

Course Number	AE8112
Course Title	Computational Fluid Dynamics and Heat Transfer
Semester/Year	Summer/Spring 2021
Instructor	Dr. Seth Dworkin

**Problem Set 5**

Submission Date	July 29, 2021
Programing Language Used	Fortran90

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Q1a Using explicit finite difference, we can discretize the given vorticity-velocity eqn, with the below Taylor's series expansion centered at  $\phi_i$

$$\frac{\partial^2 \phi}{\partial x^2} \bigg|_i = \frac{\phi_{i-1} - 2\phi_i + \phi_{i+1}}{\Delta x^2} \quad (1.1)$$

$$\text{and } \frac{\partial \phi}{\partial x} \bigg|_i = \frac{\phi_{i+1} - \phi_{i-1}}{2\Delta x} \quad (1.2)$$

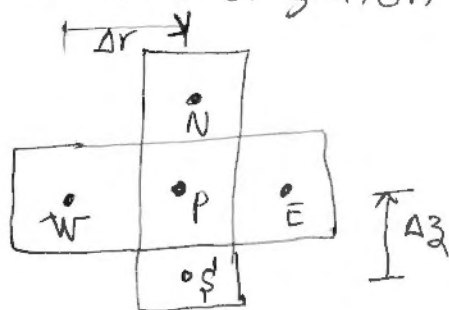
Applying this to our governing equations, we will have,  $\phi = \omega, v_r$  &  $v_z$  and for our 2-D coordinate

$$x = r \text{ \& \& } z$$

Now considering vorticity transport eqn., so that after applying the product rule, we have,

$$\rho \frac{\partial \omega}{\partial t} = \mu \frac{\partial^2 \omega}{\partial r^2} + \mu \frac{\partial^2 \omega}{\partial z^2} + \mu \left( -\frac{1}{r^2} \omega + \frac{1}{r} \frac{\partial \omega}{\partial r} \right) - \rho v_r \frac{\partial \omega}{\partial r} - \rho v_z \frac{\partial \omega}{\partial z} + \rho \frac{v_r \omega}{r}$$

So that applying the standard forward difference for the time discretization and eqn (1.1), (1.2), we get,



Q1 a) cont'd.:

$$\begin{aligned} \rho \left( \frac{\omega_p^{n+1} - \omega_p^n}{\Delta t} \right) &= \mu \left( \frac{\omega_w^n - 2\omega_p^n + \omega_E^n}{\Delta r^2} \right) + \mu \left( \frac{\omega_s^n - 2\omega_p^n + \omega_N^n}{\Delta z^2} \right) - \frac{\mu \omega_p^n}{r_p^2} \\ &+ \frac{\mu}{r_p} \left( \frac{\omega_E^n - \omega_w^n}{2\Delta r} \right) - \rho v_r^n \left( \frac{\omega_E^n - \omega_w^n}{2\Delta r} \right) - \rho v_z^n \left( \frac{\omega_N^n - \omega_s^n}{2\Delta z} \right) \\ &+ \frac{\rho v_r^n \omega_p^n}{r_p} \end{aligned}$$

now re-arranging for  $\omega_p^{n+1}$  since we will be computing it's values from an we get, initial state  $\omega_p^0$ .

$$\begin{aligned} \omega_p^{n+1} &= \frac{\mu \Delta t}{\rho \Delta r^2} (\omega_w^n - 2\omega_p^n + \omega_E^n) + \frac{\mu \Delta t}{\rho \Delta z^2} (\omega_s^n - 2\omega_p^n + \omega_N^n) - \frac{\mu \Delta t \omega_p^n}{\rho r_p^2} \\ &+ \frac{\mu \Delta t}{2\rho r_p \Delta r} (\omega_E^n - \omega_w^n) - \frac{v_r^n \Delta t}{2\Delta r} (\omega_E^n - \omega_w^n) - \frac{v_z^n \Delta t}{2\Delta z} (\omega_N^n - \omega_s^n) \\ &+ \frac{v_r^n \omega_p^n \Delta t}{r_p} + \omega_p \end{aligned}$$

expanding and grouping like terms, assuming  $\mu, \rho = \text{const}$

we get,

$$\begin{aligned} \omega_p^{n+1} &= \left( \frac{\mu \Delta t}{\rho \Delta r^2} - \frac{\mu \Delta t}{2\rho r_p \Delta r} + \frac{v_r^n \Delta t}{2\Delta r} \right) \omega_w^n + \left( \frac{\mu \Delta t}{\rho \Delta z^2} + \frac{v_z^n \Delta t}{2\Delta z} \right) \omega_s^n \\ &- \left( \frac{2\mu \Delta t}{\rho \Delta r^2} + \frac{2\mu \Delta t}{\rho \Delta z^2} + \frac{\mu \Delta t}{\rho r_p^2} - \frac{v_r^n \Delta t}{r_p} - 1 \right) \omega_p^n \\ &+ \left( \frac{\mu \Delta t}{\rho \Delta r^2} + \frac{\mu \Delta t}{2\rho r_p \Delta r} - \frac{v_r^n \Delta t}{2\Delta r} \right) \omega_E^n + \left( \frac{\mu \Delta t}{\rho \Delta z^2} - \frac{v_z^n \Delta t}{2\Delta z} \right) \omega_N^n \end{aligned}$$

Q1a | cont'd.!

For ease of computational setup, let have the above as,

$$\omega_p^{n+1} = a_w^n \omega_w^n + a_s^n \omega_s^n - a_p^n \omega_p^n + a_E^n \omega_E^n + a_w^n \omega_w^n \quad \text{--- (2.2)}$$

This can be solved directly from given initial  $\omega_p^0$

→ Now we discretize the poisson radial velocity eqn in similar way,

$$\frac{\partial^2 V_r}{\partial r^2} + \frac{\partial^2 V_r}{\partial z^2} = \frac{\partial \omega}{\partial z} + \frac{V_r}{r^2} - \frac{1}{r} \frac{\partial V_r}{\partial r} \quad \text{--- (2.3)}$$

we can apply eqs (1.1) & (1.2) to discretize as,

$$\frac{V_{rw}^{n+1} - 2V_{rp}^{n+1} + V_{re}^{n+1}}{\Delta r^2} + \frac{V_{rs}^{n+1} - 2V_{rp}^{n+1} + V_{rn}^{n+1}}{\Delta z^2} = \frac{\omega_N^{n+1} - \omega_s^{n+1}}{2\Delta z} + \frac{V_{rp}^{n+1}}{r_p^2}$$

multiplying by  $r_p^2$  and re-arranging we ge,

$$\left( \frac{r_p^2}{\Delta r^2} - \frac{r_p}{2\Delta r} \right) V_{rw}^{n+1} + \frac{r_p^2}{\Delta z^2} V_{rs}^{n+1} - \left( \frac{2r_p^2}{\Delta r^2} + \frac{2r_p^2}{\Delta z^2} + 1 \right) V_{rp}^{n+1} + \left( \frac{r_p^2}{\Delta r^2} + \frac{r_p}{2\Delta r} \right) V_{re}^{n+1} + \frac{r_p^2}{\Delta z^2} V_{rn}^{n+1} = r_p^2 \left( \frac{\omega_N^{n+1} - \omega_s^{n+1}}{2\Delta z} \right)$$

Assuming that  $\omega_p^{n+1}$  from eqn (2.2), we can re-write the above as, (for ease of computational setup).

