Quantum Monte Carlo Simulation

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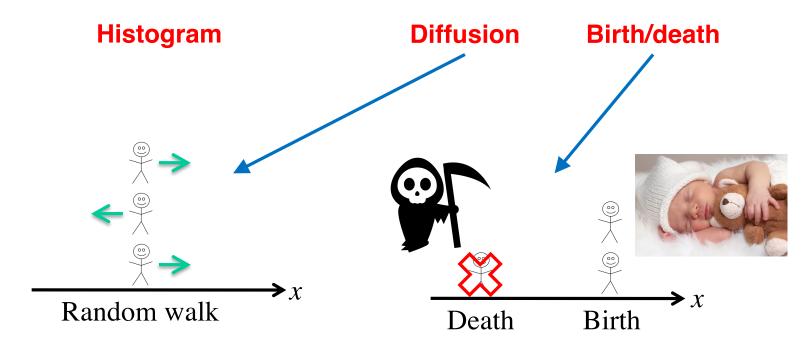
Random walk to solve quantum eigenvalue problem



Imaginary-time Schrödinger Equation

•
$$\tau = it$$

$$\frac{\partial}{\partial \tau} \psi(x, \tau) = -\frac{H}{\hbar} \psi(x, \tau) = \frac{\hbar}{2m} \frac{\partial^2}{\partial x^2} \psi(x, \tau) - \frac{V(x)}{\hbar} \psi(x, \tau)$$



$$D_{\text{QMC}} = \frac{\hbar}{2m} = 58 \text{ m}^2/\text{s} \gg D_{\text{water}} = 10^{-9} \text{ m}^2/\text{s}$$

Ground-State Filter

$$\psi(x,\tau) = \exp(-H\tau/\hbar)\psi(x,0) = \sum_{n} c_{n}\phi_{n}(x)\exp(-E_{n}\tau/\hbar) = 1$$

$$c_{n} = \int dx\phi_{n}(x)\psi(x,0)$$
Completeness: $|\psi\rangle = \sum_{n} |n\rangle\langle n|\psi\rangle$

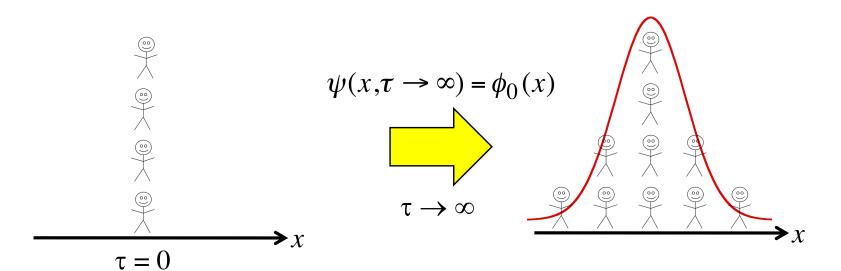
$$\langle 0|H|0\rangle = E_{0} \langle E_{1} \langle \cdots$$

$$\lim_{\tau \to \infty} \psi(x, \tau) = c_0 \phi_0(x) \exp(-E_0 \tau/\hbar)$$

• Energy estimate

It's not
$$\psi^*(x,\tau) \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right) \psi(x,\tau)$$

$$E_0 = \langle V \rangle = \frac{\int_{-\infty}^{\infty} V(x) \psi(x,\tau) dx}{\int_{-\infty}^{\infty} \psi(x,\tau) dx} \cong \frac{1}{N} \sum_{i=1}^{N} V(x_i)$$



Reference Energy

Energy measured relative to reference energy to achieve stationary population

$$\frac{\partial}{\partial \tau} \psi(x,\tau) = \frac{\hbar}{2m} \frac{\partial^2}{\partial x^2} \psi(x,\tau) - \frac{V(x) - V_{\text{ref}}}{\hbar} \psi(x,\tau)$$

$$\int_{-\infty}^{\infty} \frac{\partial}{\partial \tau} \psi(x,\tau) dx = \left[\frac{\hbar}{2m} \frac{\partial}{\partial x} \psi(x,\tau) \right]_{-\infty}^{\infty} - \int \frac{V(x)}{\hbar} \psi(x,\tau) dx + \frac{V_{\text{ref}}}{\hbar} \int \psi(x,\tau) dx$$

$$= -\int \frac{V(x)}{\hbar} \psi(x,\tau) dx + \frac{V_{\text{ref}}}{\hbar} \int \psi(x,\tau) dx$$

$$\int_{-\infty}^{\infty} dx \frac{\partial^2}{\partial x^2} \psi(x) = \left[\frac{\partial}{\partial x} \psi(x) \right]_{-\infty}^{\infty} = 0$$

