

# Swarming dynamics under different orientation (chiral) coupling mechanisms

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Synchronization and swarming are two typical self-organization behaviors representing the intrinsic dynamical and spatial ordering of units in complex systems. Here, we introduce a model that combines the Kuramoto and Vicsek models, which describes the synchronization and swarming of oscillators with different natural frequencies, chiral properties, and local coupling. The model exhibits five states: Disorder, Ring, Swarm (Multi-Clusters), Swarm (Double-Clusters), and Quick Sync. We obtain the analytic solutions of the model, explore the phase diagram of it and analytically explain the diverse emergence behaviors driven by the model. This study provides a more comprehensive exploration and integration of the interdisciplinary field encompassing synchronization dynamics and cluster dynamics.

## I. INTRODUCTION

Synchronization and swarming are two typical self-organization behaviors representing the intrinsic dynamical and spatial ordering of units in complex systems. The study of self-organized coordination of intrinsic degrees of freedom/rhythms within complex systems has deepened our understanding of self-organization in complex systems, including various levels of synchronization, partial synchronization, extended states, and chimera states<sup>1–6</sup>. In terms of synchronization theory, at the micro level, various analysis methods have been established, such as the master stability function (MSF)<sup>7</sup>, which has played an important role in the study of multi-oscillator synchronization problems. At the statistical and macro levels, successful approaches include the self-consistent equation method by Kuramoto et al.<sup>8</sup>, the Watanabe-Strogatz transformation<sup>9,10</sup>, and the Ott-Antonsen ansatz proposed by Ott and Antonsen<sup>11</sup>. Cluster behavior research focuses on the self-organized coordination of spatial movement in complex systems, contributing to our understanding of collective behavior. Spatial cluster behavior can emerge at different scales, such as animal clusters<sup>12–14</sup>, bacterial directed motion<sup>15,16</sup>, and other phenomena<sup>17–19</sup>. Numerous studies based on diverse cluster models and extensive experimental investigations<sup>20–26</sup> have revealed various collective spatial interactions and patterns in complex systems.

The Kuramoto model<sup>27</sup> and the Vicsek model<sup>20</sup> are two well-known models that describe synchronization and swarming, respectively. The Kuramoto model describes the synchronization of oscillators with a phase variable, and the Vicsek model describes the alignment of self-propelled particles with a velocity variable. Both models have been widely studied and applied in various fields, such as physics<sup>28,29</sup>, biology<sup>30–33</sup>, and engineering<sup>34–36</sup>. In recent years, there has been a notable surge in the research progress concerning the interplay between synchronization dynamics and cluster dynamics<sup>37–41</sup>,

which has led to the development of models that combine the Kuramoto and Vicsek models.

This paper continues this line of research by proposing a model that combines the Kuramoto and Vicsek models, aiming to undertake a more comprehensive exploration and integration of the interdisciplinary field encompassing synchronization dynamics and cluster dynamics. In Ceron et al.'s<sup>40</sup> and Levis et al.'s<sup>39</sup> studies, the behavior of oscillators which are non-identical (oscillators have different natural frequencies), chiral (oscillators have inherent clockwise and counter-clockwise circular orbits), and locally coupled (oscillators can only couple to motion and phase of neighbors within a given radius) was numerically investigated. Here, we introduce a model that combines these three characteristics and study the synchronization and cluster behaviors of it. We obtain the analytic solutions of the model and explore the phase diagram.

The paper is organized as follows. In Section II, we introduce the model and the numerical methods used to study it. In Section III, we present the behavior of the model and discuss the states it exhibits. In Section IV, we introduce order parameters to distinguish the states and present the phase diagram of the model. In Section V, we derive the critical lines of the transitions between states and analytically explain diverse emergence behaviors driven by the model. Finally, in Section VI, we summarize the results and discuss the implications of the model.

## II. MODEL

Oscillators have a spatial position  $\mathbf{r}_i = (x_i, y_i)$  and an internal phase  $\theta_i$  which evolve according to equations:

$$\dot{x}_i = v \cos \theta_i , \quad (1)$$

$$\dot{y}_i = v \sin \theta_i , \quad (2)$$

$$\dot{\theta}_i = \omega_i + \lambda \sum_{j=1}^N A_{ij} \sin(\theta_j - \theta_i) \quad (3)$$

for  $i = 1, 2, \dots, N$ , where  $N$  is the number of oscillators. As per Eq. (1) and (2), each oscillator moves with a constant speed  $v$

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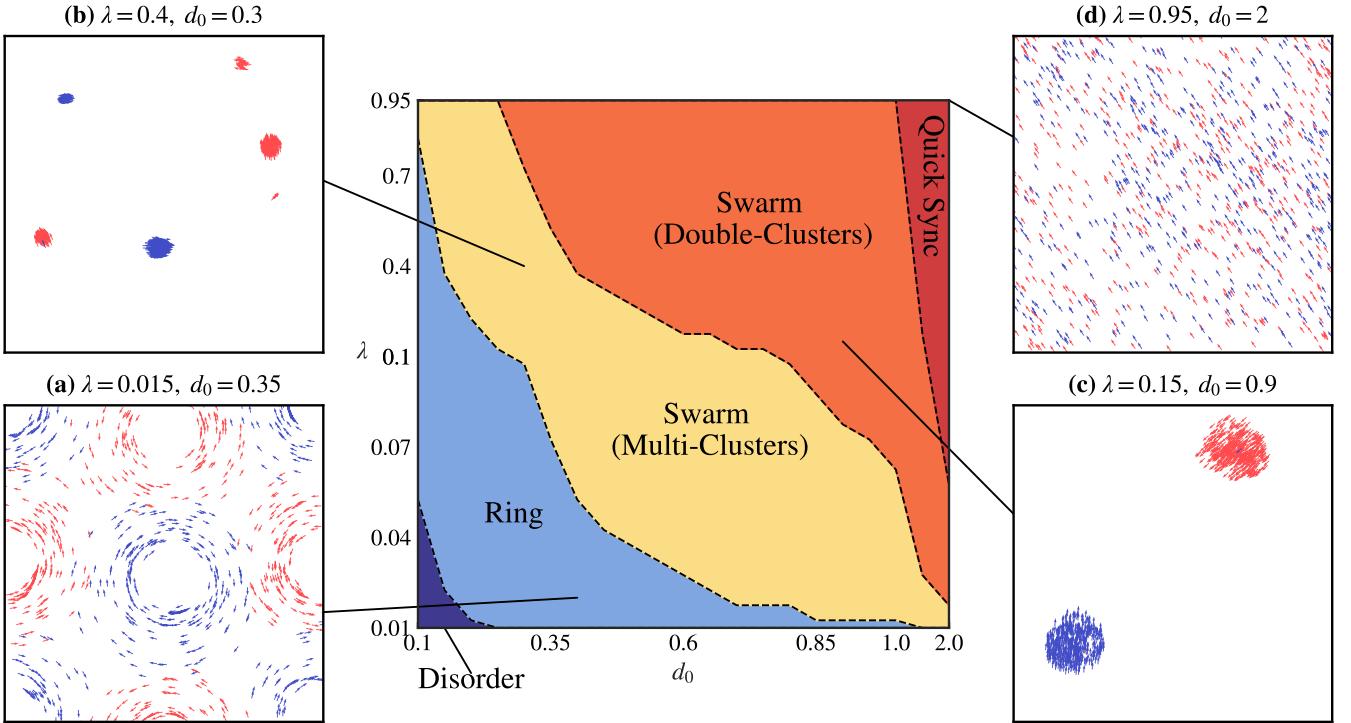


FIG. 1. Phase diagram of model Eq. (1)-(3) in the  $(\lambda-d_0)$  plane. The boundaries between states are analytical approximations given by Subsection V A. For the sake of clarity, the scale of  $\lambda$  and  $d_0$  is non-linear (For  $\lambda$  in  $[0.01, 0.1]$  and  $[0.1, 1]$ , step length is 0.1 and 0.05, respectively. For  $d_0$  in  $[0.1, 1]$  and  $[1, 2]$ , step length is 0.05 and 0.5, respectively). (a), The snapshots of Ring ( $\lambda = 0.015, d_0 = 0.35$ ). (b), Swarm (Multi-Clusters) ( $\lambda = 0.4, d_0 = 0.3$ ). (c), Swarm (Double-Clusters) ( $\lambda = 0.15, d_0 = 0.9$ ). (d), Quick Sync ( $\lambda = 0.95, d_0 = 2$ ). Two types of chiral oscillators are represented by red ( $\omega_i > 0$ ) and blue ( $\omega_i < 0$ ) arrows, respectively.

in the direction of its current phase  $\theta_i$ . The phase  $\theta_i$  evolves according to Eq. (3), where  $\omega_i$  is the natural frequency of the  $i$ th oscillator,  $\lambda$  is the coupling strength, and  $A$  is the adjacency matrix of the network, with  $A_{ij} = 1$  if there is a connection from  $i$ th to  $j$ th oscillator, and  $A_{ij} = 0$  otherwise. We can consider Eq. (1)-(3) as a generalization of the Kuramoto model and the Vicsek model in the sense that it includes both the phase and the spatial position of the oscillators.

