



Pendolo  $\begin{cases} \sum F_x = -w(t) \text{ disturbo} \\ \sum F_y = -mg \text{ gravità.} \end{cases}$

→ RIFORMLAZIONE LAGRANGIANA NSL

$$\Gamma \frac{\partial}{\partial t} \frac{\partial \mathcal{L}}{\partial \dot{q}_r} - \frac{\partial \mathcal{L}}{\partial q_r} = \sum_{i=1}^m \left( F_{ix} \frac{\partial \dot{q}_i}{\partial q_r} + F_{iy} \frac{\partial \dot{q}_i}{\partial q_r} \right) \quad \text{con } q_r, r=3,2 \text{ coordinate generalizzate.}$$

Pendolo  $\begin{cases} x_p(t) = f_1(x_c(t); \theta(t)) = x_c(t) - l \sin \theta(t) \\ y_p(t) = f_2(x_c(t); \theta(t)) = R + l \cos \theta(t). \end{cases}$

• Definiamo l'energia cinetica del modello complesso.

$$K(x_c(t); \theta(t)) = \frac{1}{2} m \dot{x}_c^2(t) + \frac{1}{2} m \left( \dot{x}_c(t) - l \cos \theta(t) \dot{\theta}(t) \right)^2 + \frac{1}{2} m \left( -l \sin \theta(t) \dot{\theta}(t) \right)^2 + \frac{1}{2} I_p \dot{\theta}^2(t)$$

determinazione le relazioni,

$$\begin{aligned} \frac{\partial}{\partial \dot{\theta}} K &= m (\dot{x}_c(t) - l \cos \theta(t) \dot{\theta}(t)) (-l \sin \theta(t)) + m (-l \sin \theta(t) \dot{\theta}(t)) (-l \sin \theta(t)) + I_p \dot{\theta}(t) = \\ &= -ml \cos \theta(t) \dot{x}_c(t) + ml^2 \cos^2 \theta(t) \dot{\theta}(t) + ml^2 \sin^2 \theta(t) \dot{\theta}(t) + I_p \dot{\theta}(t) = \\ &= (ml^2 + I_p) \dot{\theta}(t) - ml \cos \theta(t) \dot{x}_c(t). \end{aligned}$$

$$\frac{\partial}{\partial t} \frac{\partial}{\partial \dot{\theta}} K = (ml^2 + I_p) \ddot{\theta}(t) - ml \cos \theta(t) \ddot{x}_c(t) + ml \sin \theta(t) \dot{\theta}(t) \dot{x}_c(t)$$

$$\begin{aligned} \frac{\partial}{\partial \theta} K &= m (\dot{x}_c(t) - l \cos \theta(t) \dot{\theta}(t)) (l \sin \theta(t) \dot{\theta}) + m (-l \sin \theta(t) \dot{\theta}(t)) (-l \cos \theta(t) \dot{\theta}(t)) \\ &= ml \sin \theta(t) \dot{\theta}(t) \dot{x}_c(t) - ml^2 \sin \theta \cos \theta \dot{\theta}^2(t) + ml^2 \sin \theta \cos \theta \dot{\theta}^2(t) = \\ &= ml \sin \theta(t) \dot{\theta}(t) \dot{x}_c(t). \end{aligned}$$

• Cinematico Pendolo  

$$\begin{cases} x_p(t) = x_c(t) - l \sin \theta(t) \\ y_p(t) = R + l \cos \theta(t), \end{cases}$$

deriviamo rispetto il tempo,

$$\begin{cases} \dot{x}_p(t) = \dot{x}_c(t) - l \cos \theta(t) \dot{\theta}(t) \\ \dot{y}_p(t) = -l \sin \theta(t) \dot{\theta}(t) \end{cases}$$

Identifichiamo le forze sui singoli componenti.

