

Q.1. HW - WEEK II ASHMITA PANDA 1811042
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~~for light, from special relativity.~~

$$ds^2 = c^2 dt^2 - a^2(t) [dr^2 + r^2 d\Omega^2]$$

for expansion of universe,

photon geodesic would be

$$ds=0 \text{ in } ds^2 = c^2 dt^2 - a^2(t) [dr^2 + r^2 d\Omega^2]$$

⊗ & $d\Omega=0$ for (assuming photon follows a path of constant θ & ϕ)

$$\Rightarrow |c dt| = a(t) dr$$

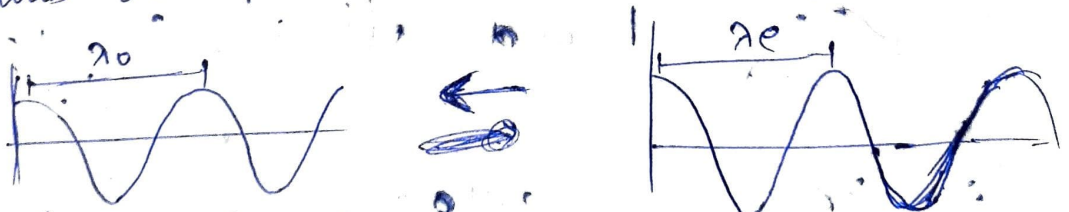
$$\Rightarrow dr = \frac{1}{c} \frac{dt}{a(t)} \quad c=1$$

$$\Rightarrow dr = \frac{dt}{a(t)} \text{ for photons.}$$

∴ for a photon emitted at $(t_0, r_0, \theta_0, \phi_0)$ & observed at $(t_e, r_e, \theta_0, \phi_0)$

$$\therefore \Delta r = r_e - r_0 = \int_{t_e}^{t_0} \frac{dt}{a(t)}$$

applying this on two wave crests / troughs.



$$\Delta r = r_0 - r_e = \int_{t_e}^{t_0} \frac{dt}{a(t)} = \int_{t_e + \frac{\lambda_e}{c}}^{t_0 + \frac{\lambda_0}{c}} \frac{dt}{a(t)}$$

now, $\int_{t_e}^{t_0} \frac{dt}{a(t)} = \int_{t_e + \frac{\lambda_e}{c}}^{t_0 + \frac{\lambda_0}{c}} \frac{dt}{a(t)}$

$$\Rightarrow \int_{t_e}^{t_0} \frac{dt}{a(t)} = \int_{t_e}^{t_0} \frac{dt}{a(t)} - \int_{t_e}^{t_e + \frac{\lambda_e}{c}} \frac{dt}{a(t)} + \int_{t_0}^{t_0 + \frac{\lambda_0}{c}} \frac{dt}{a(t)}$$

$$\Rightarrow \int_{t_e}^{t_e + \frac{\lambda_e}{c}} \frac{dt}{a(t)} = \int_{t_0}^{t_0 + \frac{\lambda_0}{c}} \frac{dt}{a(t)}$$

\therefore the time $\frac{\lambda_e}{c}$ or $\frac{\lambda_0}{c}$ (between two successive wavecrests)

is extremely small, $a(t)$ can be taken to be constant.

$$\Rightarrow \frac{1}{a(t_e)} \int_{t_e}^{t_e + \frac{\lambda_e}{c}} dt = \frac{1}{a(t_0)} \int_{t_0}^{t_0 + \frac{\lambda_0}{c}} dt$$

$$\Rightarrow \frac{1}{a(t_e)} \frac{\lambda_e}{c} = \frac{1}{a(t_0)} \frac{\lambda_0}{c} \Rightarrow \frac{\lambda_e}{\lambda_0} = \frac{a(t_e)}{a(t_0)}$$

now, using $z = \frac{\lambda_0}{\lambda_e} - 1$ and $a(t_0) = 1$ (observing at current time)

$$\Rightarrow \frac{\lambda_0}{\lambda_e} = 1 + z = \frac{a(t_0)}{a(t_e)} \Rightarrow a(t_e) = \frac{a(t_0)}{1 + z}$$

$$\text{or, } a(t_e) = \frac{1}{1 + z}$$