

Crank-Nicolson-Julia-1

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1 Derivation of the Crank–Nicolson Method

This notebook cell outlines a step-by-step derivation of the Crank–Nicolson scheme for the 1D heat equation:

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < L, \quad t > 0.$$

1.1 1. The Heat Equation

We consider the one-dimensional heat (diffusion) equation on $([0, L])$. We partition the spatial domain into (N) intervals of width $\Delta x = \frac{L}{N}$, so

$$x_i = i \Delta x, \quad i = 0, 1, \dots, N.$$

We partition the time domain into steps of size Δt , so

$$t^n = n \Delta t, \quad n = 0, 1, 2, \dots$$

We denote

$$u_i^n \approx u(x_i, t^n).$$

1.2 2. Discretize the Time Derivative

Approximate the time derivative by a forward difference:

$$\frac{\partial u}{\partial t}(x_i, t^n) \approx \frac{u_i^{n+1} - u_i^n}{\Delta t}.$$

1.3 3. Discretize the Space Derivative

Approximate the second spatial derivative by a central difference:

$$\frac{\partial^2 u}{\partial x^2}(x_i, t^n) \approx \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2}.$$

1.4 4. Crank–Nicolson: Averaging Explicit and Implicit

Crank–Nicolson is formed by **averaging** the spatial derivatives at time (n) (explicit) and (n+1) (implicit). So,

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \frac{\alpha}{2} \left[\frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{(\Delta x)^2} + \frac{u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}}{(\Delta x)^2} \right].$$

1.5 5. Introduce the Mesh Ratio

Define the mesh ratio

$$r = \frac{\alpha \Delta t}{(\Delta x)^2}.$$

Multiply both sides of the Crank–Nicolson equation by Δt and collect like terms:

$$u_i^{n+1} - u_i^n = \frac{r}{2} \left[(u_{i+1}^n - 2u_i^n + u_{i-1}^n) + (u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}) \right].$$

After rearranging, we get:

$$-\frac{r}{2} u_{i-1}^{n+1} + (1+r) u_i^{n+1} - \frac{r}{2} u_{i+1}^{n+1} = \frac{r}{2} u_{i-1}^n + (1-r) u_i^n + \frac{r}{2} u_{i+1}^n.$$

1.6 6. Tridiagonal System Form

For $i = 1, 2, \dots, N-1$, this defines a **tridiagonal linear system** in the unknowns $\{u_i^{n+1}\}$. Denote the left-hand side as $A\mathbf{u}^{n+1}$ and the right-hand side as $B\mathbf{u}^n$. We can solve it at each time step using the **Thomas algorithm**:

$$A\mathbf{u}^{n+1} = B\mathbf{u}^n.$$

1.7 7. Final Scheme

Hence, the Crank–Nicolson scheme for the 1D heat equation is:

$$-\frac{r}{2} u_{i-1}^{n+1} + (1+r) u_i^{n+1} - \frac{r}{2} u_{i+1}^{n+1} = \frac{r}{2} u_{i-1}^n + (1-r) u_i^n + \frac{r}{2} u_{i+1}^n,$$

subject to your chosen boundary conditions (Dirichlet, Neumann, etc.) at $(i = 0)$ and $(i = N)$.

1.8 Matrix Form

We can write this entire system for $(i = 1, 2, \dots, N-1)$ in **matrix form** as:

$$A \mathbf{u}^{n+1} = B \mathbf{u}^n,$$

where \mathbf{u}^n is the vector $[u_1^n, u_2^n, \dots, u_{N-1}^n]^T$,
and (A) and (B) are **tridiagonal** matrices of size $(N-1) \times (N-1)$.

1.8.1 The Matrix (A)

$$A = \begin{bmatrix} 1+r & -\frac{r}{2} & 0 & \cdots & 0 \\ -\frac{r}{2} & 1+r & -\frac{r}{2} & \ddots & \vdots \\ 0 & -\frac{r}{2} & 1+r & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & -\frac{r}{2} \\ 0 & \cdots & 0 & -\frac{r}{2} & 1+r \end{bmatrix}.$$

1.8.2 The Matrix (B)

$$B = \begin{bmatrix} 1-r & \frac{r}{2} & 0 & \cdots & 0 \\ \frac{r}{2} & 1-r & \frac{r}{2} & \ddots & \vdots \\ 0 & \frac{r}{2} & 1-r & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & \frac{r}{2} \\ 0 & \cdots & 0 & \frac{r}{2} & 1-r \end{bmatrix}.$$

Hence, the **Crank–Nicolson update** for each time step is:

$$\mathbf{u}^{n+1} = A^{-1} B \mathbf{u}^n,$$

though in practice we do **not** invert (A) directly; rather, we solve the linear system

$$A \mathbf{u}^{n+1} = B \mathbf{u}^n$$

efficiently using the **Thomas algorithm** (a specialized $(O(N))$ Gaussian elimination for tridiagonal matrices).