Q1a) Cont'd.:

$$a_w v_{rw}^{n+1} + a_s v_{rs}^{n+1} - a_p v_{rp}^{n+1} + a_e v_{re}^{n+1} + a_v v_{rv}^{n+1} = b^{n+1} \quad (2.4)$$

we will use Preconditioned Bi-CGSTAB to solve the above pentadiagonal linear system for  $v_r^{n+1}$ .

→ Now let's discretize the continuity eqn, applying the product rule we have,

$$\frac{\partial v_z}{\partial z} = -\frac{1}{r} v_r - \frac{r}{r} \frac{\partial v_r}{\partial r}$$

using first order backward difference for  $\frac{\partial v_z}{\partial z}$  and eqn (1.2) for  $\frac{\partial v_r}{\partial r}$ , we get,

$$\frac{v_{zp} - v_{zs}}{\Delta z} = -\frac{v_{rp}^{n+1}}{r_p} - \left( \frac{v_{re}^{n+1} - v_{rw}^{n+1}}{2\Delta r} \right)$$

re-arranging for  $v_{zp}^{n+1}$ , since  $v_{zs}^{n+1}$  can be known from our B.C., we will have,

$$v_{zp}^{n+1} = v_{zs}^{n+1} - \frac{\Delta z}{r_p} v_{rp}^{n+1} - \frac{\Delta z}{2\Delta r} v_{re}^{n+1} + \frac{\Delta z}{2\Delta r} v_{rw}^{n+1}$$

for ease of computational setup, we re-write the above as,

$$v_{zp}^{n+1} = v_{zs}^{n+1} - a_p v_{rp}^{n+1} - a_e v_{re}^{n+1} + a_v v_{rv}^{n+1} \quad (2.5)$$

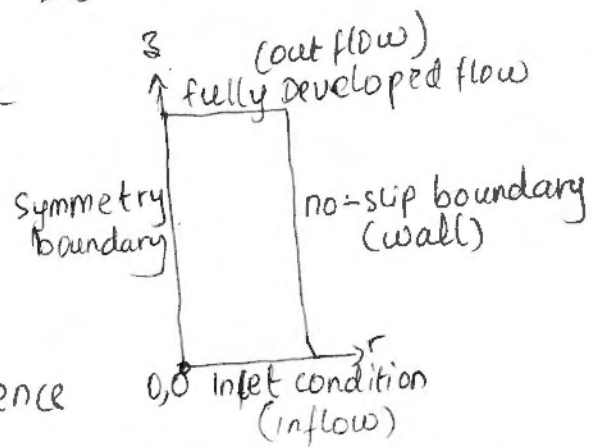
This can be solved directly for the axial velocity  $v_{zp}^{n+1}$

Q1a) cont'd.

>> Next, in order to apply the Boundary conditions (B.Cs), let's consider the four boundaries of our 2-D model as,

we have the vorticity for axisymmetric cylindrical coordinate as,

$$\omega = \frac{\partial v_r}{\partial z} - \frac{\partial v_z}{\partial r}$$



using first order one-sided difference

for  $\frac{\partial v_r}{\partial z}$ , and a second-order centered difference for  $\frac{\partial v_z}{\partial r}$ ,

the above can be discretized as,

$$\omega_p^{n+1} = \frac{v_{rn}^{n+1} - v_{rp}^{n+1}}{\Delta z} - \left( \frac{v_{ze}^{n+1} - v_{zw}^{n+1}}{2\Delta r} \right) \quad \text{--- (2.6)}$$

and we have the half parabola eqn as,

$$v_z = CLVZ - CLVZ(r^2) \quad \text{--- (2.7)}$$

So that for:

Inflow;

$$v_{zp} = CLVZ - CLVZ(r_p^2)$$

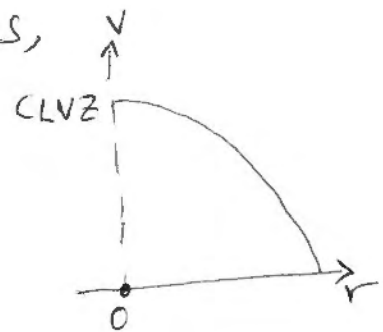
re-written as  $v_{zp} = k_1$  (for computational ease)

$$v_{rp} = 0$$

using eqn (2.6), we have,

$$\omega_p = \frac{v_{rn} - v_{rp}}{\Delta z} - \left( \frac{v_{ze} - v_{zw}}{2\Delta r} \right)$$

re-written,  $\omega_p = k_2$  (for computational ease)



Q1a) cont'd.!

wall;

$$v_{zp} = 0$$

$$v_{rp} = 0$$

Since  $\frac{\partial v_r}{\partial z} = 0$ , we will have

$$\omega_p = \frac{v_{zw}}{2\Delta r}$$

Symmetry;

$$v_{rp} = 0$$

$$\omega_p = 0$$

since  $\frac{\partial v_z}{\partial z} = 0$ , discretizing gives,

$$v_{zp} = v_{zs}$$

outflow;

$$\text{we are given } \frac{\partial \omega}{\partial z} = \frac{\partial v_r}{\partial z} = \frac{\partial v_z}{\partial z} = 0$$

and discretizing gives,

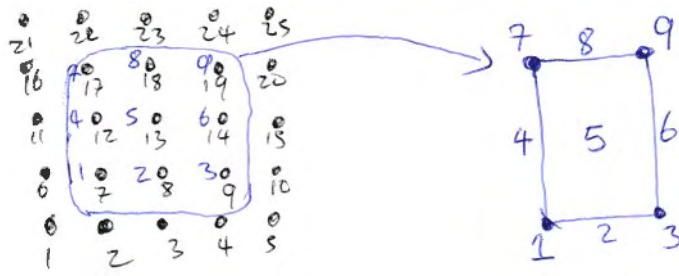
$$\omega_p = \omega_s$$

$$v_{rp} = v_{rs}$$

$$v_{zp} = v_{zs}$$

## Q1a | cont'd |

>> Now because of the unique nature for solving our  $U_r$  using a linear solver, let's assume we have a sample of grid mesh,



so that considering these 9 regions we will have the following boundary for  $U_r$ ,

① [1] The bottom left region:

applying both the inflow and symmetry B.C.s to eqn (2.4)

$$-a_p U_{rp} + a_E U_{rE} + a_N U_{rN} = b$$

② [2] The bottom region:

we apply only the inflow conditions to eqn (2.4),

$$a_w U_{rw} - a_p U_{rp} + a_E U_{rE} + a_N U_{rN} = b$$

③ [3] The bottom right region:

applying both the inflow and wall B.C.s. to eqn (2.4),

$$a_w U_{rw} - a_p U_{rp} + \cancel{a_E U_{rE}} + a_N U_{rN} = b$$

④ [4] The left region:

we will apply only the symmetry condition to eqn (2.4)



