

TNS and non-equilibrium dynamics

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Miguel Frías-Pérez, Luca Tagliacozzo (IFF)

PRB 106, 115117 (2022)
arXiv:2308.04291



MAX PLANCK INSTITUTE
OF QUANTUM OPTICS

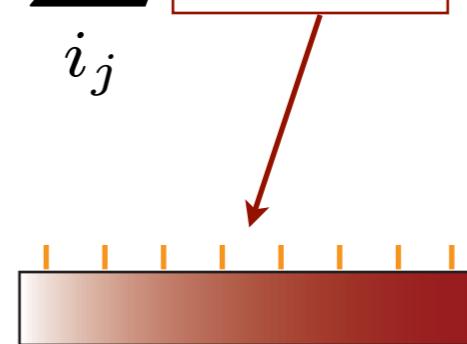


TNRG 31.10.2023

TNS in quantum many-body systems

A general state of the N-body Hilbert space has exponentially many coefficients

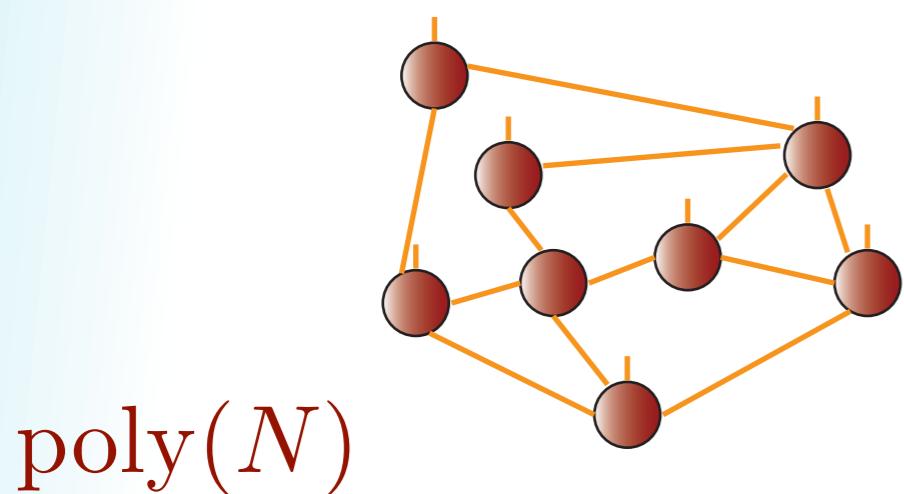
$$|\Psi\rangle = \sum_{i_j} [c_{i_1 \dots i_N}] |i_1 \dots i_N\rangle$$



N-legged tensor

ATNS has only a polynomial number of parameters

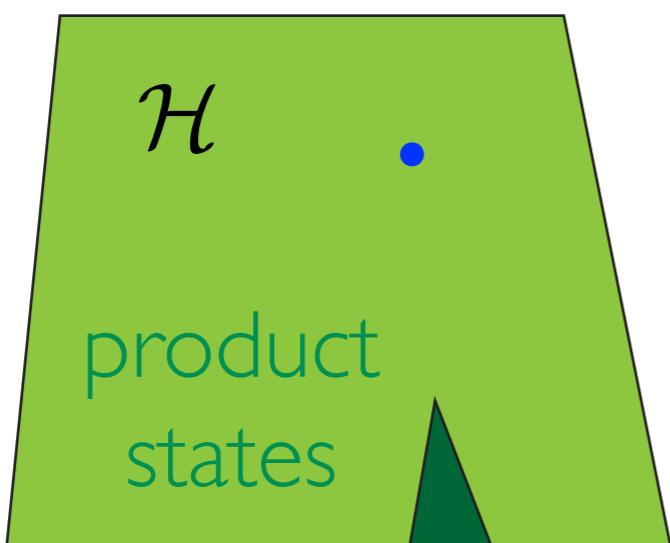
d^N



WHY SHOULD TNS BE USEFUL?

States appearing in Nature are peculiar

State at random from
Hilbert space is not
close to product

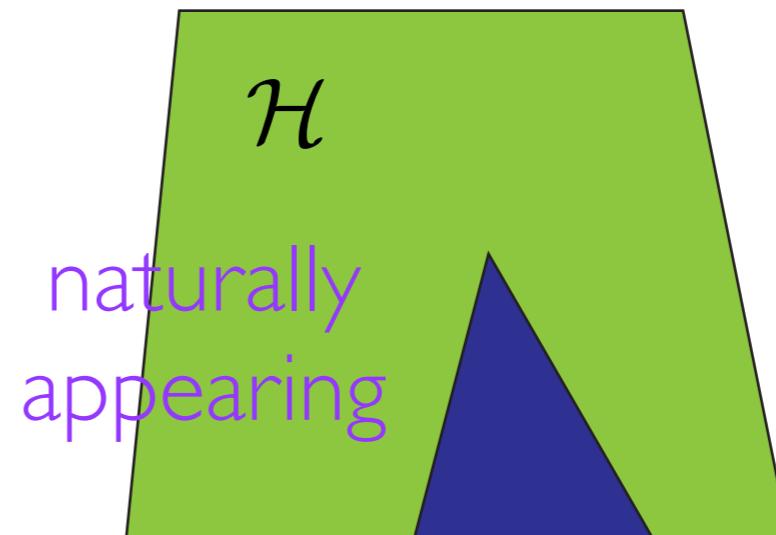
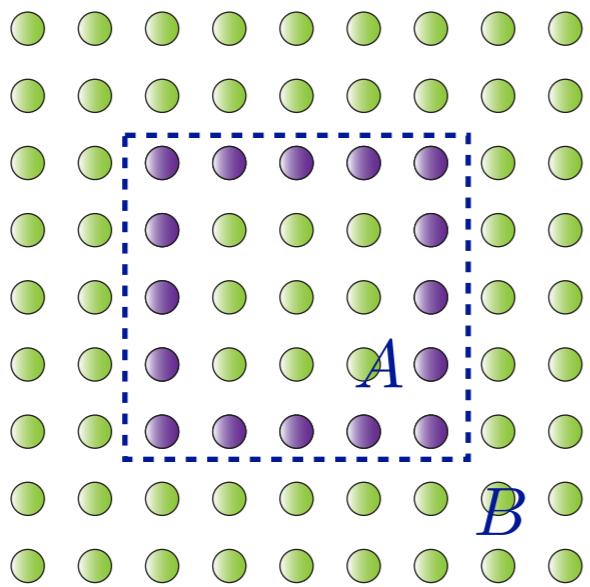


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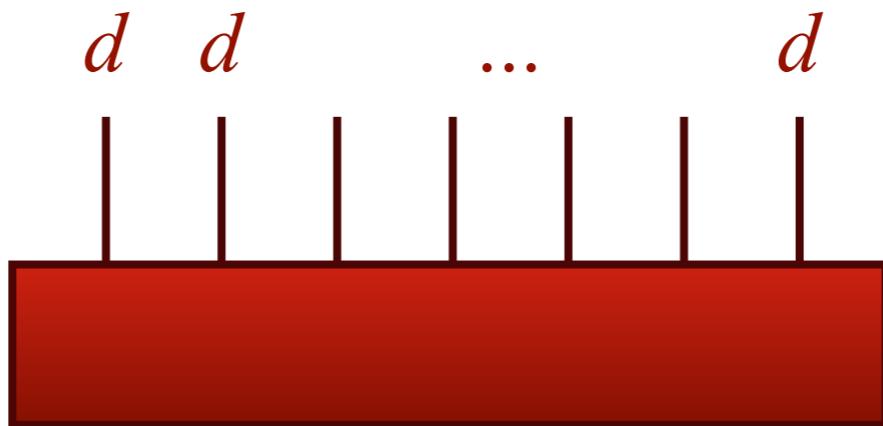
area law



We look for states with
little entanglement

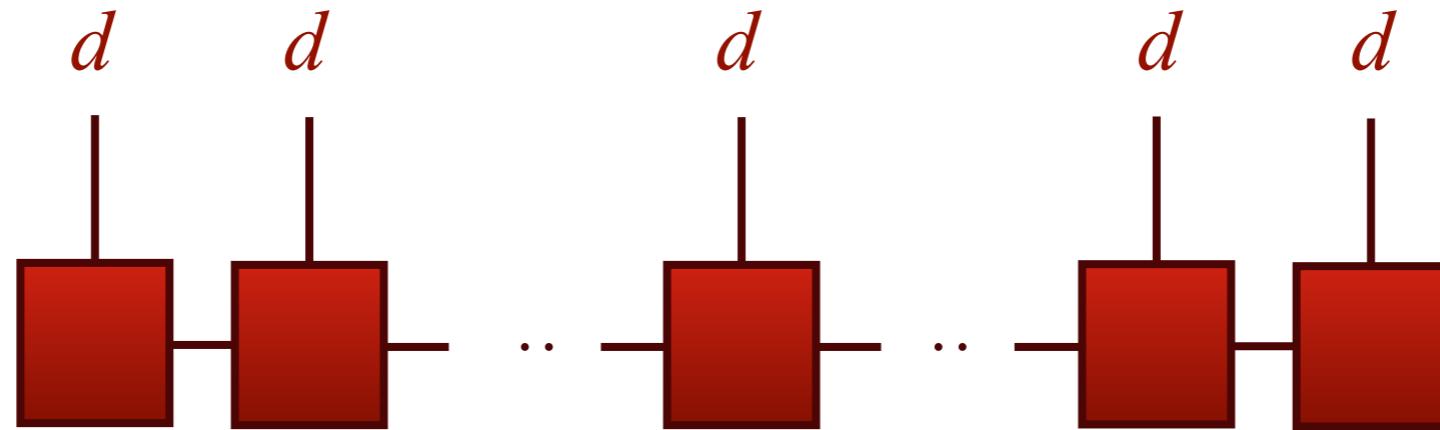
TNS = entanglement based ansatz

MPS

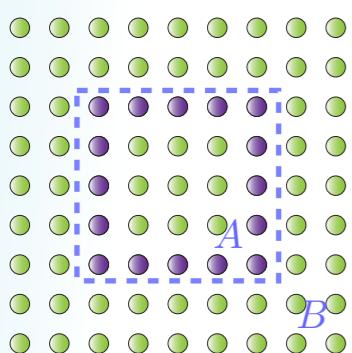


$$|\Phi\rangle = \sum_{s_1, \dots, s_N=1}^d c_{s_1, \dots, s_N} |s_1, \dots, s_N\rangle$$

MPS



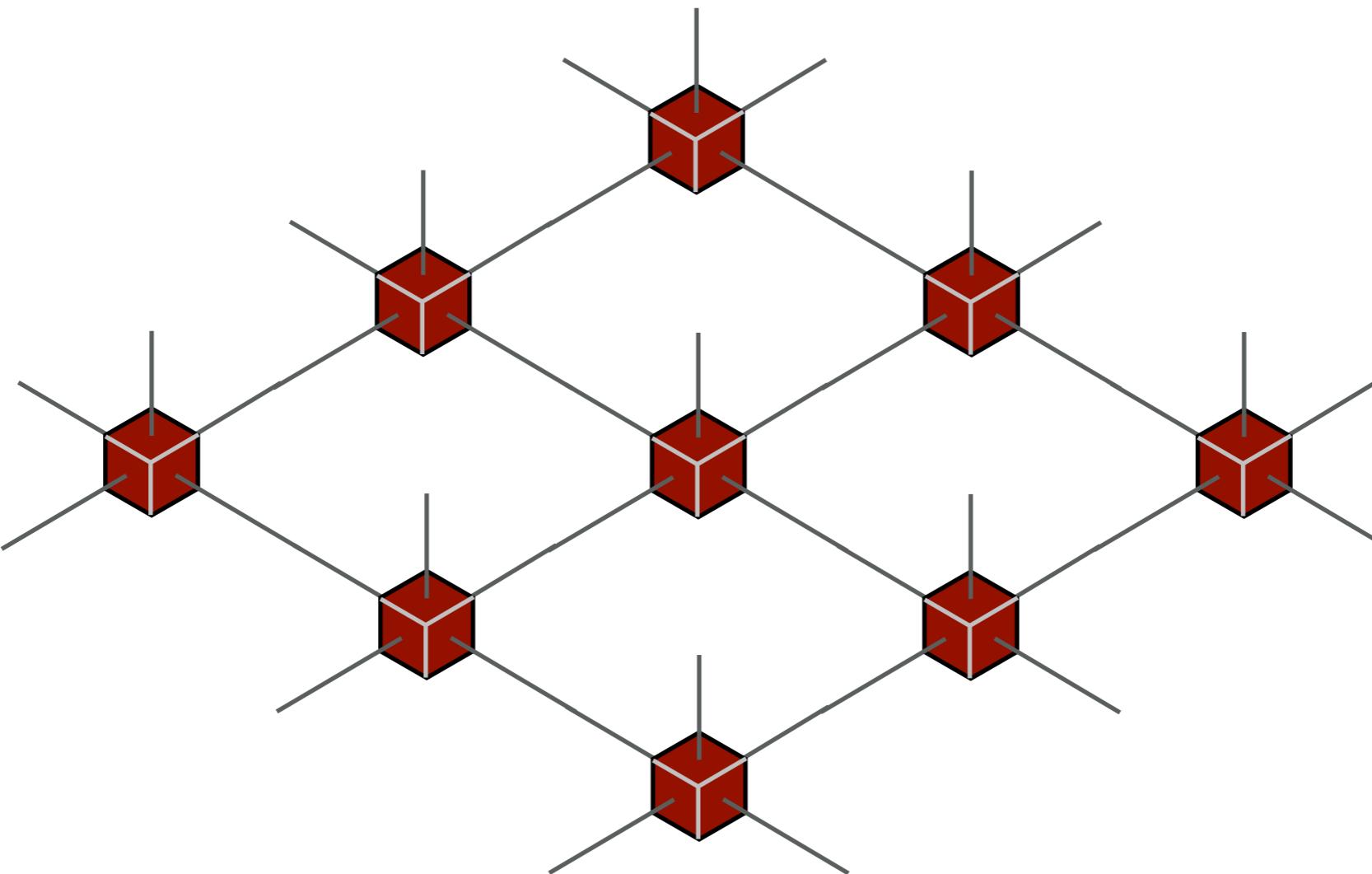
$$|\Phi_D\rangle = \sum_{s_1, \dots, s_N=1}^d \text{tr} (A^{s_1}[1] \dots A^{s_1}[N]) |s_1, \dots, s_N\rangle$$



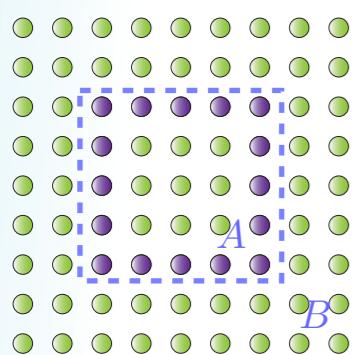
area law by construction

$$S(L/2) \leq \log D$$

PEPS

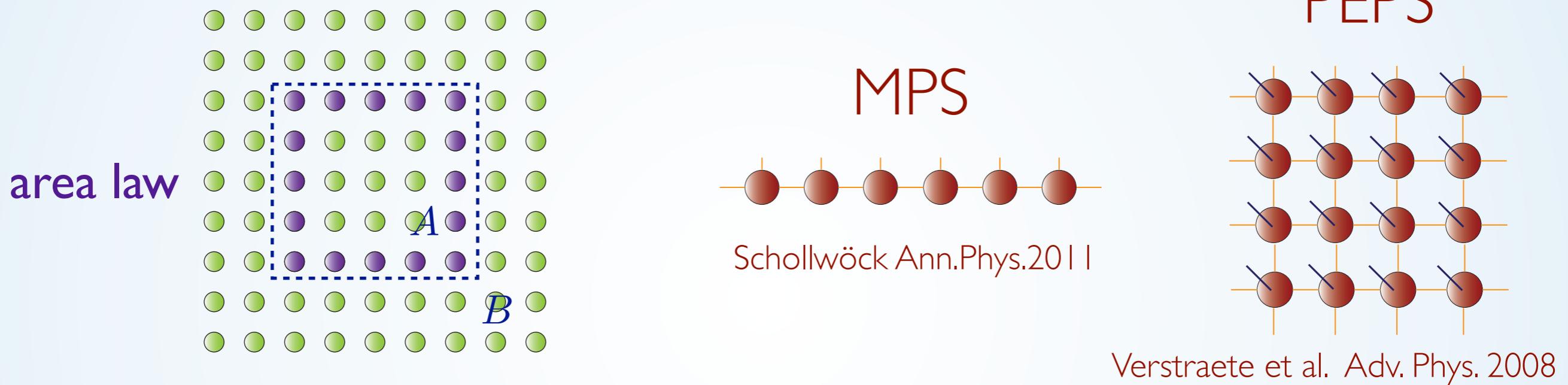


area law by construction



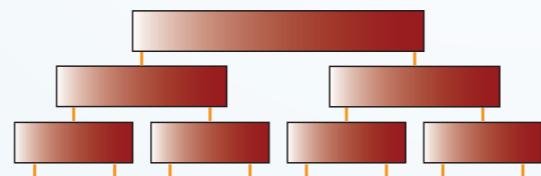
Verstraete, Cirac, 2004

TNS = entanglement based ansatz



other TNS

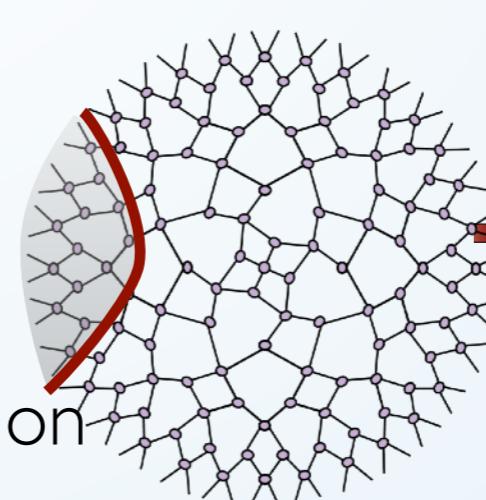
TTN



Shi et al PRA 2006

suggested connection
to AdS/CFT

Vidal PRL 2007 MERA



Swingle PRD 2012
Molina JHEP 2013
Nozaki et al JHEP 2012
Bao et al PRD 2015

TNS are very useful in the quantum many-body context...

works for GS, low energy, thermal equilibrium...

Verstraete, Cirac, PRB 2006 Hastings PRB 2006
Hastings J. Stat. Phys 2007 Molnar *et al.* PRB 2015

area laws

but not for high energy eigenstates, quenches...

