

MATHEMATICAL METHODS OF QUANTUM MECHANICS SUMMARY

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Abstract.

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Postulates of Quantum Mechanics

<i>Quantum mechanical system</i>	Hilbert space \mathcal{H}
<i>State</i>	$\psi \in \mathcal{H}, \ \psi\ = 1$
<i>Observables</i>	Self-adjoint operators on \mathcal{H}
<i>Expected Value</i> of an observable A in the state ψ	$\langle \psi, A\psi \rangle$
<i>Variance</i> of an observable A in the state ψ	$\Delta A_\psi := \langle \psi, A^2\psi \rangle - \langle \psi, A\psi \rangle^2$

Lemma 1.1 (Heisenberg Uncertainty Principle). *Let A and B two self-adjoint operators on a Hilbert space \mathcal{H} . Then for any state ψ*

$$\Delta A_\psi \Delta B_\psi \geq \frac{1}{4} |\langle \psi, [A, B]\psi \rangle|^2.$$

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Unbounded Operators

Definition 1.1 (Linear Operator). Let \mathcal{H} be a Hilbert space. A **(linear) operator on \mathcal{H}** is simply a linear map $A : D(A) \rightarrow \mathcal{H}$, where $D(A)$ is a linear subspace of \mathcal{H} .

Definition 1.2 (Closed Operator). An operator A on \mathcal{H} is said to be **closed**, iff Γ_A is closed in $\mathcal{H} \times \mathcal{H}$.

Definition 1.3 (Closable Operator). An operator A on \mathcal{H} is said to be **closable**, iff $\bar{\Gamma}_A$ is a linear graph, i.e. $(0, y) \in \Gamma_A$ implies $y = 0$. The corresponding operator associated to $\bar{\Gamma}_A$ is denoted by \bar{A} and called the **closure** of A . Clearly $A \subseteq \bar{A}$.

Definition 1.4 (Adjoint). Let A be a densely defined operator on \mathcal{H} . Set

$$D(A^*) := \{\psi \in \mathcal{H} : \exists \eta \in \mathcal{H} \text{ s.t. } \forall \varphi \in D(A) \langle A\varphi, \psi \rangle = \langle \varphi, \eta \rangle\},$$

and $A^*\psi := \eta$. The operator A^* is called the **adjoint** of A .

Theorem 1.1. Let A be a densely defined operator on \mathcal{H} . Then:

- (a) A^* is closed.
- (b) A is closable if and only if $D(A^*)$ is dense.
- (c) If A is closable, then $(\bar{A})^* = A^*$.

Definition 1.5 (Symmetric Operator). A densely defined operator A is said to be **symmetric**, iff $A \subseteq A^*$.

Definition 1.6 (Self-adjoint Operator). A densely defined operator is said to be **self-adjoint**, iff $A = A^*$.

Definition 1.7 (Essentially Self-adjoint Operator). A symmetric operator A is said to be **essentially self-adjoint**, iff \bar{A} is self-adjoint.

The Spectral Theorem

Projection Valued Measures.

Definition 1.8 (Projection Valued Measure). Let \mathcal{H} be a Hilbert space. A function $P : \mathcal{B}(\mathbb{R}) \rightarrow \mathcal{L}(\mathcal{H})$ is said to be a **projection valued measure**, iff

- (i) For all $\Omega \in \mathcal{B}(\mathbb{R})$, $P(\Omega)$ is an orthogonal projection, i.e.

$$P(\Omega)^2 = P(\Omega) = P(\Omega)^*.$$

- (ii) $P(\mathbb{R}) = \text{id}_{\mathcal{H}}$.

- (iii) If $(\Omega_n)_{n \in \omega}$ is a sequence of pairwise disjoint elements of $\mathcal{B}(\mathbb{R})$, then

$$P(\Omega)\psi = \sum_{n \in \omega} P(\Omega_n)\psi,$$

for all $\psi \in \mathcal{H}$.

Definition 1.9 (Resolution of the Identity). Let \mathcal{H} be a Hilbert space and $P : \mathcal{B}(\mathbb{R}) \rightarrow \mathcal{L}(\mathcal{H})$ a projection valued measure. The function $p : \mathbb{R} \rightarrow \mathcal{L}(\mathcal{H})$ defined by

$$p(\lambda) := P(-\infty, \lambda],$$

is called the *resolution of the identity associated to a projection valued measure*.

The Spectral Theorem.

Theorem 1.2 (Spectral Theorem). Let A be a self-adjoint operator. Then there exists a unique projection valued measure P^A such that $D(A) = \{\psi \in \mathcal{H} : \int |\lambda|^2 d\mu_\psi^A(\lambda)\}$ and

$$A = \int \lambda dp^A(\lambda).$$

The Schrödinger Equation.

Theorem 1.3. Let \mathcal{H} be a Hilbert space and $H : D(H) \rightarrow \mathcal{H}$ be self adjoint. Moreover, set $U(t) := \exp(-iHt)$ for $t \in \mathbb{R}$. Then:

- (a) $U(t)$ is a strongly continuous one parameter unitary group.
- (b) The limit

$$\lim_{t \rightarrow 0} \frac{U(t) - 1}{t} \psi$$

exists if and only if $\psi \in D(H)$. Then

$$\lim_{t \rightarrow 0} \frac{U(t) - 1}{t} \psi = -iH\psi.$$

- (c) $U(t)D(H) = D(H)$ and $[U(t), H] = 0$ on $D(H)$.
- (d) Let $\psi_0 \in D(H)$. Then $U(t)\psi_0$ uniquely solves the initial value problem

$$\begin{cases} i \partial_t \psi(t) = H\psi(t) \\ \psi(0) = \psi_0, \end{cases} \quad (1)$$

called the *Schrödinger equation*.