• For $\tau \to \infty$,

$$\psi(x,\tau) \approx c_0 \phi_0(x) \exp\left(-(E_0 - V_{\text{ref}}) \tau/\hbar\right)$$

$$\therefore \frac{\partial}{\partial \tau} \psi(x,\tau) \approx \frac{V_{\text{ref}} - E_0}{\hbar} \psi(x,\tau)$$

$$\int \frac{V_{\text{ref}} - E_0}{\hbar} \psi(x,\tau) dx = -\int \frac{V(x)}{\hbar} \psi(x,\tau) dx + \frac{V_{\text{ref}}}{\hbar} \int \psi(x,\tau) dx$$

$$\therefore \frac{E_0}{\hbar} \int \psi(x,\tau) dx = \int \frac{V(x)}{\hbar} \psi(x,\tau) dx \qquad \Longrightarrow \qquad E_0 = \frac{\int V(x) \psi(x,\tau) dx}{\int \psi(x,\tau) dx} = \langle V \rangle$$

Choice of Reference Energy

• Number of random walkers: $N(\tau) = \int_{-\infty}^{\infty} \psi(x, \tau) dx$

$$\frac{d}{d\tau}N = \frac{V_{\text{ref}} - \langle V \rangle}{\hbar} N(\tau)$$

$$\because \int_{-\infty}^{\infty} \frac{\partial}{\partial \tau} \psi(x, \tau) dx = -\int \frac{V(x)}{\hbar} \psi(x, \tau) dx + \frac{V_{\text{ref}}}{\hbar} \int_{-\infty}^{\infty} \psi(x, \tau) dx$$

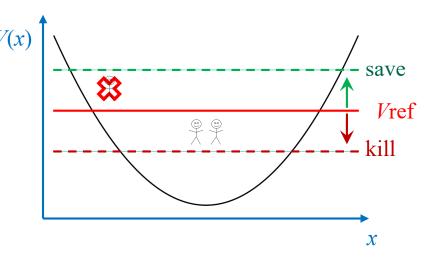
• Choose new $V_{
m ref}$ to counter-balance the current population growth

$$-\frac{d}{d\tau}N|_{\text{current}} = \frac{V_{\text{ref}}^{\text{new}} - \langle V \rangle}{\hbar}N(\tau)$$

$$\frac{V_{\text{ref}}^{\text{new}}}{\hbar} = \frac{\langle V \rangle}{\hbar} - \frac{dN/d\tau}{N(\tau)}$$

$$\therefore V_{\text{ref}} = \langle V \rangle - \frac{N(\tau + \Delta \tau) - N(\tau)}{N(\tau)\Delta \tau}$$

 $\hbar = 1$ (atomic unit)



For numerical stability, replace N by $N(\tau = 0) = N_0$

Quantum Monte Carlo Algorithm

- 1. Place N_0 walkers at the initial set of positions x_i .
- 2. Compute the reference energy, $V_{\text{ref}} = \sum_{i} V(x_i)/N_0$.
- 3. For each walker,

$$D = \frac{\hbar}{2m} = \frac{1}{2} (a. u.) = \frac{(\Delta s)^2}{2\Delta \tau}$$

- a. Randomly move the walker to the right or left by a fixed step length Δs .
- b. Compute $\Delta V = V(x) V_{\text{ref}}$ & a random number $r \in [0, 1]$. If $\Delta V > 0$ & $r < \Delta V \Delta \tau$ ($\Delta \tau = \Delta s^2$), then remove the walker. If $\Delta V < 0$ & $r < -\Delta V \Delta \tau$, then add another walker at x. Otherwise, just leave the walker at x.
- 4. Compute the mean potential energy $\langle V \rangle = \sum_i V(x_i)/N$ & the actual number of random walkers N. The new reference potential is given by