Each oscillator  $i$  is connected to all the oscillators within a action radius  $d_0$  of its position. The adjacency matrix  $A$  is defined as:

$$A_{ij} = \begin{cases} 1, & |\mathbf{r}_i - \mathbf{r}_j| \leq d_0 \\ 0, & |\mathbf{r}_i - \mathbf{r}_j| > d_0 \end{cases} \quad (4)$$

where  $|\mathbf{r}_i - \mathbf{r}_j|$  is the Euclidean distance between oscillators.

For simplicity, we consider oscillators are initially distributed uniformly in a two-dimensional square with side length  $L$  and periodic boundary conditions. Their positions  $\mathbf{r}_i(t) = (x_i(t), y_i(t))$  at given time  $t$  are given by:

$$\begin{aligned} x_i(t + \Delta t) &= x_i(t) + v \cos \theta_i(t) \Delta t \bmod L, \\ x_i(t + \Delta t) &= x_i(t) + v \cos \theta_i(t) \Delta t \bmod L, \end{aligned} \quad (5)$$

where  $\Delta t$  is the discrete time step. When two oscillators are on opposite sides of the square, the absolute value of the diff-

ference between one of their coordinates is greater than  $L/2$ . In this case, we take the minimum distance between them, which is the distance between the two points in the periodic boundary conditions. For a given pair of points  $\mathbf{r}_i$  and  $\mathbf{r}_j$ , the distance between them is  $|\mathbf{r}_i - \bar{\mathbf{r}}_j|$ , where  $\bar{\mathbf{r}}_j = (\bar{x}_j, \bar{y}_j)$  is the adjusted position of the  $j$ th oscillator, given by:

$$\bar{x}_j = \begin{cases} x_j, & |x_i - x_j| \leq L/2 \\ x_j + L, & x_i - x_j > L/2 \\ x_j - L, & x_j - x_i > L/2 \end{cases}, \quad (6)$$

$$\bar{y}_j = \begin{cases} y_j, & |y_i - y_j| \leq L/2 \\ y_j + L, & y_i - y_j > L/2 \\ y_j - L, & y_j - y_i > L/2 \end{cases}. \quad (7)$$

$|\mathbf{r}_i - \bar{\mathbf{r}}_j|$  can be proved to be the minimum distance between  $\mathbf{r}_i$  and  $\mathbf{r}_j$  in the periodic boundary conditions (see the proof in Appendix A).

Finally, we consider that the natural frequencies  $\omega_i$  are distributed in two symmetric uniform distributions, representing two types of chirality. Exactly half of the oscillators have natural frequencies in the range  $[\omega_{\min}, \omega_{\max}]$  ( $\omega_i \sim U(\omega_{\min}, \omega_{\max}), i = 1, 2, \dots, N/2$ ) and the other half in the range  $[-\omega_{\max}, -\omega_{\min}]$  ( $\omega_i \sim U(-\omega_{\max}, -\omega_{\min}), i = N/2 + 1, N/2 + 2, \dots, N$ ).

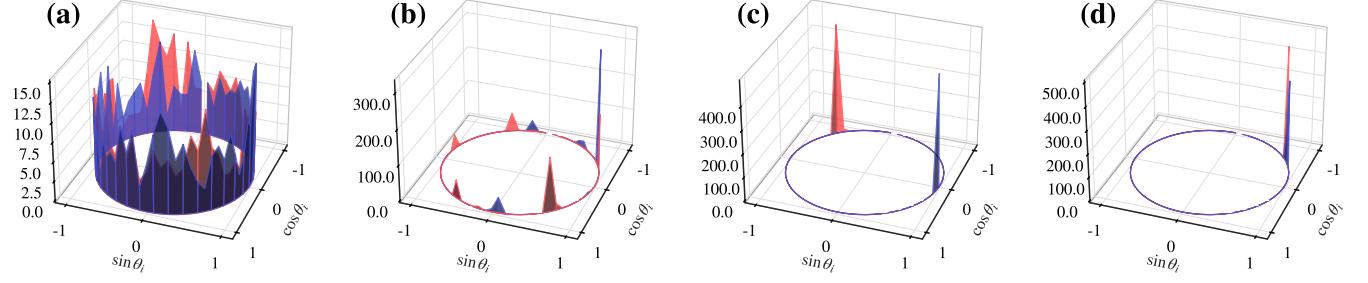


FIG. 2. Histogram of the oscillators' phases. (a), Ring state ( $\lambda = 0.015$ ,  $d_0 = 0.35$ ). (b), Swarm state (Multi-Clusters,  $\lambda = 0.8$ ,  $d_0 = 0.2$ ). (c), Swarm state (Double-Clusters,  $\lambda = 0.15$ ,  $d_0 = 0.9$ ). (d), Quick Sync state ( $\lambda = 0.95$ ,  $d_0 = 2$ ). The histograms are calculated with 70 bins.

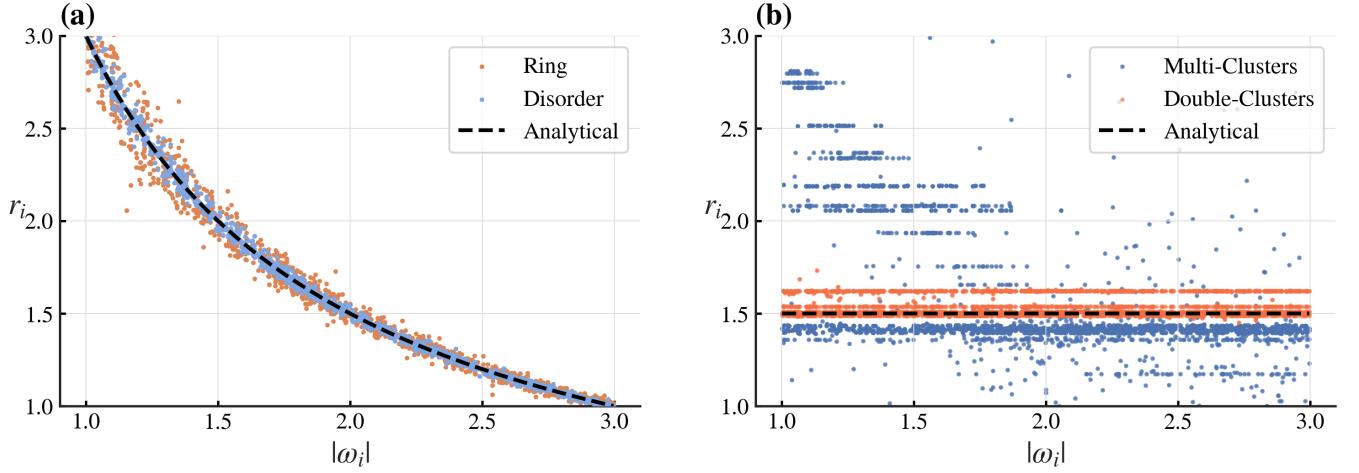


FIG. 3. The real-time and analytical rotational radius. (a), Radius for the Disorder ( $d_0 = 0.1$ ,  $\lambda = 0.01 : 0.06$ ) and Ring ( $d_0 = 0.1$ ,  $\lambda = 0.06 : 0.1$ ). The real-time rotational radius is almost constant and close to  $v/\omega_i$  for each oscillator. (b), Radius for Swarm (Multi-Clusters,  $d_0 = 0.15 : 0.25$ ,  $\lambda = 0.95$ ) and (Double-Clusters,  $d_0 = 2$ ,  $\lambda = 0.02 : 0.05$ ). Analytical line is only for Double-Clusters. All the above simulations are calculated at  $t = 60000$ .

