2^+ \psi_{i,j} + \varepsilon_2^- \delta_{\text{mod}(i,2),1} \psi_{i+1,j} \\ \left(U_x^{(\text{full})}\psi\right)_{i,j} &= \varepsilon_1^- \delta_{\text{mod}(i,2),1} \psi_{i-1,j} + \varepsilon_1^+ \psi_{i,j} + \varepsilon_1^- \delta_{\text{mod}(i,2),0} \psi_{i+1,j}\end{aligned}$$

```
/* WRK|PSI holds the new|old wave function */
for (sx=1; sx<=NX; sx++)
  for (sy=1; sy<=NY; sy++) {
    wr=al[d][t][0]*psi[sx][sy][0]-al[d][t][1]*psi[sx][sy][1];
    wi=al[d][t][0]*psi[sx][sy][1]+al[d][t][1]*psi[sx][sy][0];
    if (d==0) {
      wr+=(blx[t][sx][0]*psi[sx-1][sy][0]-blx[t][sx][1]*psi[sx-1][sy][1]);
      wi+=(blx[t][sx][0]*psi[sx-1][sy][1]+blx[t][sx][1]*psi[sx-1][sy][0]);
      wr+=(bux[t][sx][0]*psi[sx+1][sy][0]-bux[t][sx][1]*psi[sx+1][sy][1]);
      wi+=(bux[t][sx][0]*psi[sx+1][sy][1]+bux[t][sx][1]*psi[sx+1][sy][0]);}
    else if (d==1) {
      wr+=(bly[t][sy][0]*psi[sx][sy-1][0]-bly[t][sy][1]*psi[sx][sy-1][1]);
      wi+=(bly[t][sy][0]*psi[sx][sy-1][1]+bly[t][sy][1]*psi[sx][sy-1][0]);
      wr+=(buy[t][sy][0]*psi[sx][sy+1][0]-buy[t][sy][1]*psi[sx][sy+1][1]);
      wi+=(buy[t][sy][0]*psi[sx][sy+1][1]+buy[t][sy][1]*psi[sx][sy+1][0]);}
    wrk[sx][sy][0]=wr;
    wrk[sx][sy][1]=wi;}
/* Copy the new wave function back to PSI */
for (sx=1; sx<=NX; sx++)
  for (sy=1; sy<=NY; sy++)
    for (s=0; s<=1; s++) psi[sx][sy][s]=wrk[sx][sy][s];
```

Initial Wave Function

- Gaussian wave packet

$$\psi(x, t = 0) = C \exp\left(-\frac{(x - x_0)^2}{4\sigma^2}\right) \exp(ik_0 x) \sin\left(\frac{\pi y}{L_y}\right)$$

Symbol	Variable in qd.c
x_0 (packet center)	X0
σ (packet spread)	S0
$k_0^2/2$ (energy)	E0

Quantum Dynamics—II

One Dimensional System

Aiichiro Nakano

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University of Southern California*

Email: anakano@usc.edu



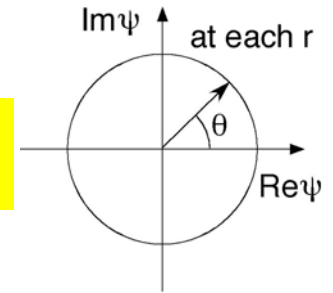
Goal: Understand qd1.c



Wave Equation

- **Complex wave function**

$$\psi(x, t) = \text{Re}\psi(x, t) + i\text{Im}\psi(x, t) \in \mathbb{C} \quad (i = \sqrt{-1})$$



- **Normalization**

$$\int dx |\psi(x, t)|^2 = 1$$

$$\psi^* \psi = (\psi_0 - i\psi_1)(\psi_0 + i\psi_1) = \psi_0^2 + \psi_1^2 \geq 0$$

- **Schrödinger equation (in atomic unit)**

$$i \frac{\partial}{\partial t} \psi(x, t) = H\psi(x, t)$$

- **Hamiltonian operator**

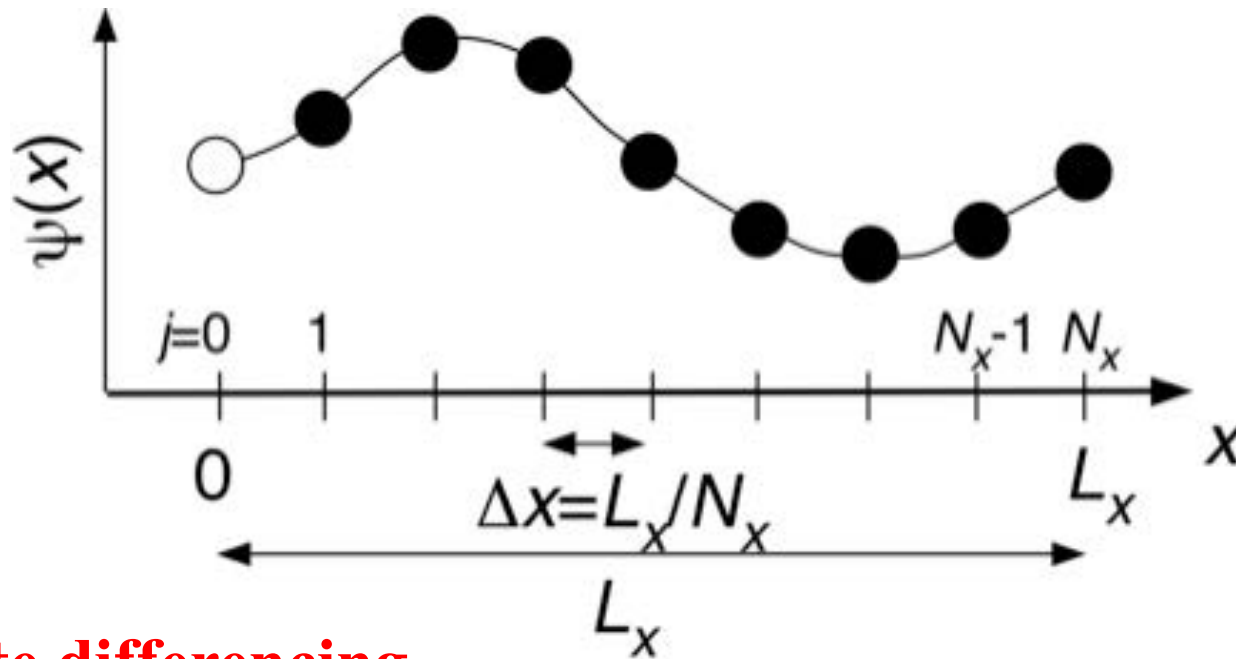
$$H = -\frac{1}{2} \frac{\partial^2}{\partial x^2} + V(x) = T_x + V$$

- **Periodic boundary condition**

$$\psi(x + L_x) = \psi(x)$$

Spatial Discretization

- **Regular 1D mesh:** $\psi_j = \psi(j\Delta x)$ ($\Delta x = L_x/N_x$)



- **Finite differencing**

$$\begin{cases} (T_x \psi)_j = -\frac{1}{2} \frac{\psi_{j-1} - 2\psi_j + \psi_{j+1}}{(\Delta x)^2} \\ (V\psi)_j = V_j \psi_j \end{cases}$$

$$\begin{aligned} \frac{\partial^2}{\partial x^2} \psi(x, t) &= \frac{\partial \psi(x+\Delta/2, t)/\partial x - \partial \psi(x-\Delta/2, t)/\partial x}{\Delta} \\ &= \frac{\frac{\psi(x+\Delta, t) - \psi(x, t)}{\Delta} - \frac{\psi(x, t) - \psi(x-\Delta, t)}{\Delta}}{\Delta} \\ &= \frac{\psi(x+\Delta, t) - 2\psi(x, t) + \psi(x-\Delta, t)}{\Delta^2} \end{aligned}$$

Temporal Propagation

- Formal solution to the Schrödinger equation:** $\frac{\partial}{\partial t}\psi(t) = -iH\psi(t)$

$$\psi(t + \Delta t) = \exp(-iH\Delta t)\psi(t) \quad \exp(-i\hat{H}t) = \sum_{n=0}^{\infty} \frac{(-it)^n}{n!} \hat{H}^n$$

- Split-operator method: unitary!**

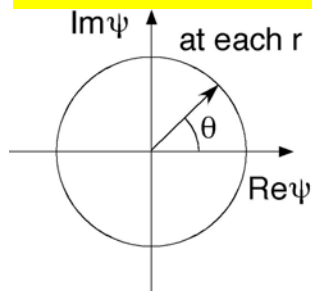
$$\begin{aligned} \psi(t + \Delta t) &= \exp(-i(T_x + V)\Delta t)\psi(t) \\ &= \exp(-iV\Delta t/2)\exp(-iT_x\Delta t)\exp(-iV\Delta t/2)\psi(t) + O([\Delta t]^3) \end{aligned}$$

Split in a way each operator is easily exponentiated

- Potential propagator (mesh point-by-point complex-number multiplications)**

$$(u_0 + iu_1)(\psi_0 + i\psi_1) = (u_0\psi_0 - u_1\psi_1) + i(u_1\psi_0 + u_0\psi_1)$$

$$(\exp(-iV\Delta t/2)\psi)_j = \exp(-iV_j\Delta t/2)\psi_j$$



Rotation

$$\begin{aligned} &= [\cos(V_j\Delta t/2) - i \sin(V_j\Delta t/2)][\text{Re}\psi_j + i \text{Im}\psi_j] \\ &= [\cos(V_j\Delta t/2)\text{Re}\psi_j + \sin(V_j\Delta t/2)\text{Im}\psi_j] \\ &\quad + i[-\sin(V_j\Delta t/2)\text{Re}\psi_j + \cos(V_j\Delta t/2)\text{Im}\psi_j] \end{aligned}$$

$$\begin{aligned} \exp(ia) &= \sum_{n=0}^{\infty} \frac{(ia)^n}{n!} = \left(1 - \frac{a^2}{2!} + \frac{a^4}{4!} - \dots\right) + i\left(a - \frac{a^3}{3!} + \frac{a^5}{5!} - \dots\right) \\ &= \cos(a) + i\sin(a) \end{aligned}$$

$$\begin{bmatrix} \cos(V_{jk}\Delta t/2) & \sin(V_{jk}\Delta t/2) \\ -\sin(V_{jk}\Delta t/2) & \cos(V_{jk}\Delta t/2) \end{bmatrix} \begin{bmatrix} \text{Re}\psi_{jk} \\ \text{Im}\psi_{jk} \end{bmatrix}$$

Kinetic Propagator: It's a Matrix!

- **Mesh-point coupling**

$$T_x \psi_j = b\psi_{j-1} + 2a\psi_j + b\psi_{j+1}$$

- **Tridiagonal matrix representation**

$$T_x = \begin{bmatrix} 2a & b & & & & & b \\ b & 2a & b & & & & \\ & b & 2a & b & & & \\ & & \ddots & \ddots & \ddots & & \\ & & & b & 2a & b & \\ & & & & b & 2a & b \\ & & & & & b & 2a \end{bmatrix}$$

Note the periodic boundary condition

$$\begin{cases} a = 1/2(\Delta x)^2 \\ b = -1/2(\Delta x)^2 \end{cases}$$



Space Splitting Method (SSM)

- 2×2 block-diagonal decomposition & split-operator exponentiation

$$T_x = \begin{bmatrix} 2a & b & & & & & b \\ b & 2a & b & & & & \\ & b & 2a & b & & & \\ & & \ddots & \ddots & \ddots & & \\ & & & b & 2a & b & \\ & & & & b & 2a & b \\ b & & & & b & 2a & b \end{bmatrix}$$

Block-by-block exponentiation

$$\exp \begin{bmatrix} \blacksquare & & \\ & \blacksquare & \\ & & \blacksquare \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} \begin{bmatrix} \blacksquare & & \\ & \blacksquare & \\ & & \blacksquare \end{bmatrix}^n$$

$$= \sum_{n=0}^{\infty} \frac{1}{n!} \begin{bmatrix} \blacksquare^n & & \\ & \blacksquare^n & \\ & & \blacksquare^n \end{bmatrix} = \begin{bmatrix} e^{\blacksquare} & & \\ & e^{\blacksquare} & \\ & & e^{\blacksquare} \end{bmatrix}$$