Osborne, PRL 2006 Vidmar *et al.*, PRL 2017
Schuch *et al.*, NJP 2008

volume law

entanglement growth in non-equilibrium scenarios limits the applicability of MPS

fundamental questions: thermalization, ETH...

global quench in 1D

entanglement
barrier

$$D_{\min}(t) \sim e^{\alpha t}$$

Osborne, PRL 2006

Schuch et al., NJP 2008

$$S(t) \propto t$$

$t = 0$

product state



easy to write as MPS

local
observables

TNS challenge:
getting around this
limitation

some recent progress

Dubail JPhysA 2017
Leviatan et al. 2017
White et al PRB 2018
Surace et al. 2018
Kvornig et al 2021
Rakovzsky et al 2022



tools to get
dynamical
properties

$t = \infty$

thermal states



well approximated as MPO

alternative: give up description of the full state

① light-cone TN for
non-equilibrium
evolution of local
observables

**M. Frías-Pérez, MCB,
PRB 106, 115117 (2022)**

② transforming long-range
entanglement into mixture

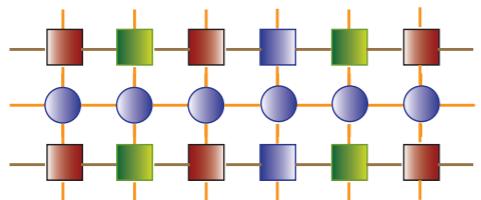
**M. Frías-Pérez, L. Tagliacozzo, MCB,
arXiv:2308.04291**

③ spectral properties of
the QMB Hamiltonian

Yang et al. PRL 124, 100602 (2020),
Lu, MCB, Cirac, PRX Quantum 2, 02032 (2021)
Yang, Cirac, MCB, PRB 106, 024307 (2022)

give up description of the full state: (local) operators

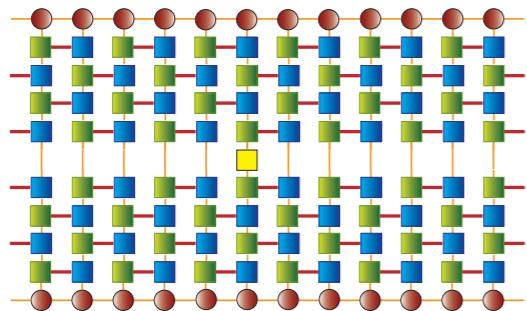
evolving operators: Heisenberg picture Hartmann et al, PRL 2009



also for mixed states
operator space entanglement

Prosen Pizorn, PRL 2008

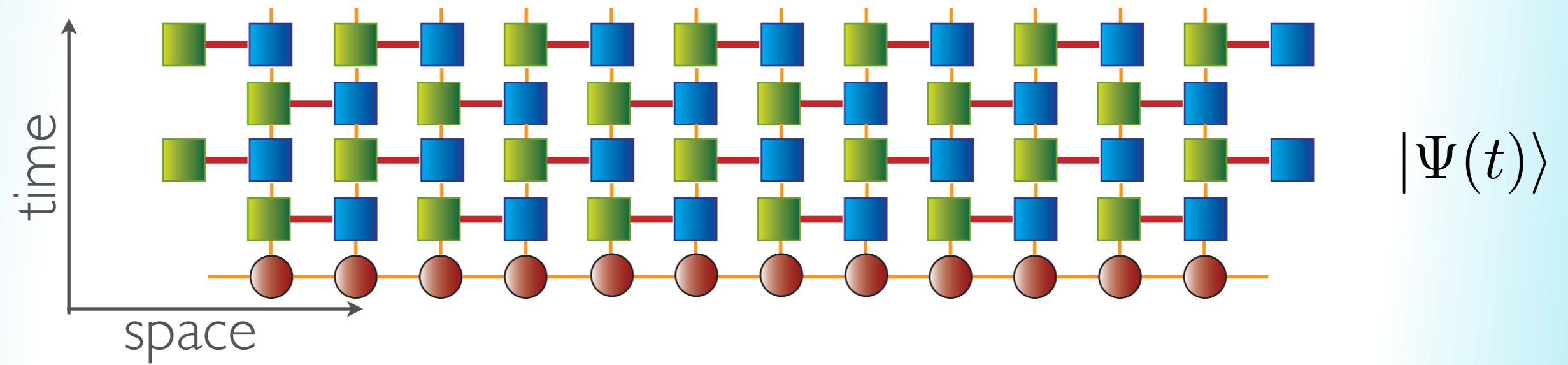
observables as TN to contract



different entanglement quantities

MCB, Hastings, Verstraete, Cirac, PRL 2009
Müller-Hermes et al., NJP 2012
Hastings, Mahajan 2014
Frías-Pérez, MCB PRB 2022

time-dependent observable as a TN

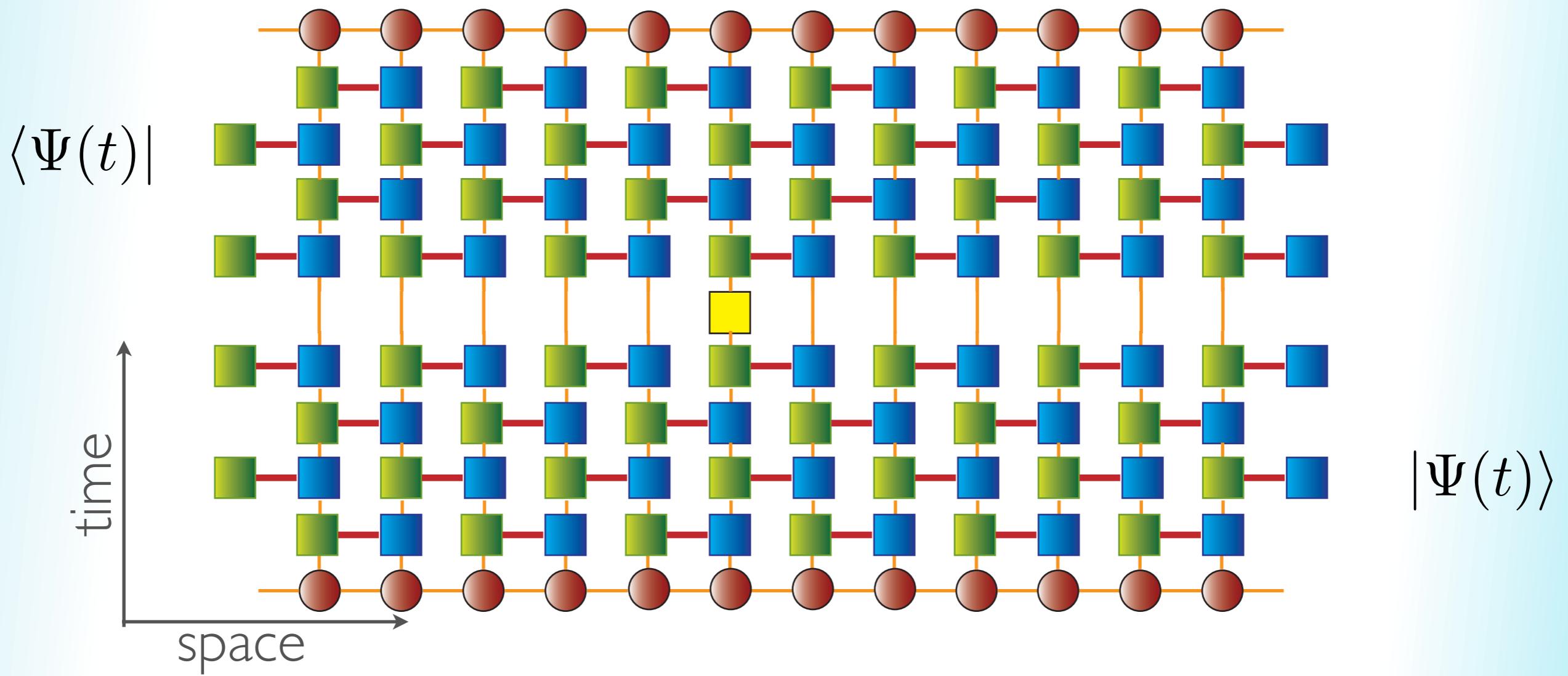


time-dependent observable as a TN

TN describe
observables, not
states

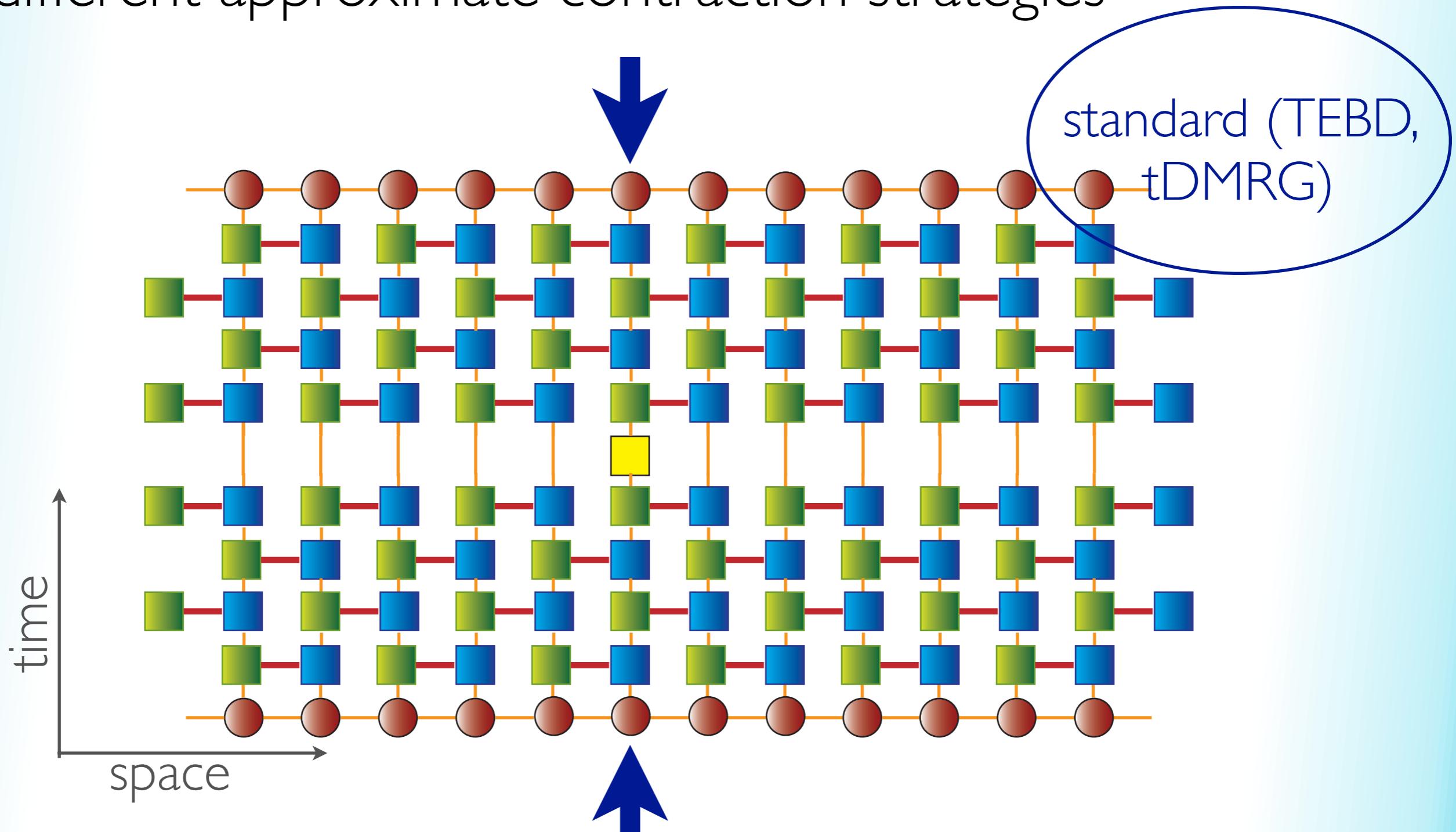
$$\langle \Psi(t) | O | \Psi(t) \rangle$$

exact contraction
not possible
 $\#P$ complete



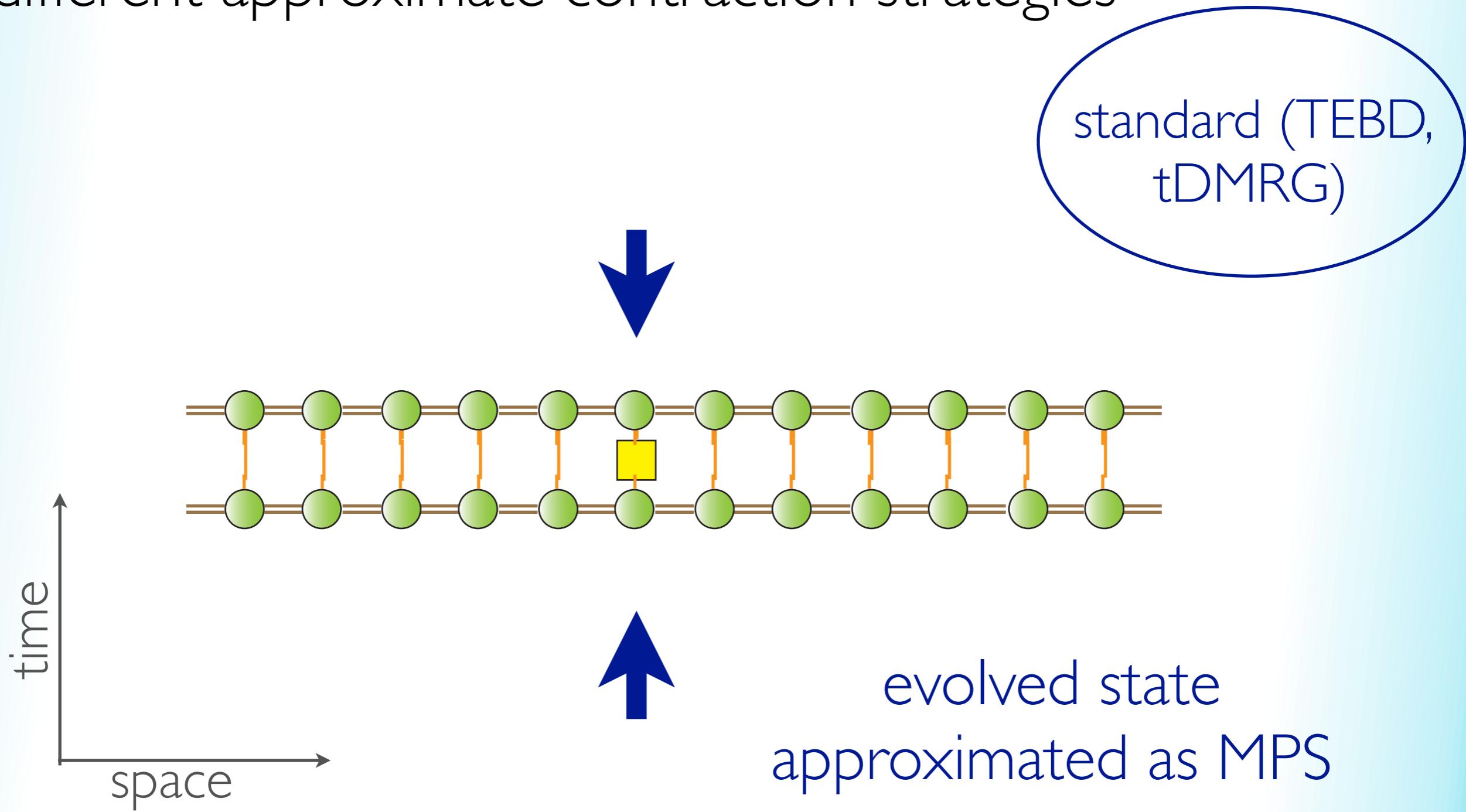
time-dependent observable as a TN

different approximate contraction strategies



time-dependent observable as a TN

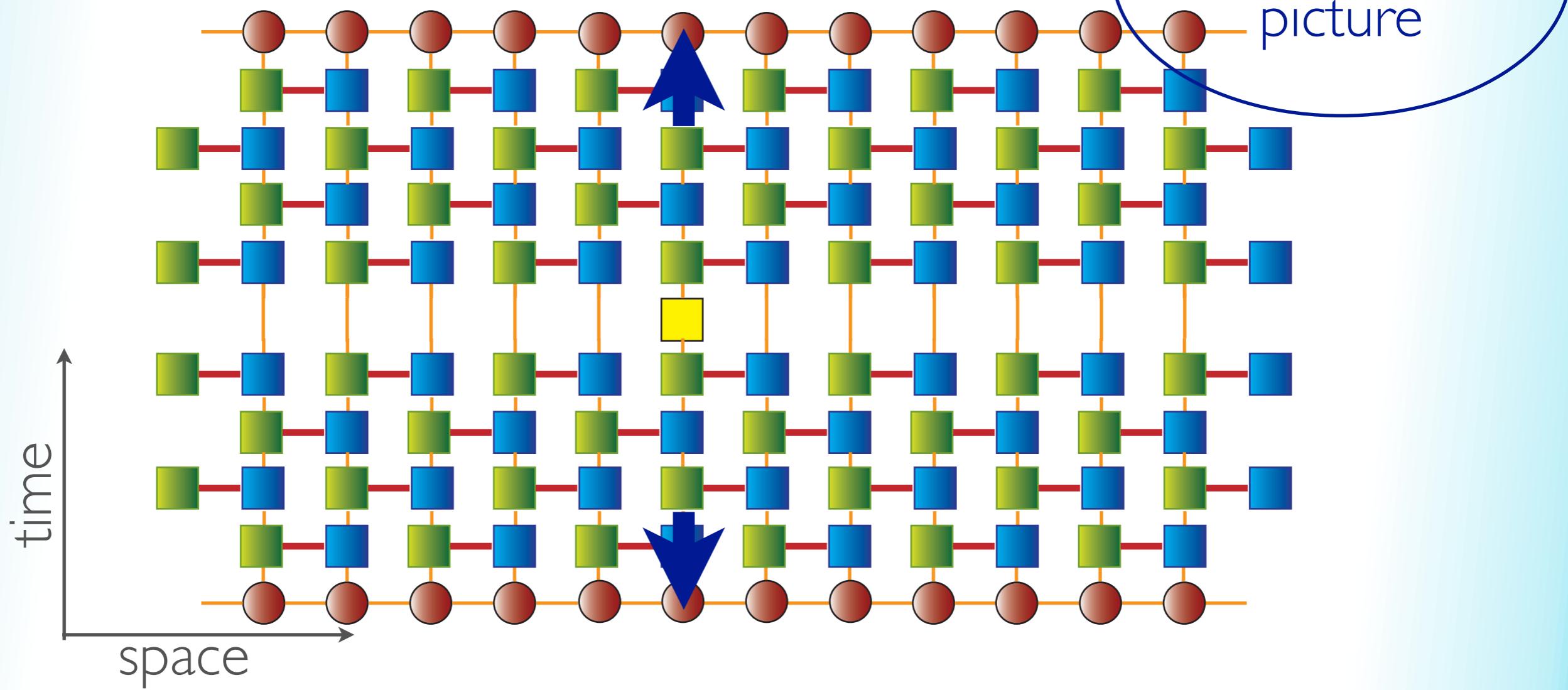
different approximate contraction strategies