### III. BEHAVIOR

We performed numerical simulations of the model to probe the behavior of its solutions (see Appendix B for details on the numerical methods).  $N = 1000$  oscillators were distributed uniformly in the square of length  $L = 10$  and their natural frequencies were distributed in the range  $[\omega_{\min}, \omega_{\max}] = [1, 3]$  and  $[-\omega_{\max}, -\omega_{\min}] = [-3, -1]$ . Two-parameter of coupling strength  $\lambda$  and action radius  $d_0$  are presented in the phase diagram in Fig. 1. We found the system settles into five states: **Disorder**, **Ring**, **Swarm** (which can be further divided into **Multi-Clusters** and **Double-Clusters**), and **Quick Sync**. In Fig. 1 we show the snapshots of the last three states and where these states are located in the phase diagram. The Disorder state is shown in Fig. 4a. We next discuss these five states.

#### A. Disorder State

Disorder state occurs when both  $\lambda$  and  $d_0$  are small. In this state, the oscillators are not asynchronous (phase histogram is

uniform, like Fig. 2a) and move in a way which similar to uncoupled oscillators ( $\lambda = 0$ ), as shown in Fig. 4a. According to Eq. (1)-(3), when  $\lambda = 0$ , the equations of oscillators' motion

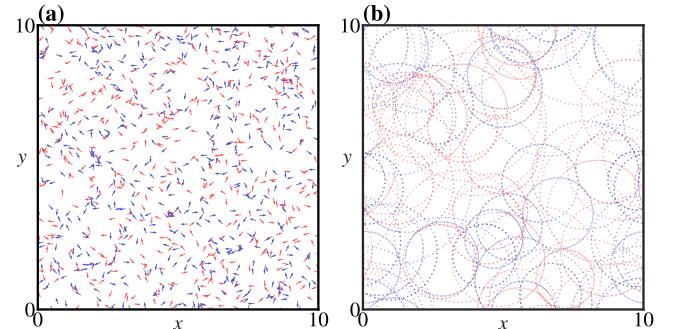


FIG. 4. Key properties of the Disorder state. (a), The snapshot of the Disorder state ( $\lambda = 0.01$ ,  $d_0 = 0.1$ ,  $T = 60000$ ). (b), The scatter plot of last 100 time steps of 20 positive chirality oscillators and 20 negative chirality oscillators.

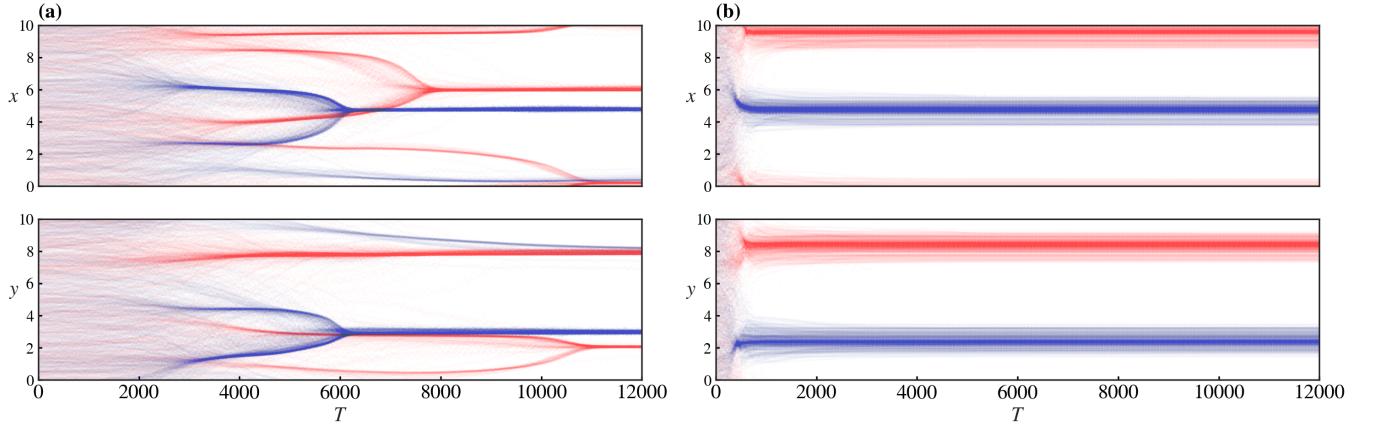


FIG. 5. Scatter plot of the real-time centers position. (a), centers position of Ring ( $\lambda = 0.02$ ,  $d_0 = 0.4$ ). As time goes on, the centers of oscillators with the same chirality converge. (b), centers position of Swarm ( $\lambda = 0.01$ ,  $d_0 = 2$ ). Unlike Ring, the centers converge quickly. The centers position are estimated with method in Fig. 6 and Eq. (11).

can be written as:

$$x_i(t) = x_i^0 + \frac{v}{\omega_i} \sin [\theta_i(0) + \omega_i t] , \quad (8)$$

$$y_i(t) = y_i^0 - \frac{v}{\omega_i} \cos [\theta_i(0) + \omega_i t] . \quad (9)$$

Then we have

$$(x_i - x_i^0)^2 + (y_i - y_i^0)^2 = \left( \frac{v}{\omega_i} \right)^2 . \quad (10)$$

In such a setup, oscillators move in a circular trajectory with radius  $v/\omega_i$  and the phases  $\theta_i$  increase linearly with time, as shown in Fig. 4b. To calculate the real-time rotational radius, we first estimate real-time centers  $\mathbf{c}_i(t)$  of the circular trajectory with method in Fig. 6 and then solve the following linear equations:

$$\begin{aligned} \mathbf{c}_i(t_1) \cdot \mathbf{v}_i(t_1) &= \mathbf{x}_i(t_1) \cdot \mathbf{v}_i(t_1) , \\ \mathbf{c}_i(t_2) \cdot \mathbf{v}_i(t_2) &= \mathbf{x}_i(t_2) \cdot \mathbf{v}_i(t_2) , \end{aligned} \quad (11)$$

where  $\mathbf{v}_i(t_1) = (x_i(t_1), y_i(t_1))$  is the velocity of  $i$ th oscillator at  $t_1$ , and  $\mathbf{v}_i(t_1) = (\cos \theta_i(t_1), \sin \theta_i(t_1))$  is the unit vector of

the velocity. According to Eq. (1)-(3), we can calculate  $\mathbf{v}_i(t_2)$  and  $\mathbf{r}_i(t_2)$ , ( $t_2 = t_1 + \Delta t$ ).

The real-time rotational radius is  $r_i(t) = |\mathbf{c}_i(t) - \mathbf{r}_i(t)|$ . We found that the real-time rotational radius is almost constant and close to  $v/\omega_i$  for each oscillator in the Disorder state, as shown in Fig. 3a. The estimation results of four states's real-time rotational centers are shown in Fig. 13 in Appendix.

## B. Ring State

The Ring state is characterized by the oscillators forming several rings with thickness, each of which is composed of oscillators with the same chirality, as is shown in Fig. 1a. Similar to Disorder state, the oscillators in the same ring cluster move in a circular trajectory with a constant rotational radius calculated in Fig. 3a. The oscillators' phase is uniformly distributed in the range  $[-\pi, \pi]$  (cf. Fig. 2a), which leads to oscillators uniformly located on the circular trajectory. Fig. 5a shows there is a long transient time before this state is reached, and in this transient time, the trajectories of oscillators with the same chirality aggregate slowly. Conversely, the oscillators with different chirality repel each other.