## Q1a | Cont'd!

$$a_s V_{rs} - a_p V_{rp} + a_E V_{re} + a_n V_{rn} = b$$

① [5] The middle region:

Here we just use eqn (2.4) directly

② [6] The right region:

we will only apply the wall condition to eqn (2.4)

$$a_w V_{rw} + a_s V_{rs} - a_p V_{rp} + a_n V_{rn} = b$$

③ [7] The top left region

applying both the outflow and symmetry B.Cs to eqn (2.4)

$$a_s V_{rs} - a_p V_{rp} + a_E V_{re} + a_n V_{rp} = b$$

simplifying gives,

④ [8] The top region:

$$a_s V_{rs} - (a_p + a_w) V_{rp} + a_E V_{re} = b$$

we will only apply the outflow condition to eqn (2.4)

$$a_w V_{rw} + a_s V_{rs} - a_p V_{rp} + a_E V_{re} + a_n V_{rp} = b$$

simplifying gives,

$$a_w V_{rw} + a_s V_{rs} - (a_p - a_n) V_{rp} + a_E V_{re} = b$$

⑤ [9] The top right region:

applying both the outflow and wall B.Cs. to eqn (2.4)

$$a_w V_{re} + a_s V_{rs} - a_p V_{rp} + a_n V_{rp} = b$$

simplifying gives,

$$a_w V_{re} + a_s V_{rs} - (a_p - a_n) V_{rp} = b$$

Q1b Using the boxed eqns from Q1a, we can write our program in fortran 90.

And we have our numeric integrated mass flow rate as,

$$\dot{m}_j = \sum_{i=1}^{N_r} (2\pi V(i) r) \quad \dot{m}_j = \sum_{i=1}^{N_r}$$

$$\dot{m}_j = 2\pi \rho \sum_{i=1}^{N_r} V_{zi} r_i$$

### Results

- >> First I ran the simulation with  $Tol = 1e^{-5}$ , and it seemed to be taking approximately more than 7 hours and more than 5 million iterations, yet the tolerance wasn't small enough.
  - >> So I reduced the tolerance to  $Tol =$  , which sufficiently small enough for a balance in both computational cost and <sup>slight</sup> accuracy.
  - >> Also I tried adaptive timestep size (by incrementing  $dt$  by  $t$  ~~to~~ at each new iteration). This worked well at lesser number of iterations than the constant time step, only if the incrementing time ( $t$ ) is small enough. Which in my case, I choose  $t = 1e^{-7}$   
for  $dt = dt + t$
- Otherwise the solution will explode if  $t$  is not small enough.

```

!*****Begin Header*****
!This program was written by Godswill Ezeorah, Student Number: 501012886 on July 20, 2021.
!This program solves an unsteady linear/non-linear equation using finite difference method
!and was written as a solution to AE8112 PS5 q1(b,c,d)
!*****End Header*****

program unsteady_Dfinite
    implicit none
    !Variable declaration
    DOUBLE PRECISION, dimension(:), ALLOCATABLE :: omg_i, V_r, V_z, r_i, z_i
    DOUBLE PRECISION, PARAMETER :: rh=8.3D-4, RR=1, L=10, mu=1.4D-4, CLVZ=5
    DOUBLE PRECISION, PARAMETER :: Pi = 4*atan(1.0)
    DOUBLE PRECISION :: dr, dz, dt, tols, rrr, start, finish
    INTEGER :: nr, nz, jj, i1, j1, l1, nnr, nnz

    call cpu_time(start) !gets the start time, for timing purpose
    open(2, file = 'PS5_Q1w.txt', status = 'unknown')
    open(3, file = 'PS5_Q1vr.txt', status = 'unknown')
    open(4, file = 'PS5_Q1vz.txt', status = 'unknown')
    Nr=21 !Number of grid points along radial direction
    Nz=201 !Number of grid points along axial direction
    nnr=nr-2; nnz=nz-2 !Number of grid points in the interior region
    dr=RR/(nr-1)
    dz=L/(nz-1)
    dt=1D-6 !given timestep
    tols=1D-5
    ALLOCATE(omg_i(nr*nnz), V_r(nr*nnz), V_z(nr*nnz), r_i(nr), z_i(nz))

    !Initialization
    omg_i=0; V_r=0; V_z=0; r_i=0; z_i=0 !first initial consideration
    !rrr=0.75; omg_i=2*CLVZ*(rrr/RR**2); V_r=0; V_z=CLVZ*(1-(rrr/1)**2) !second initial consideration

    call unsteady(omg_i, V_r, V_z, r_i, z_i) !solves the unsteady system

    !Printing and writing results
    j1=nr*nnz-(nr-1)
    l1=nr*nnz
    jj=nz
    do i1 = j1, 1, -nr
        write(2,*) z_i(jj), omg_i(i1:l1)
        write(3,*) z_i(jj), V_r(i1:l1)
        write(4,*) z_i(jj), V_z(i1:l1)
        l1=l1-nr
        jj=jj-1
    end do
    call cpu_time(finish) !gets the end time
    write(2,*) r_i(1:nr); write(3,*) r_i(1:nr); write(4,*) r_i(1:nr)

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print *, "Computation Time = ", (finish-start)/(60*60),"hours"
close(2);close(3);close(4)

contains
!*****
subroutine unsteady(oma, Vr, Vz, ri, zi)
  DOUBLE PRECISION, dimension(nr*nz) :: oma, Vz, omo, b, Vr
  DOUBLE PRECISION, dimension(nnr*nnz):: d, e, f, g, h, bi, Vri
  DOUBLE PRECISION, dimension(nr) :: ri
  DOUBLE PRECISION, dimension(nz) :: zi, mdot, som
  DOUBLE PRECISION :: aa,aw,as,ap,ae,an, k1, k2, k3, t, rp, zp
  DOUBLE PRECISION :: Mmdot, Amdot
  integer :: i, j, k, nj, nk, ll, ij, iii
!!This subroutine solves unsteady fluid flow through a pipe problem using vorticity method

  open(1, file = 'PS5_Q1.txt', status = 'unknown')
  !Variable initialization
  d=0; e=0; f=0; g=0; h=0
  t=1D-8; rp=0; zp=0
  b=0
  Vri=0
  Mmdot=1
  ij=1
  do i = 1,nz !Loop for creating the axial and radial gridpoint dimensions
    if ( i<=nr ) then
      ri(i)=rp
      zi(i)=zp
    else
      zi(i)=zp
    end if
    rp=rp+dr
    zp=zp+dz
  end do
  do while (Mmdot>tols) !Main loop for time step
    j=nr+1
    k=nr+nr
    ll=1
    omo=oma

    !The below loop solves the vorticity at each gridpoint
    do i = 1,nr*nz
      aw=mu*dt/(rh*dr**2)-mu*dt/(2*rh*ri(ll)*dr)+Vr(i)*dt/(2*dr)
      as=mu*dt/(rh*dz**2)+Vz(i)*dt/(2*dz)
      ap=2*mu*dt/(rh*dr**2)+2*mu*dt/(rh*dz**2)+mu*dt/(rh*ri(ll)**2)-Vr(i)*dt/ri(ll) -1
      ae=mu*dt/(rh*dr**2)+mu*dt/(2*rh*ri(ll)*dr)-Vr(i)*dt/(2*dr)
      an=mu*dt/(rh*dz**2)-Vz(i)*dt/(2*dz)

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k2=(Vr(i+nr)/dz)-((Vz(i+1)-Vz(i-1))/(2*dr))
k3=Vz(i-1)/(2*dr)
!The 9 if-statements below are for the 9 regions of our CV formulation
if ( i==1 ) then !bottom left region
    oma(i)=0 !symmetry condition
    ll=ll+1
else if ( i < nr ) then !bottom region
    oma(i)=k2
    ll=ll+1
else if ( i==nr ) then !bottom right region
    oma(i)=k3
    ll=1
else if ( i==nr*nz-(nr-1) ) then !top left region
    oma(i)=0 !symmetry condition
    ll=ll+1
else if ( i==nr*nz ) then !top right region
    oma(i)=k3 !wall condition
else if ( i==j ) then !left region
    oma(i)=0 !symmetry condition
    j=j+nr
    ll=ll+1
else if ( i==k ) then !right region
    oma(i)=k3 !wall condition
    k=k+nr
    ll=1
else if ( i < (nr*nz)-nr ) then !Interior region
    oma(i)=aw*omo(i-1)+as*omo(i-nr)-ap*omo(i)+ae*omo(i+1)+an*omo(i+nr)
    ll=ll+1
else !top region
    oma(i)=oma(i-nr) !outflow condition
    ll=ll+1
end if
end do
!print *, 'Oma=', oma(1:nr*nz)