Cart.  $\begin{cases} \sum F_x = u(t) - \mu \dot{x}_c(t) \\ \sum F_y \text{ trascurte.} \end{cases}$

Cart.  $\begin{cases} x_c(t) = f_1(x_c(t); \theta(t)) = x_c(t) \\ y_c(t) = f_2(x_c(t); \theta(t)) = R + l \cos \theta(t) \end{cases}$

coordinate generalizzate.

infine, otteniamo che

$$\frac{\partial}{\partial t} \frac{\partial}{\partial \theta} K - \frac{\partial}{\partial \theta} K = (m\ell^2 + I_p) \ddot{\theta}(t) - m\ell \cos \theta(t) \dot{x}_c(t) = F_\theta$$

con  $F_\theta$  forza generalizzata.

$$F_\theta = \sum_{i=1}^2 \left( F_i x \frac{\partial f_i}{\partial \theta} + F_i y \frac{\partial g_i}{\partial \theta} \right) = -\omega(t)(-\ell \cos \theta(t)) - mg(-\ell \sin \theta(t))$$

$$\frac{\partial f_1}{\partial \theta} = \frac{\partial f_2}{\partial \theta} = 0.$$

$$\frac{\partial g_2}{\partial \theta} = -\ell \cos \theta(t), \quad \frac{\partial g_1}{\partial \theta} = -\ell \sin \theta(t)$$

Io Equazione Dinamica Sisteme

$$\bullet (m\ell^2 + I_p) \ddot{\theta}(t) - m\ell \cos \theta(t) \dot{x}_c(t) - mg\ell \sin \theta(t) = \omega(t)\ell \cos \theta(t).$$

Determiniamo le variazioni di  $K$  rispetto a  $x_c$

$$\begin{aligned} \frac{\partial K}{\partial \dot{x}_c} &= M \dot{x}_c(t) + m(\dot{x}_c(t) - \ell \cos \theta(t) \dot{\theta}(t)) = \\ &= (M+m) \dot{x}_c(t) - m\ell \cos \theta(t) \dot{\theta}(t). \end{aligned}$$

$$\frac{\partial K}{\partial x_c} = 0.$$

$$\frac{\partial}{\partial t} \frac{\partial}{\partial \dot{x}_c} K = (M+m) \ddot{x}_c(t) - m\ell \cos \theta(t) \ddot{\theta}(t) + m\ell \sin \theta(t) \dot{\theta}^2(t) = F_{x_c}$$

$$F_{x_c} \hat{=} \sum_{i=1}^2 \left( F_i x \frac{\partial f_i}{\partial \dot{x}_c} + F_i y \frac{\partial g_i}{\partial \dot{x}_c} \right) =$$

$$\frac{\partial f_1}{\partial \dot{x}_c} = 1, \quad \frac{\partial g_1}{\partial \dot{x}_c} = 0. \quad = u(t) - \mu \dot{x}_c(t) - \omega(t).$$

$$\frac{\partial f_2}{\partial \dot{x}_c} = 1, \quad \frac{\partial g_2}{\partial \dot{x}_c} = 0$$

Modelli Non Lineari Sistemi in ANALISI

$$\bullet (m\ell^2 + I_p) \ddot{\theta}(t) - m\ell \cos \theta(t) \dot{x}_c(t) - mg\ell \sin \theta(t) = \omega(t)\ell \cos \theta(t)$$

$$\bullet (M+m) \ddot{x}_c(t) - m\ell \cos \theta(t) \ddot{\theta}(t) + m\ell \sin \theta(t) \dot{\theta}^2(t) + \mu \dot{x}_c(t) = u(t) - \omega(t).$$

• Aumentazione in  $Ox$ .

Ricordiamo che premette in  $Ox$  equivale ad effettuare le seguenti approssimazioni:

$$\sin \theta(t) \approx \theta(t); \cos \theta(t) \approx 1; \dot{\theta}^2 \approx 0.$$

Il modello linearizzato in  $Ox$ , sarà:

$$1) (m\ell^2 + I_p) \ddot{\theta}(t) - m\ell \ddot{x}_c(t) - m\ell \ell \dot{\theta}(t) = \ell \cdot w(t)$$

$$2) (M+m) \ddot{x}_c(t) - m\ell \ddot{\theta}(t) + \mu \dot{x}_c(t) = u(t) - w(t).$$