time-dependent observable as a TN

different approximate contraction strategies
evolved operator as

MPO



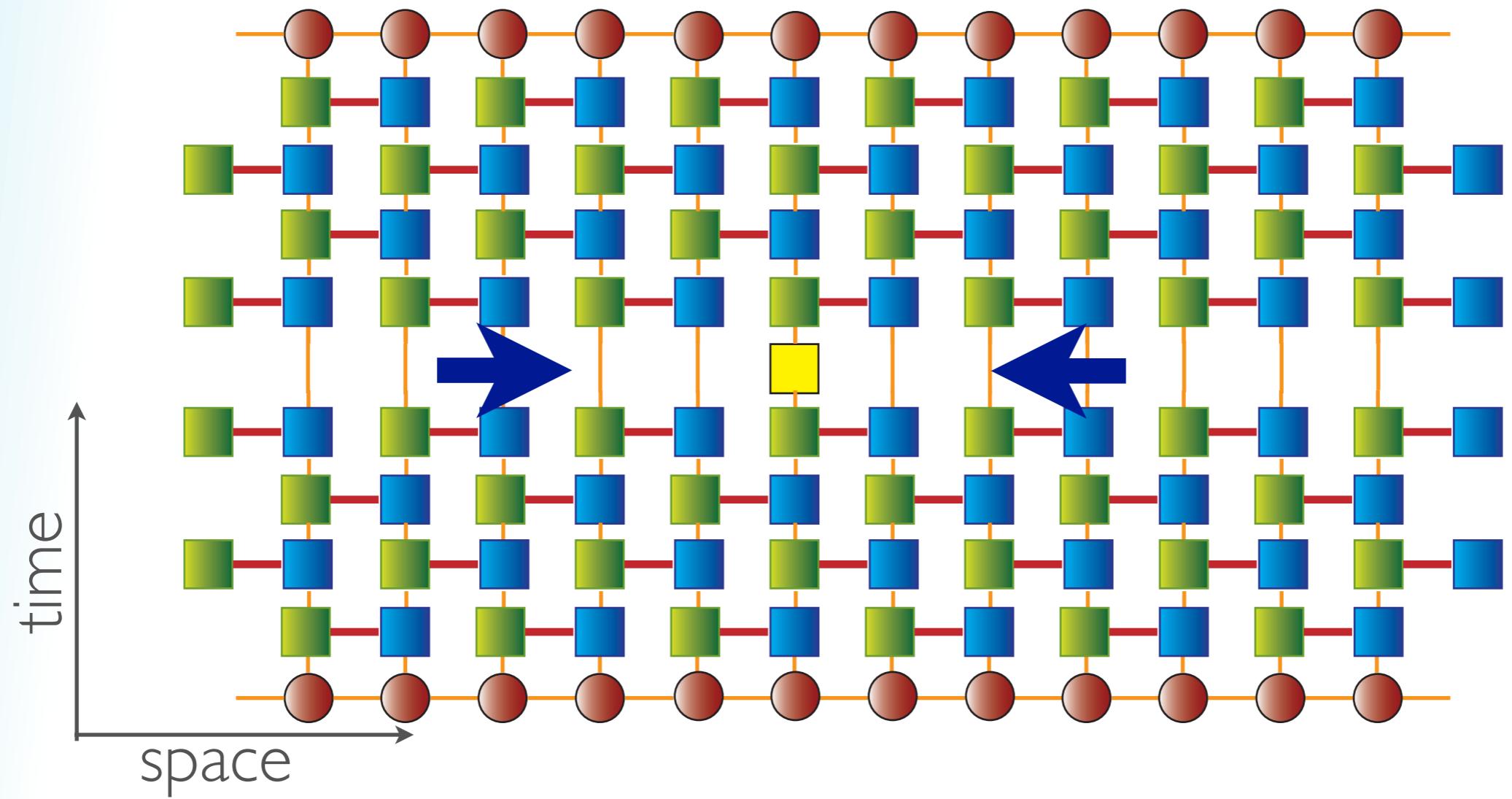
time-dependent observable as a TN

for infinite systems, transverse folding approach

MCB, Hastings, Verstraete, Cirac, PRL 2009

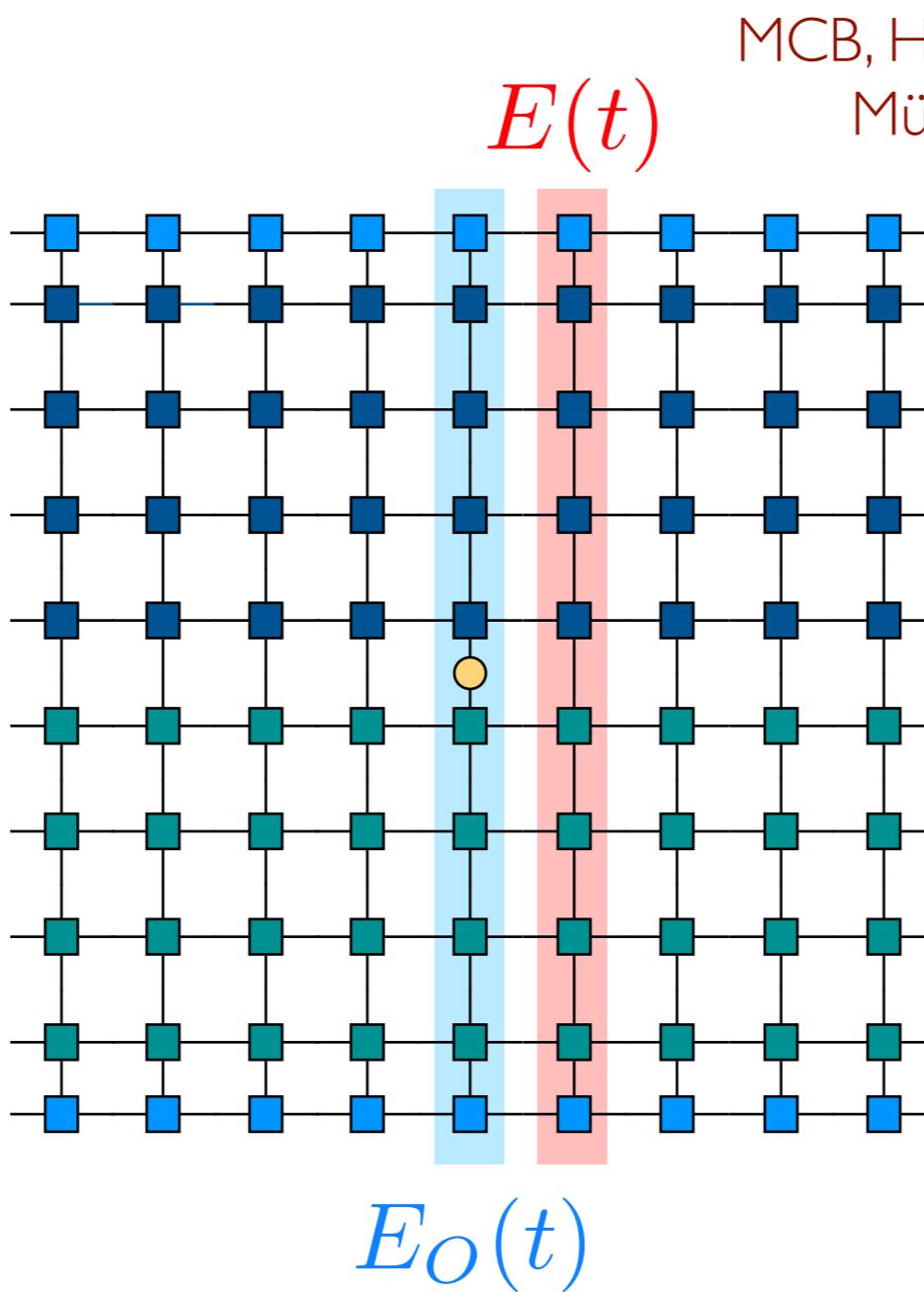
Müller-Hermes, Cirac, MCB, NJP 2012

Hastings, Mahajan 2014



transverse folding approach

for infinite systems, transverse folding approach



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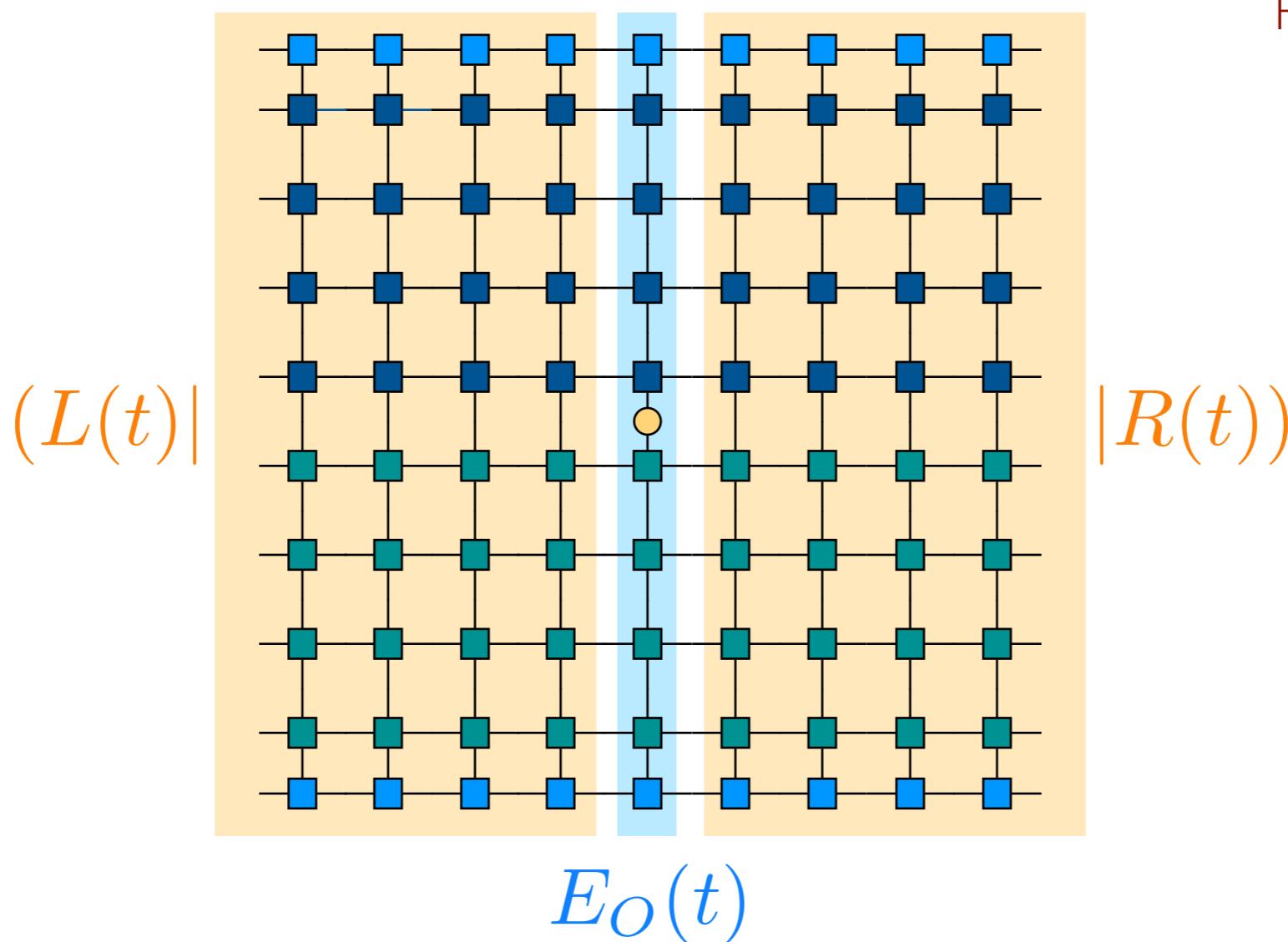
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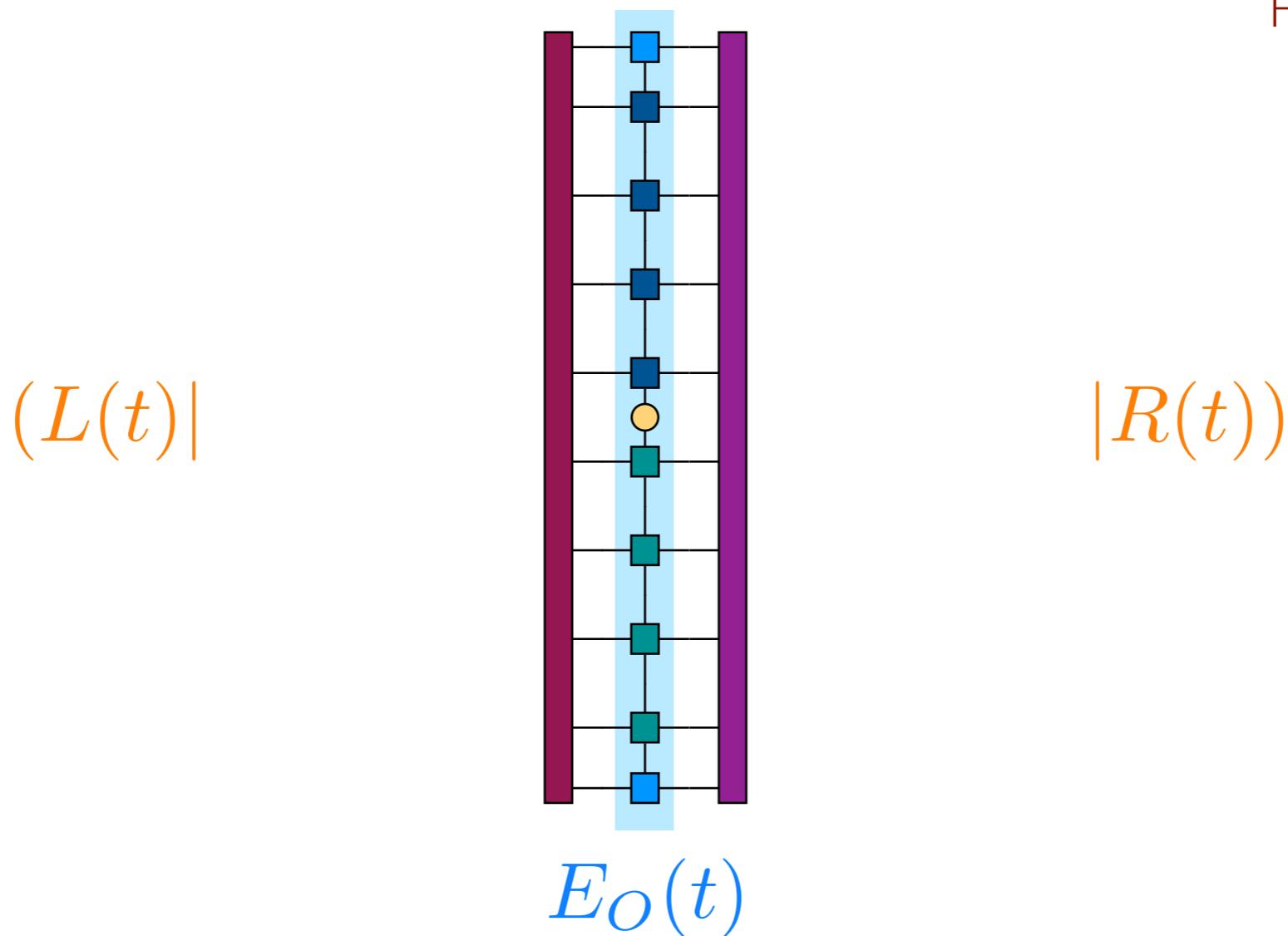
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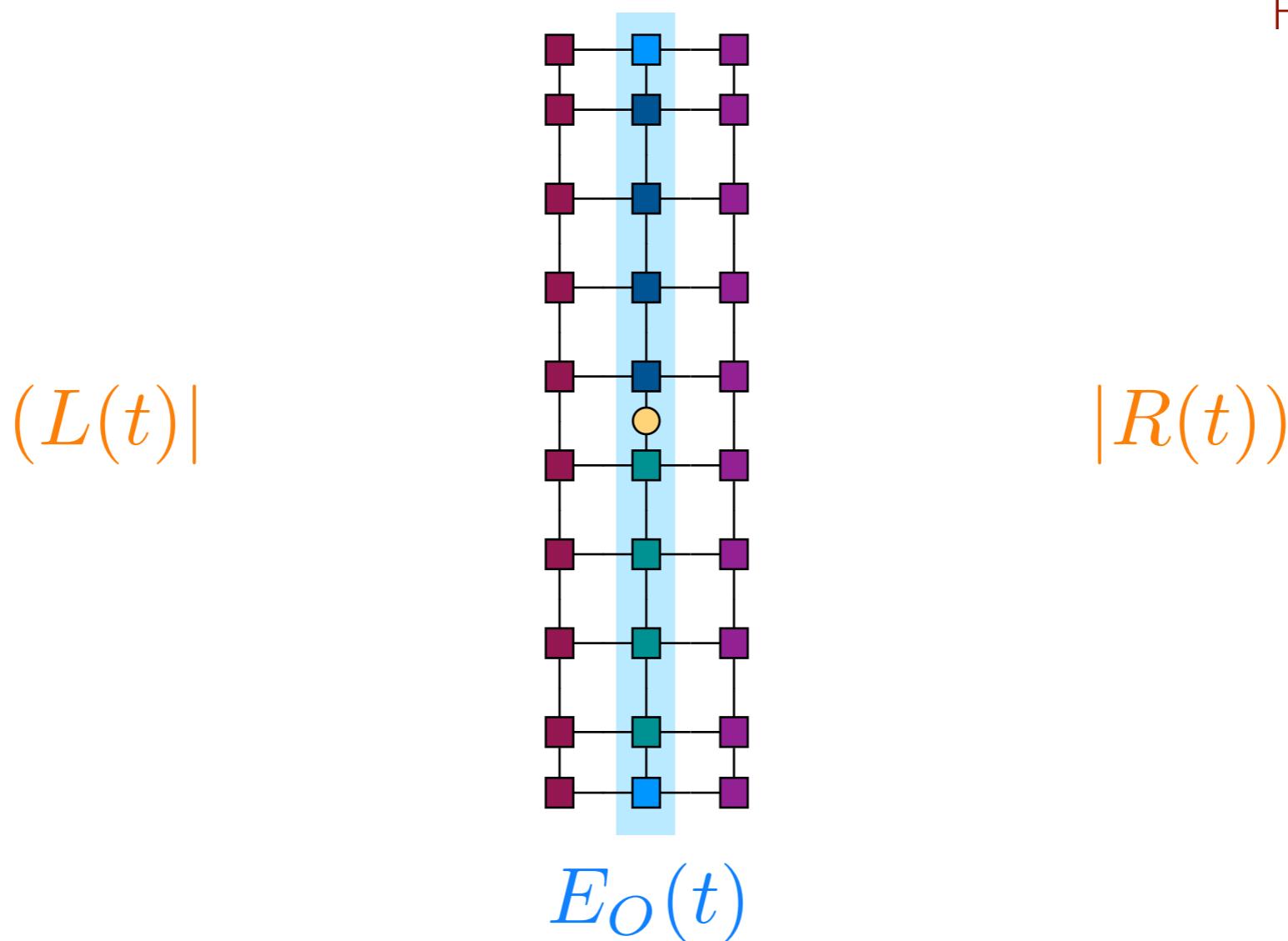
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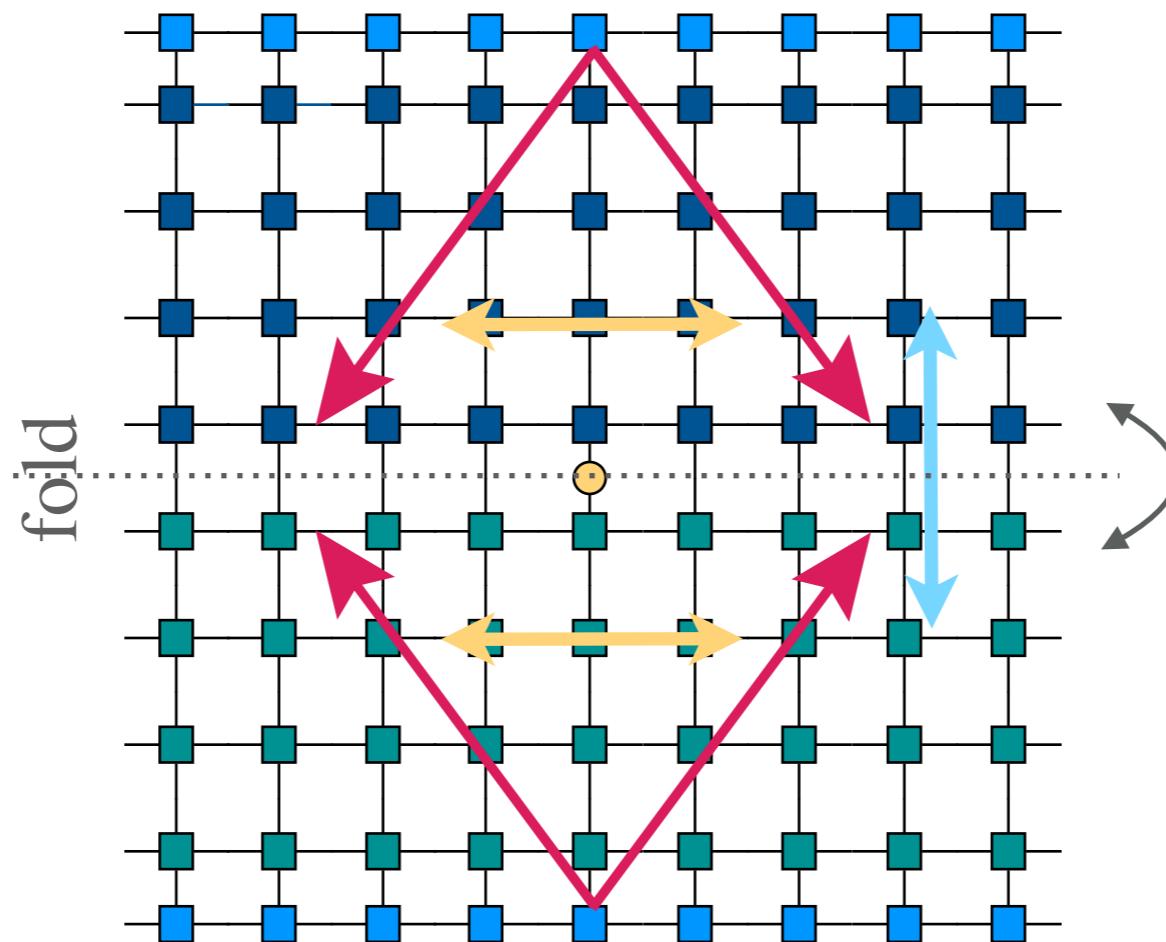
transverse folding approach