Split-operator (Trotter expansion) again

$$= \frac{1}{2} \begin{bmatrix} a & b & & & \\ b & a & & & \\ & & a & b & \\ & & b & a & \\ & & & \ddots & \\ & & & & a & b \\ & & & & b & a \end{bmatrix} + \begin{bmatrix} a & & & & b \\ & a & b & & \\ & b & a & & \\ & & & \ddots & \\ & & & & a & b \\ & & & & b & a \end{bmatrix} + \frac{1}{2} \begin{bmatrix} a & b & & & \\ b & a & & & \\ & & a & b & \\ & & b & a & \\ & & & \ddots & \\ & & & & a & b \\ & & & & b & a \end{bmatrix}$$

$$\exp(-i\Delta t T_x) = U_x^{(\text{half})} U_x^{(\text{full})} U_x^{(\text{half})} + \mathcal{O}([\Delta t]^3)$$

$$= \exp \left(-\frac{i\Delta t}{2} \begin{bmatrix} a & b & & & \\ b & a & & & \\ & & a & b & \\ & & b & a & \\ & & & \ddots & \\ & & & & a & b \\ & & & & b & a \end{bmatrix} \right) \exp \left(-i\Delta t \begin{bmatrix} a & & & & b \\ & a & b & & \\ & b & a & & \\ & & & \ddots & \\ & & & & a & b \\ & & & & b & a \end{bmatrix} \right) \exp \left(-\frac{i\Delta t}{2} \begin{bmatrix} a & b & & & \\ b & a & & & \\ & & a & b & \\ & & b & a & \\ & & & \ddots & \\ & & & & a & b \\ & & & & b & a \end{bmatrix} \right)$$

How? Block diagonal → block-by-block exponentiation

Space Splitting Method (SSM)

$$\exp(-i\Delta t T_x) = U_x^{(\text{half})} U_x^{(\text{full})} U_x^{(\text{half})} + O([\Delta t]^3)$$

$$= \exp \left(-\frac{i\Delta t}{2} \begin{bmatrix} a & b & & \\ b & a & & \\ & & a & b \\ & & b & a \\ & & & \ddots \\ & & & & a & b \\ & & & & b & a \end{bmatrix} \right) \exp \left(-i\Delta t \begin{bmatrix} a & & & & & b \\ & a & b & & & \\ & b & a & & & \\ & & & \ddots & & \\ & & & & a & b \\ & & & & b & a \end{bmatrix} \right) \exp \left(-\frac{i\Delta t}{2} \begin{bmatrix} a & b & & \\ b & a & & \\ & & a & b \\ & & b & a \\ & & & \ddots \\ & & & & a & b \\ & & & & b & a \end{bmatrix} \right)$$

$$= \begin{bmatrix} \varepsilon_2^+ & \varepsilon_2^- & & & & \\ \varepsilon_2^- & \varepsilon_2^+ & & & & \\ & & \varepsilon_2^+ & \varepsilon_2^- & & \\ & & \varepsilon_2^- & \varepsilon_2^+ & & \\ & & & \ddots & & \\ & & & & \varepsilon_2^+ & \varepsilon_2^- \\ & & & & \varepsilon_2^- & \varepsilon_2^+ \end{bmatrix} \begin{bmatrix} \varepsilon_1^+ & & & & & \\ & \varepsilon_1^+ & \varepsilon_1^- & & & \\ & \varepsilon_1^- & \varepsilon_1^+ & & & \\ & & & \ddots & & \\ & & & & \varepsilon_1^+ & \varepsilon_1^- \\ & & & & \varepsilon_1^- & \varepsilon_1^+ \end{bmatrix} \begin{bmatrix} \varepsilon_1^- & & & & & \\ \varepsilon_2^+ & \varepsilon_2^- & & & & \\ \varepsilon_2^- & \varepsilon_2^+ & & & & \\ & & \varepsilon_2^+ & \varepsilon_2^- & & \\ & & \varepsilon_2^- & \varepsilon_2^+ & & \\ & & & \ddots & & \\ & & & & \varepsilon_2^+ & \varepsilon_2^- \\ & & & & \varepsilon_2^- & \varepsilon_2^+ \end{bmatrix}$$

$$\begin{cases} \varepsilon_n^+ = \frac{1}{2} \left[\exp \left(-\frac{i\Delta t}{n} (a + b) \right) + \exp \left(-\frac{i\Delta t}{n} (a - b) \right) \right] \\ \varepsilon_n^- = \frac{1}{2} \left[\exp \left(-\frac{i\Delta t}{n} (a + b) \right) - \exp \left(-\frac{i\Delta t}{n} (a - b) \right) \right] \end{cases} \quad \text{Just need } 2 \times 2 \text{ exponentiation}$$

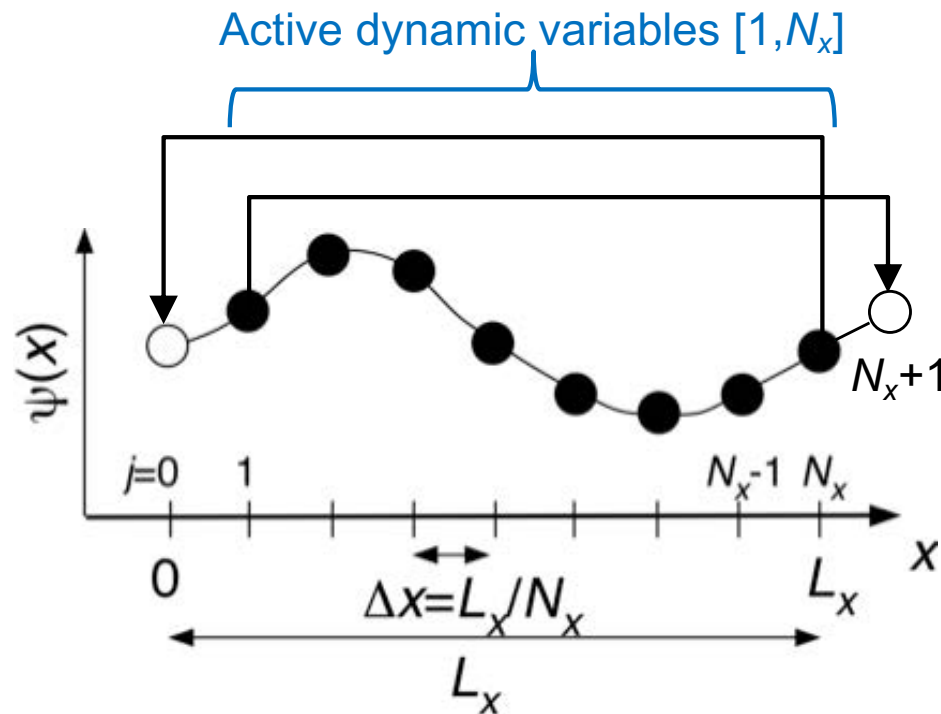
$$\exp \left(-\frac{i\Delta}{(2)} \begin{bmatrix} a & b \\ b & a \end{bmatrix} \right)$$

Use eigen-decomposition & telescoping

Data Structures in Program qd1.c

- **Wave function:** `psi[NX+2][2]` $\psi[j][0|1] = (\text{Re}|\text{Im})\psi_{j\Delta x}$
- **Periodic boundary condition by auxiliary elements**

```
for (s=0; s<=1; s++) {  
    psi[0][s] = psi[NX][s];  
    psi[NX+1][s] = psi[1][s];  
}
```



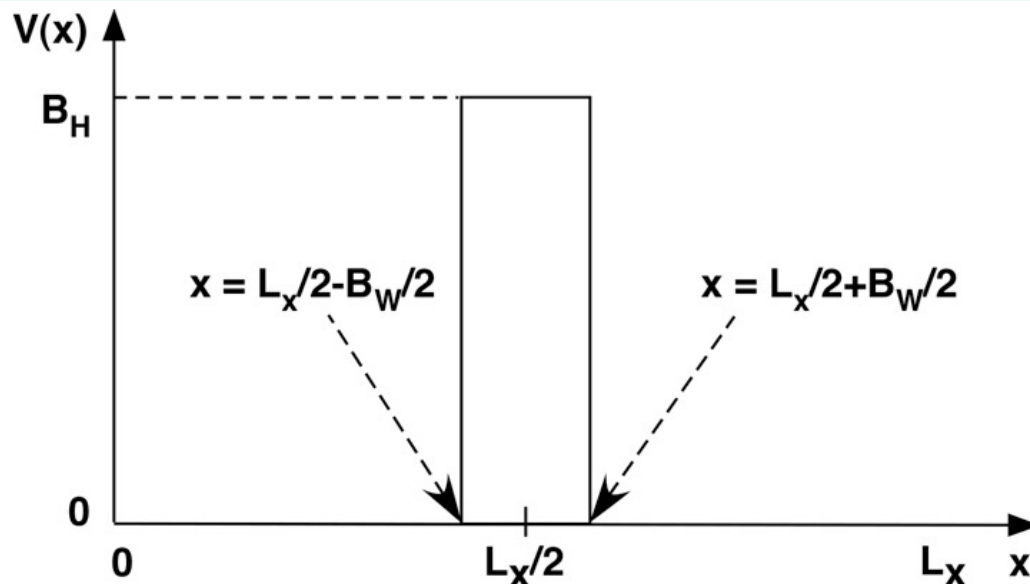
Potential Propagator in qd1.c

- **Potential barrier:** $v[NX+2]$
- **Potential propagator:** $\exp(-iV\Delta t/2)$, $u[NX+2][2]$
- **Potential propagation:** $\psi \leftarrow \exp(-iV\Delta t/2) \psi$

```
for (sx=1; sx<=NX; sx++)  
    wr=u[sx][0]*psi[sx][0]-u[sx][1]*psi[sx][1];  
    wi=u[sx][0]*psi[sx][1]+u[sx][1]*psi[sx][0];  
    psi[sx][0]=wr;  
    psi[sx][1]=wi;  
}
```

$$\begin{cases} u[j][0] = \cos\left(-\frac{\Delta}{2}V_j\right) \\ u[j][1] = \sin\left(-\frac{\Delta}{2}V_j\right) \end{cases}$$

$$\exp\left(-\frac{iV_j\Delta}{2}\right)\psi_j \equiv u\psi = (u_0 + iu_1)(\psi_0 + i\psi_1) = \overbrace{(u_0\psi_0 - u_1\psi_1)}^{\text{new } \psi_0} + i \overbrace{(u_0\psi_1 + u_1\psi_0)}^{\text{new } \psi_1}$$



Kinetic Propagator in qd1.c

$$\begin{aligned} \left(U_x^{(\text{half})} \psi \right)_i &= \varepsilon_2^- \delta_{\text{mod}(i,2),0} \beta_{\text{half}}^{(\text{low})} \psi_{i-1} + \varepsilon_2^+ \alpha_{\text{half}} \psi_i + \varepsilon_2^- \delta_{\text{mod}(i,2),1} \beta_{\text{half}}^{(\text{up})} \psi_{i+1} \\ \left(U_x^{(\text{full})} \psi \right)_i &= \varepsilon_1^- \delta_{\text{mod}(i,2),1} \beta_{\text{full}}^{(\text{low})} \psi_{i-1} + \varepsilon_1^+ \alpha_{\text{full}} \psi_i + \varepsilon_1^- \delta_{\text{mod}(i,2),0} \beta_{\text{full}}^{(\text{up})} \psi_{i+1} \end{aligned}$$

```
for (sx=1; sx<=NX; sx++) { // wrk[][]|psi[][] holds new|old wave function
    wr=al[t][0]*psi[sx][0]-al[t][1]*psi[sx][1]; // al[0|1][]:  $\alpha_{\text{half|full}}$ 
    wi=al[t][0]*psi[sx][1]+al[t][1]*psi[sx][0];
    wr+=(bl[t][sx][0]*psi[sx-1][0]-bl[t][sx][1]*psi[sx-1][1]); // bl[0|1][]:  $\beta_{\text{half|full}}^{(\text{low})}$ 
    wi+=(bl[t][sx][0]*psi[sx-1][1]+bl[t][sx][1]*psi[sx-1][0]);
    wr+=(bu[t][sx][0]*psi[sx+1][0]-bu[t][sx][1]*psi[sx+1][1]); // bu[0|1][]:  $\beta_{\text{half|full}}^{(\text{high})}$ 
    wi+=(bu[t][sx][0]*psi[sx+1][1]+bu[t][sx][1]*psi[sx+1][0]);}
    wrk[sx][0]=wr;
    wrk[sx][1]=wi;}
for (sx=1; sx<=NX; sx++) // Copy new wave function back to psi
    for (s=0; s<=1; s++)
        psi[sx][s]=wrk[sx][s];
```

$$\psi_j \leftarrow \beta_l \psi_{j-1} + \alpha \psi_j + \beta_u \psi_{j+1}$$

$$\begin{cases} \varepsilon_n^+ = \frac{1}{2} \left[\exp \left(-\frac{i\Delta t}{n} (a+b) \right) + \exp \left(-\frac{i\Delta t}{n} (a-b) \right) \right] \\ \varepsilon_n^- = \frac{1}{2} \left[\exp \left(-\frac{i\Delta t}{n} (a+b) \right) - \exp \left(-\frac{i\Delta t}{n} (a-b) \right) \right] \end{cases}$$

$$\exp(-i\Delta t T_x) = U_x^{(\text{half})} U_x^{(\text{full})} U_x^{(\text{half})} + O([\Delta t]^3)$$