```

```

!The below Loop composes the radial velocity at the boundary
j=nr+1
k=nr+nr
ll=1
do i = 1,nr*nz
    !The 9 if-statements below are for the 9 regions of our formulation
    if ( i==1 ) then !bottom left region
        Vr(i)=0 !inflow or symmetry condition
        b(i)=(ri(ll)**2)*(0-0)/(2*dz)
        ll=ll+1
    else if ( i < nr ) then !bottom region

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    Vr(i)=0 !inflow condition
    b(i)=(ri(ll)**2)*(oma(i+nr)+((Vz(i+1)-Vz(i-1))/(2*dr)))/(2*dz)
    ll=ll+1
else if ( i==nr ) then          !bottom right region
    Vr(i)=0 !wall conditions
    b(i)=(ri(ll)**2)*(oma(i+nr)-(Vz(i-1)/(2*dr)))/(2*dz)
    ll=1
else if ( i==nr*nz-(nr-1) ) then !top left region
    Vr(i)=0 !symmetry condition
    b(i)=ri(ll)**2*(0-0)/(2*dz)
    ll=ll+1
else if ( i==nr*nz ) then      !top right region
    Vr(i)=0
    b(i)=(ri(ll)**2)*(oma(i)-oma(i-nr-nr))/(2*dz)
else if ( i==j ) then          !left region
    Vr(i)=0 !symmetry conditions
    b(i)=ri(ll)**2*(0-0)/(2*dz)
    j=j+nr
    ll=ll+1
else if ( i==k ) then          !right region
    Vr(i)=0 !wall conditions
    b(i)=(ri(ll)**2)*(oma(i+nr)-oma(i-nr))/(2*dz)
    k=k+nr
    ll=1
else if ( i < (nr*nz)-nr ) then !Interior region
    b(i)=(ri(ll)**2)*(oma(i+nr)-oma(i-nr))/(2*dz)
    ll=ll+1
else
    !top region
    b(i)=(ri(ll)**2)*(oma(i)-oma(i-nr-nr))/(2*dz)
    ll=ll+1
end if
end do
nj=nnr+1; nk=nnr+nnr
j=nr+1; k=1
ll=2
do i = 1,nnr*nnz !This loop composes the interior radial velocity pentadiagonal vectors
    aw=(ri(ll)**2)/(dr**2) - ri(ll)/(2*dr)
    aa=(ri(ll)**2)/(dz**2)
    ap=(2*ri(ll)**2)/(dr**2) + (2*ri(ll)**2)/(dz**2) + 1
    ae=(ri(ll)**2)/(dr**2) + ri(ll)/(2*dr)
    !The 9 if-statements below are for the 9 regions of our formulation
    if ( i==1 ) then !bottom left region
        f(i)=-ap
        g(i+1)=ae
        h(i+nnr)=aa
        bi(i)=b(j+k)

```

```

    k=k+1
    ll=ll+1
else if ( i < nnr ) then          !bottom region
    e(i-1)=aw
    f(i)=-ap
    g(i+1)=ae
    h(i+nnr)=aa
    bi(i)=b(j+k)
    k=k+1
    ll=ll+1
else if ( i==nnr ) then          !bottom right region
    e(i-1)=aw
    f(i)=-ap
    h(i+nnr)=aa
    bi(i)=b(j+k)
    j=j+nr
    k=1
    ll=2
else if ( i==nnr*nnz-(nnr-1) ) then !top left region
    d(i-nnr)=aa
    f(i)=-ap+aa
    g(i+1)=ae
    bi(i)=b(j+k)
    k=k+1
    ll=ll+1
else if ( i==nnr*nnz ) then      !top right region
    d(i-nnr)=aa
    e(i-1)=aw
    f(i)=-ap+aa
    bi(i)=b(j+k)
else if ( i==nj ) then           !left region
    d(i-nnr)=aa
    f(i)=-ap
    g(i+1)=ae
    h(i+nnr)=aa
    bi(i)=b(j+k)
    nj=nj+nnr
    k=k+1
    ll=ll+1
else if ( i==nk ) then           !right region
    d(i-nnr)=aa
    e(i-1)=aw
    f(i)=-ap
    h(i+nnr)=aa
    bi(i)=b(j+k)
    nk=nk+nnr

```

```

        j=j+nr
        k=1
        ll=2
    else if ( i < (nnr*nnz)-nnr ) then      !Interior region
        d(i-nnr)=aa
        e(i-1)=aw
        f(i)=-ap
        g(i+1)=ae
        h(i+nnr)=aa
        bi(i)=b(j+k)
        k=k+1
        ll=ll+1
    else                                     !top region
        d(i-nnr)=aa
        e(i-1)=aw
        f(i)=-ap+aa
        g(i+1)=ae
        bi(i)=b(j+k)
        k=k+1
        ll=ll+1
    end if
end do
!For computational efficiency, the, A penta-diagonal matrix is split to 5 vectors
call Bi_CGSTAB_P(d,e,f,g,h,bi,Vri) !Solves the linear system at each time-step
j=nr+1
k=nr+nr
ll=1
do i = 1,nr*nz!This loop adds the Bi_CGSTAB solved interior to the boundary radial velocity

    if ( i>nr+1 .and. i/=j .and. i/=k ) then
        Vr(i)=Vri(ll)
        ll=ll+1
    else if ( i==nr*nz-(nr-1) ) then !interior region
        ll=nnr*nnz-(nnr-1)
    else if ( i==j ) then             !left region
        j=j+nr
    else if ( i==k ) then             !right region
        k=k+nr
    end if
end do
!print "(a5,25f10.7)", 'Vr=', Vr(1:nr*nz)

!The below loop solves the axial velocity at each grid point
j=nr+1
k=nr+nr
ll=1

