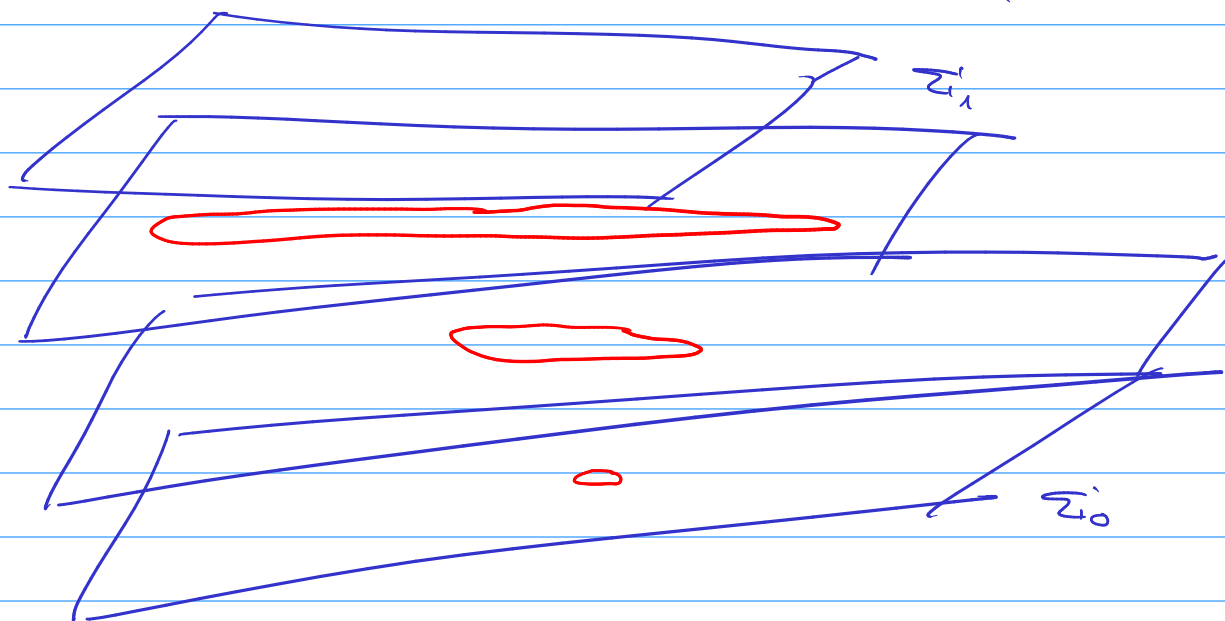


## Blue Sky Catastrophes

- Let  $(M, d\lambda)$  be an exact symplectic manifold and  $H \in C^\infty(M \times [0, 1])$  such that
- (i) 0 is a regular value of  $H_\sigma = H(\cdot, \sigma) \in C^\infty(M)$  for  $\sigma \in [0, 1]$ .
  - (ii)  $H^{-1}(0)$  is compact.
  - (iii) (Contact Condition) Set  $\Sigma_\sigma := H_\sigma^{-1}(0)$ ,  
 $dH_\sigma(X)|_{\Sigma_\sigma} > 0 \quad \forall \sigma \in [0, 1]$ ,  
 where  $X \in \mathfrak{X}(M)$  denotes the Liouville vector field defined by  $i_X d\lambda = \lambda$ .  
 (this just means that each  $(\Sigma_\sigma, \lambda|_{\Sigma_\sigma})$  is a hypersurface of restricted contact type).



Suppose we are given a smooth family  $(\gamma_\sigma, \tau_\sigma)_{\sigma \in [0, 1]}$  of parametrised periodic orbits, that is,

- (i)  $(\gamma_\sigma, \tau_\sigma) \in \mathcal{L}\Sigma_\sigma \times (0, +\infty)$  for all  $\sigma \in [0, 1]$ .
- (ii)  $\dot{\gamma}_\sigma(t) = \tau_\sigma X_{H_\sigma}(\gamma_\sigma(t)) \quad \forall t \in \mathbb{S}^1$ .

$$\left[ \begin{array}{l} \mathcal{L}\Sigma_\sigma := C^\infty(\mathbb{S}^1, \Sigma_\sigma) \\ \text{Can define } \gamma_\tau(t) = \gamma(t/\tau), \quad t \in \tau\mathbb{S}^1 \\ \text{The } \gamma_\tau \text{ is an integral curve of } X_H. \end{array} \right]$$

Ex. Define  $\bar{H} \in C^\infty(\mathbb{R} \times [0,1])$  by energy shift of a fixed Hamiltonian  $H \in C^\infty(\mathbb{R})$ :  
$$\bar{H}(x, \sigma) := H(x) - \sigma.$$

If you have a nondegenerate periodic orbit on  $T_0$ , you can extend it to a family of periodic orbits (locally).