intuition: model free propagating excitations

MCB, Hastings, Verstraete, Cirac, PRL 2009

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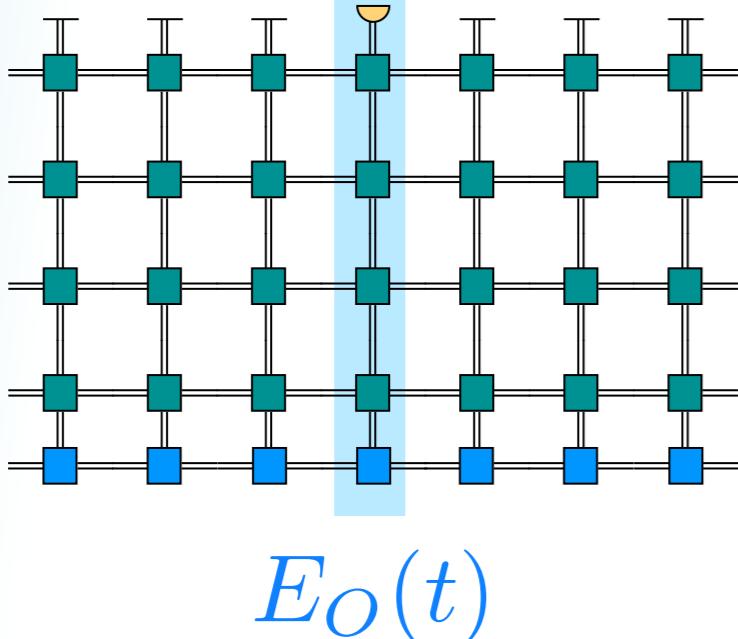
transverse folding approach

free propagating excitations

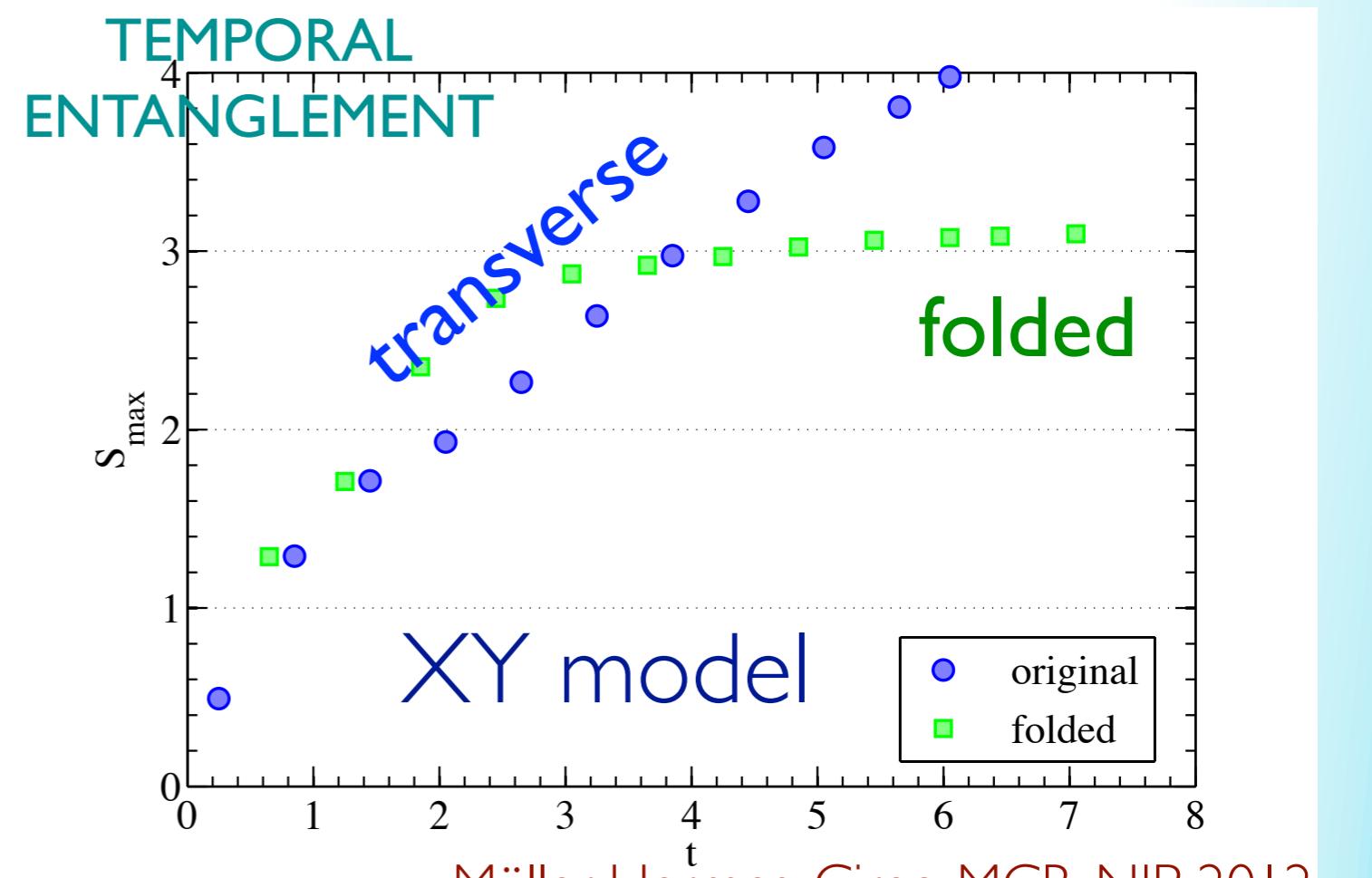
recently: influence functional
Sonner et al, Ann. Phys 2021

Lerose et al. PRX 2021
Ye, Chan, J. Chem. Phys. 2021

see also Carignano 2307.11649



closest real case: global quench
in free fermionic models



Müller-Hermes, Cirac, MCB, NJP 2012
see also Giudice et al., PRL 128, 220401 (2022)

transverse folding + light cone = TLCC

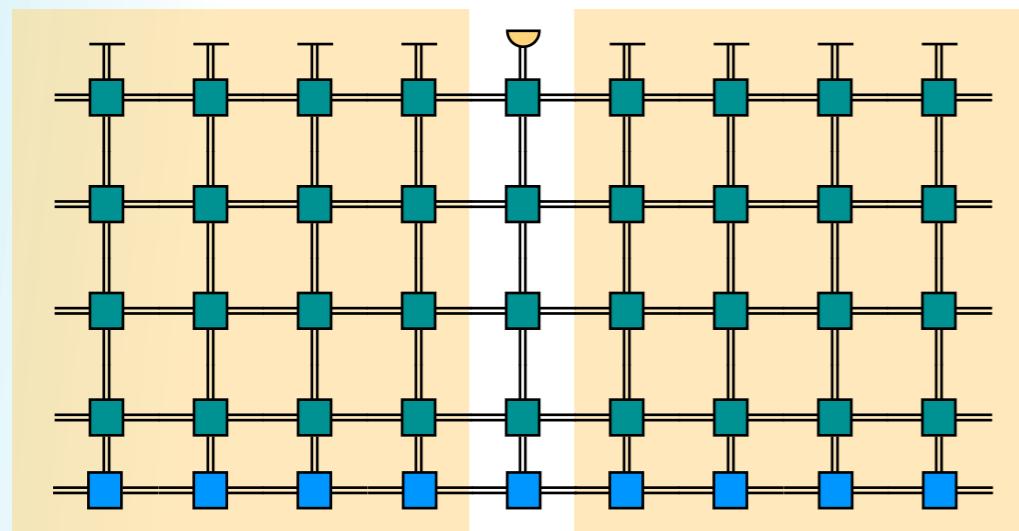
cancelling local unitaries

gain in efficiency

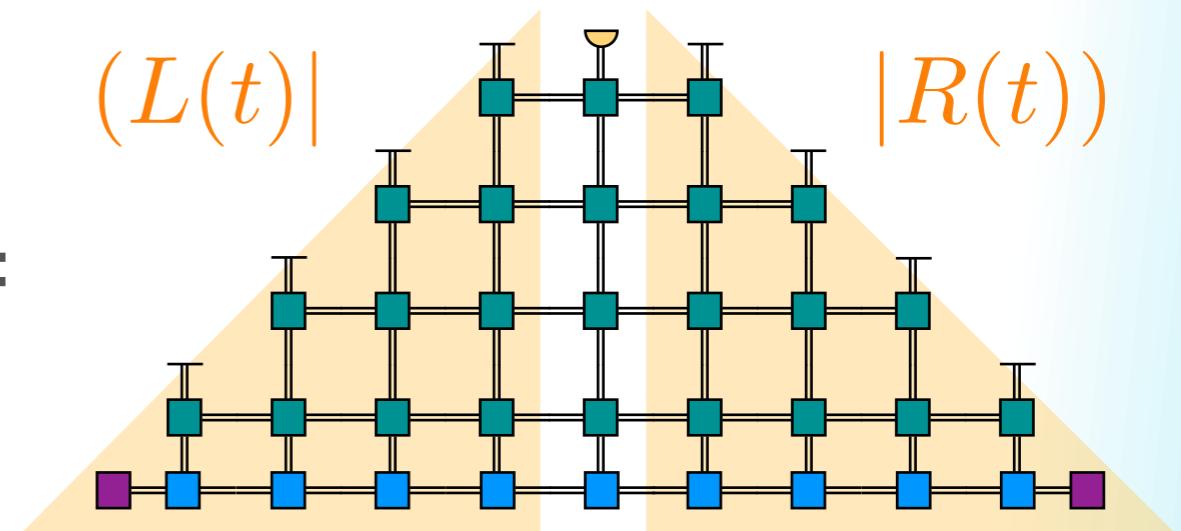
systematic increment of t

improved convergence with
Hastings' truncation

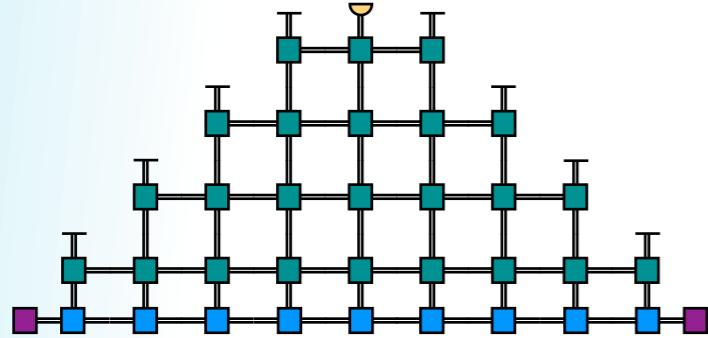
Hastings, Mahajan 2014



=



transverse folding + light cone = TLCC



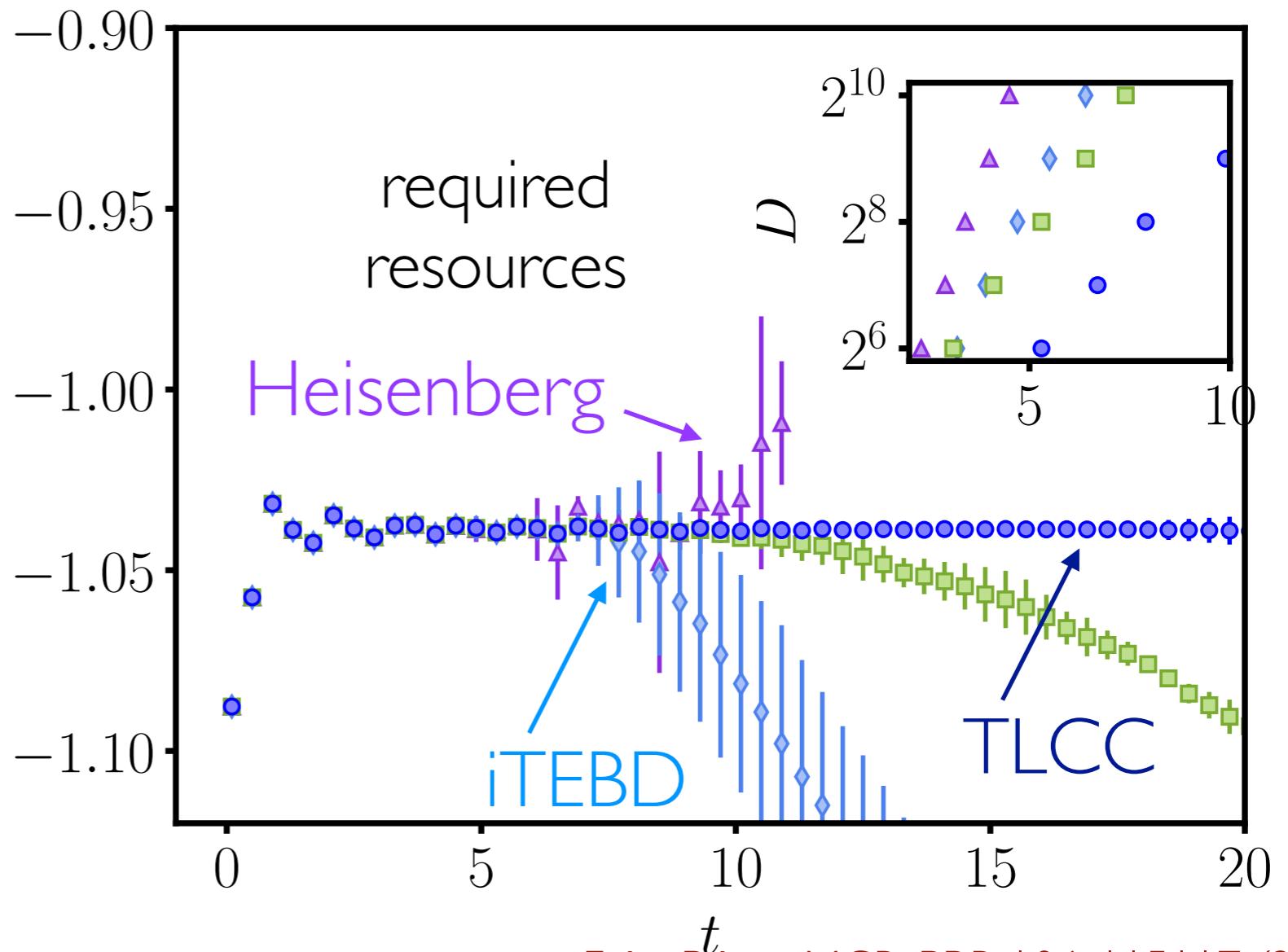
$(J, g, h) =$

$(1, -1.05, 0.5)$

$\langle H \rangle$

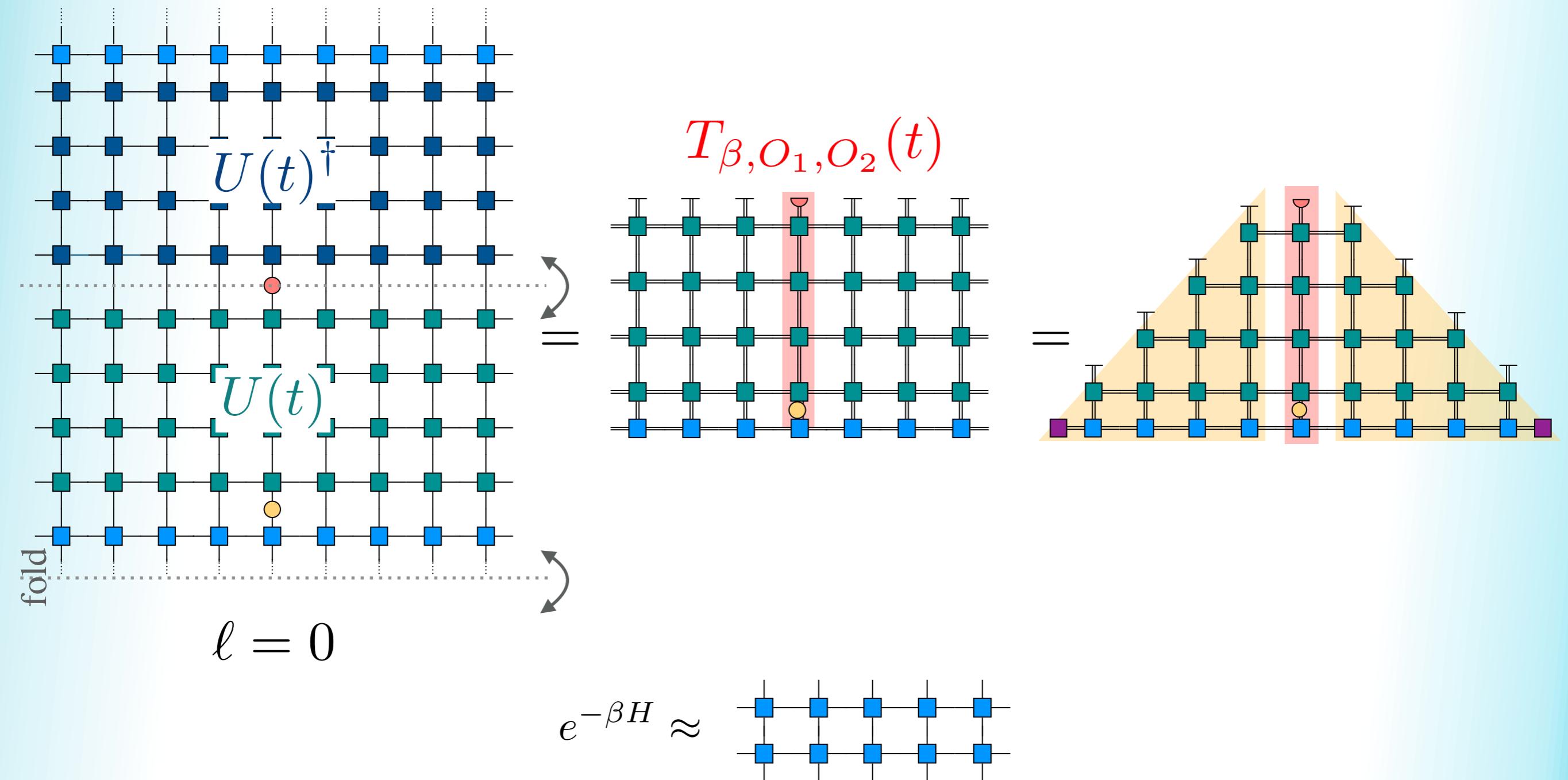
quench from $|X+\rangle$

$$H_{\text{Ising}} = J \sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i \sigma_x^{[i]} + h \sum_i \sigma_z^{[i]}$$