Key points: - This is **second-order accurate** in both time and space. - **Unconditionally stable** for linear diffusion problems. - Solved via a tridiagonal linear system at each timestep.

2 Thomas Algorithm: Step-by-Step Derivation

We want to solve a **tridiagonal linear system** of n equations:

$$\begin{aligned}
b_1 x_1 + c_1 x_2 &= d_1, \\
a_2 x_1 + b_2 x_2 + c_2 x_3 &= d_2, \\
a_3 x_2 + b_3 x_3 + c_3 x_4 &= d_3, \\
&\vdots \\
a_n x_{n-1} + b_n x_n &= d_n.
\end{aligned}$$

Here: - a_i are the subdiagonal entries ($i = 2, \dots, n$; often we take $a_1 = 0$ by convention), - b_i are the main diagonal entries ($i = 1, \dots, n$), - c_i are the superdiagonal entries ($i = 1, \dots, n-1$; often $c_n = 0$ by convention), - d_i are the constants on the right-hand side.

Our goal is to find x_1, x_2, \dots, x_n in $\mathcal{O}(n)$ time. The procedure involves two phases:

1. **Forward Elimination** – to remove the subdiagonal a_i .
 2. **Back Substitution** – to solve the resulting upper-triangular system.
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2.1 1. Forward Elimination

We sequentially eliminate a_2, a_3, \dots, a_n by modifying each row in turn.

2.1.1 Step 1 (Row 1)

The first equation is

$$b_1 x_1 + c_1 x_2 = d_1.$$

We **normalize** this row by dividing through by b_1 (assuming $b_1 \neq 0$):

$$x_1 + \frac{c_1}{b_1} x_2 = \frac{d_1}{b_1}.$$

Define the **modified** superdiagonal and right-hand side for row 1:

$$c'_1 = \frac{c_1}{b_1}, \quad d'_1 = \frac{d_1}{b_1}.$$

Hence the first row (in **modified** form) is:

$$x_1 + c'_1 x_2 = d'_1.$$

2.1.2 Step 2 (Row 2)

The second equation in the original system is

$$a_2 x_1 + b_2 x_2 + c_2 x_3 = d_2.$$

We want to **eliminate** x_1 from this equation using the (already normalized) Row 1.

1. Multiply the (modified) Row 1 by a_2 and **subtract** from Row 2:

$$a_2(x_1 + c'_1 x_2) = a_2 d'_1.$$

Subtracting, we get:

$$(b_2 - a_2 c'_1) x_2 + c_2 x_3 = d_2 - a_2 d'_1.$$

2. Define

$$\begin{aligned} b'_2 &= b_2 - a_2 c'_1, \\ d'_2 &= d_2 - a_2 d'_1. \end{aligned}$$

So the second equation becomes

$$b'_2 x_2 + c_2 x_3 = d'_2.$$

3. **Normalize** this row by dividing through by b'_2 :

$$x_2 + \frac{c_2}{b'_2} x_3 = \frac{d'_2}{b'_2}.$$

Define:

$$c'_2 = \frac{c_2}{b'_2}, \quad d'_2 = \frac{d'_2}{b'_2}.$$

Hence Row 2 is now:

$$x_2 + c'_2 x_3 = d'_2.$$

2.1.3 Step 3 (General Row i)

For $i = 3, 4, \dots, n$, we proceed **similarly**:

1. The i -th equation (before modification) is

$$a_i x_{i-1} + b_i x_i + c_i x_{i+1} = d_i.$$

2. We have already normalized Row $i - 1$, so we can subtract a_i times that row to eliminate x_{i-1} . Symbolically:

$$b'_i = b_i - a_i c'_{i-1}, \quad d'_i = d_i - a_i d'_{i-1}.$$

3. **Normalize** Row i by dividing through by b'_i :

$$c'_i = \frac{c_i}{b'_i}, \quad d'_i = \frac{d_i}{b'_i}.$$

After processing row n , the subdiagonal entries are effectively **eliminated**. The system is now **upper-triangular**, represented by:

$$x_i + c'_i x_{i+1} = d'_i, \quad (i = 1, 2, \dots, n-1).$$

For the last equation ($i = n$), it simply becomes:

$$x_n = d'_n.$$

2.2 2. Back Substitution

Now we solve for the x_i **from the bottom up**:

1. **Initialize** the solution at the last row:

$$x_n = d'_n.$$

2. **For** $i = n-1, n-2, \dots, 1$:

$$x_i = d'_i - c'_i x_{i+1}.$$

This recovers all the unknowns x_1, \dots, x_n in a single **backward** pass.