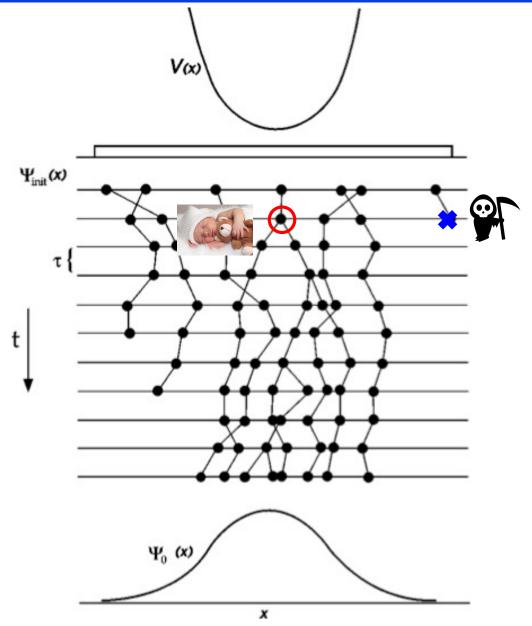
$$V_{\text{ref}} = \langle V \rangle - \frac{N(\tau + \Delta \tau) - N(\tau)}{N(\tau)\Delta \tau} N_0$$

$$\Delta \psi(x, \tau) = -\frac{\Delta V(x)}{\hbar} \Delta \tau \cdot \psi(x, \tau)$$

The average $\langle V \rangle$ is an estimate of the ground state energy.

5. Repeat steps 3–4 until the estimates of the ground state energy $\langle V \rangle$ have reached a steady state value with only random fluctuations. Average $\langle V \rangle$ over many Monte Carlo steps to compute the ground state energy.

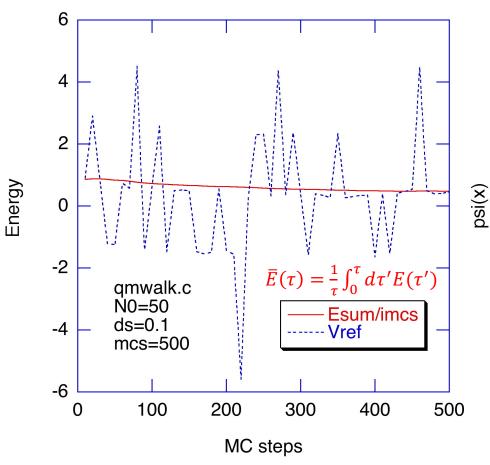
Random Walk + Birth/Death



W. M. C. Foulkes et al., Rev. Mod. Phys. 73, 33 ('01)

Numerical Example

• $V(x) = x^2/2$ (in atomic unit)



(after $0.4 \times 500 = 200$ MC steps of thermalization)

$$D = \frac{\hbar}{2m} = \frac{1}{2} (a. u.) = \frac{(\Delta s)^2}{2\Delta \tau}$$

$$\frac{1}{\tau} \int_{-\infty}^{+\infty} \frac{\Delta^2}{2} \varphi(\Delta) d\Delta = D$$

A. Einstein, Ann. Phys. 17, 549 (1905)

Where to Go from Here

• Types of QMC for many electrons: diffusion MC & variational MC (zero temperature); path-integral MC (finite temperature)

Interacting Electrons, Martin, Reining & Ceperley (Cambridge Univ. Press, '16)

• QMC for realistic materials

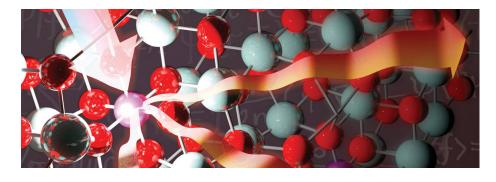
Quantum Monte Carlo simulations of solids, Foulkes et al., Rev. Mod. Phys. 73, 33 ('01)

QMC software

QMCPACK: http://qmcpack.org

QWalk: http://qwalk.github.io/mainline

CASINO: https://vallico.net/casinoqmc



Neural networks to represent many-electron wave functions (~ variational QMC)

NetKet: https://www.netket.org