computing response functions

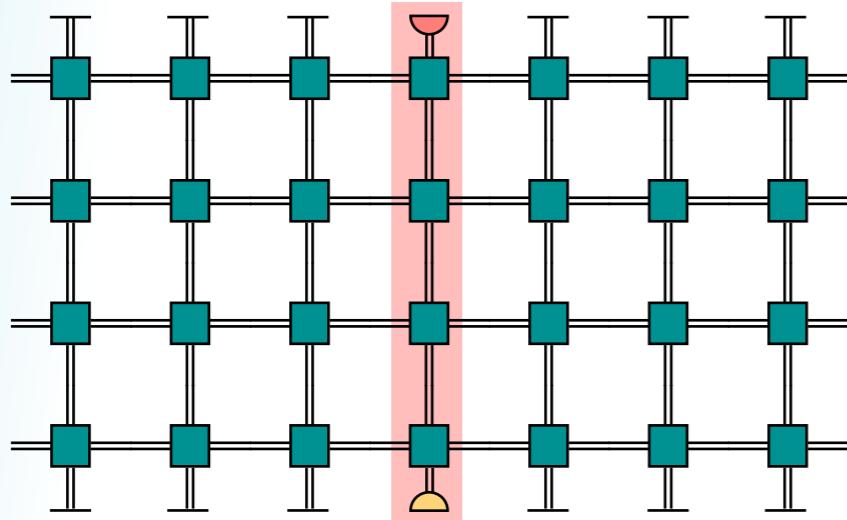
$$C_{1,2}(t, \ell, \beta) = \text{tr}(\rho_\beta O_2^{[\ell]}(t) O_1^{[0]}(0))$$



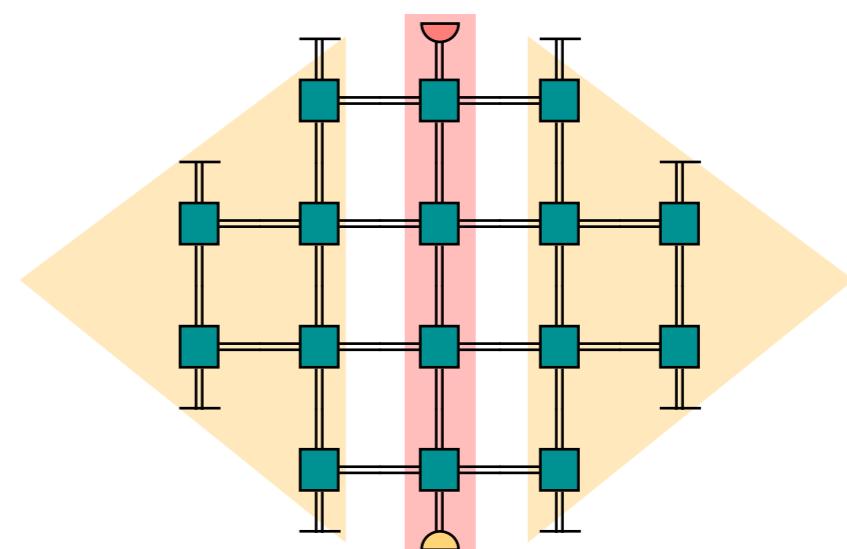
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$$C_{1,2}(t, \ell, \beta) = \text{tr}(\rho_\beta O_2^{[\ell]}(t) O_1^{[0]}(0))$$

$\ell = 0$



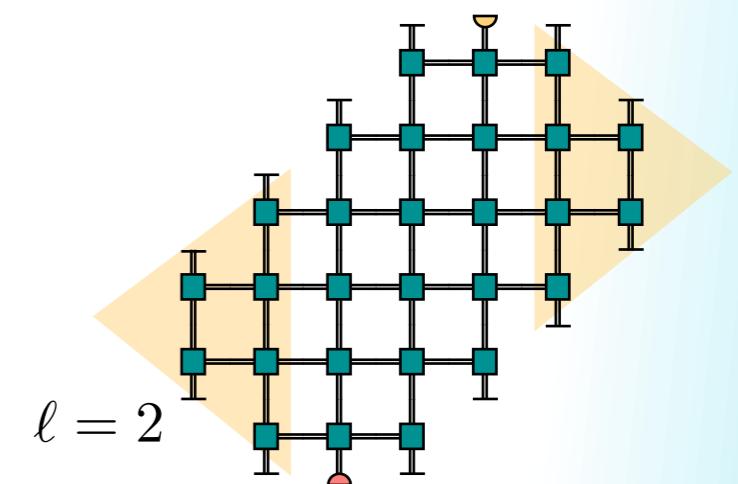
=



$T_{\beta=0, O_1, O_2}(t)$

infinite temperature

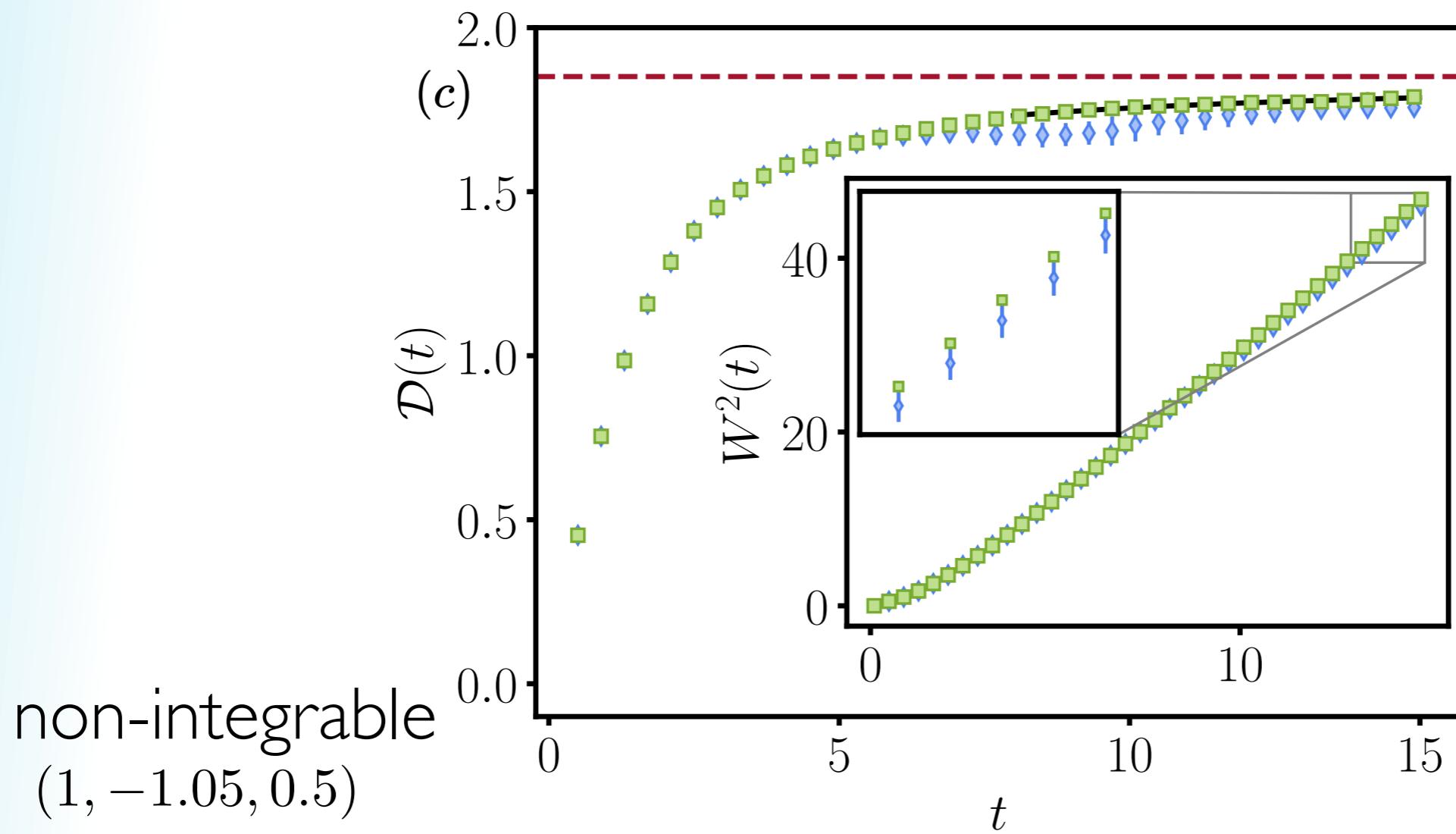
$\beta = 0$



computing response functions

tilted Ising
energy density
infinite temperature

$$\begin{aligned} O_E^{[i]} := & J\sigma_i^z\sigma_{i+1}^z + \frac{g}{2}(\sigma_i^x + \sigma_{i+1}^x) \\ & + \frac{h}{2}(\sigma_i^z + \sigma_{i+1}^z) \end{aligned}$$



diffusion
constant

alternative: give up description of the full state

① light-cone TN for
non-equilibrium
evolution of local
observables

**M. Frías-Pérez, MCB,
PRB 106, 115117 (2022)**

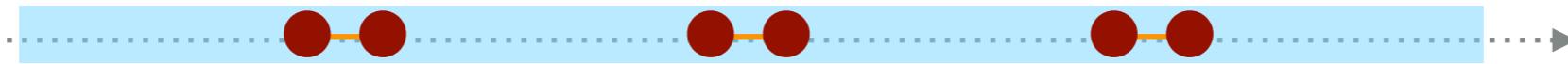
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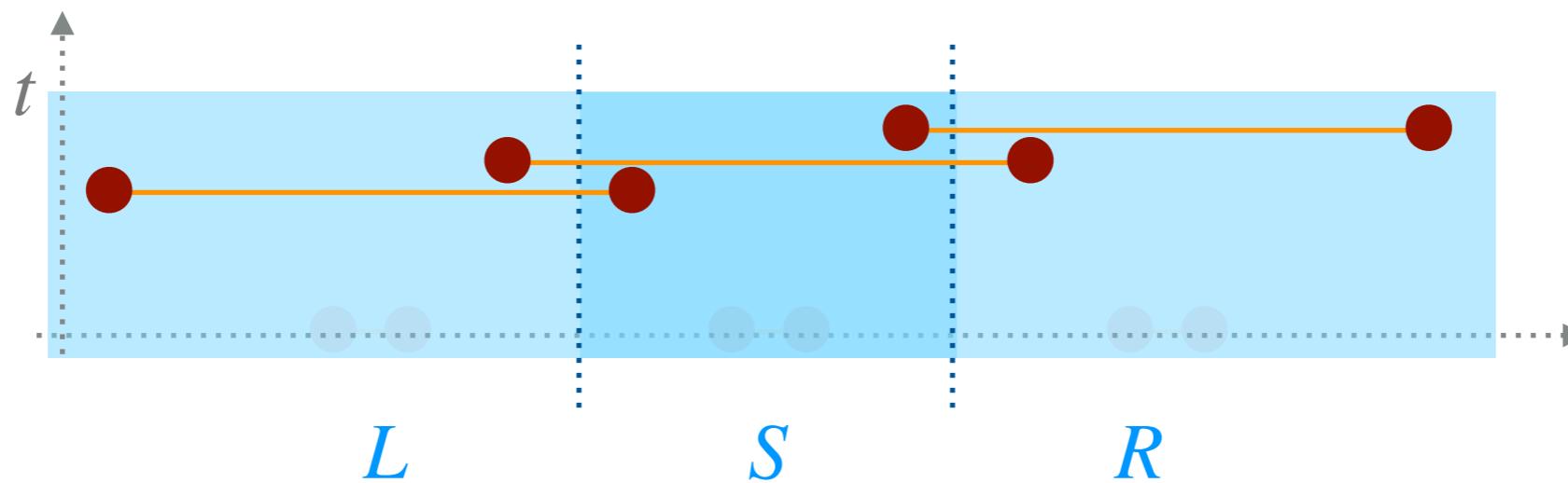
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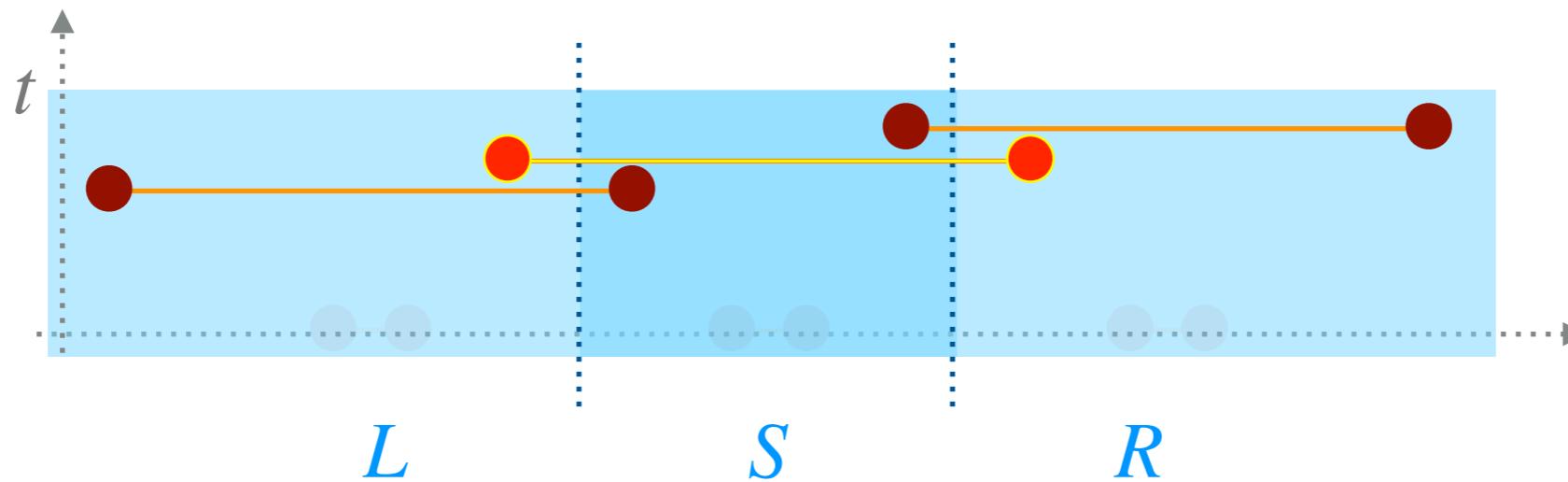
back to quasiparticle intuition



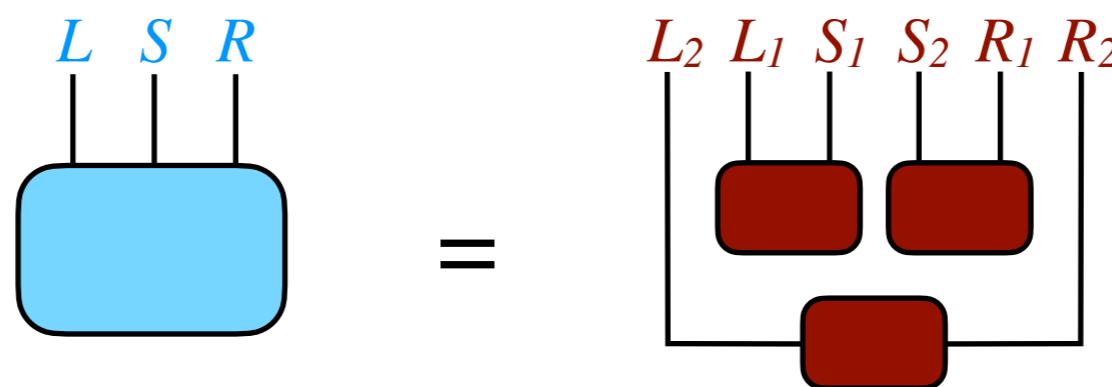
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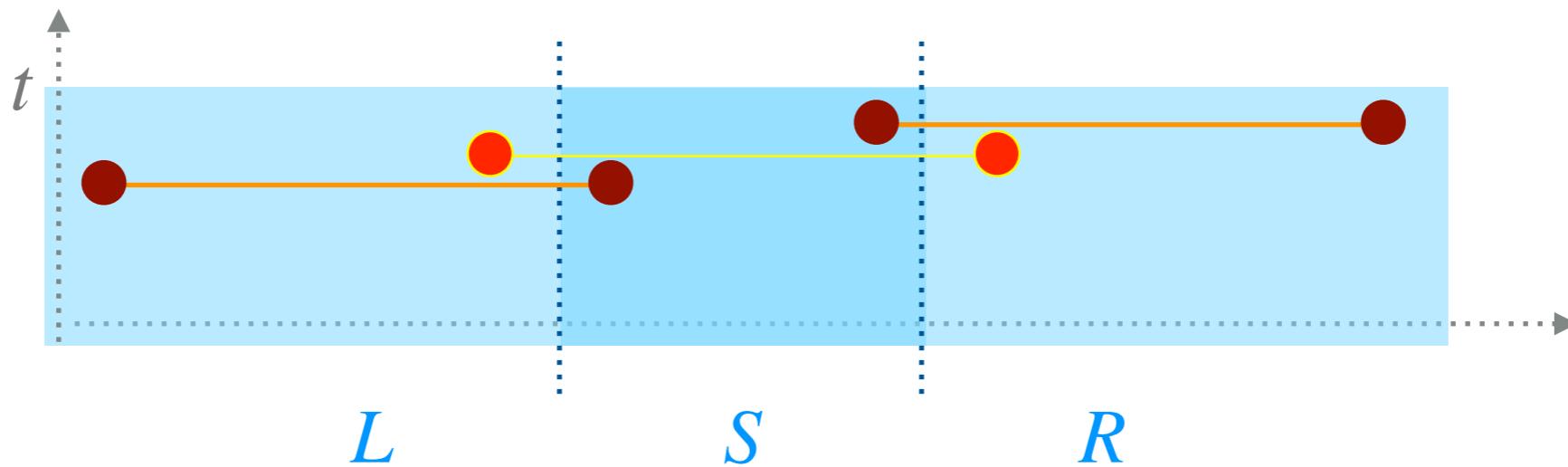
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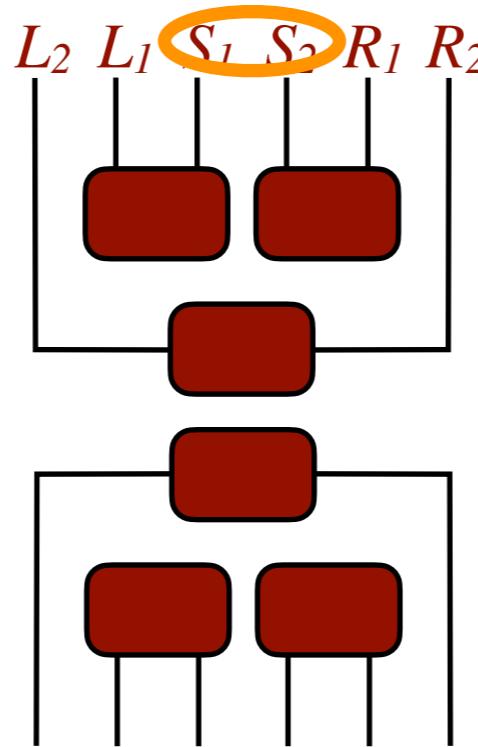
$$|\Psi_{LSR}\rangle = |\phi_{L_1S_1}\rangle \otimes |\phi_{S_2R_1}\rangle \otimes |\phi_{L_2R_2}\rangle$$