## C. Swarm State

Swarm State is a state where the oscillators form spatial clusters and align into several clusters [Fig. 1b, c and Fig. 5b]. When  $\lambda$  and  $d_0$  increases, the number of clusters decreases by 2, which is named by Double-Clusters state, and other states are named by Multi-Clusters state. The clusters are composed of oscillators with the same chirality, and the phase  $\theta_i$  of the oscillators in the same cluster is synchronized as seen in Fig. 2b and 2c, which means that the oscillators in the cluster move with the same velocity  $\mathbf{v}_i = (\cos \theta_s, \sin \theta_s)$  and rotational radius  $r_i = v/\theta_s$ , where  $\theta_s$  is the oscillators' phase in

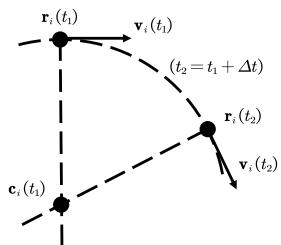


FIG. 6. Estimation for real-time centers.  $\mathbf{c}_i(t_j)$ ,  $\mathbf{r}_i(t_j)$  and  $\mathbf{v}_i(t_j)$  are the center of the circular trajectory, the position and the velocity of the  $i$ th oscillator at time  $t_j$ , respectively. The line from  $\mathbf{c}_i(t_j)$  to  $\mathbf{r}_i(t_j)$  is perpendicular to  $\mathbf{v}_i(t_j)$ .

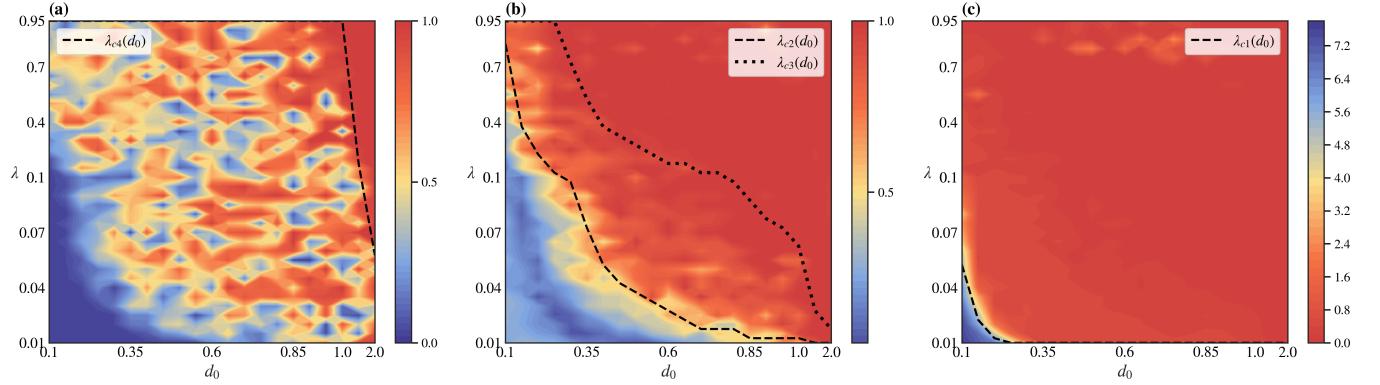


FIG. 7. Order parameter heatmaps of  $(\lambda, d_0)$  plane and the critical lines of the transitions between states. (a), order parameter  $R$  and critical line of  $\lambda_{c4}$ . (b), order parameter  $R_s$  and critical lines of  $\lambda_{c2}, \lambda_{c3}$ . (c), order parameter  $\Delta\Omega$  and critical lines of  $\lambda_{c1}$ . All order parameters are calculated at  $t = 60000$ .

the cluster. Based on this property, we can calculate  $\theta_s$  and  $r_i$  with Eq. (3):

$$\begin{aligned} N_s \omega_s &= \sum_{i=1}^{N_s} \left( \omega_i + \lambda \sum_{j=1}^{N_s} A_{ij} \sin(\theta_j - \theta_i) \right) \\ \omega_s &= \frac{1}{N_s} \sum_{i=1}^{N_s} \omega_i + \frac{\lambda}{N_s} \sum_{i=1}^{N_s} \sum_{j=1}^{N_s} A_{ij} \sin(\theta_j - \theta_i) \quad (12) \\ &= \frac{1}{N_s} \sum_{i=1}^{N_s} \omega_i, \end{aligned}$$

where  $N_s$  is the number of oscillators in the cluster. As  $\omega_i \sim U(\omega_{\min}, \omega_{\max})$  and  $\omega_i \sim U(-\omega_{\max}, -\omega_{\min})$  for two types of chirality, we can calculate  $\theta_i$ ,  $\omega_s$  and  $r_i$  with  $\omega_i$  for Double-Clusters state:

$$\begin{aligned} \theta_i &= \omega_s = \begin{cases} (\omega_{\max} + \omega_{\min})/2, & i = 1, 2, \dots, N/2 \\ -(\omega_{\max} + \omega_{\min})/2, & i = N/2 + 1, \dots, N \end{cases} \\ r_i &= \frac{v}{|\omega_s|}, \quad (13) \end{aligned}$$

as shown in Fig. 3b. But for Multi-Clusters, due to which oscillators are synchronized within each cluster is not accurately known, we can only calculate the real-time rotational radius of the them. As seen in Fig. 3b, similar to Double-Clusters, some local platforms appear in the real-time rotational radius due to synchronization.

#### D. Quick Sync State

Quick Sync state is a simple state where total oscillators are synchronized quickly, as shown in Fig. 1d. and 2d. The oscillators are synchronized in an extremely short time, which leads them have no time to form clusters (can also be considered as a special case of Swarm state). Due to the two types of chirality oscillators are synchronized and the distributions

of them is symmetric, the phase velocities of total oscillators are close to zero according to Eq. (12).

#### IV. ORDER PARAMETER

Having described the four states of our system, we next discuss how to distinguish them. We use the order parameter  $R$  to measure global synchronization. The order parameter  $R$  is defined as:

$$R = \left| \frac{1}{N} \sum_{j=1}^N e^{i\theta_j} \right|. \quad (14)$$

The order parameter  $R$  is the absolute value of the mean of the complex numbers  $e^{i\theta_j}$ , which can be interpreted as the mean direction of the oscillators. When  $R = 1$ , the oscillators are completely synchronized, and when  $R = 0$ , the oscillators are completely desynchronized. Fig 7a shows the order parameter  $R$  in the parameter plane. The order parameter  $R$  is close to 1 in the Quick Sync state, close to 0 in the Disorder state and most of the Ring state, and between 0 and 1 in other states. In these states, we see that the order parameter  $R$  changes non-monotonically in the sense that phases in these states are not globally synchronized, and each cluster's phase velocity  $\omega_s \neq 0$ . When the phases of different clusters are exactly equal, the order parameter  $R$  is close to 1, and when they are exactly opposite,  $R$  is close to 0.

Having realized that the order parameter  $R$  is not enough to distinguish the states with clusters, we next define the following order parameter  $R_s$  to metric the local synchronization,

$$R_s = \frac{1}{N_c} \sum_{k=1}^{N_c} \left| \frac{1}{|C_k|} \sum_{j \in C_k} e^{i\theta_j} \right|, \quad (15)$$

where  $N_c$  is the number of clusters,  $C_k$  is the  $k$ th cluster, and  $|C_k|$  is the number of oscillators in the  $k$ th cluster. We can

consider  $R_s$  as an order parameter that studies the spatial position and internal phase simultaneously. To determine the classification of clusters, we use the following method: we first calculate the relative center distance matrix  $D_{ij} = |\mathbf{c}_i - \bar{\mathbf{c}}_j|$ , where  $\bar{\mathbf{c}}_j = (\bar{x}_j, \bar{y}_j)$  is the adjusted position of the  $j$ th oscillator's rotational center calculated by Eq. (6), (7) and (11). The reason of using the distance between centers instead of the distance between oscillators' positions is that the oscillators in the Ring state are uniformly distributed on the circular trajectory, and the distance between them is much larger than the distance between their centers. Then we use the DBSCAN algorithm to cluster the oscillators. The DBSCAN algorithm is a density-based clustering algorithm, which can find clusters of arbitrary shapes and sizes. We set the minimum number of oscillators in a cluster to be 5 and the maximum distance between two oscillators in the same cluster to be 0.3 (see Appendix C for details on the determination of these parameters). One example of the classification of clusters is shown in Fig. 8. We then calculate the order parameter  $R_s$  for each cluster. The order parameter  $R_s$  is close to 1 in the Swarm state ( $R_s$  of Double-Clusters state is closer to 1 than Multi-Clusters) and Quick Sync state, and close to 0 in Disorder state and most of the Ring state, between 0 and 1 in other Ring states with local clusters, as shown in Fig. 7b.