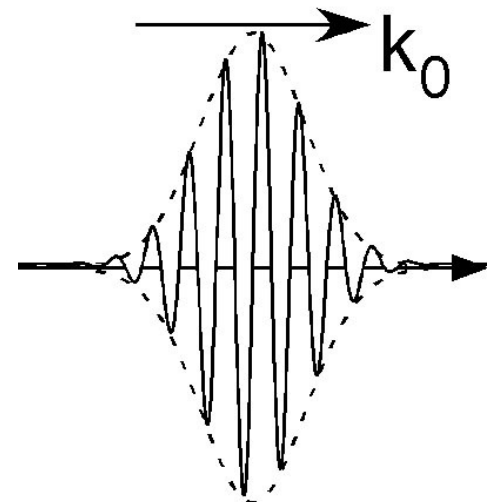
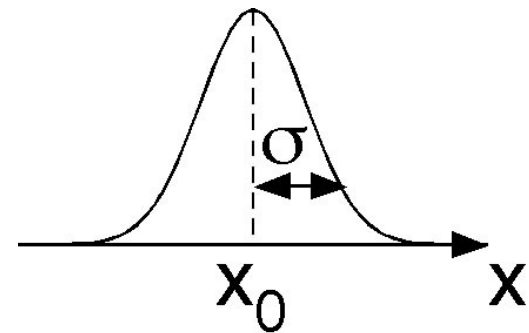
$$= \begin{bmatrix} \varepsilon_2^+ & \varepsilon_2^- & & & & \\ \varepsilon_2^- & \varepsilon_2^+ & & & & \\ & & \varepsilon_2^+ & \varepsilon_2^- & & \\ & & \varepsilon_2^- & \varepsilon_2^+ & & \\ & & & & \ddots & \\ & & & & & \varepsilon_2^+ & \varepsilon_2^- \\ & & & & & \varepsilon_2^- & \varepsilon_2^+ \end{bmatrix} \begin{bmatrix} \varepsilon_1^+ & & & & & \\ & \varepsilon_1^+ & \varepsilon_1^- & & & \\ & \varepsilon_1^- & \varepsilon_1^+ & & & \\ & & & \ddots & & \\ & & & & \varepsilon_1^+ & \varepsilon_1^- \\ & & & & \varepsilon_1^- & \varepsilon_1^+ \end{bmatrix} \begin{bmatrix} \varepsilon_1^- & & & & & \\ & \varepsilon_2^+ & \varepsilon_2^- & & & \\ & \varepsilon_2^- & \varepsilon_2^+ & & & \\ & & & \varepsilon_2^+ & \varepsilon_2^- & \\ & & & \varepsilon_2^- & \varepsilon_2^+ & \\ & & & & \ddots & \\ & & & & & \varepsilon_2^+ & \varepsilon_2^- \\ & & & & & \varepsilon_2^- & \varepsilon_2^+ \end{bmatrix}$$

Initial Wave Function

- Gaussian wave packet

$$\psi(x, t = 0) = C \exp\left(-\frac{(x - x_0)^2}{4\sigma^2}\right) \exp(ik_0 x)$$

Symbol	Variable in qd1.c
x_0 (packet center)	X0
σ (packet spread)	S0
$k_0^2/2$ (energy)	E0



Quantum Dynamics—III

Spectral Method

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Goal: Understand Fourier transform in the context of the orthonormal plane-wave basis set in a vector space

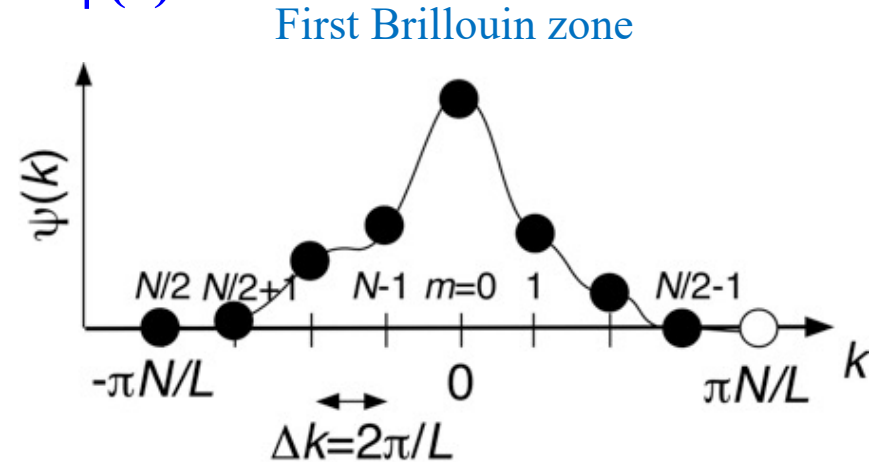
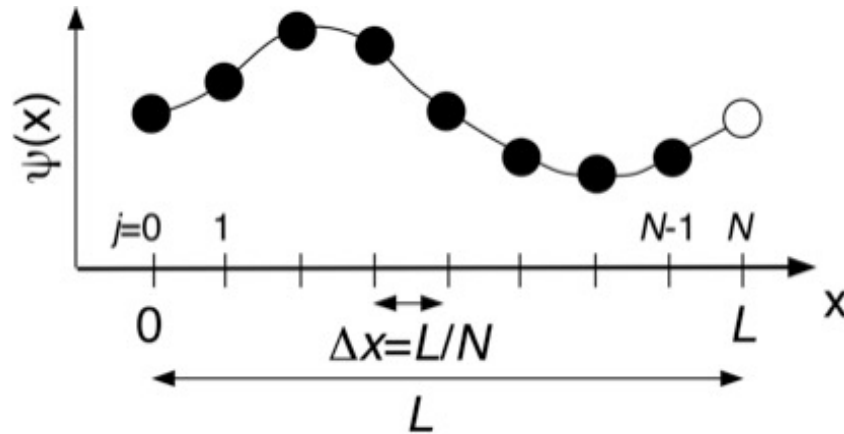
Resolution of identity:

$$1 = \sum_n |n\rangle\langle n|$$



Discrete Fourier Transform

- **Discretize** $\psi(x) \in \mathbb{C}$ ($x \in [0, L]$) on N mesh points, $x_j = j\Delta x$ ($j = 0, \dots, N-1$), with equal mesh spacing, $\Delta x = L/N$
- **Periodic boundary condition:** $\psi(x + L) = \psi(x)$

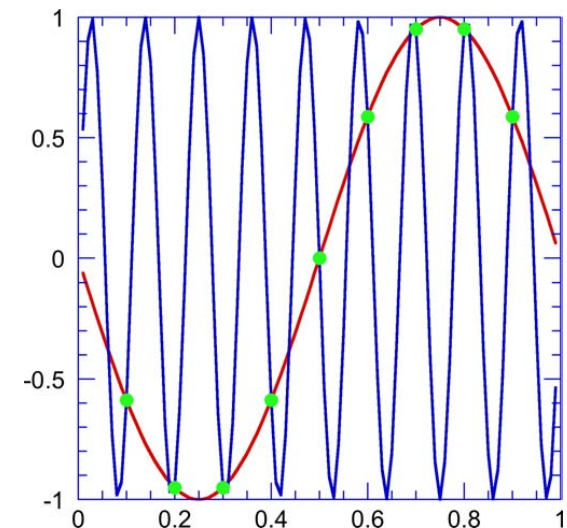


- **Discrete Fourier transform:** Represents $\psi(x)$ as a linear combination of $\exp(ikx) = \cos(kx) + i \sin(kx)$, with different wave numbers, k

$$\psi_j = \sum_{m=0}^{N-1} \tilde{\psi}_m \exp(ik_m x_j)$$

$$k_m = \begin{cases} 2\pi m/L & (m = 0, 1, \dots, N/2 - 1) \\ 2\pi (m - N)/L & (m = N/2, N/2 + 1, \dots, N - 1) \end{cases}$$

$$\tilde{\psi}_m = \frac{1}{N} \sum_{j=0}^{N-1} \psi_j \exp(-ik_m x_j)$$



Wave Numbers

- Periodic boundary condition, $\psi(x + L) = \psi(x)$, is guaranteed by choosing $k_m = 2\pi m/L$

$$e^{i\left(\frac{2\pi m}{L}(x+L)\right)} = e^{i\left(\frac{2\pi m}{L}x + 2\pi m\right)} = e^{i\frac{2\pi m}{L}x} \underbrace{(e^{i2\pi})^m = 1^m = 1}_{e^{i2\pi m}} = e^{i\frac{2\pi m}{L}x}$$

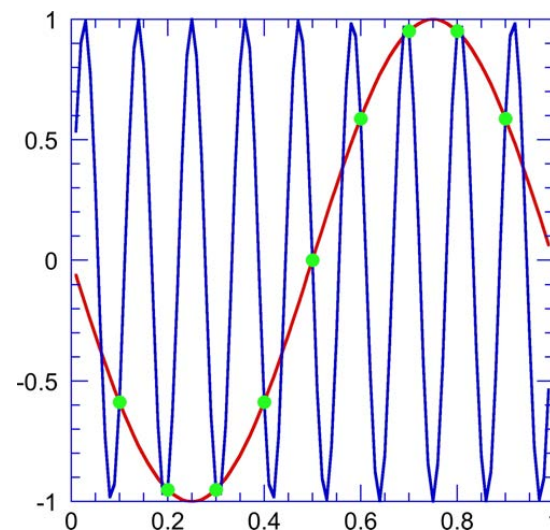
- Folding back the latter half of wave numbers by $\frac{2\pi N}{L} = 2\pi/\Delta x$ (cf. first Brillouin zone):

$$\text{From } \left[0, \frac{2\pi N}{L}\right] = \left[0, \frac{2\pi}{\Delta x}\right] \text{ to } \left[-\frac{2\pi}{L} \cdot \frac{N}{2}, \frac{2\pi}{L} \cdot \frac{N}{2}\right] = \left[-\frac{\pi}{\Delta x}, \frac{\pi}{\Delta x}\right]$$

$$\Delta x = \frac{L}{N}$$

The shift won't change wave-function value on any grid point

$$\begin{aligned} & e^{i\left(k_m - \frac{2\pi N}{L}\right)x_j} \\ &= e^{ik_mx_j} e^{-i\frac{2\pi N}{L} \cdot \frac{L}{N}j} \\ &= e^{ik_mx_j} e^{-i2\pi j} \\ &= e^{ik_mx_j} \end{aligned}$$



Orthonormal Basis Set

- **N-dimensional vector space:** $|\psi\rangle = (\psi_0, \psi_1, \dots, \psi_{N-1})$
- **Plane-wave basis set:** $\{|m\rangle = b_m(x_j) = \frac{1}{\sqrt{N}} \exp(ik_m x_j) \mid m = 0, 1, \dots, N-1\}$
- **Orthonormality:** $\langle m|n\rangle = \sum_{j=0}^{N-1} b_m^*(x_j) b_n(x_j) = \delta_{m,n} = \begin{cases} 1 & (m = n) \\ 0 & (m \neq n) \end{cases}$

$$\begin{aligned} \therefore \langle m|n\rangle &= \frac{1}{N} \sum_{j=0}^{N-1} \exp(i(k_n - k_m)x_j) = \frac{1}{N} \sum_{j=0}^{N-1} \exp\left(i \frac{2\pi}{N} (n - m)j\right) \\ &= \begin{cases} \frac{1}{N} \frac{\exp(i2\pi(n-m)) - 1}{\exp(i \frac{2\pi}{N} (n-m)) - 1} = 0 & (m \neq n) \\ \frac{1}{N} \cdot N = 1 & (m = n) \end{cases} \end{aligned}$$

$(k_n - k_m)x_j = \frac{2\pi(n-m)}{L} \cdot \frac{L}{N}j$
Geometric series
 $S = \sum_{j=0}^{N-1} \left(e^{i \frac{2\pi}{N} (n-m)} \right)^j = 1 + \dots + r^{N-1}$
 $rS = r + \dots + r^N$
 $\therefore (r-1)S = (r^N - 1)$

$|\psi\rangle = \sum_m c_m |m\rangle$
 $\langle n| \times \Downarrow$

- **Completeness:** $|\psi\rangle = \sum_{m=0}^{N-1} |m\rangle \langle m|\psi\rangle$ **or** $1 = \sum_{m=0}^{N-1} |m\rangle \langle m|$ $\langle n|\psi\rangle = \sum_m c_m \overbrace{\langle n|m\rangle}^{\delta_{nm}} = c_n$