back to quasiparticle intuition

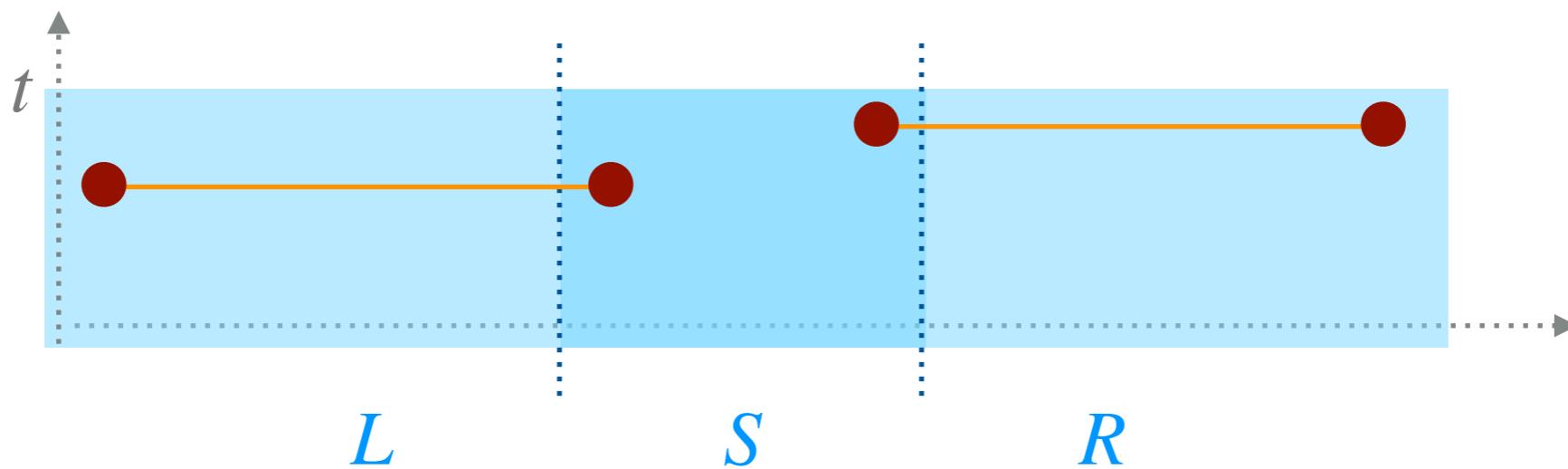


local observables on
any S block

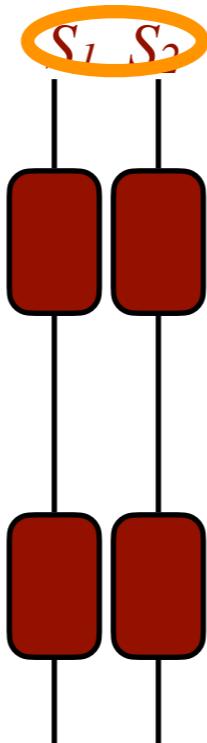


$|\Psi_{LSR}\rangle\langle\Psi_{LSR}|$

back to quasiparticle intuition

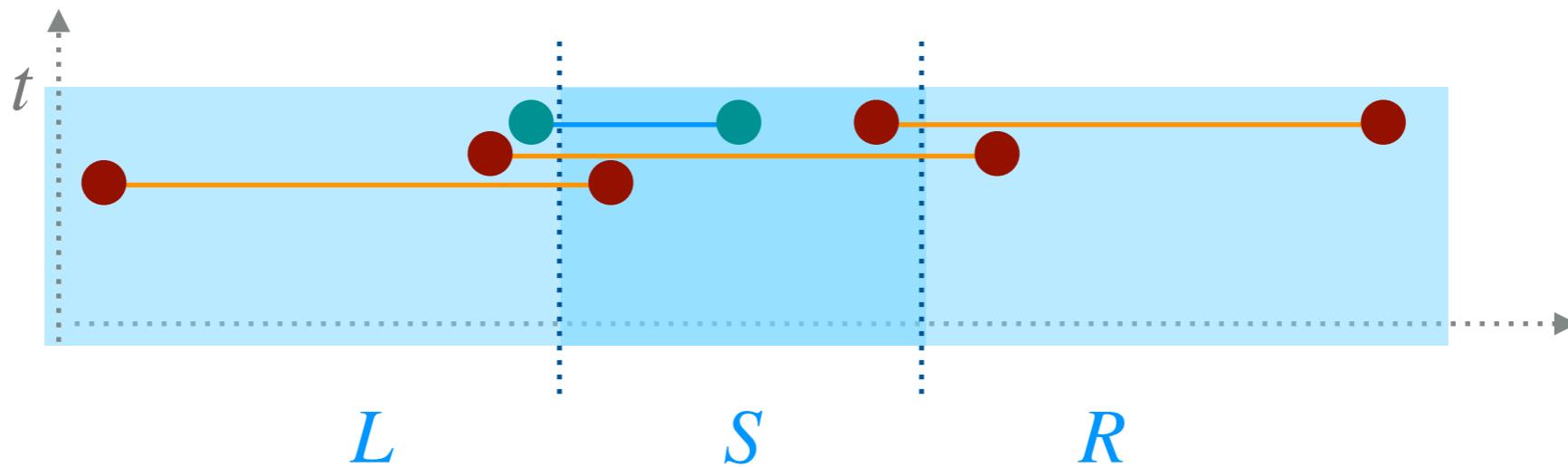


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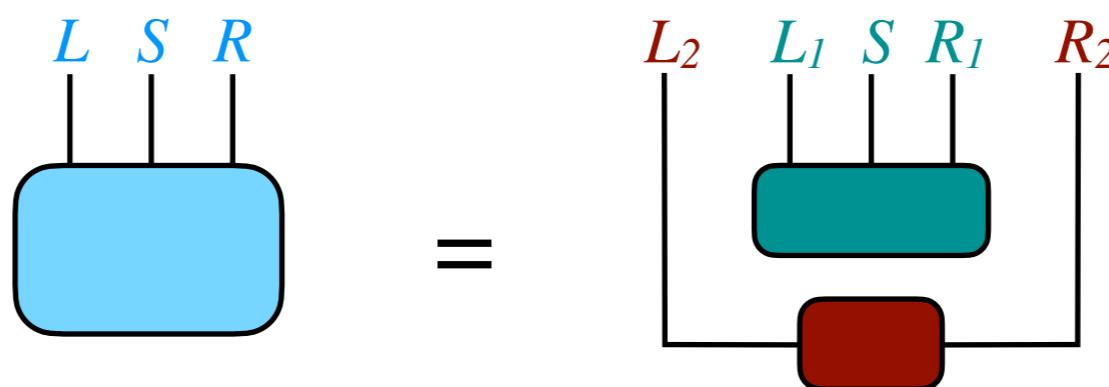


long-range
entanglement
contributes as
statistical mixture

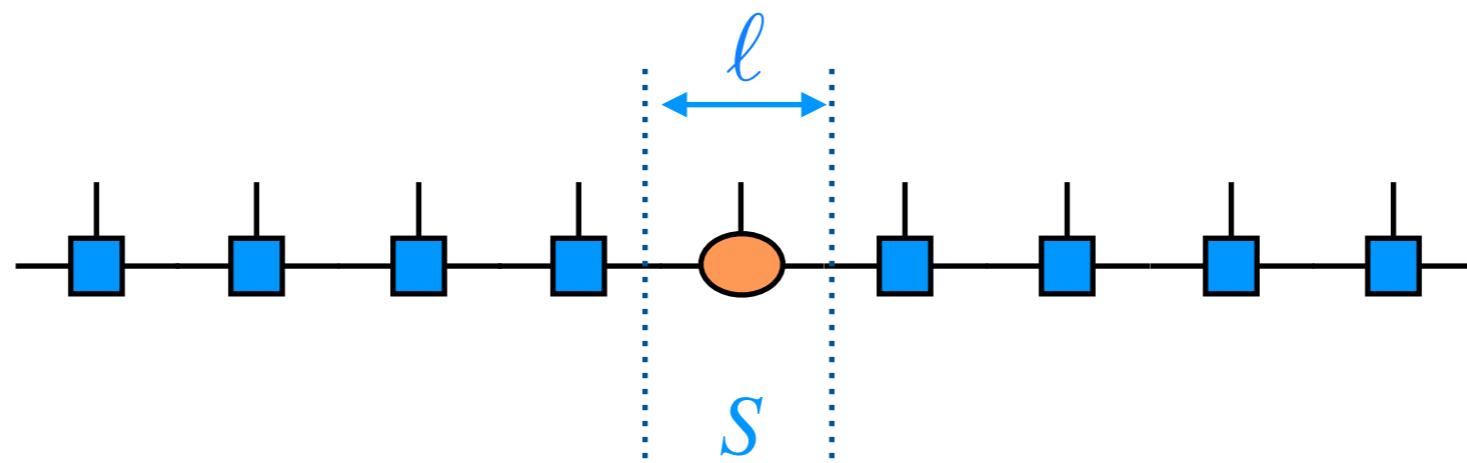
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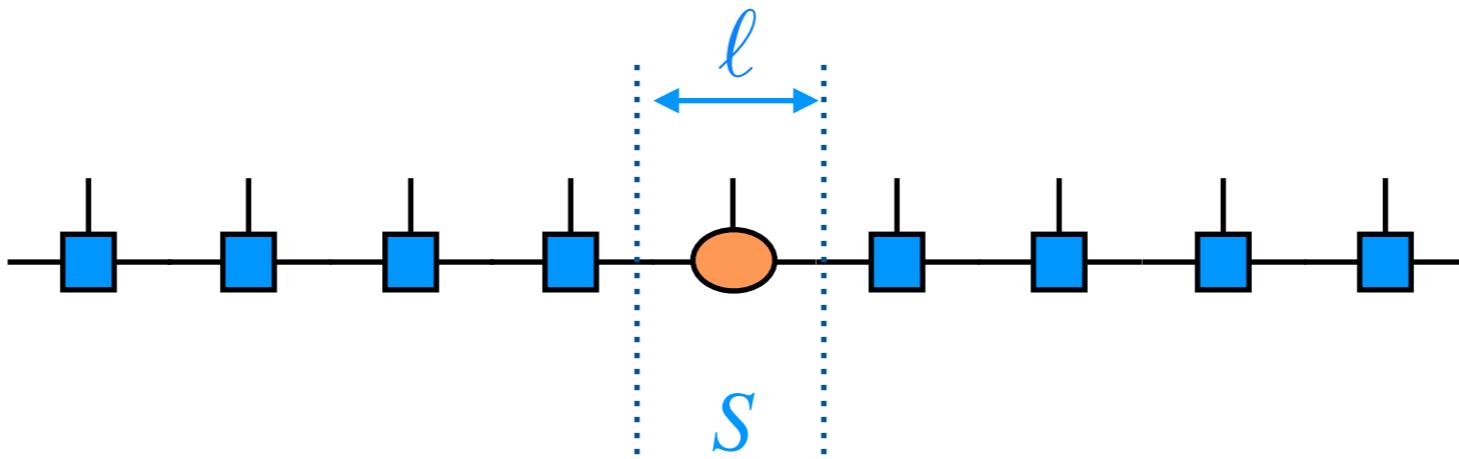
$$|\Psi_{LSR}\rangle = |\phi_{L_1SR_1}\rangle \otimes |\phi_{L_2R_2}\rangle$$



TN picture



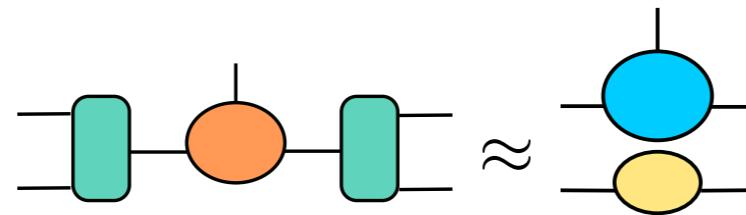
TN picture



$$\min \left\| \text{---} \circlearrowleft \text{---} - \text{---} \text{---} \text{---} \right\|^2$$

The right-hand side of the equation shows a tensor network component. It consists of three tensors represented by colored circles: an orange circle on the left, a blue circle in the middle, and a yellow circle on the right. They are connected sequentially by lines, forming a chain of three tensors.

long-range entanglement in TN

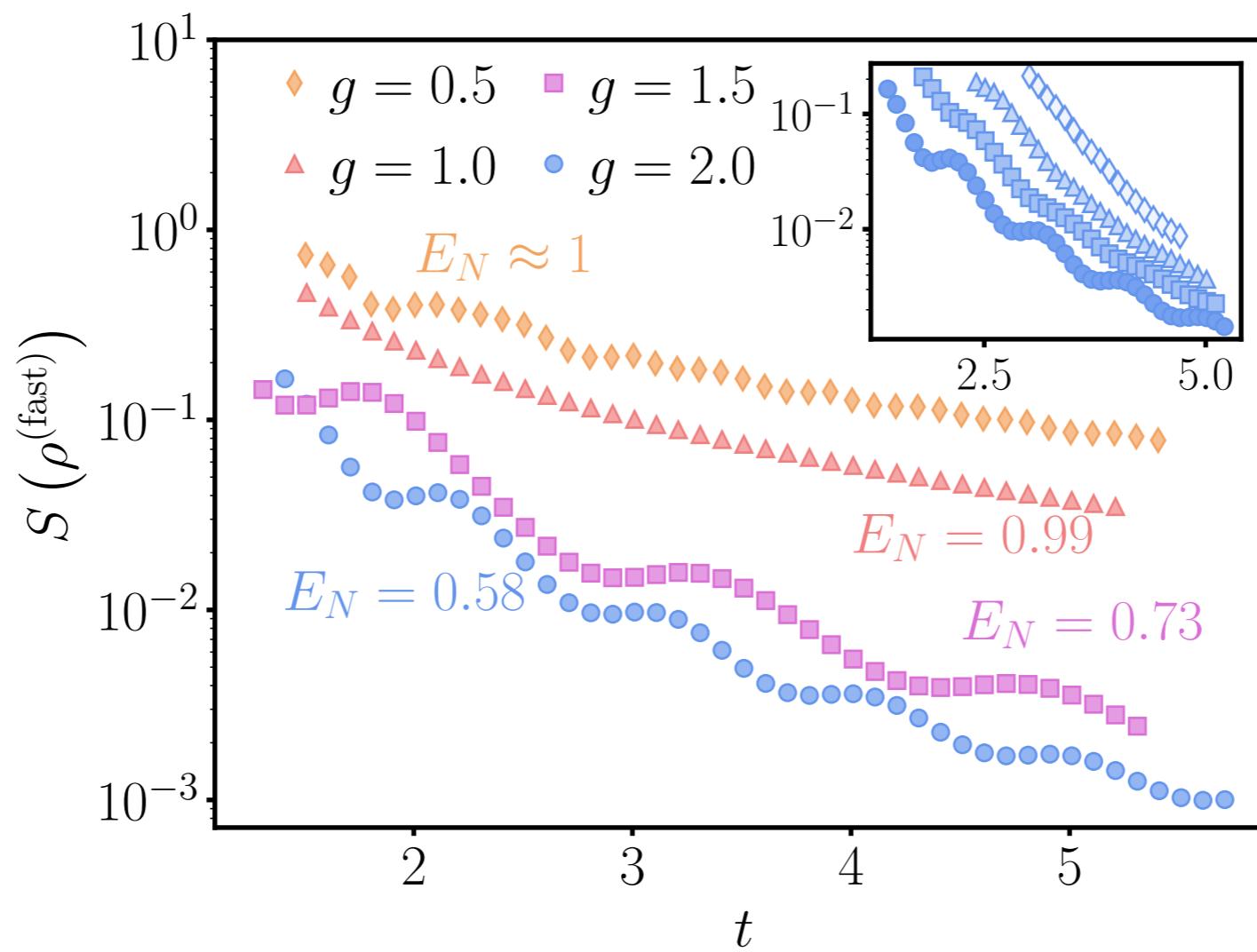
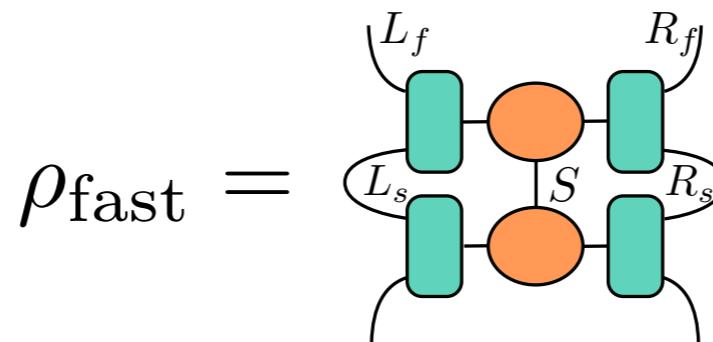


$$S_{\text{fast-slow}} \dashrightarrow \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array}$$

$$\rho_{\text{fast}} = \begin{array}{c} \text{Diagram showing two coupled systems. Each system has two green rectangular components (L_s and R_s) connected by a dashed vertical line. Between them is an orange circular component (S). A dashed red vertical line connects the S components of both systems. Labels L_f and R_f are at the top and bottom respectively.} \\ E_N(\rho_{\text{fast}}) \end{array}$$

