Lecture 10: The Poincaré Recurrence Theorem

Yannis Bähni

University of Augsburg yannis.baehni@math.uni-augsburg.de

June 17, 2021

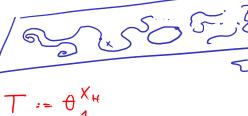
Ergodic Theory

the study of measure theoretical do

Definition (Dynamical System)

A *dynamical system* on a probability space (X, A, μ) is a measurable transformation $T: X \to X$ such that $T_*\mu = \mu$.

Ot to denois



Let $T: X \to X$ be a dynamical system on a probability space (X, A, μ) . Given $A \in A$, set

$$E := \bigcap_{n=0}^{\infty} E_n$$
 where $E_n := \bigcup_{k=n}^{\infty} T^{-k} A$.

Then $\mu(A \cap E) = \mu(A)$. here it is crucial that P preserves μ .

Proof A good exercise in mensue theory

We have that x EARE; I and only if the exists a sequere (ky) EN s.t |

ky -> 00 and Tkn(x) EA

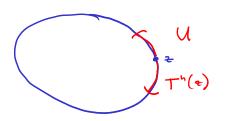
Y we N.

Theorem (Poincant Recurrence the one

Let $T: X \to X$ be a dynamical system on a probability space (X, A, μ) . Let $A \in A$ be measurable with $\mu(A) > 0$. Then almost every point of A returns infinitely often to A under T.

 $T_{\theta} \colon S^{1} \to S^{1}$ · fe Q, that is B - w/n, so Th(2) = 2.

· D € Dr, the the west {T"(2): N ∈ N } ∈ 5° is dense by the clove Theore.



The most be some iterate in U because otherwise U, $\pm U$, $\pm 3U$, \dots nello are disjoint. But as $\mu(U) > 0$, $\rightarrow \mu(S^{2}) = \infty$.

Regular Energy Surfaces

nition (Regular Energy Surface)

Definition (Regular Energy Surface)

A regular energy surface in a Hamiltonian system (M, ω, H) is defined to be an embedded hypersurface $\Sigma = H^{-1}(0)$ such that $Crit(H) \cap \Sigma = \emptyset$. Any such function is called a defining Hamiltonian function for Σ .

Since 0 is a regular value of 11, the preimage H-1(0) is an embedded hyperriface value by the by the implicit fruition theorem. Moreover, Tx Zi = leardy x + ye Zi. Clin. Xx is to quet to Z. $\mathsf{Aff}^{\times}(X^{\mathsf{H}}(X)) = -\omega(X^{\mathsf{H}}(X)^{\mathsf{I}}X^{\mathsf{H}}(X)) - O.$ for all x = Z. .

Every regular energy surface is orientable.

Prof. let in be a Riemannia metic on M. The consider the "normalised" gradient

X = grad + H = EX(U)

where UEH is on open neighborhood of Zi 5.1. All + O (this on 'be done if I is ampant, your is defend on 2). Volume on 2: ixing.

dH. + + = dH (qulu + · + x) I quelut + + 1/2 ¥ { = (= =). -> Hopx = +

Let Σ be an embedded hypersurface in a symplectic manifold (M, ω) . Then

$$\ker \omega|_{\Sigma} \to \Sigma$$

is a line bundle, called the **characteristic line bundle of** Σ .

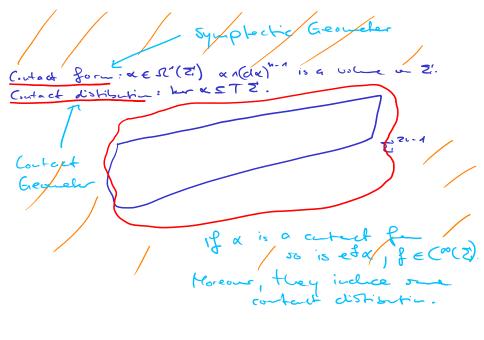
Ex. If I is a regular energy hypersurface, then XHI spans the characteristic line builte. Indeed, for xEI and vETx I we came $\omega(\chi_{\mathcal{H}}(x), v) = -dH_{\times}(v) - O.$ -> XHIT belongs to the characteristic l'a bulle . F. A XH never varishes on so it does incled open it. But why line distillation? with has lever runk, so there exists upe ler ul 7/50%. The <u> = ler w/t, | Seconse du TxZ" + din TxZ = din, M

Let $H, \widetilde{H} \in C^{\infty}(M)$ be two defining Hamiltonian functions for a regular energy surface Σ in a symplectic manifold (M, ω) . Then there exists a <u>nowhere-vanishing</u> function $f \in C^{\infty}(\Sigma)$ such that

$$X_{\widetilde{H}}|_{\Sigma} = fX_{H}|_{\Sigma}.$$

Front. We know from previous discussion text both Xff and XH span the line distriction her why.