Idea Study limit behaviour of the family  $(\gamma_\sigma, \tau_\sigma)_{\sigma \in [0,1]}$  as  $\sigma \rightarrow 1$ .

For example, it might be the case that  
$$\tau_\sigma \rightarrow +\infty \quad \text{as} \quad \sigma \rightarrow 1.$$

This is referred to as a blue sky catastrophe.  
A convenient way for studying such a limit behavior is by introducing the  $\omega$ -limit set  $\omega(\gamma_\sigma, \tau_\sigma)$  of the family  $(\gamma_\sigma, \tau_\sigma)$ , consisting of all

$$(\gamma_1, \tau_1) \in \mathbb{R} \times (0, +\infty)$$

such that there exists a sequence  $(\sigma_k) \in [0,1]$  with

$$\sigma_k \rightarrow 1 \quad \text{and} \quad (\gamma_{\sigma_k}, \tau_{\sigma_k}) \xrightarrow{C^0} (\gamma_1, \tau_1)$$

as  $k \rightarrow \infty$ .

Theorem (Belbrun/Frauenfelder/Koert) The  $\omega$ -limit set  $\omega(\gamma_\sigma, \tau_\sigma)$  is nonempty, compact and connected.

Proof. Compactness and connectedness are straight forward but tedious to prove. Note, that in order to prove  $\omega(\gamma_\sigma, \tau_\sigma) \neq \emptyset$ , it is enough to show that  $\tau_\sigma$  is uniformly bounded from below and above, i.e. there exists

$\epsilon > 0$  such that

$$0 < \frac{1}{\epsilon} \leq \tau_\sigma \leq \epsilon \quad \forall \sigma \in [0, 1].$$

Proof of  $\square$ .

Idea Recall that  $(\gamma_\sigma, \tau_\sigma) \in \text{Crit } \mathcal{A}^{\#_\sigma} \quad \forall \sigma \in [0, 1]$ .

Note, without loss of generality, we may assume

$$(*) \quad X_{\#_\sigma}|_{\Sigma_\sigma} = R_\sigma \quad \forall \sigma \in [0, 1].$$

Indeed, note that both  $X_{\#_\sigma}$  and  $R_\sigma$  belong to the characteristic distribution  $\ker \omega|_{\Sigma_\sigma}$ .

One can check that

$$X_{\#_\sigma}|_{\Sigma_\sigma} = d\mathcal{H}_\sigma(X)|_{\Sigma_\sigma} R_\sigma \quad \forall \sigma \in [0, 1].$$

So replace  $\mathcal{H}$  by  $f \circ \mathcal{H}$ , where  $f_\sigma$  is an extension of  $1/d\mathcal{H}_\sigma(X)$ ,  $f \in C^\infty(\mathcal{H}, (0, +\infty))$ .

Hence we have the period-action equality

$$\mathcal{A}^{\#_\sigma}(\gamma_\sigma, \tau_\sigma) = \int_0^1 \gamma_\sigma^* \lambda - \tau_\sigma \int_0^1 \mathcal{H}_\sigma \circ \gamma_\sigma \quad - 0$$

$$= \int_0^1 \gamma_\sigma^* \lambda$$

$$= \int_0^1 \lambda(\dot{\gamma}_\sigma)$$

$$\begin{aligned} (*) \quad &= \tau_\sigma \int_0^1 \lambda(R_\sigma \circ \gamma_\sigma) \\ &= \tau_\sigma \end{aligned}$$

Reeb fields satisfy  
 $iR_\sigma d\lambda = 0$   
 and  $\lambda(R_\sigma) = 1$ .