- **Fourier transform:** $\psi_j = \sum_{m=0}^{N-1} \exp(ik_m x_j) \left[\frac{1}{N} \sum_{l=0}^{N-1} \exp(-ik_m x_l) \psi_l \right] \tilde{\psi}_m$

Spectral Method

- Kinetic-energy operator is diagonal in the momentum space:** $\tilde{\psi}_m \xrightarrow{T} \frac{k_m^2}{2} \tilde{\psi}_m$

$$-\frac{1}{2} \frac{\partial^2}{\partial x^2} \sum_{m=0}^{N-1} \tilde{\psi}_m \exp(ik_m x) = \sum_{m=0}^{N-1} \frac{k_m^2}{2} \tilde{\psi}_m \exp(ik_m x)$$

- Potential-energy operator is diagonal in the real space:** $\psi_j \xrightarrow{V} V_j \psi_j$
- Split-operator technique & spectral method**

$$\psi(t + \Delta t) = \exp\left(-\frac{iV\Delta t}{2}\right) \xrightarrow{F} \exp(-iT\Delta t) \xrightarrow{F^{-1}} \exp\left(-\frac{iV\Delta t}{2}\right) \psi(t) + O((\Delta t)^3)$$

$$1. \quad \psi_j \xrightarrow{\exp(-iV\Delta t/2)} \exp(-iV_j \Delta t/2) \psi_j$$

$$2. \quad \psi_j \xrightarrow{F^{-1}} F^{-1} \psi_j = \tilde{\psi}_m = \frac{1}{N} \sum_{j=1}^N \psi_j \exp(-ik_m x_j)$$

$$3. \quad \tilde{\psi}_m \xrightarrow{\exp(-iT\Delta t)} \exp(-ik_m^2 \Delta t/2) \tilde{\psi}_m \quad \exp\left(\frac{i\Delta t}{2} \frac{\partial^2}{\partial x^2}\right) \sum_m \tilde{\psi}_m e^{-ik_m x} = \sum_m \exp\left(-\frac{i\Delta t k_m^2}{2}\right) \tilde{\psi}_m e^{-ik_m x}$$

$$4. \quad \tilde{\psi}_m \xrightarrow{F} F \tilde{\psi}_m = \psi_j = \sum_{m=1}^N \tilde{\psi}_m \exp(ik_m x_j)$$

$$5. \quad \psi_j \xrightarrow{\exp(-iV\Delta t/2)} \exp(-iV_j \Delta t/2) \psi_j$$

Exact exponentiation!

Numerical Recipes FFT: four1()

- Spectral method requires

$$\psi_j \xrightarrow{F^{-1}} F^{-1}\psi_j = \tilde{\psi}_m = \frac{1}{N} \sum_{j=1}^N \psi_j \exp(-ik_m x_j)$$

$$\tilde{\psi}_m \xrightarrow{F} F\tilde{\psi}_m = \psi_j = \sum_{m=1}^N \tilde{\psi}_m \exp(ik_m x_j)$$

- `four1(double data[], unsigned long nn, int isign)`

On input, the `data[]` array contains $2*nn$ elements that represent nn complex function values, such that `data[2*j-1]` & `data[2*j]` ($j = 1, \dots, nn$) are the real & imaginary parts of the function value on the j -th grid point

- On output `data[]` is replaced by:

— `isign = 1`

$$data_j \leftarrow \sum_{m=1}^N data_m \exp(i2\pi mj / N)$$

— `isign = -1`

$$data_m \leftarrow \sum_{j=1}^N data_j \exp(-i2\pi mj / N)$$

$$k_m x_j = \frac{2\pi m}{L} \times \frac{L}{N} j = \frac{2\pi mj}{N}$$

- Note that `four1()` does not perform the division by N in F^{-1}

See `four1.c` in the class home page

Using four1()

- Define double psi[2*N], where psi[2*j] & psi[2*j+1] (j = 0, ..., N-1) are the real & imaginary parts of ψ_j

```
/* Fourier transform */
four1(psi-1, (unsigned long) N, 1);

/* Inverse Fourier transform */
four1(psi-1, (unsigned long) N, -1);
for (int j=0; j<2*N; j++)
    psi[j] /= N;
```

- Note that four1() assumes 1 offset for the first argument but psi[] is 0 offset

$$\psi_j \xrightarrow{F^{-1}} F^{-1}\psi_j = \tilde{\psi}_m = \frac{1}{N} \sum_{j=1}^N \psi_j \exp(-ik_m x_j)$$

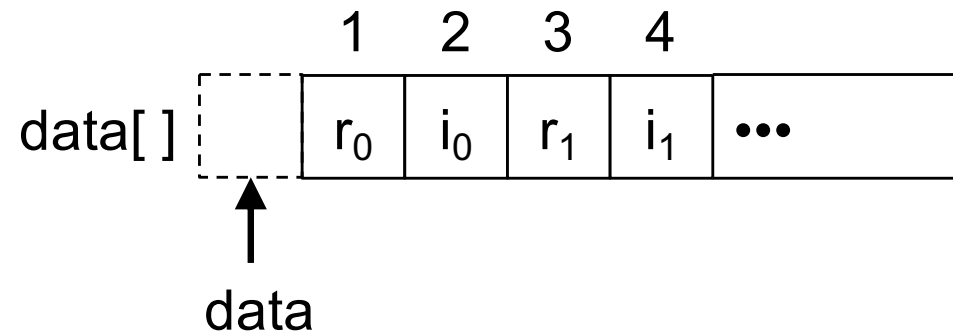
$$\tilde{\psi}_m \xrightarrow{F} F\tilde{\psi}_m = \psi_j = \sum_{m=1}^N \tilde{\psi}_m \exp(ik_m x_j)$$

	0	1	2	3		2N-2	2N-1
psi[]	Re ψ_0	Im ψ_0	Re ψ_1	Im ψ_1	...	Re ψ_{N-1}	Im ψ_{N-1}

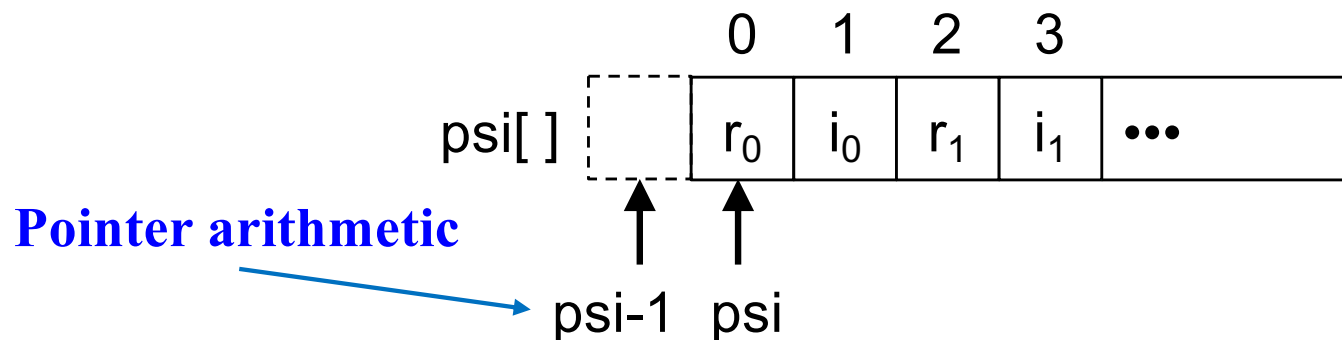
Array Offset

```
four1(psi-1, (unsigned long) N, 1);
```

- four1() assumes 1-offset (because of its Fortran origin)**



- But `psi[]` uses 0-offset (C convention)**



In C language, array name is a pointer to the zero-th element

Energy

$$\langle H \rangle = \langle T \rangle + \langle V \rangle$$

$$= \int dx \psi^*(x) \left(-\frac{1}{2} \frac{\partial^2}{\partial x^2} \right) \psi(x) + \int dx \psi^*(x) V(x) \psi(x)$$

discretize $\cong dx \sum_{j=0}^{N-1} \psi_j^* \left(-\frac{1}{2} \frac{\partial^2}{\partial x^2} \right) \psi_j + dx \sum_{j=0}^{N-1} \psi_j^* V_j \psi_j$

$$= dx N \sum_{m=0}^{N-1} \frac{k_m^2}{2} |\tilde{\psi}_m|^2 + dx \sum_{j=0}^{N-1} V_j |\psi_j|^2 \quad \text{weighted sums}$$

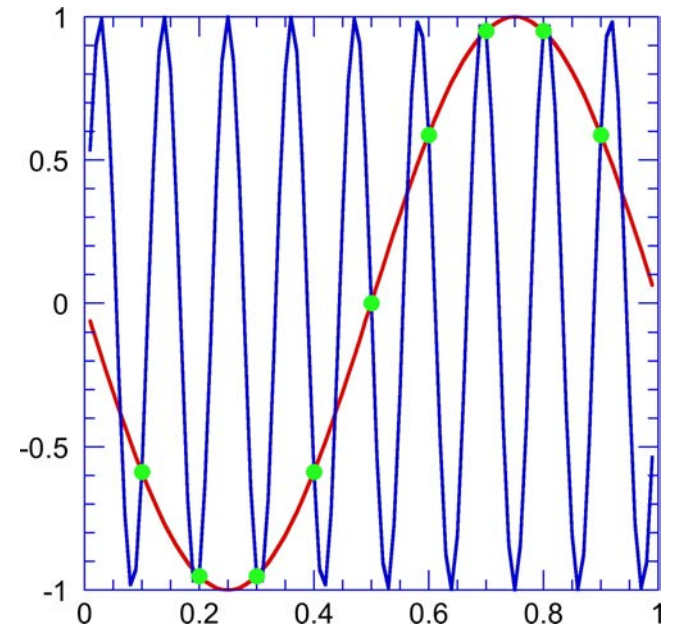
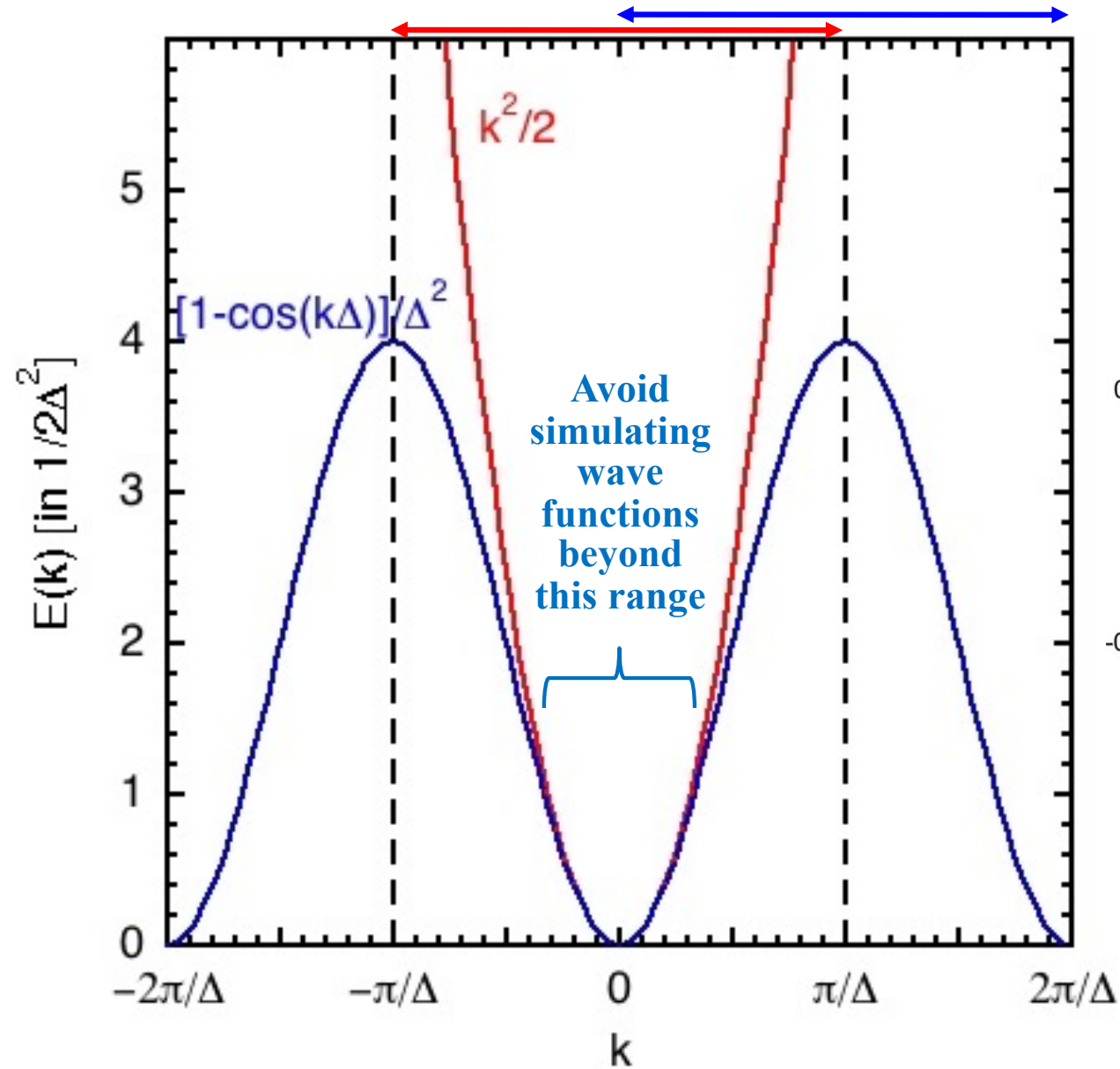
In calc_energy():