```



```

do i = 1,nr*nz
  aa=dz/(2*dr)
  ap=dz/ri(ll)
  k1=CLVZ-CLVZ*ri(ll)
  !The 9 if-statements below are for the 9 regions of our CV formulation
  if ( i==1 ) then !bottom left region
    Vz(i)=k1 !inflow condition
    ll=ll+1
  else if ( i < nr ) then !bottom region
    Vz(i)=k1
    ll=ll+1
  else if ( i==nr ) then !bottom right region
    Vz(i)=0
    ll=1
  else if ( i==nr*nz-(nr-1) ) then !top left region
    Vz(i)=Vz(i-nr) !symmetry condition
    ll=ll+1
  else if ( i==nr*nz ) then !top right region
    Vz(i)=0 !wall conditions
  else if ( i==j ) then !left region
    Vz(i)=Vz(i-nr) !symmetry conditions
    j=j+nr
    ll=ll+1
  else if ( i==k ) then !right region
    Vz(i)=0 !wall conditions
    k=k+nr
    ll=1
  else if ( i < nr*nz-nr ) then !Interior region
    Vz(i)=Vz(i-nr)-ap*Vr(i)-aa*Vr(i+1)+aa*Vr(i-1)
    ll=ll+1
  else !top region
    Vz(i)=Vz(i-nr) !outflow condition
    ll=ll+1
  end if
end do
!print "(a5,25f8.5)", 'Vz=', Vz(1:nr*nz)
k=0
do j = 1, nz !this loop solves the mass flow-rate along each z-gridpoint
  mdot(j) = 0.d0
  do i = 1, nr
    k=k+1
    mdot(j) = mdot(j) + 2.d0*pi*rh*Vz(k)*ri(i)
  end do
end do

do i = 1,Nz

```

```

        som(i)=abs(mdot(i)-mdot(1))
    end do
    !print "(a5,25f8.3)", 'som=', som(1:nz)
    !print "(a5,25f8.3)", 'mdot=', mdot(1:nz)
    Mmdot=maxval(som)/mdot(1)
    Amdot=(sum(som)/nz)/mdot(1)

    !for printing results
    print "(i9,E12.3,E12.3)", ij, Mmdot, Amdot
    if (ij == 1) then
        Mmdot=1
        iii=101
    elseif (ij <= 100) then
        write(1,*) ij, Mmdot, Amdot
    elseif (ij == iii) then
        print *, achar(27)//"[2J" !clears the console
        iii=iii+5
    elseif (ij >= 3D6) then !set the maximum number of iteration
        Mmdot=tols
    end if

    !dt=dt+t !Adaptive time-step
    ij=ij+1
end do

close(1)
end subroutine unsteady
!*****
!*****
subroutine Bi_CGSTAB_P(d,e,f,g,h,b,x)
    implicit none
    DOUBLE PRECISION, DIMENSION(nnr*nnz), INTENT(OUT) :: x
    DOUBLE PRECISION, dimension(nnr*nnz), INTENT(IN):: d, e, f, g, h, b
    DOUBLE PRECISION, DIMENSION(nnr*nnz) :: d1,e1,g1,h1,d2,e2,g2,h2
    DOUBLE PRECISION, DIMENSION(nnr*nnz) :: p, r, r0, y, z, v, s, t, k, ks, kt
    DOUBLE PRECISION :: rho0, rho, w, alpha, beta, r_check, tol, nan
    INTEGER :: i
    !!This subroutine solves a penta-diagonal linear system using
    !!The Preconditioned Bi_CGSTAB Algorithm by Van Der Vorst

    !Variable initialization
    d2=0; e2=0; g2=0; h2=0
    !since we are dealing with vectors, I have done the multiplication of two vectors,
    !using this pattern
    !this multiples each vectors with x, for A*x
    d1=d(1:nnr*nnz)*x(1:nnr*nnz);e1=e(1:nnr*nnz)*x(1:nnr*nnz);

```

```

g1=g(1:nnr*nnz)*x(1:nnr*nnz);h1=h(1:nnr*nnz)*x(1:nnr*nnz);
!This circularly shifts the vectors to the appropriate position, before addition
d2(nnr+1:nnr*nnz)=d1(1:nnr*nnz-nnr);e2(2:nnr*nnz)=e1(1:nnr*nnz-1);
g2(1:nnr*nnz-1)=g1(2:nnr*nnz)
h2(1:nnr*nnz-nnr)=h1(nnr+1:nnr*nnz)
!Hence  $A*x = d2+e2+f(1:nnr*nnz)*x(1:nnr*nnz)+g2+h2$ 
r0=b-(d2+e2+f(1:nnr*nnz)*x(1:nnr*nnz)+g2+h2)
r=r0
w=1; alpha=1; rho0=1; r_check=1
v=0; p=0; k=0
tol=10E-8
!For the inverse of K
do i = 1, nnr*nnz
    k(i)=1/f(i)
end do
i=0
!main algorithm loop
do while (r_check>tol)
    i=i+1
    rho=dot_product(r0,r)
    beta=(rho/rho0)*(alpha/w)
    rho0=dot_product(r0,r)
    p=r+beta*(p-w*v)
    y=k(1:nnr*nnz)*p(1:nnr*nnz)
    d1=d(1:nnr*nnz)*y(1:nnr*nnz);e1=e(1:nnr*nnz)*y(1:nnr*nnz);
    g1=g(1:nnr*nnz)*y(1:nnr*nnz);h1=h(1:nnr*nnz)*y(1:nnr*nnz);
    d2(nnr+1:nnr*nnz)=d1(1:nnr*nnz-nnr);e2(2:nnr*nnz)=e1(1:nnr*nnz-1);
    g2(1:nnr*nnz-1)=g1(2:nnr*nnz)
    h2(1:nnr*nnz-nnr)=h1(nnr+1:nnr*nnz)
    v=(d2+e2+f(1:nnr*nnz)*y(1:nnr*nnz)+g2+h2)
    alpha=rho/dot_product(r0,v)
    if (rho==0) then
        alpha=0
    end if
    s=r-alpha*v
    z=k(1:nnr*nnz)*s(1:nnr*nnz)
    d1=d(1:nnr*nnz)*z(1:nnr*nnz);e1=e(1:nnr*nnz)*z(1:nnr*nnz);
    g1=g(1:nnr*nnz)*z(1:nnr*nnz);h1=h(1:nnr*nnz)*z(1:nnr*nnz);
    d2(nnr+1:nnr*nnz)=d1(1:nnr*nnz-nnr);e2(2:nnr*nnz)=e1(1:nnr*nnz-1);
    g2(1:nnr*nnz-1)=g1(2:nnr*nnz)
    h2(1:nnr*nnz-nnr)=h1(nnr+1:nnr*nnz)
    t=(d2+e2+f(1:nnr*nnz)*z(1:nnr*nnz)+g2+h2)
    d1=d(1:nnr*nnz)*s(1:nnr*nnz);e1=e(1:nnr*nnz)*s(1:nnr*nnz);
    g1=g(1:nnr*nnz)*s(1:nnr*nnz);h1=h(1:nnr*nnz)*s(1:nnr*nnz);
    d2(nnr+1:nnr*nnz)=d1(1:nnr*nnz-nnr);e2(2:nnr*nnz)=e1(1:nnr*nnz-1);
    g2(1:nnr*nnz-1)=g1(2:nnr*nnz)