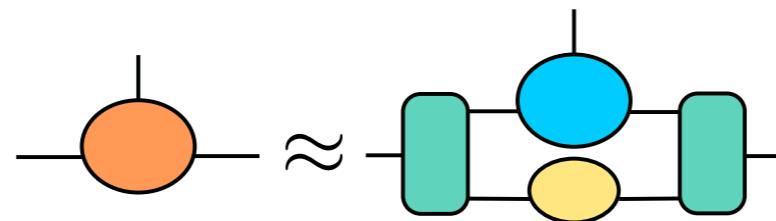
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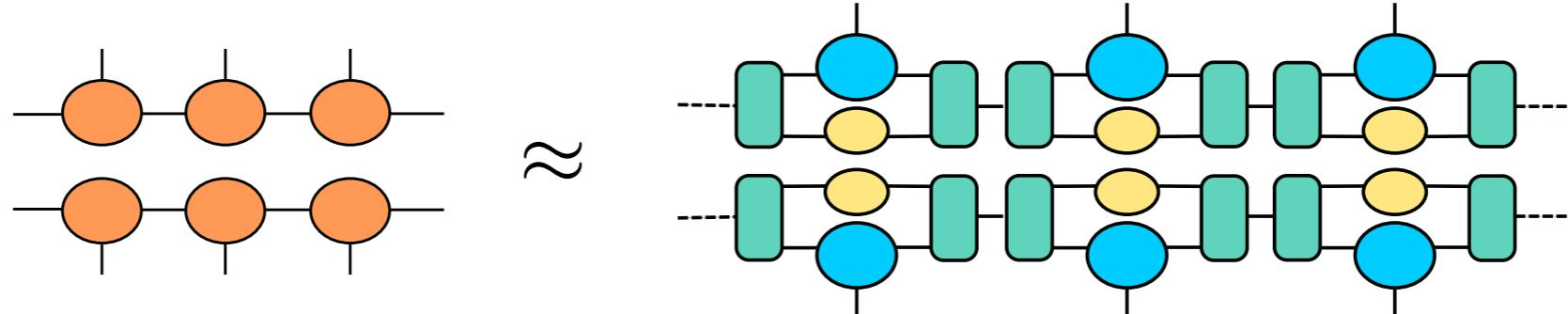


aTN algorithm

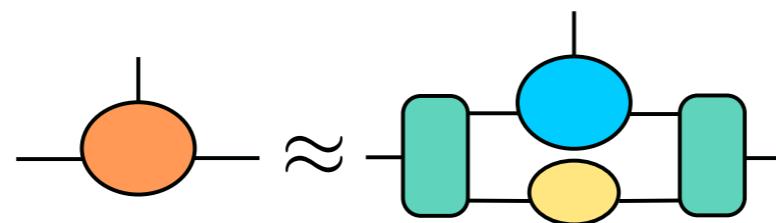
effective TN description



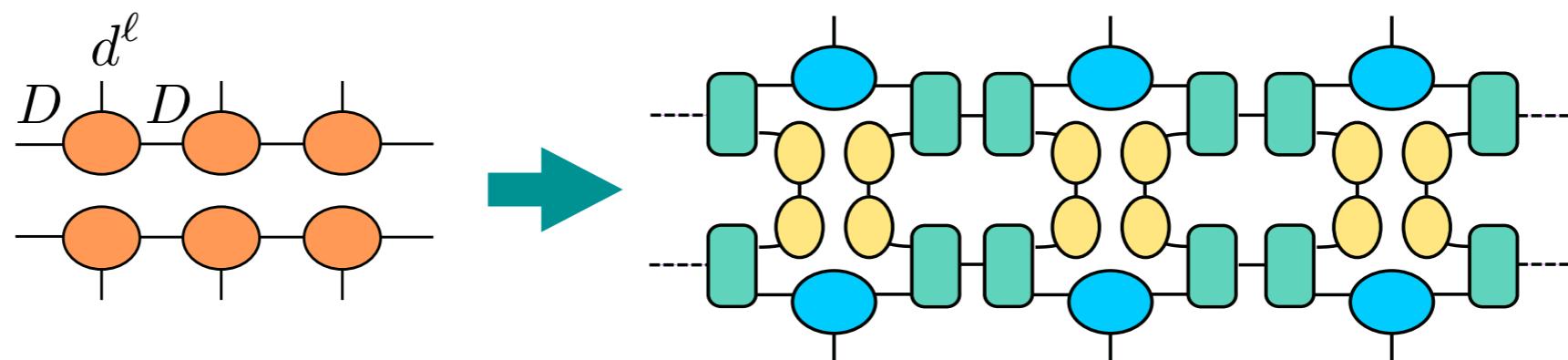
contribute as mixture to local
observables in neighbouring blocks



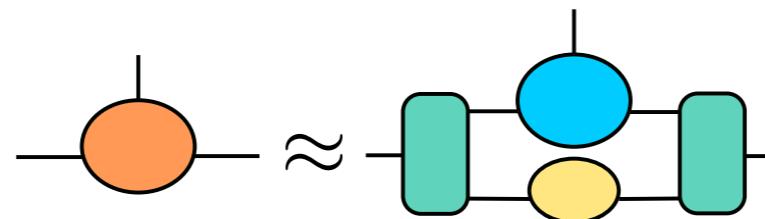
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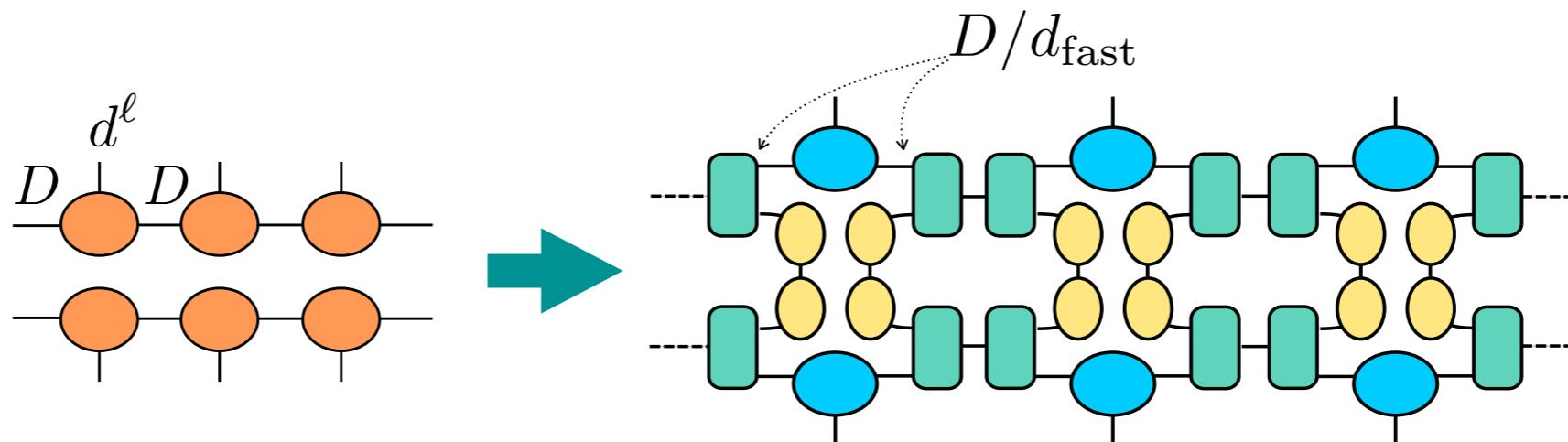
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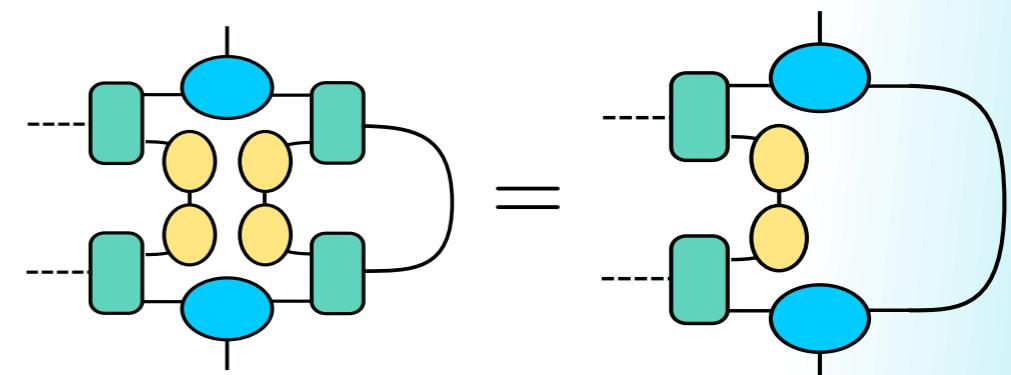
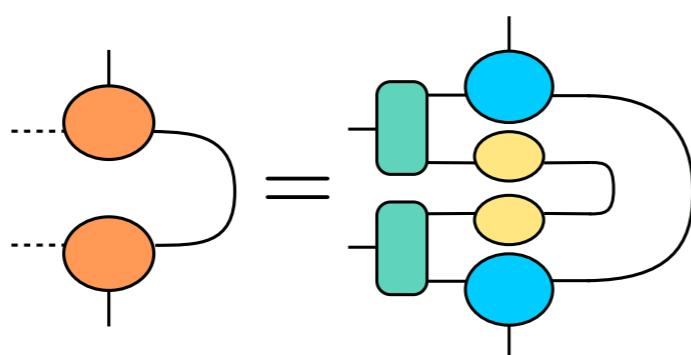
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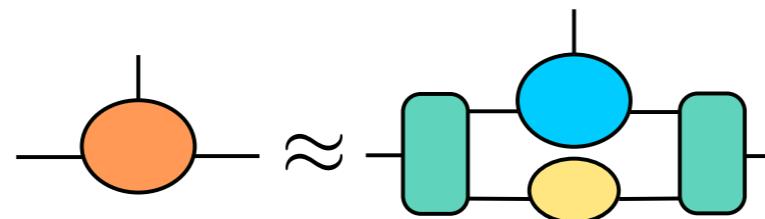
contribute as mixture to local
observables in neighbouring blocks



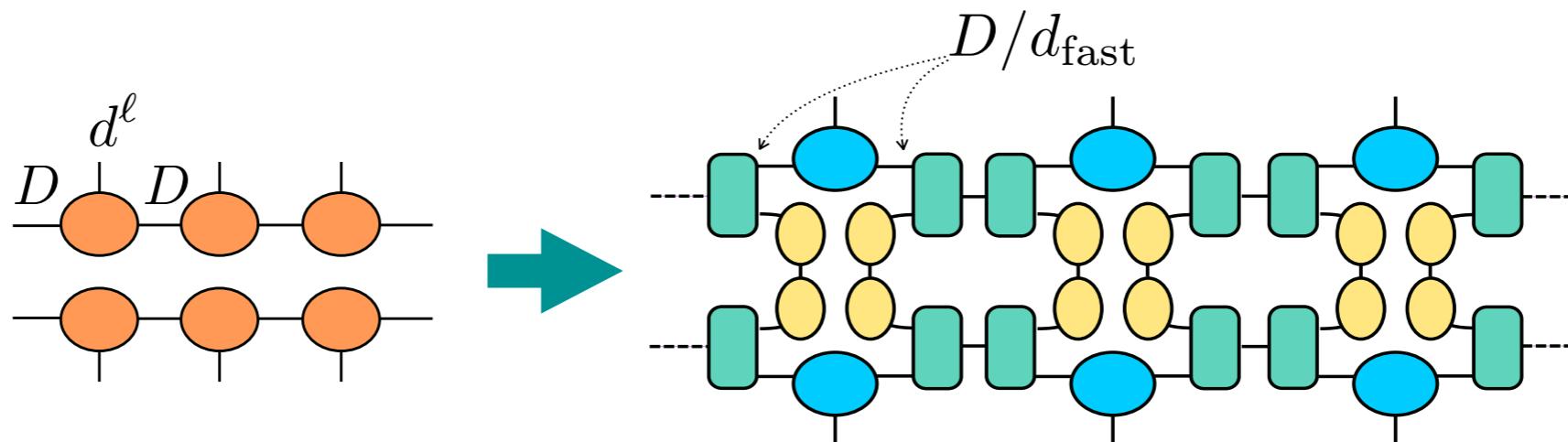
preserves
reduced density
matrices



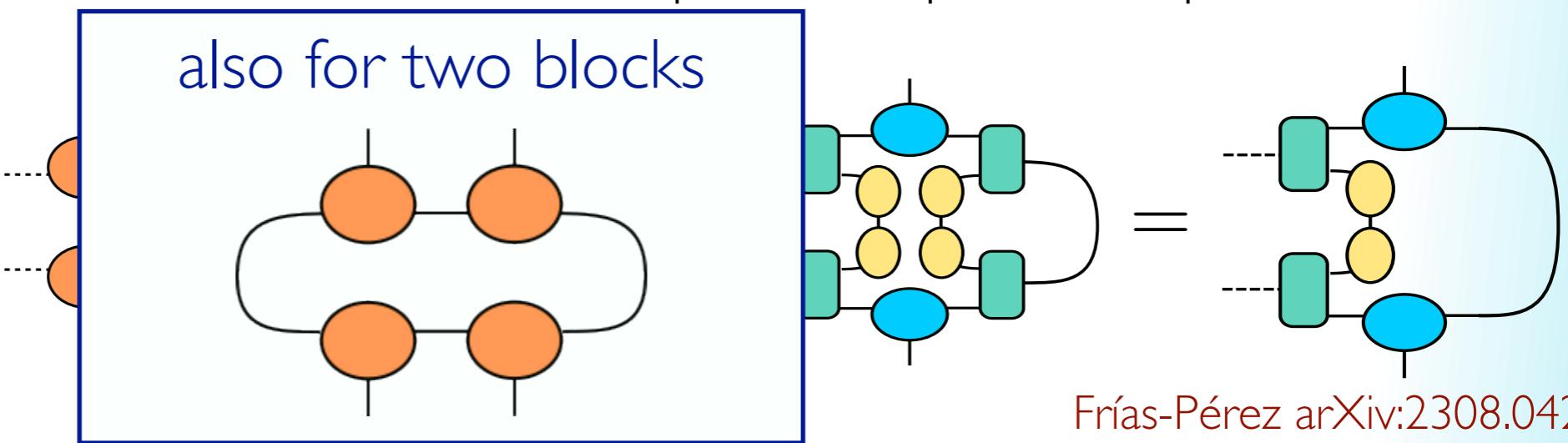
effective TN description



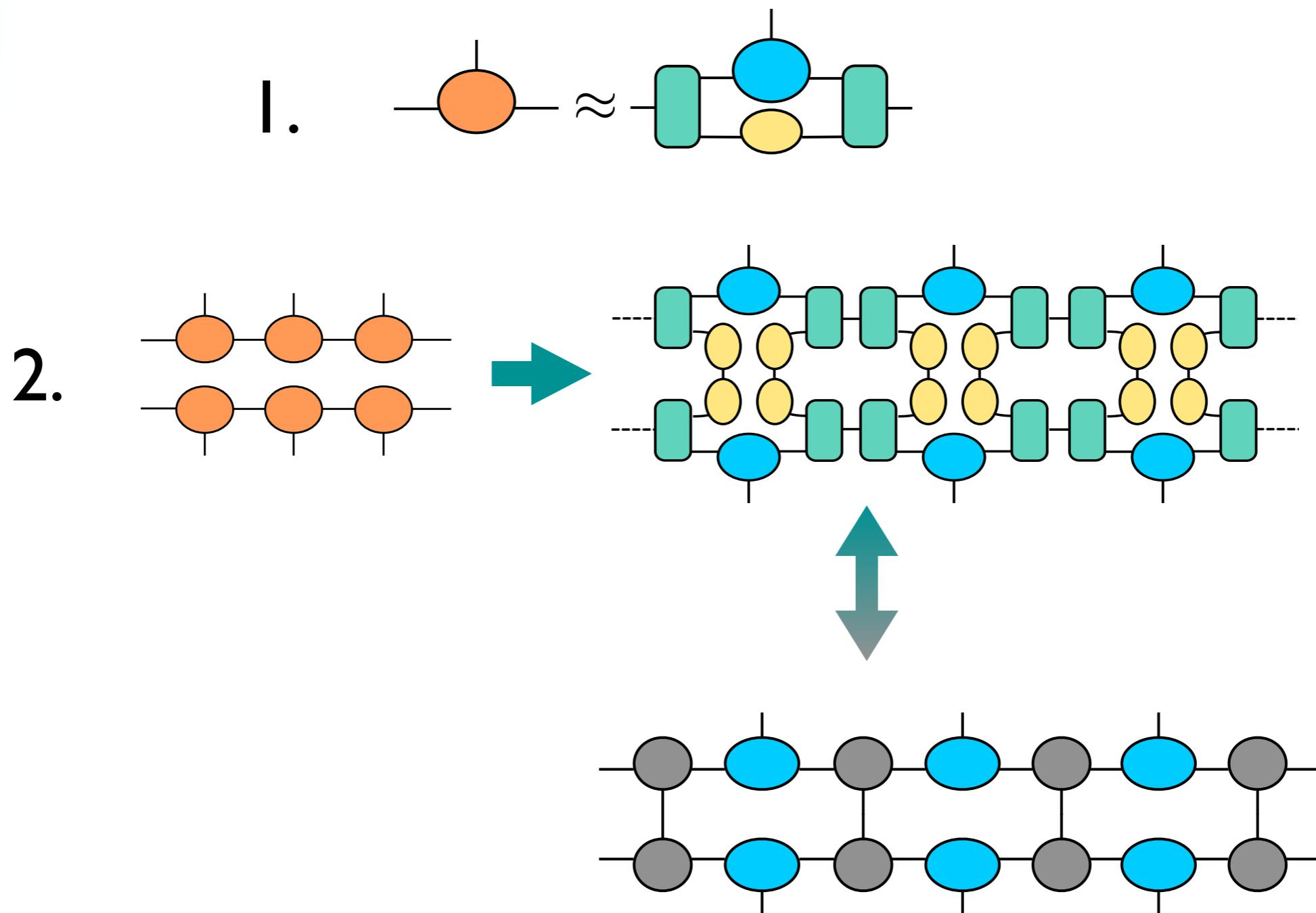
contribute as mixture to local
observables in neighbouring blocks