1. $\tilde{\psi}_m \leftarrow F^{-1}[\psi_j]$
2. $E_{\text{kin}} \leftarrow dx N \sum_{m=0}^{N-1} \frac{k_m^2}{2} |\tilde{\psi}_m|^2$
3. $\psi_j \leftarrow F[\tilde{\psi}_m]$ // don't forget to come back
4. $E_{\text{pot}} \leftarrow dx \sum_{j=0}^{N-1} V_j |\psi_j|^2$

Kinetic Energy

$$\begin{aligned}
 \langle T \rangle &= dx \sum_{j=0}^{N-1} \overbrace{\sum_{m=0}^{N-1} \tilde{\psi}_m^* \exp(-ik_m x_j)}^{\psi^*(x_j)} \left(-\frac{1}{2} \frac{\partial^2}{\partial x^2} \right) \overbrace{\sum_{n=0}^{N-1} \tilde{\psi}_n \exp(ik_n x_j)}^{\psi(x_j)} \\
 &= dx \sum_{j=0}^{N-1} \sum_{m=0}^{N-1} \tilde{\psi}_m^* \exp(-ik_m x_j) \sum_{n=0}^{N-1} \frac{k_n^2}{2} \tilde{\psi}_n \exp(ik_n x_j) \cong \frac{\frac{d^2}{dx^2} e^{ik_n x_j}}{\Delta x^2} e^{-ik_n \Delta x - 2 + e^{ik_n \Delta x}} e^{ik_n x_j} \\
 &= dx \sum_{m=0}^{N-1} \tilde{\psi}_m^* \sum_{n=0}^{N-1} \frac{k_n^2}{2} \tilde{\psi}_n \underbrace{\sum_{j=0}^{N-1} \exp(i(k_n - k_m)x_j)}_{N\langle m|n \rangle} = \frac{2[\cos(k_n \Delta x) - 1]}{\Delta x^2} e^{ik_n x_j} \\
 &= dx \sum_{m=0}^{N-1} \tilde{\psi}_m^* \sum_{n=0}^{N-1} \frac{k_n^2}{2} \tilde{\psi}_n N \delta_{m,n} \quad \text{addition theorem} \\
 &= dx N \sum_{m=0}^{N-1} \frac{k_m^2}{2} |\tilde{\psi}_m|^2 \quad \xrightarrow{\Delta x \rightarrow 0} \frac{k_n^2}{2}
 \end{aligned}$$

Continuum vs. Discrete Kinetic Energy



Total Energy Conservation

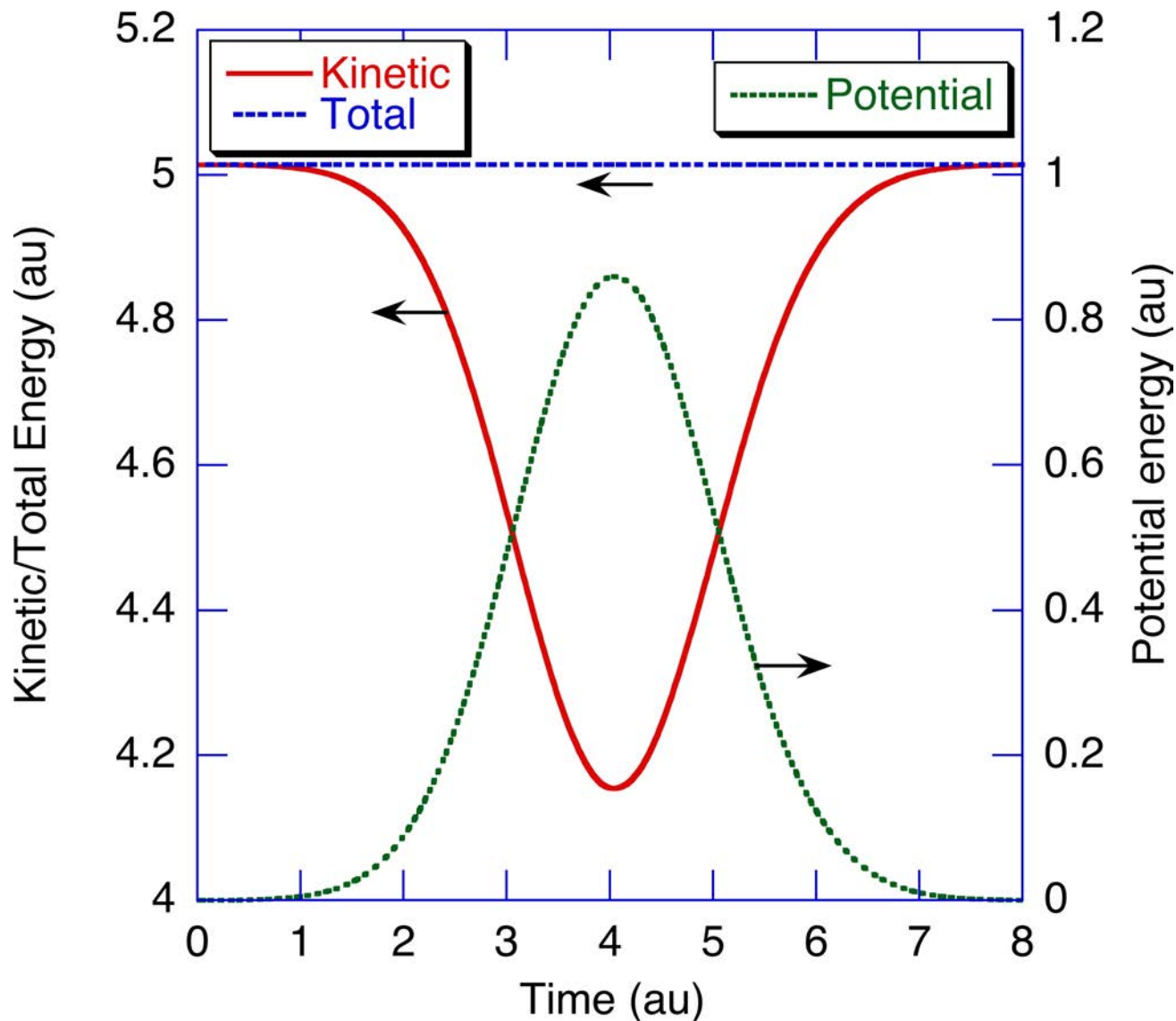
- **Energy eigenvalues & eigenvectors:** $H|n\rangle = \varepsilon_n|n\rangle \quad (n = 0, \dots, N-1)$
- **Time evolution of a wave function**

$$|\psi(t)\rangle = \exp(-iHt) \sum_{n=0}^{N-1} |n\rangle \langle n|\psi(0)\rangle = \sum_{n=0}^{N-1} \exp(-i\varepsilon_n t) |n\rangle \langle n|\psi(0)\rangle$$

- **Total energy** = 1: completeness (resolution of identity)

$$\begin{aligned} \langle \psi(t) | H | \psi(t) \rangle &= \left(\sum_{m=0}^{N-1} \langle \psi(0) | m \rangle \exp(i\varepsilon_m t) \langle m | \right) H \left(\sum_{n=0}^{N-1} \exp(-i\varepsilon_n t) | n \rangle \langle n | \psi(0) \rangle \right) \\ &= \left(\sum_{m=0}^{N-1} \langle \psi(0) | m \rangle \exp(i\varepsilon_m t) \langle m | \right) \left(\sum_{n=0}^{N-1} \exp(-i\varepsilon_n t) \varepsilon_n | n \rangle \langle n | \psi(0) \rangle \right) \\ &= \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} \exp(i(\varepsilon_m - \varepsilon_n)t) \varepsilon_n \langle \psi(0) | m \rangle \langle n | \psi(0) \rangle \underbrace{\langle m | n \rangle}_{\delta_{m,n}} \\ &= \sum_{n=0}^{N-1} \varepsilon_n |\langle n | \psi(0) \rangle|^2 = \text{constant} \end{aligned}$$

Energy Conservation for 1D Square Barrier



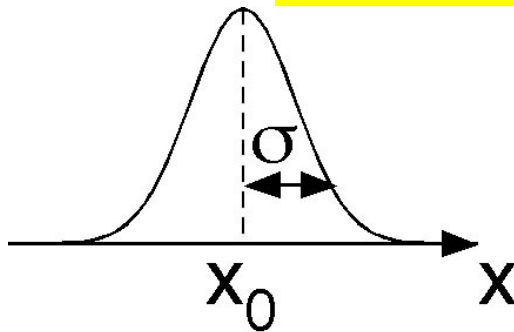
Energy conservation: good program verification

Initial Wave Function

- Gaussian wave packet

$$\psi(x, t = 0) = C \exp\left(-\frac{(x - \overset{x_0}{x_0})^2}{4\sigma^2}\right) \exp(ik_0 x)$$

σ $E_0 = k_0^2/2$



$$i \frac{\partial}{\partial t} \exp\left(ik_0 x - i \frac{k_0^2}{2} t\right) = \frac{k_0^2}{2} \exp\left(ik_0 x - i \frac{k_0^2}{2} t\right)$$

Free-space solution

||

$$\left[-\frac{1}{2} \frac{d^2}{dx^2} + 0\right] \exp\left(ik_0 x - i \frac{k_0^2}{2} t\right) = \frac{k_0^2}{2} \exp\left(ik_0 x - i \frac{k_0^2}{2} t\right)$$

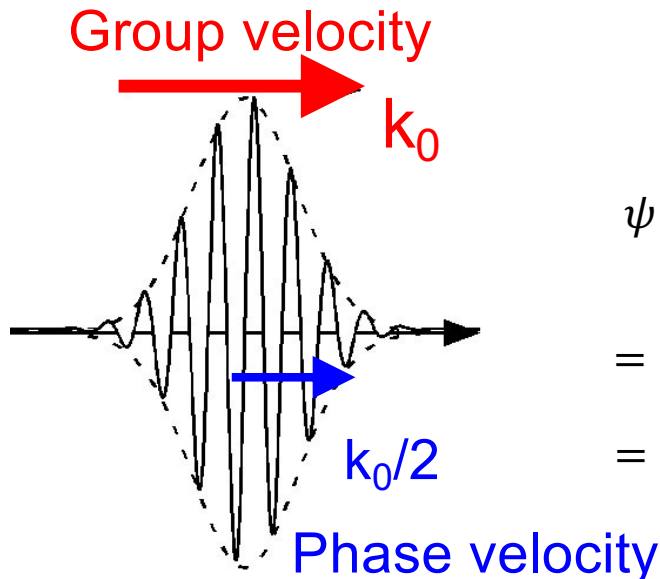
$$\psi(x, t) = \int dk \tilde{\psi}(k) \exp(ikx - i\omega(k)t) \quad \text{Here, } \omega(k) = k^2/2$$

$$= \int dk \tilde{\psi}(k) \exp\left(i \overset{k_0 + k - k_0}{\tilde{k}} x - i \overset{\cong \omega(k_0) + \frac{d\omega}{dk_0}(k - k_0)}{\omega(k)} t\right)$$