```

```

h2(1:nnr*nnz-nnr)=h1(nnr+1:nnr*nnz)
ks=(d2+e2+f(1:nnr*nnz)*s(1:nnr*nnz)+g2+h2)
d1=d(1:nnr*nnz)*t(1:nnr*nnz);e1=e(1:nnr*nnz)*t(1:nnr*nnz);
  g1=g(1:nnr*nnz)*t(1:nnr*nnz);h1=h(1:nnr*nnz)*t(1:nnr*nnz);
d2(nnr+1:nnr*nnz)=d1(1:nnr*nnz-nnr);e2(2:nnr*nnz)=e1(1:nnr*nnz-1);
  g2(1:nnr*nnz-1)=g1(2:nnr*nnz)
h2(1:nnr*nnz-nnr)=h1(nnr+1:nnr*nnz)
kt=(d2+e2+f(1:nnr*nnz)*t(1:nnr*nnz)+g2+h2)
nan=dot_product(kt,ks)
w=nan/dot_product(kt,kt)
if (nan==0) then
  w=0
end if
x=x+alpha*y+w*z
r=s-w*t
r_check=norm2(r)
end do
end subroutine Bi_CGSTAB_P
!*****

end program unsteady_Dfinite

```

### Q1 b) Cont'd!

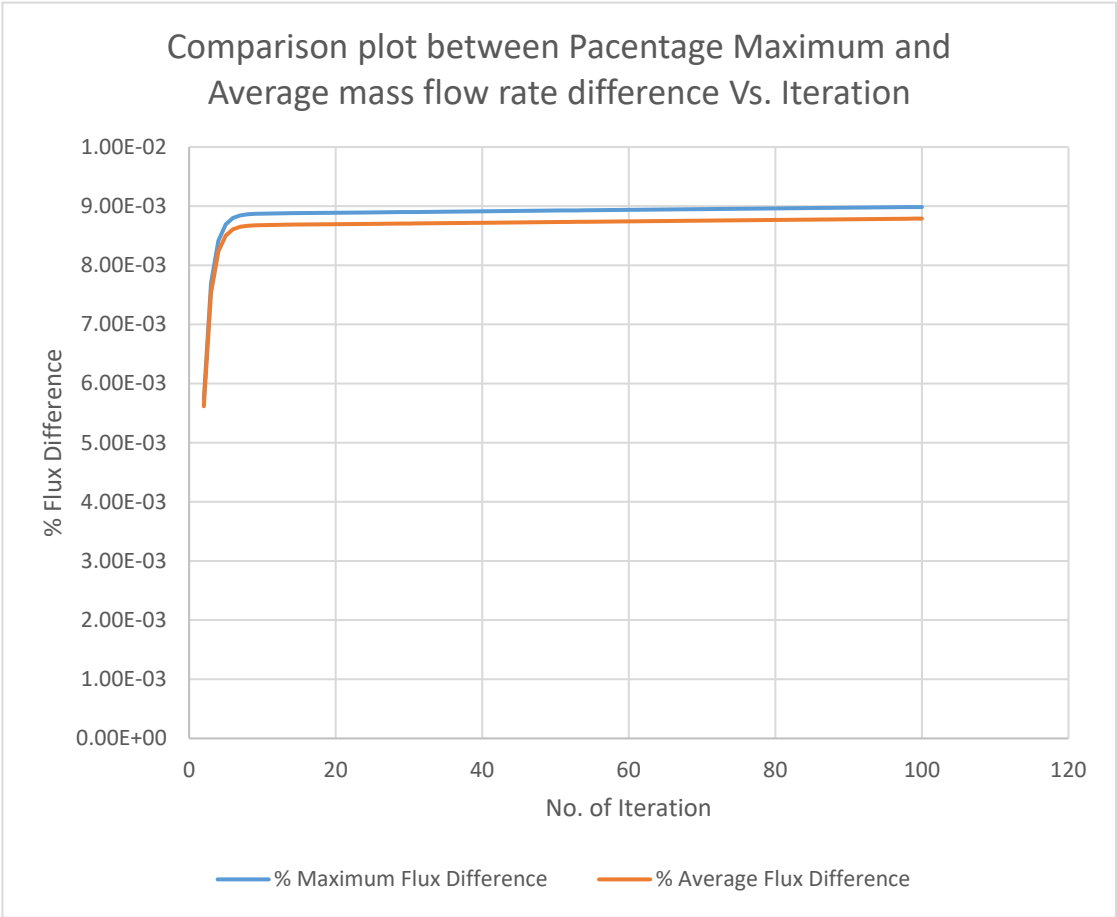
>> For computation time, I observed that slight adjustment to the tolerance value ~~greatly~~ affects the computation time. ~~For instance~~ if adaptive timestep is not used. Even after 7 hours of simulation on a constant timestep ( $\Delta t$ ), the  $\text{tol} = 10^{-5}$  is not still reached. But I was able to simulate for 3 million iterations at just 3.088 hours using vector only Bi-CGSTAB instead of the conventional matrix type. This improved my computational time greatly.

>> The accuracy which is a direct function of our Percentage maximum/average flow rate difference (i.e based on ~~the~~ Conservation law), will be as accurate as we want it to be, in the expense of our computational time.