Now we can compute further because  $(\gamma_\sigma, \tau_\sigma)$  is a critical point of  $\mathcal{A}^{\#_\sigma}$ .

$$\partial_\sigma \tau_\sigma = \partial_\sigma (\mathcal{A}^{\#_\sigma}(\gamma_\sigma, \tau_\sigma))$$

$$= (\partial_\sigma \mathcal{A}^{\#_\sigma})(\gamma_\sigma, \tau_\sigma) + d\mathcal{A}^{\#_\sigma}|_{(\gamma_\sigma, \tau_\sigma)}(\partial_\sigma(\gamma_\sigma, \tau_\sigma))$$

$$= -\tau_\sigma \int_0^1 (\partial_\sigma \mathcal{H}_\sigma) \circ \gamma_\sigma.$$

In particular

$$|\partial_\sigma T_\sigma| \leq k T_\sigma \quad \forall \sigma \in [0, 1),$$

where  $k$  is such that

$$\|\partial_\sigma H_\sigma\|_{L^\infty(H^{-1}(0))} \leq k.$$

Integrating / Gronwall's inequality yields

$$T_0 e^{-k} \leq T_\sigma \leq T_0 e^k \quad \forall \sigma \in [0, 1).$$

### Local Rabinowitz-Floer Homology

One can apply local Rabinowitz-Floer homology to the  $\omega$ -limit set  $\omega(\gamma_0, T_0)$  as it is compact and connected in the contact (stable) case.

To simplify the discussion, assume that  $H \in C^\infty(H \times [0, 1])$  is parametrised by energy.

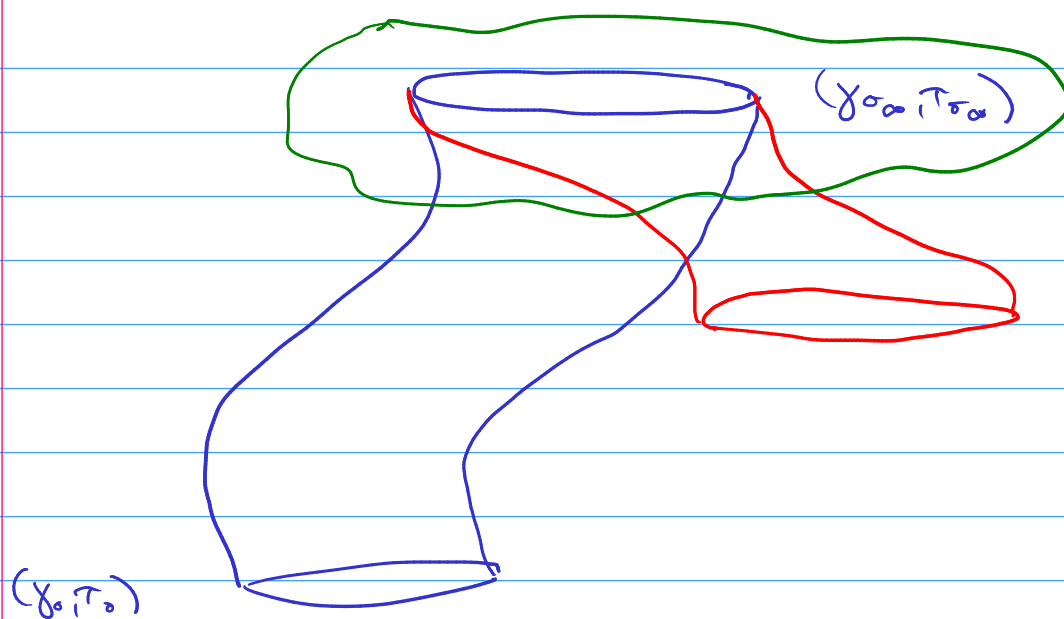
Assume  $(\gamma_0, T_0)$  is a nondegenerate parametrised periodic orbit on  $\Sigma_0$ . By regular orbit cylinder theorem, there exists  $0 < \sigma_\infty \leq 1$  and a smooth family of nondegenerate orbits  $(\gamma_\sigma, T_\sigma)_{\sigma \in [0, \sigma_\infty)}$ .



Roughly speaking, there are two options for the family.

- ① the family extends across  $\sigma_\infty$ .
- ②  $(\gamma_{\sigma_\infty}, T_{\sigma_\infty})$  is degenerate and there exists another family of parametrised periodic orbits with the same  $\omega$ -limit set.

↑  
Forcing Theorem



For a construction of local  $R\#H$  see the thesis by Kathrin Naef ( $R\#H$ ), 2018.  
 There must exist such a second family "killing" the first one as

$$R\#H^{\text{loc}}(\Sigma_\infty, \lambda|_{\Sigma_\infty}, w(\gamma_0, \tau_0)) = 0.$$

Floer chain complex is generated by 1-periodic orbits

Rabinowitz chain complex is generated by  $\tau$ -periodic orbits on fixed energy hypersurface.