Combining the order parameter  $R$  and  $R_s$ , we can find only the distinction between Ring and Disorder states has not been resolved. Except the study for synchronization, we also define an order parameter  $\Delta\Omega$  to metric the phase locking of the oscillators:

$$\Delta\Omega = \frac{1}{N_c} \sum_{k=1}^{N_c} \left[ \frac{1}{|C_k|^2} \sum_{i,j \in C_k} (\langle \dot{\theta}_i \rangle - \langle \dot{\theta}_j \rangle)^2 \right], \quad (16)$$

where  $\langle \dot{\theta}_i \rangle$  is the average of the phase velocity of the  $i$ th cluster, which can be calculated by

$$\langle \dot{\theta}_i \rangle = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{t_0}^{t_0+T} \dot{\theta}_i(t) dt. \quad (17)$$

We estimate  $\langle \dot{\theta}_i \rangle$  by the average of the phase velocity of the oscillators in the  $i$ th cluster at the last 1000 time steps. The order parameter  $\Delta\Omega > 0$  in the Disorder state, and  $\Delta\Omega = 0$  in other states, as shown in Fig. 7c.

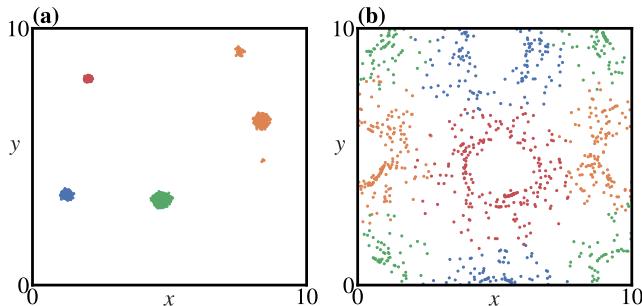


FIG. 8. Two examples of classification results. (a):  $(\lambda = 0.4, d_0 = 0.3)$ . (b):  $(\lambda = 0.02, d_0 = 0.4)$ .

TABLE I. Order parameter values in each state

State	$R$	$R_s$	$\Delta\Omega$
Disorder	$= 0$	$= 0$	$> 0$
Ring	$= 0$	$= 0$	$= 0$
Swarm (Multi-Clusters)	$> 0$	$\rightarrow 1^a$	$= 0$
Swarm (Double-Clusters)	$> 0$	$\Rightarrow 1^a$	$= 0$
Quick Sync	$= 1$	$= 1$	$= 0$

<sup>a</sup> Note that the  $R_s$  of Double-Clusters state is closer to 1 than that of Multi-Clusters state.

To sum up, using  $R$ ,  $R_s$  and  $\Delta\Omega$  in combination allows us to discern all the equilibrium state of our system. The order parameter values in each state are summarized in Table I.

## V. ANALYSIS

### A. Analytical approximations of the critical boundaries

In this section we derive the analytical approximations of the boundaries between the states. The boundaries between the states are determined by the critical lines  $\lambda_{c1}$ ,  $\lambda_{c2}$ ,  $\lambda_{c3}$ , and  $\lambda_{c4}$ , which are the critical values of  $\lambda$  for the transitions between the states with given  $d_0$ .

#### 1. $\lambda_{c1}$ : Disorder to Ring

We first consider the transition between the Disorder and Ring states. The oscillators in both states move in a circular trajectory (cf. Fig. 4b), and the difference between them is that the oscillators in the Ring state are phase-locked. Therefore, oscillators in Disorder state have a very small probability of being randomly distributed on the ring, but without phase locking. The critical coupling strength  $\lambda_{c1}$  can be calculated as the critical value of  $\lambda$  for the phase locking of oscillators on the same ring. We consider the following synchronous dynamics of oscillators on the same ring:

$$\dot{\theta}_1 = \omega_1 + \lambda \sum_{j=2}^{N_c} A_{1j} \sin(\theta_j - \theta_1), \quad (18)$$

$$\dot{\theta}_j = \omega_j, \quad (19)$$

for  $j = 2, 3, \dots, N_c$ , where  $N_c$  is the number of oscillators on the same ring, and  $\theta_1$  is the phase of the oscillator about to be phase-locked. Introducing the phase difference  $\Delta\theta_j = \theta_j - \theta_1$  and  $\Delta\omega_j = \omega_j - \omega_1$ , we have:

$$\Delta\dot{\theta}_j = \Delta\omega_j + \lambda \sum_{k=2}^{N_c} A_{jk} \sin \Delta\theta_k. \quad (20)$$

Each oscillator only needs to be phase-locked with neighboring oscillators (minimum  $|\Delta\omega_j|$ ) to achieve phase locking of entire ring due to the symmetry. The minimum  $|\Delta\omega_j|$  in a ring is  $|\Delta\omega_j| = (\omega_{\max} - \omega_{\min})/N_c$ . When  $\lambda \sum_{k=2}^{N_c} A_{jk} \geq |\Delta\omega_j|$ ,

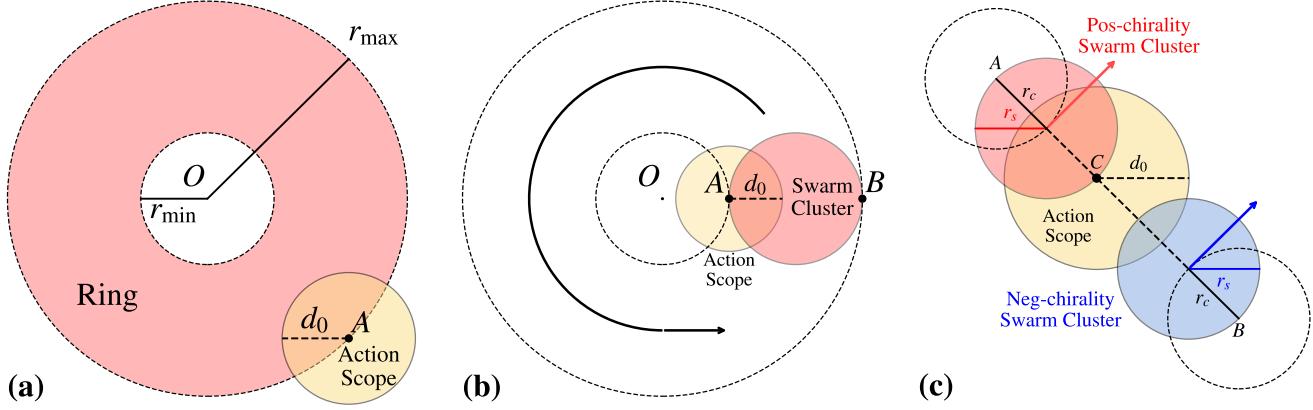


FIG. 9. The schematic plot of the analytical approximations. (a), The 1st oscillator is at point A which is on the outer edge of the ring, and the overlapping area of the yellow circle (action scope) and the red ring is  $S_1(d_0)$ . (b), The 1st oscillator is at point A which is on the inner edge of the ring, and the overlapping area of the yellow circle (action scope) and the red ring is  $S_2(d_0)$ . (c), The 1st oscillator is at point C which is on the edge of red circle, and the overlapping area of the yellow circle (action scope) and the blue circle is  $S_3(d_0)$ .

Eq. (20) has fixed point solutions, and the oscillators are phase-locked. Therefore, the critical coupling strength  $\lambda_{c1}$  is:

$$\lambda_{c1} = \frac{\omega_{\max} - \omega_{\min}}{N_c \sum_{j=2}^{N_c} A_{1j}}, \quad (21)$$

where  $\sum_{j=2}^{N_c} A_{1j}$  is the number of oscillators within the action scope of the 1st oscillator on the ring. Obviously, this is a function of  $d_0$ . We define it as

$$N_1(d_0) = N_c \frac{S_1(d_0)}{S_R} = \frac{N_c S_1(d_0)}{\pi(r_{\max}^2 - r_{\min}^2)}, \quad (22)$$

where  $S_1$  is the overlapping area of the action scope of the 1st oscillator and the ring,  $S_R$  is the area of the ring, and  $r_{\max} = v/\omega_{\min}$  and  $r_{\min} = v/\omega_{\max}$  are the outer and inner radius of the ring, respectively. In order to achieve phase locking of all oscillators on the ring, we need to consider the minimum value of  $N_1(d_0)$ . As shown in Fig. 9a, the 1st oscillator is at point A which is on the outer edge of the ring, and the overlapping area of the yellow circle (action scope) and the red ring is  $S_1(d_0)$ . Elementary geometry gives

$$\begin{cases} S_1(d_0) = d_0^2 \frac{\alpha}{2} + r_{\max}^2 \frac{\beta}{2} - r_{\max} d_0 \sin \frac{\alpha}{2} \\ \beta = 2 \arccos \left( 1 - \frac{d_0^2}{2r_{\max}^2} \right) \\ \alpha = \pi - \frac{\beta}{2} \end{cases}, \quad (23)$$

where  $\alpha$  and  $\beta$  are the angles of sectors from two overlapping circles' centers. Overall, we have

$$\begin{cases} \lambda_{c1} = \frac{\pi(r_{\max}^2 - r_{\min}^2)(\omega_{\max} - \omega_{\min})}{N_c^2 \left( d_0^2 \frac{\alpha}{2} + r_{\max}^2 \frac{\beta}{2} - r_{\max} d_0 \sin \frac{\alpha}{2} \right)} \\ \beta = 2 \arccos \left( 1 - \frac{d_0^2}{2r_{\max}^2} \right) \\ \alpha = \pi - \frac{\beta}{2} \end{cases}. \quad (24)$$