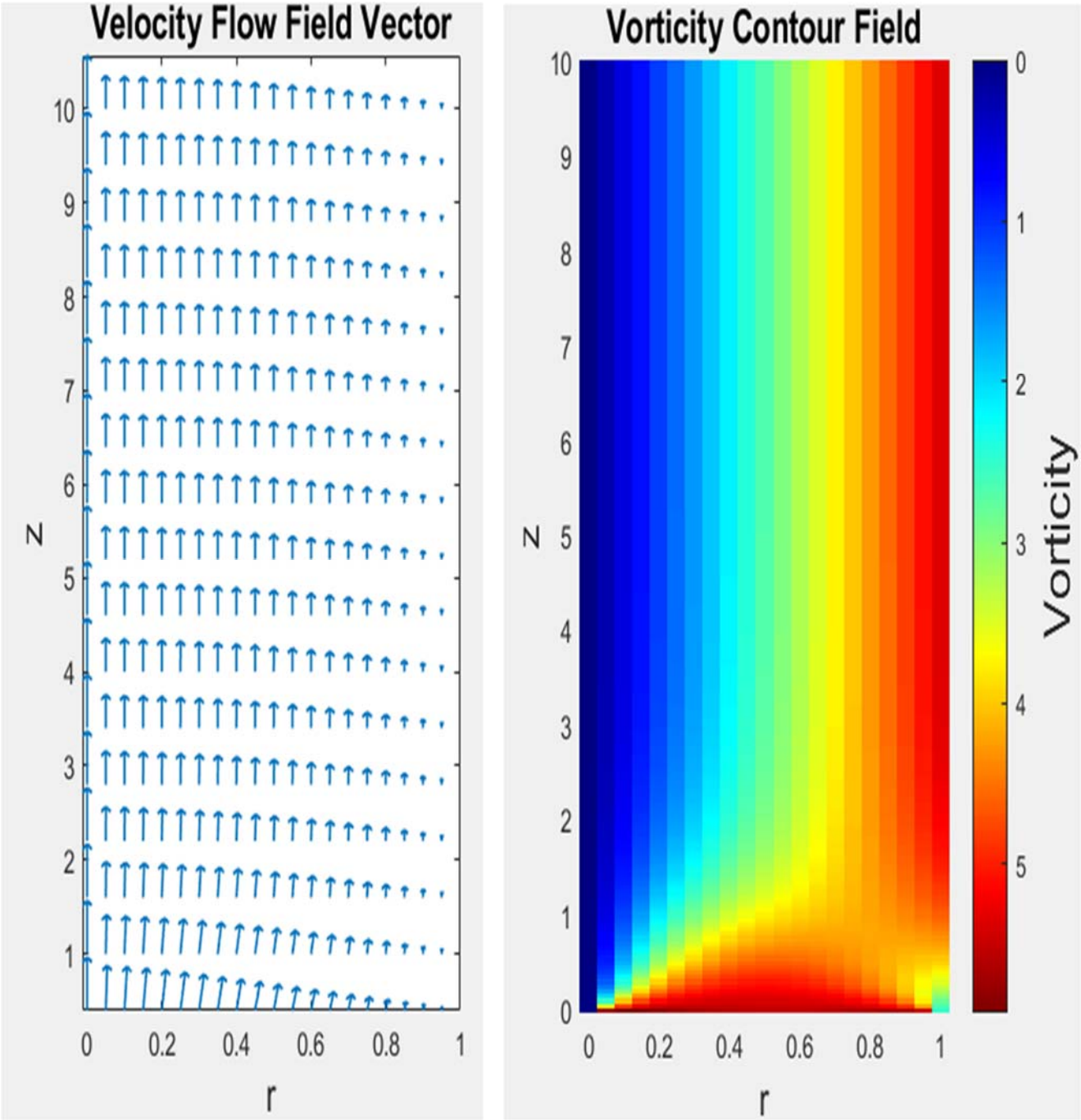
But what I've been able to deduce from our simulation, is that the solution gets ~~more~~ slightly inaccurate as we ~~increas~~ at new iterations upto about 50,000 iterations. And then it starts to get accurate again and on-ward in expense of computation time.

Q1C  $\Rightarrow$  From below plot, we can see that both percentage Maximum and average flow rate difference have similar behaviour against the number of iteration. And thus either one of them can be used to check for the accuracy of our solution. (This was taken from samples of just 100 iterations).

$\Rightarrow$  And setting our tolerance to  $\text{tol} = 8e^{-3}$  is small enough for us to understand the behaviour of the solution (ie. after it begins to converge again, from about 50,000 iterations). With this we got about 3 million iterations. And the vector plot for the flow field, as well as the contour plot for the vorticity is shown below;



Q1c] cont'd.]



Computational Time = 3.0881987847222221 hours



Q1d | If we assume an initial guess of;

$$v_r = 0, \quad v_z = CLVZ \left( 1 - \left( \frac{r_{\text{guess}}}{R} \right)^2 \right), \quad \omega = 2(CL VZ) \left( \frac{r_{\text{guess}}}{R^2} \right)$$

lets assume  $R$  to be our overall radius  $R = 1\text{cm}$   
and let's play with  $0 \leq r \leq 1\text{cm}$  (we can call this  $r_{\text{guess}}$ ).

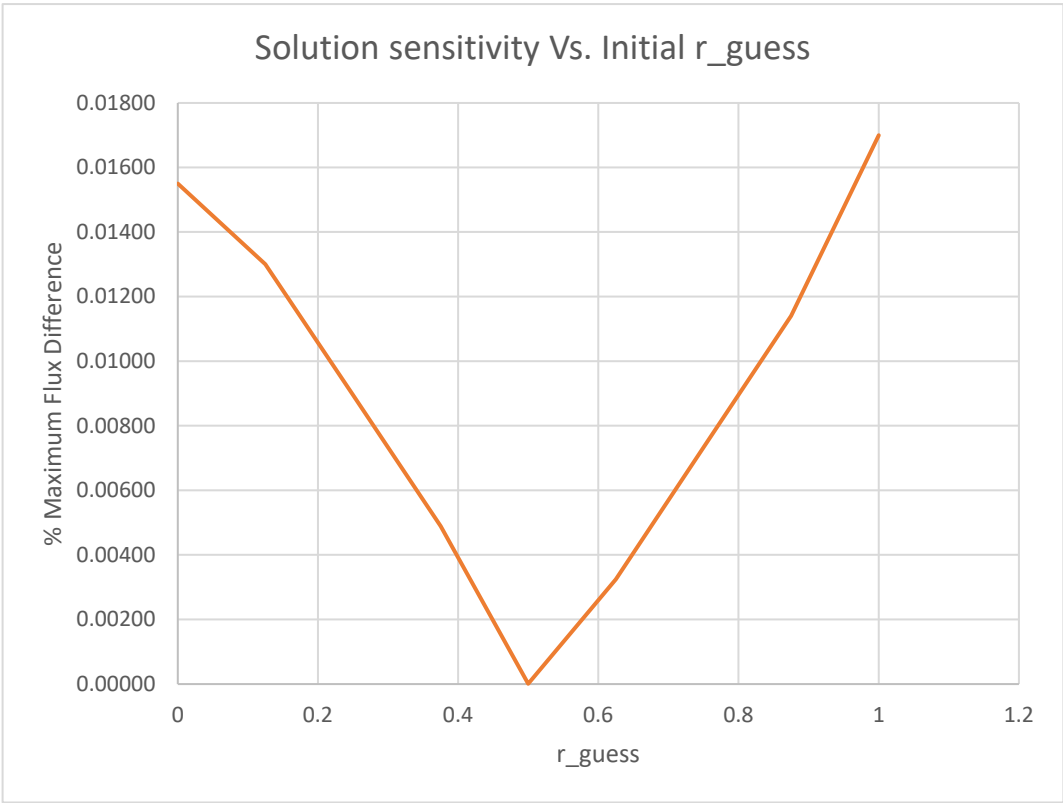
>> The sensitivity of varying this  $r_{\text{guess}}$  is shown in the table below. Which at  $r_{\text{guess}} = 0.5$  we can see the solution converging very at just 9 iteration.