This means that O'ff is just a reparametisation of OXH! In perfector, imparametrised periodic orbits coincide



Let Σ be a compact regular energy surface in a Hamiltonian system (M, ω, H) . Denote by θ the flow of X_H on Σ . Then there exists a unique regular θ -invariant probability measure μ_{Σ} on Σ , that is, we have that

$$\theta_t^* \mu_{\Sigma} = \mu_{\Sigma} \qquad \forall t \in \mathbb{R}.$$

Let (M, Ω) be a compact oriented smooth manifold of positive dimension and suppose that $\varphi \in \mathrm{Diff}(M)$ such that $\underline{\varphi}^*\Omega = \underline{\Omega}$. Then there exists a unique regular $\underline{\varphi}$ -invariant probability measure $\underline{\mu}_{\Omega}$ such that

$$\int_{M} f\Omega = \int_{M} f d\mu_{\Omega} \qquad \forall f \in C^{0}(M).$$

Froug. Clear from lecture
$$g$$
, as we emptre

$$\int_{\mathcal{H}} f \Omega = \int_{\mathcal{H}} \varphi^*(f \Omega) = \int_{\mathcal{H}} (f \circ \varphi) \varphi^* \Omega$$

$$= \int_{\mathcal{H}} (f \circ \varphi) d \rho \Omega$$

$$= \int_{\mathcal{H}} f \partial_{\varphi} (\varphi_* \rho_{\varphi}).$$

Let M be a smooth manifold. Suppose that $\eta \in \Omega^1(M)$ is nowhere-vanishing and $\xi \in \Omega^k(M)$. Then $\eta \wedge \xi = 0$ if and only if there exists $\zeta \in \Omega^{k-1}(M)$ such that $\xi = \eta \wedge \zeta$.

Ford:
$$w'' = dH \wedge x$$
.

If $\beta \in \Omega^{2n-1}(U)$ is another onch $\beta = 0$.

Here
$$dH \wedge (x - \beta) = 0$$

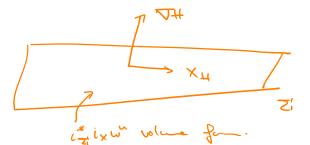
Find $\gamma \in \Omega^{2n-2}(U)$ s.t
$$x - \beta = dH \wedge \gamma$$
.

$$t_{\overline{\alpha}} = t_{\overline{\alpha}} (dH \wedge \gamma) + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \beta = d(H \circ t_{\overline{\alpha}}) \wedge t_{\overline{\alpha}} \gamma + t_{\overline{\alpha}} \gamma +$$

$$0 = i \times (qH \vee r) = \{xqH\}^{r} - qH \vee i \times r$$

$$V = \frac{\|\Delta H\|_{S}}{\Delta H} = V$$

$$V = \frac{\|\Delta H\|_{S}}{\Delta H} = V$$



Hoper + Zehnder: Humittonian Dynamics and Str

Theorem (Poincaré's Recurrence Theorem)

Let Σ be a compact regular energy surface in a Hamiltonian system (M, ω, H) . Then for almost every $x \in \Sigma$, with respect to the probability measure μ_{Σ} , there exists a sequence $(t_k) \subseteq \mathbb{R}$ such that

$$t_k \to +\infty$$
 and $\lim_{k \to \infty} \theta_{t_k}^{X_H}(x) = x$.

Proof. This is up to details the Poinceré recurrent in Zi is a recurrent theorem from enjoire theory. In Zi is a recurrent point.

this concially uses that I can cover bypronface with a countable number of metric balls (I adils a netic Structure of Renamin man fold)