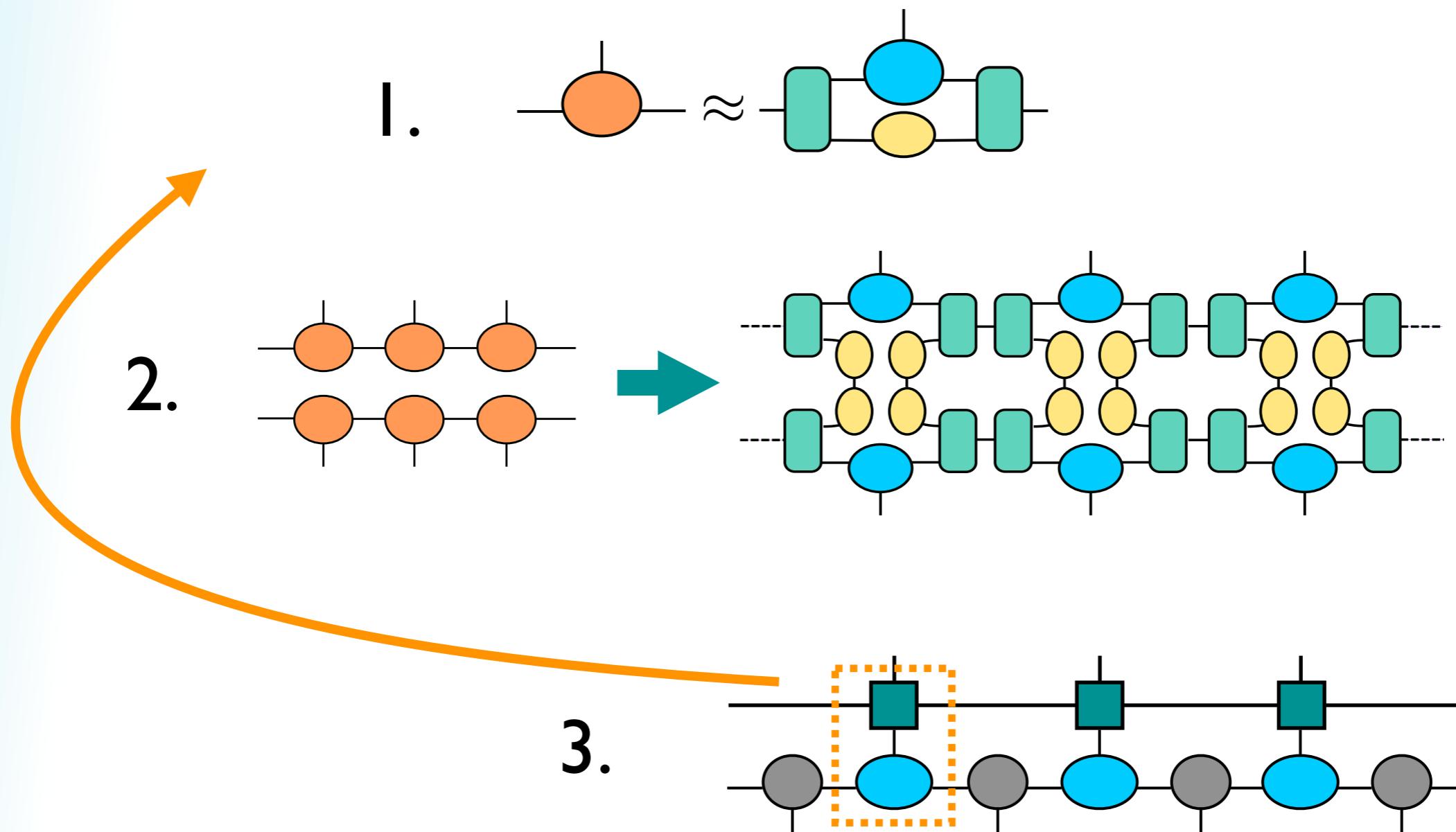
preserves
reduced density
matrices



TN algorithm



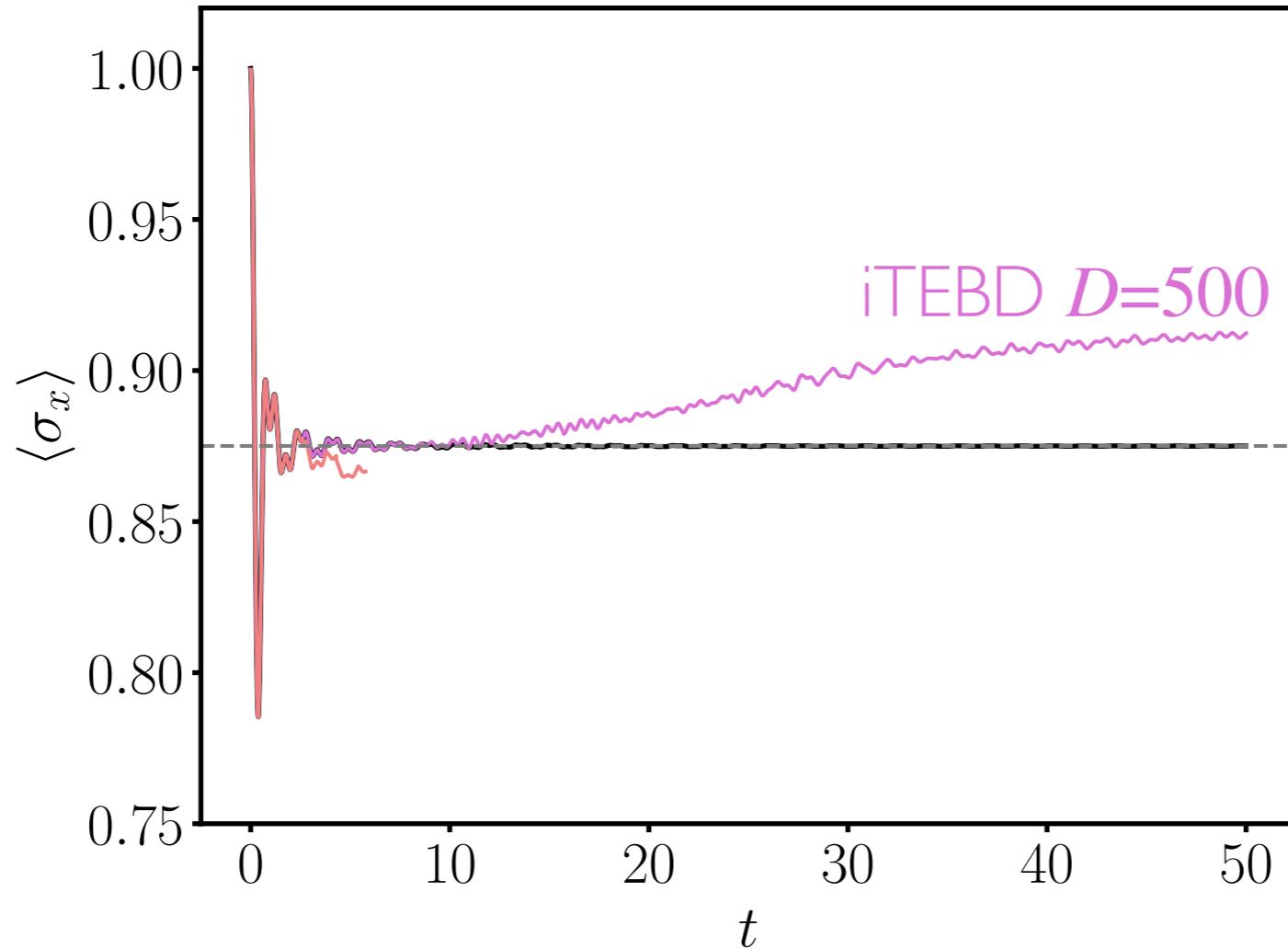
TN algorithm



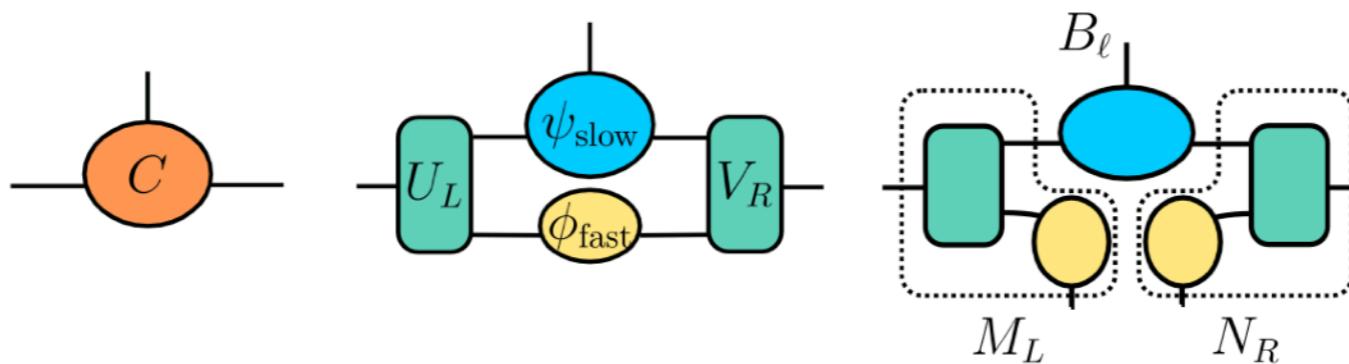
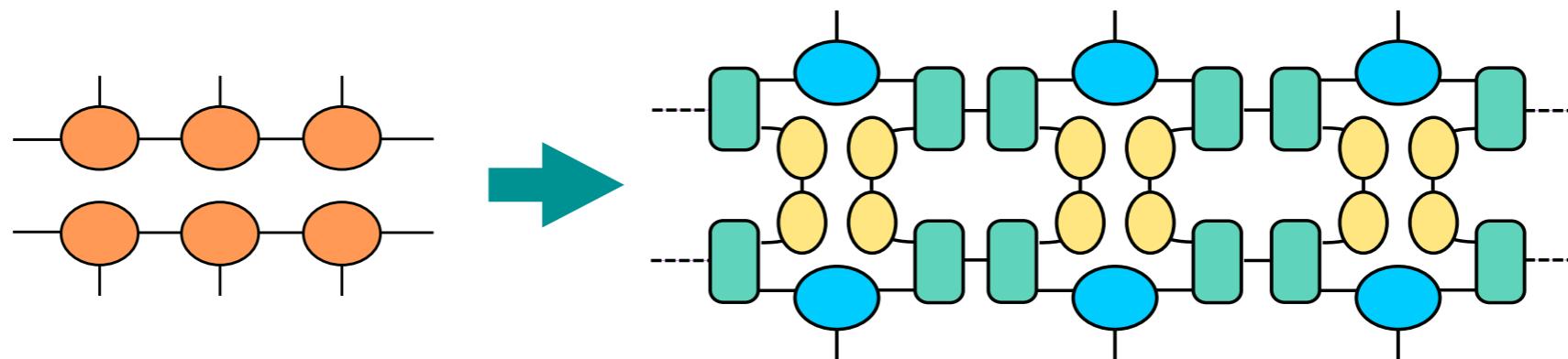
TN algorithm

$$H_{\text{Ising}} = -J \left(\sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i \sigma_x^{[i]} \right)$$

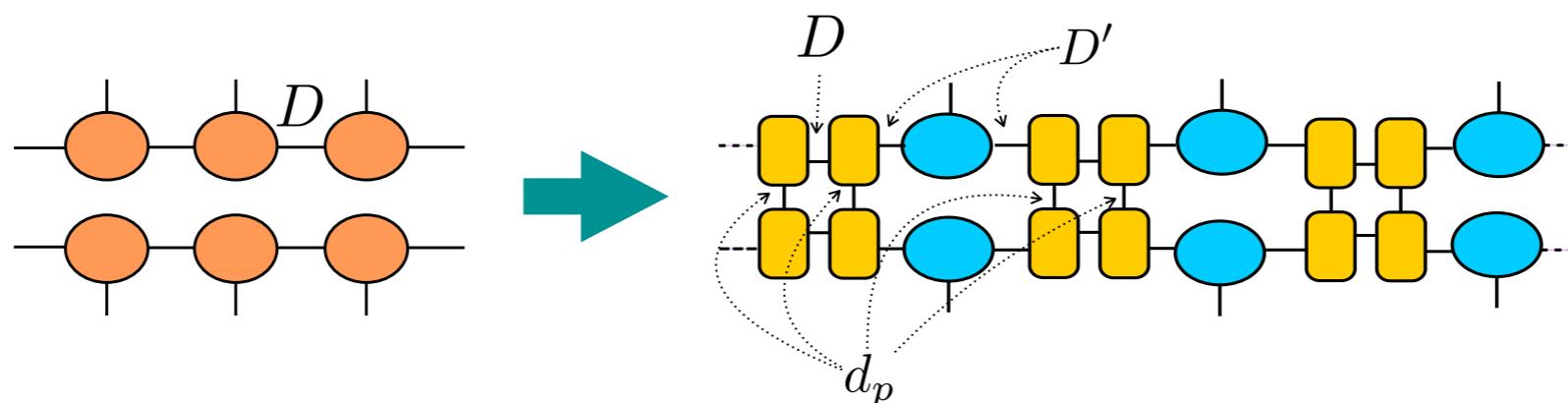
$J=1$
 $g=2$



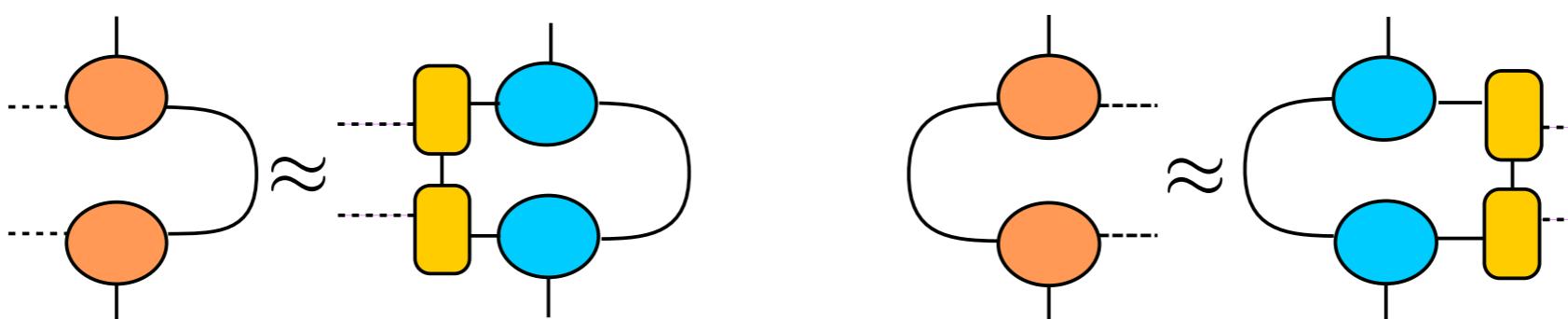
improved heuristic TN algorithm



improved heuristic TN algorithm



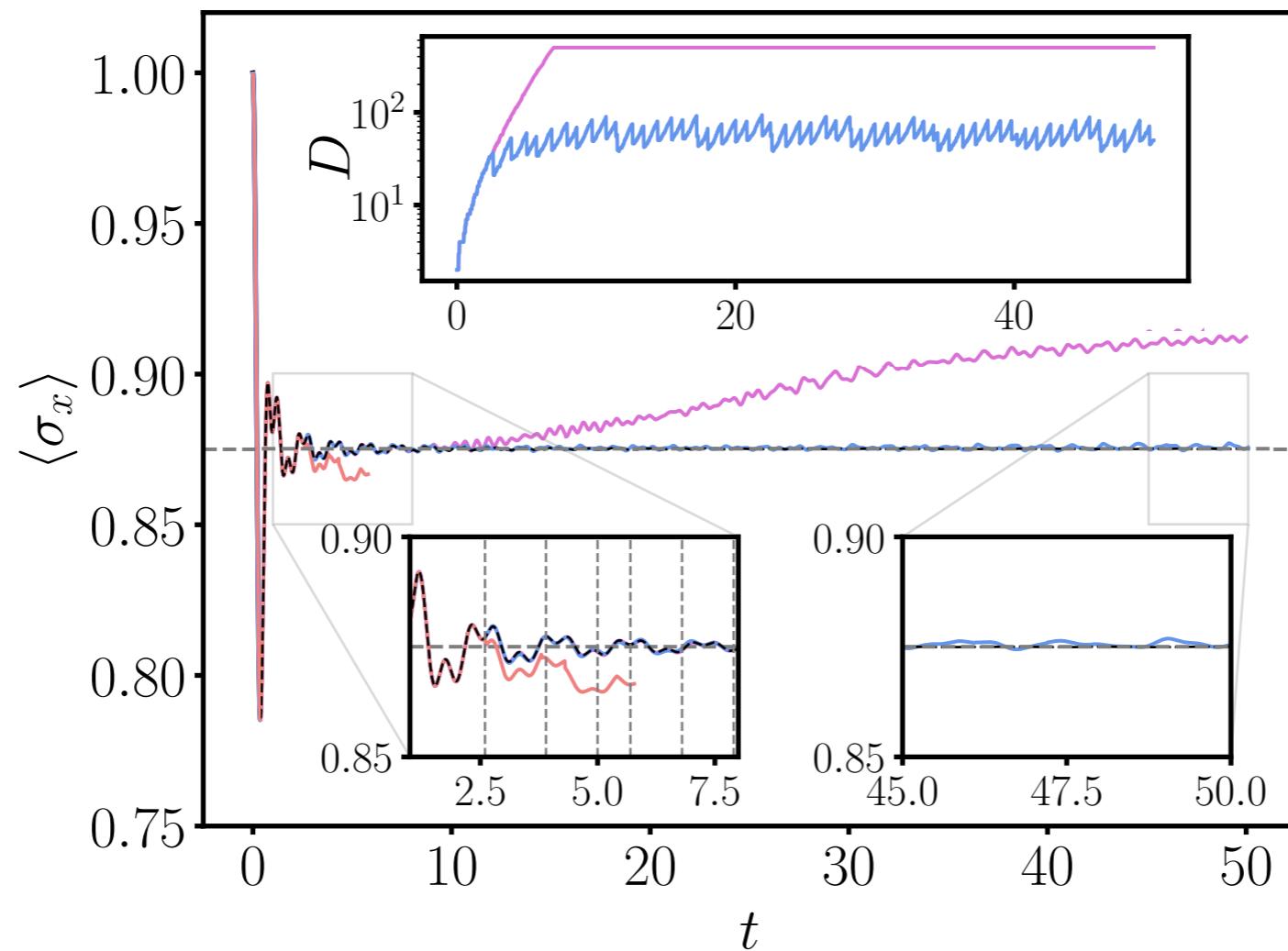
such that condition on rdms holds



TN algorithm

$$H_{\text{Ising}} = -J \left(\sum_{i=1}^{N-1} \sigma_z^{[i]} \sigma_z^{[i+1]} + g \sum_i \sigma_x^{[i]} \right)$$

$J=1$
 $g=2$



To conclude



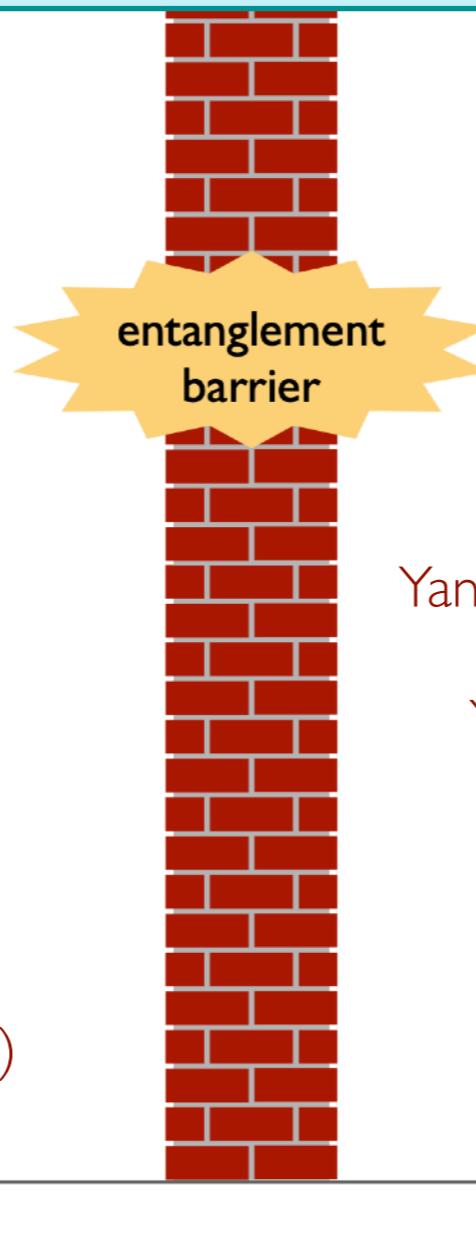
energy filters & TNS can provide other
(classical and quantum) tools to get dynamical
properties

changing entanglement
perspective:
transforming long-range
entanglement into
mixture

Frías-Pérez, Tagliacozzo, MCB,
arXiv:2308.04291

light-cone TN contraction
for local observables

Frías-Pérez, MCB, PRB 106, 115117 (2022)



spectral properties of a
QMB Hamiltonian

Yang, Iblisdir, Cirac, MCB, PRL 124, 100602 (2020)
Lu, PRX Quantum 2, 020321 (2021)
Yang, Cirac, MCB, PRB 106, 024307 (2022)