Week 4: Theory Solutions

Solution 1 We know that the Hadamard Gate is the QFT for N=2. Accordingly, the second root of unity is $\omega=\exp(\frac{2i\pi}{2})=e^{i\pi}=-1$. For N=2, we have two basis vectors $|0\rangle$ and $|1\rangle$. Therefore for $|0\rangle$ we have

$$|0\rangle \xrightarrow{QFT} \frac{1}{\sqrt{2}} \sum_{y=0}^{1} \omega_2^{-0y} |y\rangle$$
$$= \frac{1}{\sqrt{2}} \sum_{y=0}^{1} |y\rangle$$
$$= \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

and for $|1\rangle$ we have

$$\begin{aligned} |1\rangle &\xrightarrow{QFT} \frac{1}{\sqrt{2}} \sum_{y=0}^{1} \omega_2^{-1y} |y\rangle \\ &= \frac{1}{\sqrt{2}} (-1^{-1 \times 0} |0\rangle) + \frac{1}{\sqrt{2}} (-1^{-1 \times 1} |1\rangle) \\ &= \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \end{aligned}$$

Hence the required Fourier basis states are $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ and $\frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$. These are sometimes known as the $|+\rangle$ and $|-\rangle$ states.

Solution 2 We have the two basic equations $A|\psi\rangle = \lambda |\psi\rangle$ and $AA^* = I$. Now,

$$A|\psi\rangle = \lambda|\psi\rangle$$

$$|\psi\rangle^*A^* = \lambda^*|\psi\rangle^*$$

$$\implies |\psi\rangle^*A^*A|\psi\rangle = \lambda^*|\psi\rangle^*\lambda|\psi\rangle$$

$$\implies |\psi\rangle^*I|\psi\rangle = \lambda^*\lambda|\psi\rangle^*|\psi\rangle$$

$$\implies |\psi\rangle^*|\psi\rangle = \lambda^*\lambda|\psi\rangle^*|\psi\rangle$$

$$\implies |\psi\rangle|^2 = |\lambda|^2||\psi\rangle|^2$$

$$\implies |\lambda|^2 = 1$$

$$\implies |\lambda| = 1$$

Solution 3 We know that the MS qubit has the form

$$\frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i 0.y_n}|1\rangle\right) = \frac{1}{\sqrt{2}}\left(|0\rangle + e^{2\pi i y_n/2}|1\rangle\right) = \frac{1}{\sqrt{2}}\left(|0\rangle + e^{\pi i y_n}|1\rangle\right)$$

Since $e^{\pi i}=-1$, depending on the value of y_n (bit can be 0 or 1) we have the phase of $|1\rangle$ as +1 or -1. Hence we get a form of $\frac{1}{\sqrt{2}}\left(|0\rangle+|1\rangle\right)$ or $\frac{1}{\sqrt{2}}\left(|0\rangle-|1\rangle\right)$. This is nothing but a Hadamard transformation.

Solution 4

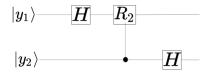


Figure 1: QFT circuit for 2 qubits

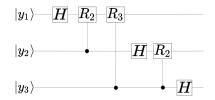


Figure 2: QFT circuit for 3 qubits

Solution 5 When $2^n \phi \neq a$, $\beta \neq 0$.

$$\Pr[a] = \frac{1}{N^2} \left| \sum_{k=0}^{N-1} e^{(2\pi i k \beta)} \right|^2$$

Now, using geometric series sum we get

$$= \frac{1}{N^2} \left| \frac{1 - e^{(2\pi i N \beta)}}{1 - e^{(2\pi i \beta)}} \right|^2$$

Now we use the observation that $|1 - e^{2ix}|^2 = 4 |\sin x|^2$, and substitute this in our formula

to get

$$= \frac{1}{N^2} \left| \frac{\sin \pi N \beta}{\sin \pi \beta} \right|^2$$

$$\geq \frac{1}{N^2} \left| \frac{\sin \pi N \beta}{\pi \beta} \right|^2$$

$$\geq \frac{1}{N^2} \left| \frac{2N \beta}{\pi \beta} \right|^2$$

$$\geq \frac{4}{\pi^2}$$

Second step is because $|\sin x| \le |x|, \quad \forall |x| \le \frac{1}{2\pi N}$. Third step is because $|\sin \pi N\beta| \ge |2N\beta|, \quad \forall \quad |\beta| \le \frac{1}{2N}$.

Solution 6

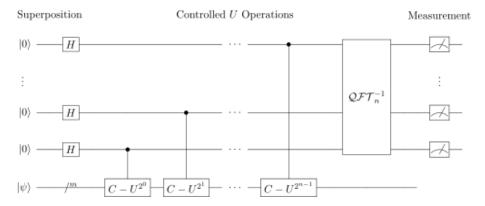


Figure 3: Circuit for Quantum Phase Estimation ¹

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