Then we can calculate the critical line  $\lambda_{c1}$  in the  $(\lambda-d_0)$  plane with  $N_c = 250$  (Ring state first appears as four rings, more details in Subsection V C), as shown in Fig. 7c, where the line

accurately divides the phase-locked and non-phase-locked regions.

## 2. $\lambda_{c2}$ : Ring to Swarm (Multi-Clusters)

According to  $R_s$  in Fig. 7b, the oscillators in the cluster are synchronized, and the oscillators with different chirality are not synchronized. Therefore,  $\lambda_{c2}$  is the critical value of  $\lambda$  for the synchronization of oscillators in the same cluster. We still consider the synchronous dynamics in Eq. (18) and (19). To achieve synchronization, the oscillators with the largest phase difference (maximum  $|\Delta\omega_j|$ ) in the cluster need to be phase locked. The maximum  $|\Delta\omega_j|$  in a cluster is  $|\Delta\omega_j| = \omega_{\max} - \omega_{\min}$ . Thus, the critical coupling strength  $\lambda_{c2}$  is:

$$\lambda_{c2} = \frac{\omega_{\max} - \omega_{\min}}{\sum_{j=2}^{N_s} A_{1j}}, \quad (25)$$

where  $N_s$  is the number of oscillators in the cluster, and  $\sum_{j=2}^{N_s} A_{1j}$  is the number of oscillators within the action scope of the 1st oscillator in the cluster, which is  $N_2(d_0)$ . We define it as

$$N_2(d_0) = N_s \frac{S_2(d_0)}{S_S} = \frac{N_s S_2(d_0)}{\pi r_s^2}, \quad (26)$$

where  $S_2$  is the overlapping area of the action scope of the 1st oscillator and the cluster,  $S_S$  is the area of the cluster. Considering that this is a phase transition between ring and cluster states, the oscillator moves with a rotational radius in the ring state until it gathers into clusters, so the radius of the cluster  $r_s$  is  $(r_{\max} - r_{\min})/2$ , as shown in Fig. 9b. Similar to Eq. (23), the minimum overlapping area  $S_2$  (overlapping area of two circles

in Fig. 9b) can be calculated as

$$\begin{cases} S_2(d_0) = d_0^2 \frac{\alpha}{2} + r_s^2 \frac{\beta}{2} - r_s d_0 \sin \frac{\alpha}{2} \\ \beta = 2\arccos \left( 1 - \frac{d_0^2}{2r_s^2} \right) \\ \alpha = \pi - \frac{\beta}{2} \end{cases} . \quad (27)$$

Therefore, we have

$$\begin{cases} \lambda_{c2} = \frac{\pi r_s^2 (\omega_{\max} - \omega_{\min})}{N_s \left( d_0^2 \frac{\alpha}{2} + r_s^2 \frac{\beta}{2} - r_s d_0 \sin \frac{\alpha}{2} \right)} \\ \beta = 2\arccos \left( 1 - \frac{d_0^2}{2r_s^2} \right) \\ \alpha = \pi - \frac{\beta}{2} \\ r_s = \frac{r_{\max} - r_{\min}}{2} \end{cases} . \quad (28)$$

Then we can calculate the critical line  $\lambda_{c2}$  in the  $(\lambda-d_0)$  plane with  $N_s = 500$  (each chirality has half of the total population  $N = 1000$ ), as shown in Fig. 7b, where the line divides whether local synchronization occurred.

### 3. $\lambda_{c3}$ : Multi-Clusters to Double-Clusters

The derivation of  $\lambda_{c3}$  is similar to that of  $\lambda_{c2}$ , but the difference is that the centers of oscillators in the Swarm (Double-Clusters) state converge quickly (cf Fig. 5b), which means oscillators do not require a long transient to form clusters. We consider that the oscillator density  $\rho$  (Number of oscillators per unit area) under initial conditions can achieve synchronization of oscillators of the same chirality.  $\rho$  can be calculated as

$$\rho = \frac{N}{L^2} . \quad (29)$$

Then we have

$$\sum_{j=2}^{N_c} A_{1j} = \rho \pi d_0^2 = \frac{N \pi d_0^2}{L^2} . \quad (30)$$

The critical coupling strength  $\lambda_{c3}$  is:

$$\lambda_{c3} = \frac{L^2 (\omega_{\max} - \omega_{\min})}{N \pi d_0^2} . \quad (31)$$

Similarly, we can calculate the critical line  $\lambda_{c3}$  in the  $(\lambda-d_0)$  plane, as shown in Fig. 7b, where the values of region above the line (Double-Clusters) is closer to 1 than below the line (Multi-Clusters).

### 4. $\lambda_{c4}$ : Swarm (Double-Clusters) to Quick Sync

The Quick Sync state can be considered as a special case of the Swarm state where the oscillators are globally synchronized. Our derivation approach comes from adiabatic parameter tuning (see Appendix D). We assume that the oscillators of

two chiralities have been synchronized in the Swarm (Double-Clusters) state, as is shown in Fig. 9c. Next, we consider the synchronization of the two chiralities. According to Eq. (13), two clusters with opposite chirality are moving on their respective circular trajectories with the same rotational radius

$$r_c = \frac{2v}{\omega_{\max} + \omega_{\min}} . \quad (32)$$

In numerical simulations, we observe that the radius  $r_s$  of clusters in Swarm state increases with the increase of  $\lambda$  and  $d_0$ . Thus, the radius of the clusters reach maximum value of  $r_c$  (see proof in Appendix E) when the  $\lambda$  and  $d_0$  reach the critical values.

Considering that the trajectories of two chiralities will repel each other before oscillators achieving global synchronization, let's take the maximum relative distance between the motion trajectories of two clusters in space, which is line  $AB$  in Fig. 9c. The distance  $l_{AB}$  is

$$l_{AB} = \frac{L}{\sqrt{2}} . \quad (33)$$

Similar to the derivation of  $\lambda_{c2}$ , we consider the synchronization dynamics between a single oscillator ( $C$  in Fig. 9c) and a cluster (blue circle in Fig. 9c), the difference is that the single oscillator and cluster here belong to different chirality. When the centers of two clusters move onto line segment  $AB$ , their distance is the shortest. At this point, due to symmetry, if the oscillators at the edges of the two clusters are synchronized, global synchronization can be achieved. Obviously, when

$$d_0 + 2r_c + 2r_s < l_{AB} , \quad (34)$$

the action scope of oscillator at point  $C$  does not overlap with the neg-chirality cluster (blue circle), so the critical coupling strength  $\lambda_{c4}$  does not exist. When  $d_0 + 2r_c + 2r_s > l_{AB}$ , the overlapping area  $S_3(d_0)$  can be calculated as

$$\begin{cases} S_3(d_0) = r_s^2 \frac{\alpha}{2} + d_0^2 \frac{\beta}{2} - d_0 r_d \sin \frac{\beta}{2} \\ r_d = \frac{L}{\sqrt{2}} - r_s - 2r_c \\ \beta = 2\arccos \frac{d_0^2 + r_d^2 - r_s^2}{2d_0 r_d} \\ \alpha = 2\arccos \frac{r_s^2 + r_d^2 - d_0^2}{2r_s r_d} \end{cases} , \quad (35)$$

where  $r_s = r_c$ . The global maximum value of  $|\Delta\omega_j|$  is  $\omega_{\max} - (-\omega_{\max}) = 2\omega_{\max}$ . Therefore, we have

$$\begin{cases} \lambda_{c4} = \frac{2\pi r_s^2 \omega_{\max}}{N_c \left( \frac{\alpha}{2} r_s^2 + \frac{\beta}{2} d_0^2 - d_0 r_d \sin \frac{\beta}{2} \right)} \\ r_d = \frac{L}{\sqrt{2}} - r_s - 2r_c \\ \beta = 2\arccos \frac{d_0^2 + r_d^2 - r_s^2}{2d_0 r_d} \\ \alpha = 2\arccos \frac{r_s^2 + r_d^2 - d_0^2}{2r_s r_d} \end{cases} . \quad (36)$$

Then we can calculate the critical line  $\lambda_{c4}$  in the  $(\lambda-d_0)$  plane, as shown in Fig. 7a, where the line divides whether global synchronization occurred.