### • RAPPRESENTAZIONE IN STATO

Possiamo scrivere le seguenti:

$$x_1(t) = \theta(t), \quad x_2(t) = \dot{\theta}(t), \quad x_3(t) = x_c(t), \quad x_4(t) = \dot{x}_c(t).$$

$$\begin{cases} \dot{x}_1(t) = x_2(t) \\ \dot{x}_2(t) = \frac{1}{(m\ell^2 + I_p)} \left\{ m\ell \ell x_1(t) + \ell w(t) + m\ell \dot{x}_4(t) \right\} \\ \dot{x}_3(t) = x_4(t) \\ \dot{x}_4(t) = \frac{1}{(M+m)} \left\{ -\mu x_4(t) + u(t) - w(t) + m\ell \dot{x}_2(t) \right\}. \end{cases}$$

Sostituendo le scritte espressioni,

$$\begin{aligned} \dot{x}_1(t) &= x_2(t) \\ \dot{x}_2(t) &= \frac{1}{(m\ell^2 + I_p)} \left\{ m\ell \ell x_1(t) + \ell w(t) + \frac{m\ell}{(M+m)} \left[ -\mu x_4(t) + u(t) - w(t) + m\ell \dot{x}_2(t) \right] \right\} \\ &= \frac{m\ell \ell}{(m\ell^2 + I_p)} x_1(t) + \frac{\ell}{(m\ell^2 + I_p)} w(t) - \frac{m\ell \mu}{(m\ell^2 + I_p)(M+m)} x_4(t) + \frac{m\ell}{(m\ell^2 + I_p)(M+m)} u(t) + \\ &\quad - \frac{m\ell}{(m\ell^2 + I_p)(M+m)} w(t) + \frac{(m\ell)^2}{(m\ell^2 + I_p)(M+m)} \dot{x}_2(t). \end{aligned}$$

onero,

$$\begin{aligned} \left[ 1 - \frac{(m\ell)^2}{(m\ell^2 + I_p)(M+m)} \right] \dot{x}_2(t) &= \frac{m\ell \ell}{(m\ell^2 + I_p)} x_1(t) - \frac{m\ell \mu}{(m\ell^2 + I_p)(M+m)} x_4(t) + \\ &\quad + \frac{m\ell}{(m\ell^2 + I_p)(M+m)} u(t) + \left[ \frac{\ell}{(m\ell^2 + I_p)} - \frac{m\ell}{(m\ell^2 + I_p)(M+m)} \right] w(t) \end{aligned}$$

poi,

$$\dot{x}_3(t) = x_4(t).$$

ed infine

$$\dot{x}_4(t) = \frac{1}{(r\ell+m)} \left\{ -\mu x_4(t) + u(t) - \omega(t) + \frac{m\ell}{(m\ell^2 + I_p)} [m\ell x_3(t) + \ell\omega(t) + m\ell x_4(t)] \right\}$$

$$= -\frac{\mu}{(r\ell+m)} x_4(t) + \frac{1}{(r\ell+m)} u(t) - \frac{1}{(r\ell+m)} \omega(t) + \frac{(m\ell)^2 f}{(m\ell^2 + I_p)(r\ell+m)} x_4(t) +$$

$$+ \frac{m\ell^2}{(m\ell^2 + I_p)(r\ell+m)} \omega(t) + \frac{(m\ell)^2}{(m\ell^2 + I_p)(r\ell+m)} \dot{x}_4(t)$$

riportando a primo membro,

$$\left[ 1 - \frac{(m\ell)^2}{(m\ell^2 + I_p)(r\ell+m)} \right] \dot{x}_4(t) = -\frac{\mu}{(r\ell+m)} x_4(t) + \frac{(m\ell)^2 f}{(m\ell^2 + I_p)(r\ell+m)} x_4(t) +$$

$$+ \frac{1}{(r\ell+m)} u(t) + \left[ \frac{m\ell^2}{(m\ell^2 + I_p)(r\ell+m)} - \frac{1}{(r\ell+m)} \right] \omega(t).$$

In forma compatta,

$$\dot{x} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ A_{11} & 0 & 0 & A_{24} \\ 0 & 0 & 0 & 1 \\ A_{41} & 0 & 0 & A_{44} \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix} + \begin{bmatrix} 0 \\ B_2 \\ 0 \\ B_4 \end{bmatrix} u(t) + \underbrace{\begin{bmatrix} 0 \\ D_2 \\ 0 \\ D_4 \end{bmatrix}}_{u(t) \text{ turba}} \omega(t)$$

$$y = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix}$$