>> But the strange part is with the vector and contour plot which is shown below for  $r_{\text{guess}} = 0.5$  and it looks much different from the solution if  $v_r = 0, v_z = 0 \ \& \ \omega = 0$ .

I also observed that  $r_{\text{guess}} = 0$  yields similar results to our previous guess of  $v_r = v_z = \omega = 0$ .

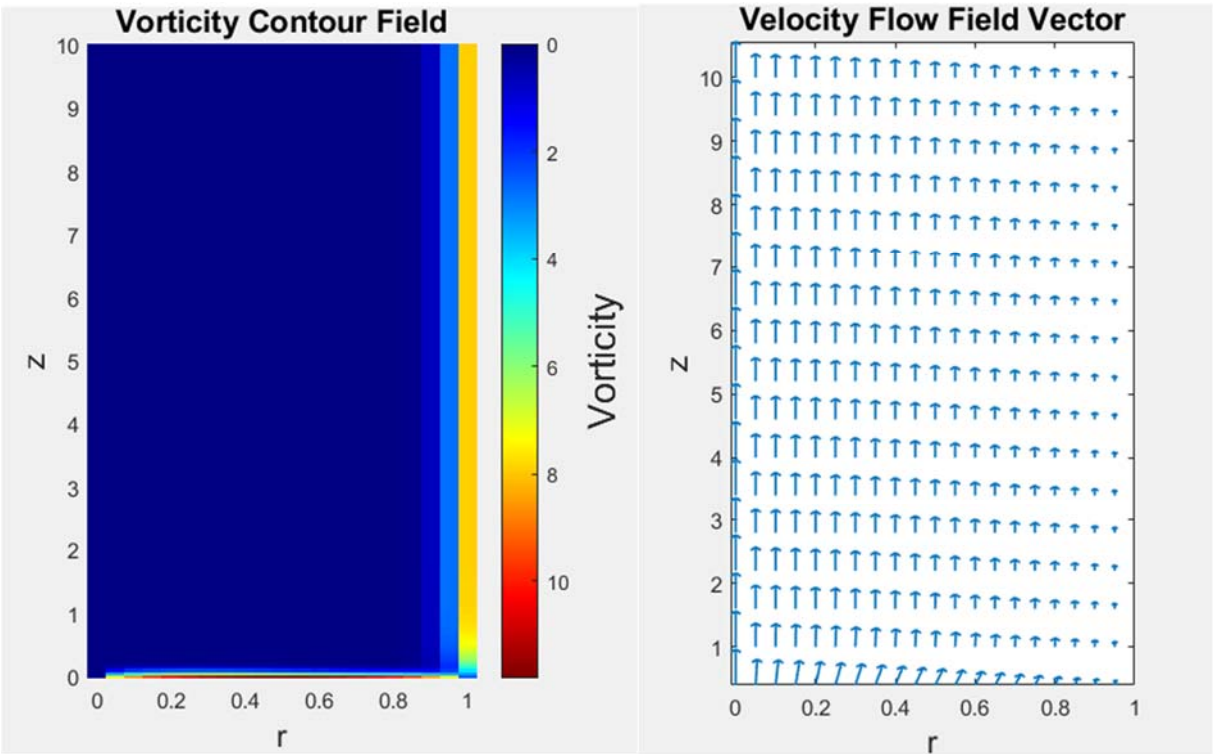
Q1d] cont'd.]

No. of Iteration	r_guess	% Maximum Flux Difference
10000	1	0.01550
10000	0.125	0.01300
10000	0.25	0.00895
10000	0.375	0.00489
9	0.5	0.00001
10000	0.625	0.00324
10000	0.75	0.00732
10000	0.875	0.01140
10000	1	0.01700

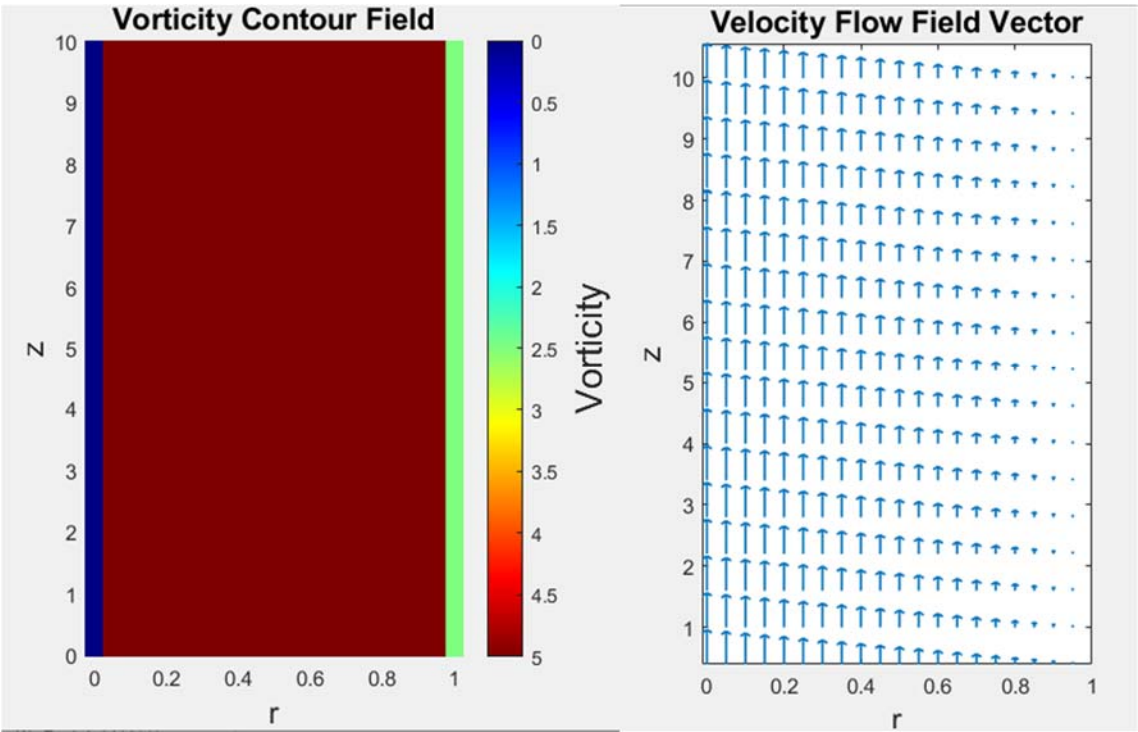


Q1d] cont'd.]

Plot when  $r_{guess} = 0$  :



Plot when  $r_{guess} = 0.5$  :



Q1d] cont'd.]

Plot when  $r_{guess} = 1$  :