## B. Repulsion-attraction mechanism

In the Ring state, the oscillators of two chiralities exhibit a spatial repulsion-attraction mechanism driven by phase coupling. Specifically, the motion trajectories of chiral oscillators of the same chirality are attracted to each other, while the motion trajectories of chiral oscillators of different chirality repel each other. Considering that the ring state occurs when the coupling strength  $\lambda$  and the action radius  $d_0$  are small and the trajectories of the oscillators in the Ring state are almost circular under uncoupled conditions, we can analyze the repulsion-attraction mechanism by considering the motion trajectories of the oscillators.

Under the assumption of no coupling, the center  $\mathbf{c}_i$  of the  $i$ th oscillator moves on a circular trajectory with radius  $r_i = v/\omega_i$  and angular velocity  $\omega_i$  according to Eq. (8). The coordinates of the rotation center can be estimated as the mean coordinates of points on the trajectory. the horizontal ordinate of the  $i$ th oscillator's rotation center is

$$\begin{aligned} C_{xi}(t_0) &= \frac{\omega_i}{2\pi} \int_{t_0}^{t_0+2\pi/\omega_i} \left\{ x_i^0(t_0) + \frac{v}{\omega_i} \sin[\theta_i(t_0) + \omega_i t] \right\} dt \\ &= x_i^0(t_0) \\ &= x_i(t_0) - \frac{v}{\omega_i} \sin \theta_i(t_0) . \end{aligned} \quad (37)$$

Similarly, the vertical ordinate of the  $i$ th oscillator's rotation center is

$$C_{yi}(t_0) = y_i(t_0) + \frac{v}{\omega_i} \cos \theta_i(t_0) . \quad (38)$$

Due to the symmetry of the ring, we can consider the interaction between only two oscillators. For the  $i$ th and  $j$ th oscillators, the distance between their rotation centers  $|\mathbf{c}_j - \mathbf{c}_i|$  is

$$|\mathbf{c}_j - \mathbf{c}_i| = \sqrt{\left( x_j - \frac{v}{\omega_j} \sin \theta_j - x_i + \frac{v}{\omega_i} \sin \theta_i \right)^2 + \left( y_j + \frac{v}{\omega_j} \cos \theta_j - y_i - \frac{v}{\omega_i} \cos \theta_i \right)^2} . \quad (39)$$

Considering the action radius  $d_0$  is small, the coordinates of the two oscillators are very close to each other when coupling occurs, which means  $(x_i, y_i) \approx (x_j, y_j)$ . Then we have

$$|\mathbf{c}_j - \mathbf{c}_i| = \sqrt{-\frac{2v^2 \cos(\theta_i - \theta_j)}{\omega_i \omega_j} + \frac{v^2}{\omega_i^2} + \frac{v^2}{\omega_j^2}} . \quad (40)$$

Furthermore, due to the symmetry of coupling from Eq. (3), The phase of the two oscillators after being coupled will have a pair of symmetrical angular offsets compared to those without coupling:

$$\begin{aligned} \hat{\theta}_i &= \theta_i + \lambda \sin(\theta_j - \theta_i) T , \\ \hat{\theta}_j &= \theta_j + \lambda \sin(\theta_i - \theta_j) T . \end{aligned} \quad (41)$$

Let's set  $\theta_s = \lambda \sin(\theta_j - \theta_i) T$  as the angular offset and  $\theta_s$  clearly satisfies

$$\theta_s \in \begin{cases} \left[0, \frac{\theta_i - \theta_j}{2}\right], & \theta_j \geq \theta_i \\ \left[\frac{\theta_i - \theta_j}{2}, 0\right], & \theta_j < \theta_i \end{cases} \quad (42)$$

Then we have the distance driven by the coupling  $|\hat{\mathbf{c}}_j - \hat{\mathbf{c}}_i|$  as

$$\begin{aligned} |\hat{\mathbf{c}}_j - \hat{\mathbf{c}}_i| &= \sqrt{-\frac{2v^2 \cos(\hat{\theta}_i - \hat{\theta}_j)}{\omega_i \omega_j} + \frac{v^2}{\omega_i^2} + \frac{v^2}{\omega_j^2}} \\ &= \sqrt{-\frac{2v^2 \cos(\theta_i - \theta_j + 2\theta_s)}{\omega_i \omega_j} + \frac{v^2}{\omega_i^2} + \frac{v^2}{\omega_j^2}} . \end{aligned} \quad (43)$$

To compare Eq. (39) and Eq. (43), we only need to compare the cosine term in the two equations. According to Inequality (42), we have

$$\theta_i - \theta_j + 2\theta_s \in \begin{cases} [0, \theta_i - \theta_j], & \theta_i - \theta_j \geq 0 \\ [\theta_i - \theta_j, 0], & \theta_i - \theta_j < 0 \end{cases} , \quad (44)$$

which means

$$\cos(\theta_i - \theta_j) \leq \cos(\theta_i - \theta_j + 2\theta_s) . \quad (45)$$

Therefore, when the coupling occurs, the distance between the rotation centers of two oscillators with the same signs of natural frequency  $\omega$  (chirality) will decrease, which means the oscillators of the same chirality are attracted to each other, while the distance between the centers with opposite signs of  $\omega$  (chirality) will increase, which means the oscillators of different chirality repel each other.

## C. Number of Rings

## VI. CONCLUSIONS

## Appendix A: PROOF OF THE ADJUSTED POSITION

In this section, we prove the distances between oscillators' adjusted position is the minimum distance in periodic boundary conditions.

*Proof.* To prove this, we only need to prove the adjusted distance  $|\mathbf{r}_i - \bar{\mathbf{r}}_j|$  is not longer than the raw distance  $|\mathbf{r}_i - \mathbf{r}_j|$ .

For  $(x_i - x_j)^2$  and  $(x_i - \bar{x}_j)^2$ , if  $x_j = \bar{x}_j$ , we have  $(x_i - x_j)^2 = (x_i - \bar{x}_j)^2$ . If  $x_j \neq \bar{x}_j$ , we have

$$\begin{aligned} (x_i - \bar{x}_j)^2 - (x_i - x_j)^2 &= (x_i \pm L - x_i)^2 - (x_j - x_x)^2 \\ &= \begin{cases} L^2 + 2L(x_j - x_x), & x_i - x_j > L/2 \\ L^2 - 2L(x_j - x_x), & x_i - x_j < L/2 \end{cases} \\ &< L^2 - L^2 \\ &= 0 \end{aligned} \quad (\text{A1})$$

Then, we have  $(x_i - \bar{x}_j)^2 \leq (x_i - x_j)^2$ . Similarly, we have  $(y_i - \bar{y}_j)^2 \leq (y_i - y_j)^2$ . Therefore, we have

$$\begin{aligned} |\mathbf{r}_i - \bar{\mathbf{r}}_j| &= \sqrt{(x_i - \bar{x}_j)^2 + (y_i - \bar{y}_j)^2} \\ &\leq \sqrt{(x_i - x_j)^2 + (y_i - y_j)^2} \\ &= |\mathbf{r}_i - \mathbf{r}_j| . \end{aligned} \quad (\text{A2})$$

## Appendix B: NUMERICAL METHODS

All the simulations of the model Eq. (1)-(3) were run on Python using Euler integration, with a time step  $\Delta t = 0.01$ , and a total time of  $T = 60000$ .