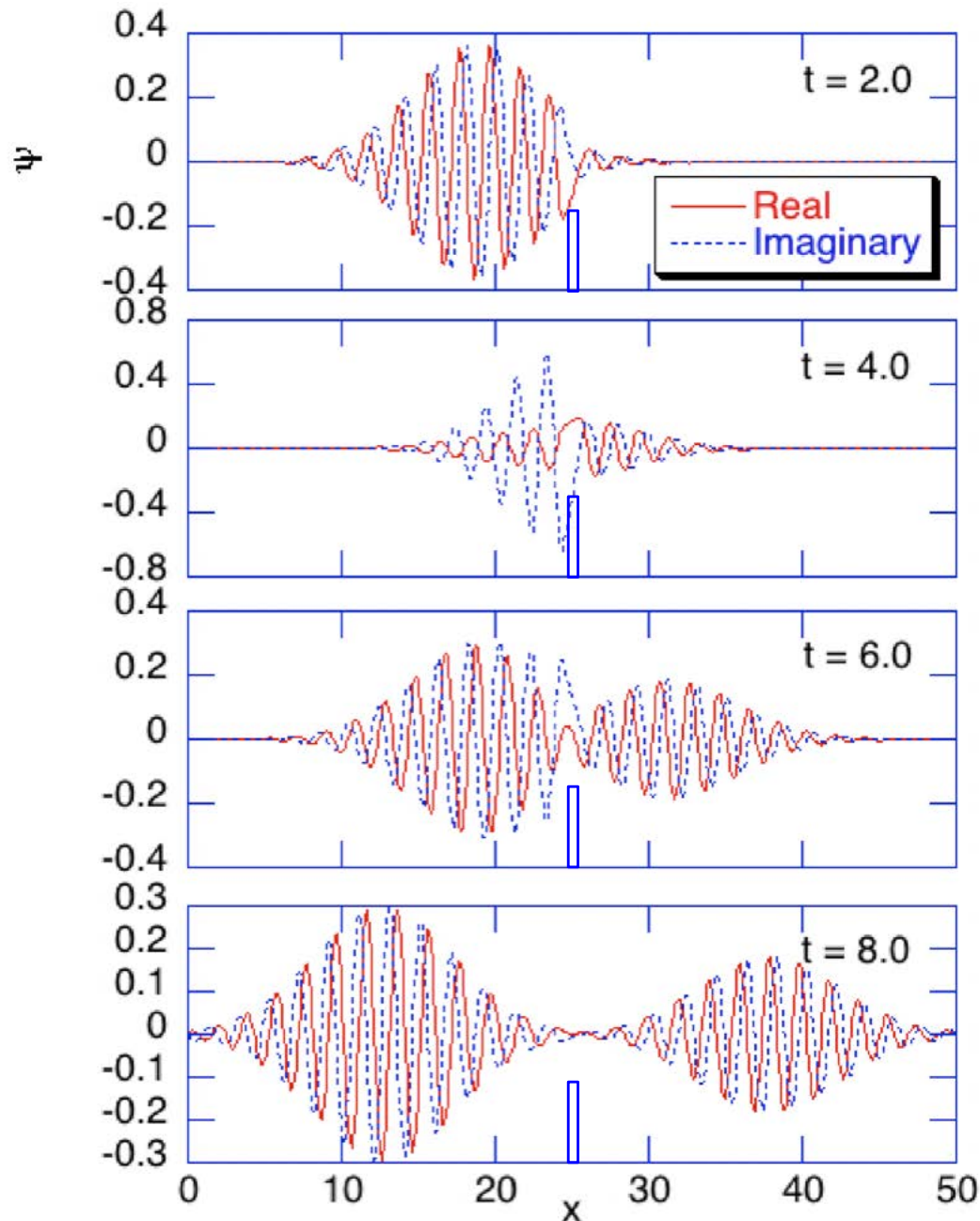
$$= \exp\left(ik_0 \left(x - \frac{\omega(k_0)}{k_0} t\right)\right) \int dk \tilde{\psi}(k) \exp\left(i(k - k_0) \left(x - \frac{d\omega}{dk_0} t\right)\right)$$

Phase velocity

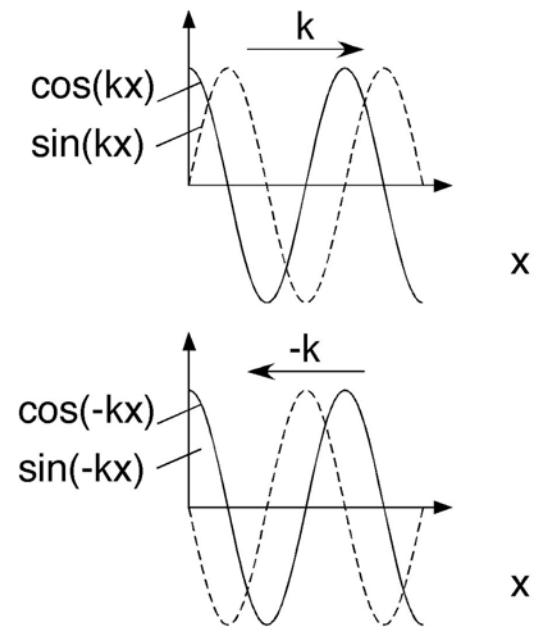
Group velocity



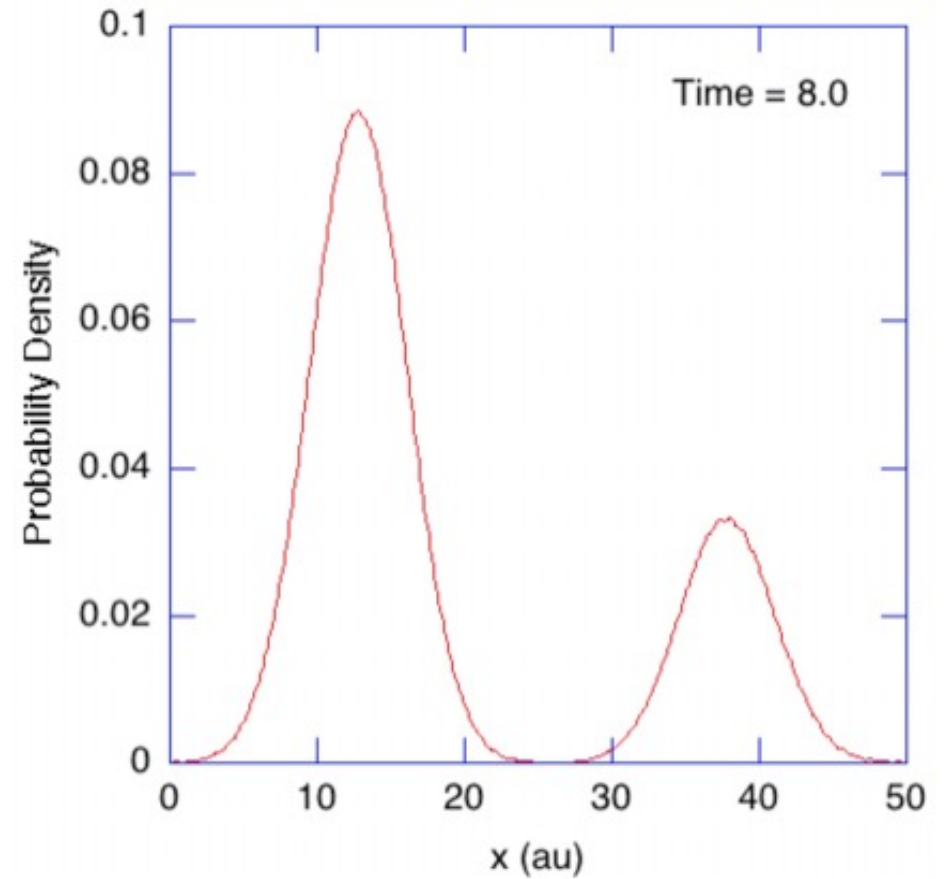
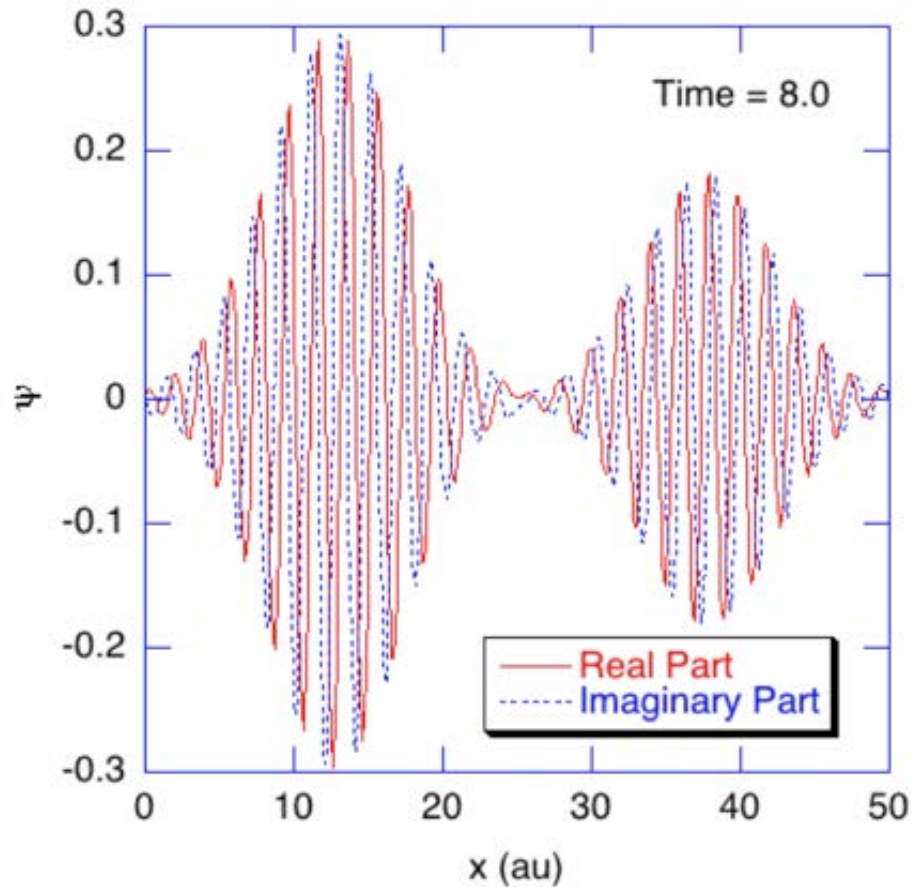
Numerical Example



$$\Delta t = 2 \times 10^{-3} \text{ au}$$



Wave Function & Probability

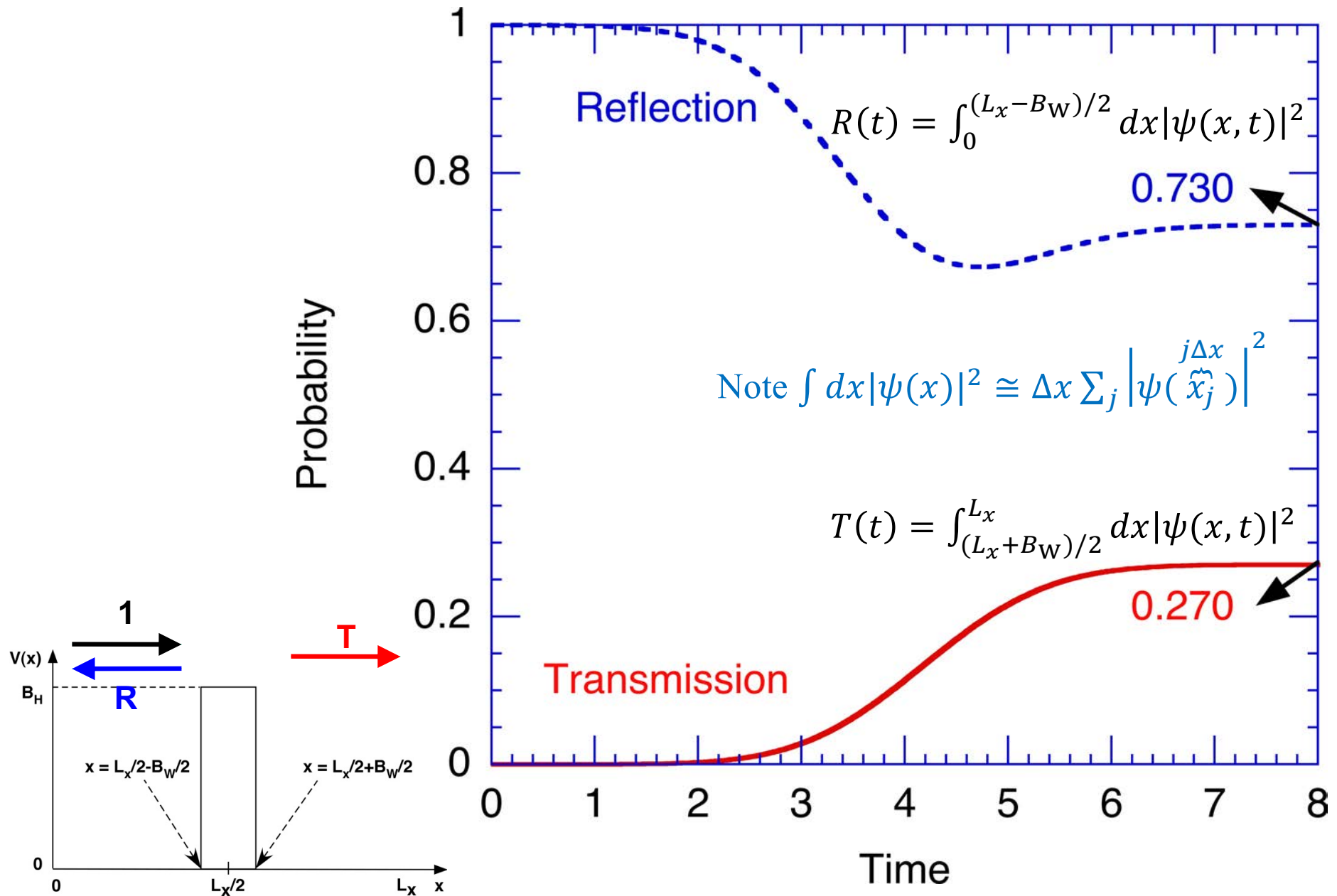


```
50.0          LX
2.0e-3        DT
4000          NSTEP
10            NECAL
12.5 3.0 5.0  X0 S0 E0
5.0 1.0       BH BW
50.0          EH
```

qd1.in

```
qd1.h  #define NX 512
```

Transmission & Reflection Coefficients



Top 10 Algorithms in History

In putting together this issue of *Computing in Science & Engineering*, we knew three things: it would be difficult to list just 10 algorithms; it would be fun to assemble the authors and read their papers; and, whatever we came up with in the end, it would be controversial. We tried to assemble the 10 algorithms with the greatest influence on the development and practice of science and engineering in the 20th century. Following is our list (here, the list is in chronological order; however, the articles appear in no particular order):

