

Partial trace

Definition:

The partial trace tr_B is a mapping from the density matrices ρ_{AB} on a composite space $\mathcal{H}_A \otimes \mathcal{H}_B$ onto density matrices ρ_A on \mathcal{H}_A . It is defined as the linear extension of the mapping

$$\text{tr}_B: S \otimes T \mapsto \text{tr}(T)S$$

for any matrix S on \mathcal{H}_A and T on \mathcal{H}_B .

Let $\{|a_i\rangle\}$ be a basis of \mathcal{H}_A , and $\{|b_i\rangle\}$ be a basis of \mathcal{H}_B . Any density matrix ρ_{AB} on $\mathcal{H}_A \otimes \mathcal{H}_B$ can then be decomposed as $\rho_{AB} = \sum_{ijkl} c_{ijkl} |a_i\rangle\langle a_j| \otimes |b_k\rangle\langle b_l|$ and the partial trace reads

$$\text{tr}_B \rho_{AB} = \sum_{ijkl} c_{ijkl} |a_i\rangle\langle a_j| \langle b_l | b_k \rangle, \quad (1)$$

which is a density matrix ρ_A on \mathcal{H}_A . Note, that $\text{tr}|b_k\rangle\langle b_l| = \langle b_l | b_k \rangle$ is a complex number and note the exchange of indices compared to the operator.

Example:

The two-qubit spin singlet $|\psi^-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$ corresponds to the density matrix

$$|\psi^-\rangle\langle\psi^-| = \frac{1}{2}(|01\rangle\langle 01| - |01\rangle\langle 10| - |10\rangle\langle 01| + |10\rangle\langle 10|),$$

on which we act with the partial trace according to the definition (1)

$$\begin{aligned} \text{tr}_B |\psi^-\rangle\langle\psi^-| &= \frac{1}{2}(|0\rangle\langle 0| \langle 1 | 1 \rangle - |0\rangle\langle 1| \langle 0 | 1 \rangle - |1\rangle\langle 0| \langle 1 | 0 \rangle + |1\rangle\langle 1| \langle 0 | 0 \rangle) \\ &= \frac{1}{2}(|0\rangle\langle 0| + |1\rangle\langle 1|), \end{aligned}$$

to obtain the reduced state $\rho_A = \frac{1}{2}\mathbb{1}$.

Formula for two qubits:

A generic two-qubit state ρ_{AB} can be expanded with respect to the orthonormal basis $\{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$ as

$$\rho_{AB} = \rho_{11}|00\rangle\langle 00| + \rho_{12}|00\rangle\langle 01| + \rho_{13}|00\rangle\langle 10| + \dots + \rho_{44}|11\rangle\langle 11|.$$

The partial trace is calculated according to (1)

$$\rho_A = (\rho_{11} + \rho_{22})|0\rangle\langle 0| + (\rho_{13} + \rho_{24})|0\rangle\langle 1| + (\rho_{31} + \rho_{42})|1\rangle\langle 0| + (\rho_{33} + \rho_{44})|1\rangle\langle 1|.$$

The factor in front of $|1\rangle\langle 0|$, say, can be derived as follows: the preimages of $|1\rangle\langle 0|$ under the partial trace are $|10\rangle\langle 00|$ and $|11\rangle\langle 01|$ and come with factors ρ_{31} and ρ_{42} in ρ_{AB} . In matrix form the formula reads

$$\text{tr}_B \begin{pmatrix} \rho_{11} & \rho_{12} & \rho_{13} & \rho_{14} \\ \rho_{21} & \rho_{22} & \rho_{23} & \rho_{24} \\ \rho_{31} & \rho_{32} & \rho_{33} & \rho_{34} \\ \rho_{41} & \rho_{42} & \rho_{43} & \rho_{44} \end{pmatrix} = \begin{pmatrix} \text{tr} \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} & \text{tr} \begin{pmatrix} \rho_{13} & \rho_{14} \\ \rho_{23} & \rho_{24} \end{pmatrix} \\ \text{tr} \begin{pmatrix} \rho_{31} & \rho_{32} \\ \rho_{41} & \rho_{42} \end{pmatrix} & \text{tr} \begin{pmatrix} \rho_{33} & \rho_{34} \\ \rho_{43} & \rho_{44} \end{pmatrix} \end{pmatrix} = \begin{pmatrix} \rho_{11} + \rho_{22} & \rho_{13} + \rho_{24} \\ \rho_{31} + \rho_{42} & \rho_{33} + \rho_{44} \end{pmatrix}.$$

Partial trace over A:

Finally, we mention that the partial trace tr_A over the subsystem A is defined in the obvious way and can be calculated similarly to tr_B . For convenience we state the formula for two qubits

$$\text{tr}_A \begin{pmatrix} \rho_{11} & \rho_{12} & \rho_{13} & \rho_{14} \\ \rho_{21} & \rho_{22} & \rho_{23} & \rho_{24} \\ \rho_{31} & \rho_{32} & \rho_{33} & \rho_{34} \\ \rho_{41} & \rho_{42} & \rho_{43} & \rho_{44} \end{pmatrix} = \begin{pmatrix} \text{tr} \begin{pmatrix} \rho_{11} & \rho_{13} \\ \rho_{31} & \rho_{33} \end{pmatrix} & \text{tr} \begin{pmatrix} \rho_{12} & \rho_{14} \\ \rho_{32} & \rho_{34} \end{pmatrix} \\ \text{tr} \begin{pmatrix} \rho_{21} & \rho_{23} \\ \rho_{41} & \rho_{43} \end{pmatrix} & \text{tr} \begin{pmatrix} \rho_{22} & \rho_{24} \\ \rho_{42} & \rho_{44} \end{pmatrix} \end{pmatrix} = \begin{pmatrix} \rho_{11} + \rho_{33} & \rho_{12} + \rho_{34} \\ \rho_{21} + \rho_{43} & \rho_{22} + \rho_{44} \end{pmatrix}.$$

It is a good exercise to convince oneself that these formulae make sense with respect to the matrix representation of the tensor product

$$\begin{pmatrix} a_{00} & a_{01} \\ a_{10} & a_{11} \end{pmatrix} \otimes \begin{pmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{pmatrix} = \begin{pmatrix} a_{00} \begin{pmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{pmatrix} & a_{01} \begin{pmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{pmatrix} \\ a_{10} \begin{pmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{pmatrix} & a_{11} \begin{pmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{pmatrix} \end{pmatrix}.$$