## Appendix C: DETERMINATION OF DBSCAN'S PARAMETERS

DBSCAN (Density-Based Spatial Clustering of Applications with Noise) is a density-based clustering algorithm, which can find clusters of arbitrary shapes and sizes. The algorithm has two parameters:  $\varepsilon$  and  $m$ .  $\varepsilon$  is the maximum distance between two samples for one to be considered as in the neighborhood of the other, and  $m$  is the minimum number of samples in a neighborhood for a point to be considered as a core point.

We traverse all values between 0.15 and 0.5 with a step length of 0.05, and for each value of  $\varepsilon$ , we calculate the number of clusters of Swarm state with  $m = 5$  (which is 0.5% of the population  $N = 1000$  of the system). Then we record the minimum counts of clusters in total states.

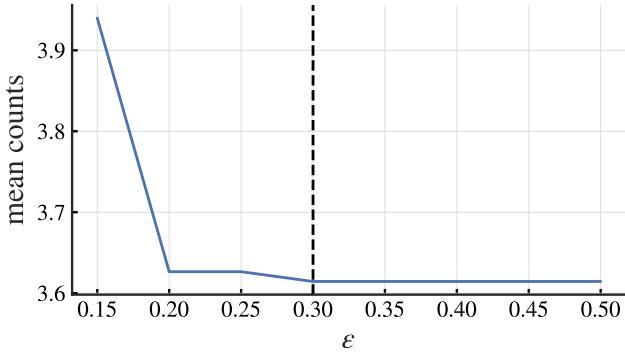


FIG. 10. The minimum counts of clusters with  $m = 0$  and different  $\varepsilon$ . The number of clusters is calculated by DBSCAN algorithm.

As shown in Fig. 10, the mean counts of clusters converges to at  $\varepsilon = 0.3$ . Then we set  $\varepsilon$  to be 0.3, and set  $m$  to be 5.

## Appendix D: ADIABATIC TUNING OF THE ACTION RADIUS

## Appendix E: PROOF OF THE MAXIMUM RADIUS OF CLUSTERS

In this section, we prove the maximum radius of cluster in the Swarm state is  $r_c$ , which is the radius of its rotation trajectory. We prove this by contradiction.

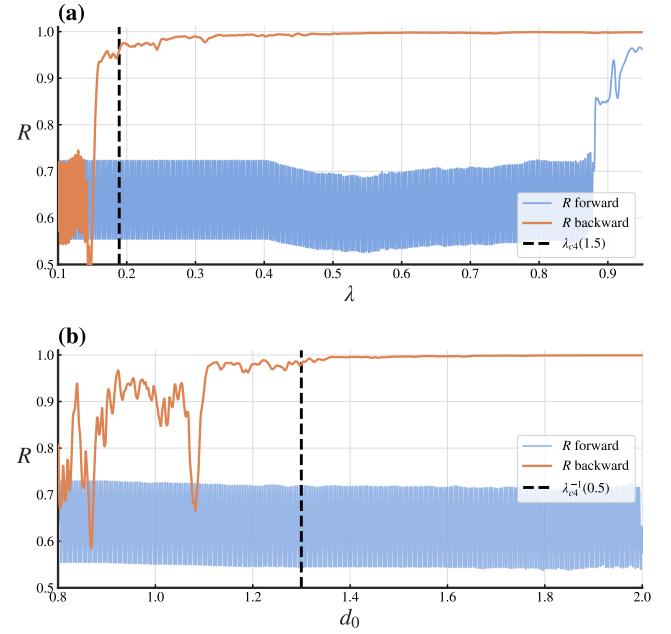


FIG. 11. The adiabatic tuning of the coupling strength  $\lambda$  and the action radius  $d_0$ . (a), The coupling strength  $\lambda$  is adiabatically tuned from 0.01 to 0.95 with fixed  $d_0 = 1.5$ . (b), The action radius  $d_0$  is adiabatically tuned from 0.8 to 2 with fixed  $\lambda = 0.5$ .

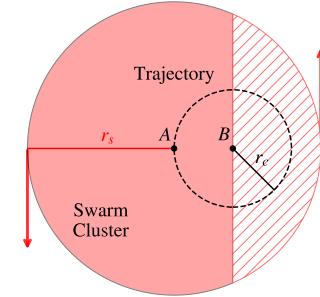


FIG. 12. The schematic plot of the proof of the maximum radius of clusters.

*Proof.* We assume that the maximum radius of the cluster is  $r_s > r_c$ . Then we have the case that the cluster's rotation center is inside the cluster circle with radius  $r_s$ . Fig. 12 shows the schematic plot of this case.  $A$  and  $B$  is the center of the cluster and trajectory circle, respectively. The cluster circle is divided into two parts by  $B$ : red solid and red // hatching. In order to enable clusters to move along trajectories, the oscillators in two parts must have contrary phase velocities. This contradicts the fact that the oscillators in the cluster are synchronized. Therefore, the maximum radius of the cluster is  $r_c$ .  $\square$

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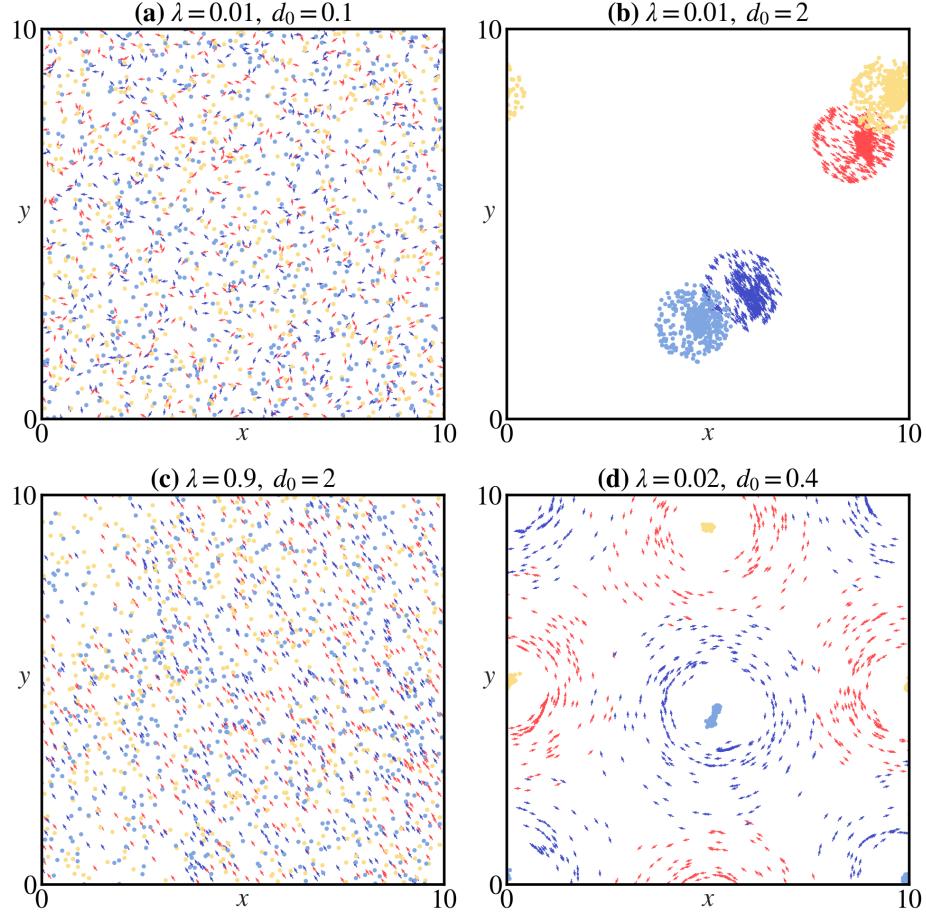


FIG. 13. Estimation results of real-time rotational centers. The centers of two types of chirality oscillators are represented by light yellow ( $\omega_i > 0$ ) and light blue ( $\omega_i < 0$ ) points, respectively. (a), Disorder state ( $\lambda = 0.01$ ,  $d_0 = 0.1$ ). (b), Swarm state, ( $\lambda = 0.01$ ,  $d_0 = 2$ ). (c), Quick Sync stat ( $\lambda = 0.9$ ,  $d_0 = 2$ ). (d), Ring state ( $\lambda = 0.02$ ,  $d_0 = 0.4$ ).

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