- Metropolis Algorithm for Monte Carlo
- Simplex Method for Linear Programming
- Krylov Subspace Iteration Methods
- The Decompositional Approach to Matrix Computations
- The Fortran Optimizing Compiler
- QR Algorithm for Computing Eigenvalues
- Quicksort Algorithm for Sorting
- Fast Fourier Transform
- Integer Relation Detection
- Fast Multipole Method

PHYS 516

CSCI 596

CSCI 653

IEEE CiSE, Jan/Feb (2000)

Fast Fourier Transform

- Danielson-Lanczos algorithm:**

$$\begin{aligned}
 \psi_j &= \sum_{m=0}^{N-1} \tilde{\psi}_m \exp(ik_m x_j) = \sum_{m=0}^{N-1} \tilde{\psi}_m \exp(i2\pi m j / N) \quad \textcolor{red}{O(N^2)!} \\
 &= \sum_{m=0}^{N/2-1} \tilde{\psi}_{2m} \exp(i2\pi (2m) j / N) + \sum_{m=0}^{N/2-1} \tilde{\psi}_{2m+1} \exp(i2\pi (2m+1) j / N) \\
 &\quad \text{even terms} \qquad \qquad \qquad \text{odd terms} \\
 &= \sum_{m=0}^{N/2-1} \tilde{\psi}_{2m} \exp(i2\pi m j / (N/2)) + \exp(i2\pi j / N) \sum_{m=0}^{N/2-1} \tilde{\psi}_{2m+1} \exp(i2\pi m j / (N/2))
 \end{aligned}$$

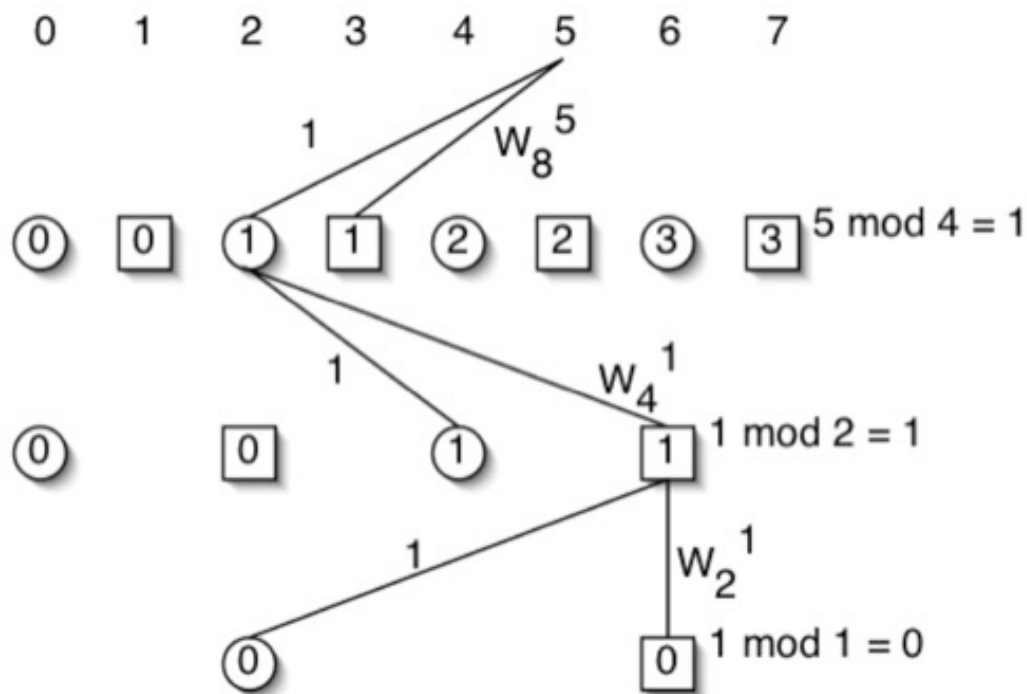
$$\psi_j = \psi_j^0 + W_N^j \psi_j^1 \quad \left\{ \begin{array}{l} \psi_j^0 = \sum_{m=0}^{N/2-1} \tilde{\psi}_{2m} \exp(i2\pi m j / (N/2)) \quad \text{subarray Fourier decompositions} \\ \psi_j^1 = \sum_{m=0}^{N/2-1} \tilde{\psi}_{2m+1} \exp(i2\pi m j / (N/2)) \\ W_N = \exp(i2\pi / N) \quad j \text{ read as } j \bmod N/2 \end{array} \right.$$

Divide-and-conquer

Fast Fourier Transform

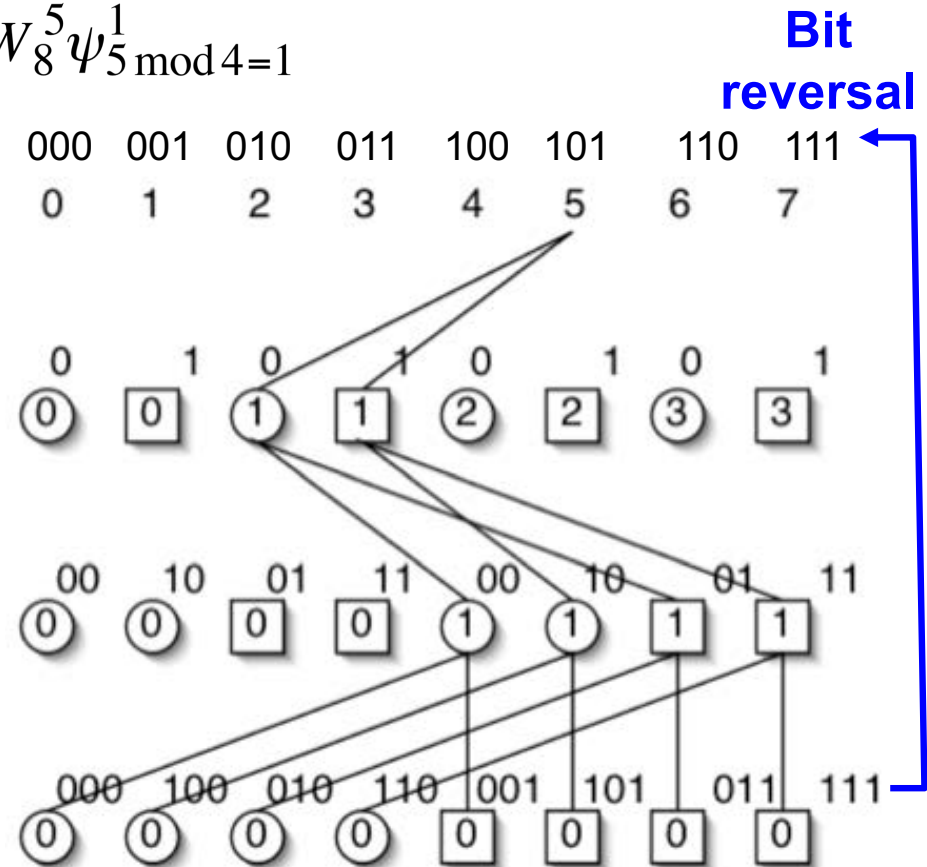
- Recursive sub-Fourier transforms: $\psi_j = \psi_j^0 + W_N^j \psi_j^1$

$$\psi_5 = \psi_{5 \bmod 4=1}^0 + W_8^5 \psi_{5 \bmod 4=1}^1$$



010 → 010 → 2 011 → 110 → 6

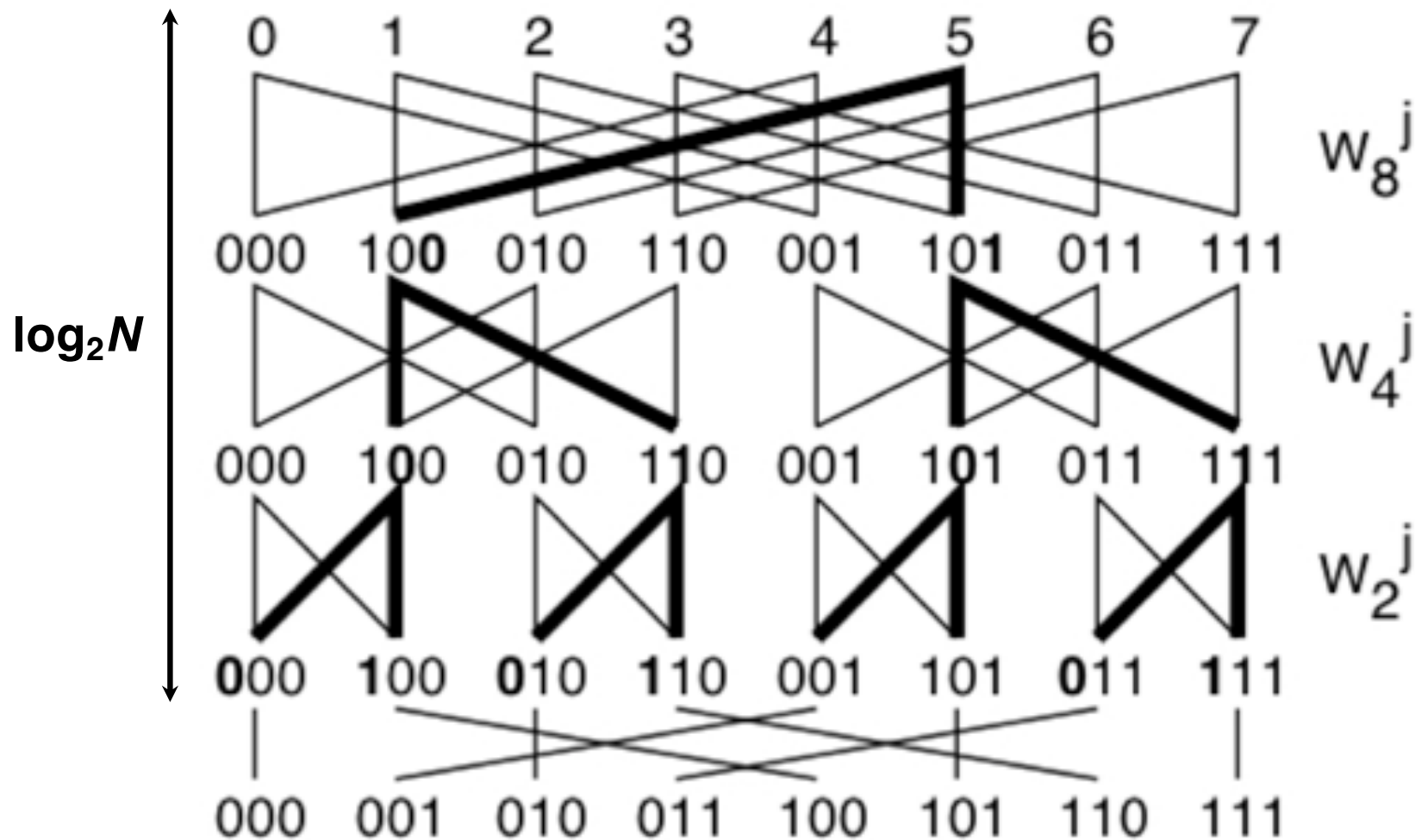
Path bit reversal = element



Recursion tree
 $O(N)$ operations per element

Fast Fourier Transform Algorithm

- Butterfly (hypercube) data exchange after bit-reversal:

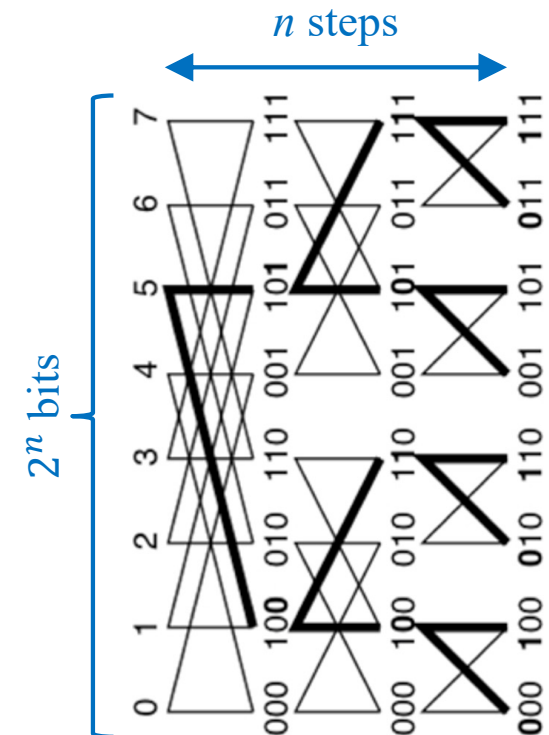
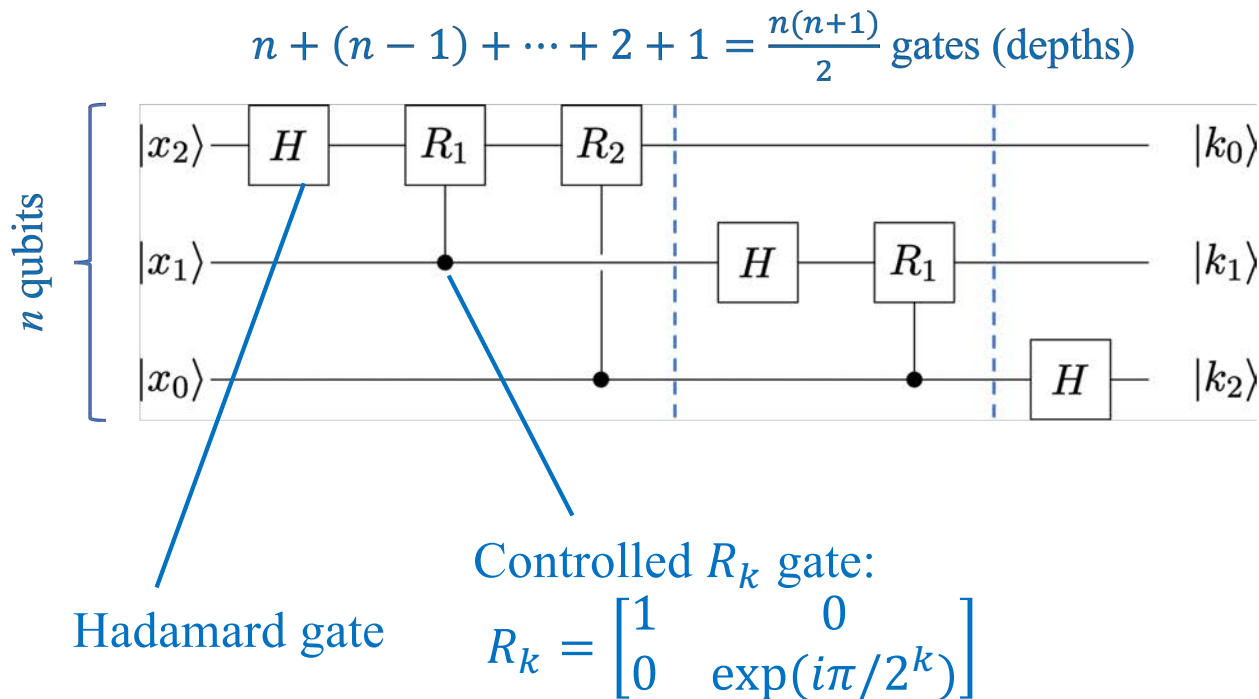


- Many computations are shared among the recursion trees
- $2N\log_2 N$ arithmetic operations

Quantum Fourier Transform (QFT)

- On a quantum computer, quantum parallelism allows Fourier transform to be performed using only n qubits & $O(n^2)$ gates for $N = 2^n$

$$\text{QFT: } |x\rangle \rightarrow \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{i\frac{2\pi kx}{N}} |k\rangle$$



- Compare QFT with $2N \log_2 N = O(2^n n)$ arithmetic operations in classical fast Fourier transform (FFT):

Exponential operation reduction: $O(2^n n/n^2) = O(2^n/n)$

Parallelizing Quantum Dynamics

Aiichiro Nakano

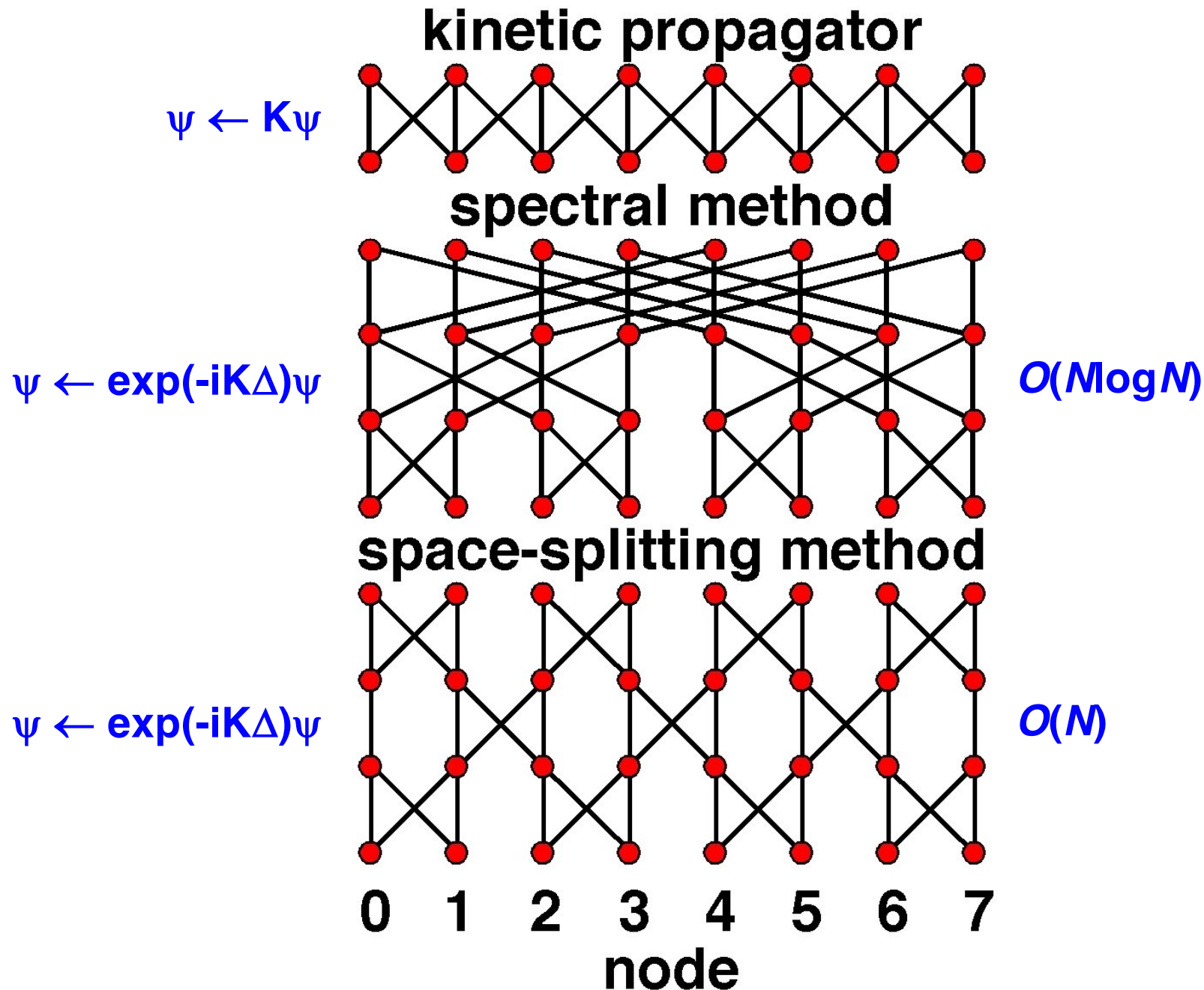
*Collaboratory for Advanced Computing & Simulations
Department of Computer Science
Department of Physics & Astronomy
Department of Chemical Engineering & Materials Science
Department of Quantitative & Computational Biology
University of Southern California*

Email: anakano@usc.edu

See <https://aiichironakano.github.io/cs596.html>

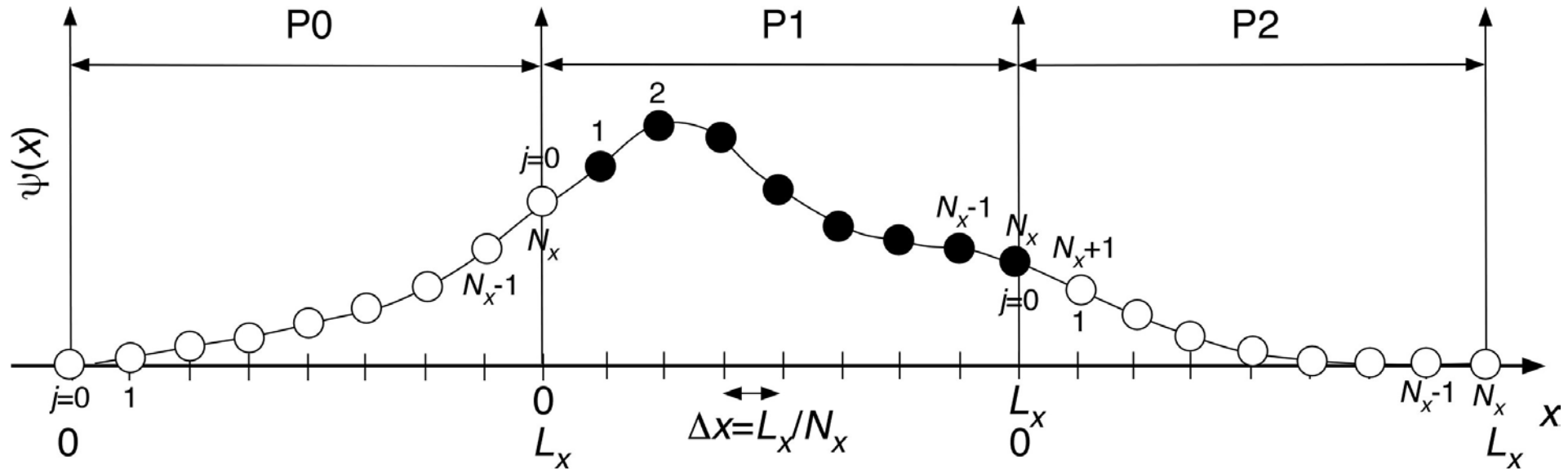


Parallel QD Communications



Parallelization of Space Splitting Method

- Self-centric spatial decomposition



- Local & global coordinates

$$\begin{cases} x_j = j\Delta x \\ x_j^{(\text{global})} = j\Delta x + pL_x \end{cases}$$

- Global coordinates only in `init_prop()` & `init_wavefn()`

Boundary Wave Function Caching

- Parallelized `periodic_bc()`

```
plw = (myid-1+nproc)%nproc; /* Lower partner process */  
pup = (myid+1)%nproc; /* Upper partner process */
```

```
/* Cache boundary wave function value at the lower end */
```

```
dbuf[0:1] ← psi[NX][0:1];
```

```
Send dbuf to pup;
```

```
Receive dbufr from plw;
```

```
psi[0][0:1] ← dbufr[0:1];
```

```
/* Cache boundary wave function value at the upper end */
```

```
dbuf[0:1] ← psi[1][0:1];
```

```
Send dbuf to plw;
```

```
Receive dbufr from pup;
```

```
psi[NX+1][0:1] ← dbufr[0:1];
```