2.3 3. Summary of the Thomas Algorithm

- **Forward Elimination:**

For $i = 1$ to n :

1. Normalize the current row by its diagonal (b'_i).
2. Use the normalized row to eliminate a_{i+1} in the next row.

- **Back Substitution:**

Start from $x_n = d'_n$ and move backward using $x_i = d'_i - c'_i x_{i+1}$.

2.3.1 Complexity

Because each step uses only a few arithmetic operations per row, the **Thomas algorithm** runs in $\mathcal{O}(n)$ time, much faster than $\mathcal{O}(n^3)$ for generic Gaussian elimination.

2.4 4. Final Formulas

Putting it all together, if we define:

$$\begin{aligned}
 \text{(Initialization)} \quad c'_1 &= \frac{c_1}{b_1}, & d'_1 &= \frac{d_1}{b_1}, \\
 \text{(Forward sweep for } i = 2, \dots, n) \quad b'_i &= b_i - a_i c'_{i-1}, & d'_i &= d_i - a_i d'_{i-1}, \\
 & c'_i = \frac{c_i}{b'_i}, & d'_i &= \frac{d'_i}{b'_i}, \\
 \text{(Back Substitution)} \quad x_n &= d'_n, & x_i &= d'_i - c'_i x_{i+1} \quad (i = n-1, \dots, 1),
 \end{aligned}$$

then $\{x_i\}$ is the unique solution of the original **tridiagonal system**.

That is the Thomas algorithm in a **step-by-step** derivation, showing how we systematically eliminate subdiagonal terms and then back-substitute.

```
[29]: using LinearAlgebra
      using PyPlot
```

```
[30]: function solve_by_thomas_algorithm(a,b,c,d)
      """
      Solve a tridiagonal system A x = d using the Thomas algorithm.
      a, b, c are the lower, main, and upper diagonals (each 1D arrays).
      d is the right-hand side array.

      Returns the solution array x.
      """
      N = length(b)
      bp = zeros(Float64,N)
      cp = zeros(Float64,N)
      dp = zeros(Float64,N)
      xs = zeros(Float64,N)
      cp[1] = c[1]/b[1]; dp[1] = d[1]/b[1]
      for i = 2:N
          bp[i] = b[i] - a[i]*cp[i-1]
          dp[i] = d[i] - a[i]*dp[i-1]
          cp[i] = c[i]/bp[i]
          dp[i] = dp[i]/bp[i]
      end
      xs[end] = dp[end]
      for i=N-1:-1:1
          xs[i] = dp[i] - cp[i]*xs[i+1]
      end
      return xs
  end
```

```

"""
params =  $\alpha$ , L, Nx, Nt,  $\Delta t$ , u_boundary, u_ini, r
"""
function solve_heat_equation(params)
    """
        Solve the 1D heat equation using the Crank-Nicolson method with Thomas_
         $\rightarrow$  algorithm.

        params = ( $\alpha$ , L, Nx, Nt,  $\Delta t$ , u_boundary, u_ini, r)
             $\alpha$           -> Thermal diffusivity
            L            -> Length of the domain
            Nx           -> Number of spatial grid points
            Nt           -> Number of time steps
             $\Delta t$        -> Time step size
            u_boundary   -> Tuple (u_left, u_right) for Dirichlet BCs
            u_ini        -> Initial temperature profile (1D numpy array of length Nx)
            r            ->  $\alpha \Delta t / (\Delta x^2)$ 

        Returns:
            ts          -> 1D array of time values
            solutions   -> 2D array of shape (Nx, Nt+1),
                        where solutions[:, k] is the solution at time step k.
    """

    # Retrieve the parameters
     $\alpha$ , L, Nx, Nt,  $\Delta t$ , u_boundary, u_ini, r = params

    # Load initial values
    u = u_ini

    # Define time points for simulation
    ts = range(0.0, length=Nt, step= $\Delta t$ )

    # Variable for storing the solution (Nt+1 as we also store initial value)
    solutions = zeros(Float64, Nx, Nt+1)

    # Time loop
    for i=1:Nt

        # Store
        solutions[:,i] = u

        # For RHS
        d = zeros(Float64, Nx)
        for i=1:Nx
            d[i] = (1-r)*u[i]

```



```

        if i>1
            d[i] += r/2*u[i-1]
        end
        if i<Nx
            d[i] += r/2*u[i+1]
        end
    end

    # Adjust Boundary conditions
    d[1] += r*u_boundary[1]
    d[end] += r*u_boundary[end]

    # For LHS
    b = (1+r)*ones(Nx)
    a = [0; -r/2*ones(Nx-1)]
    c = [-r/2*ones(Nx-1); 0]

    u = solve_by_thomas_algorithm(a,b,c,d)

    # Fix boundary values
    u[1], u[end] = u_boundary
end

# Store the last solution
solutions[:,Nt+1] = u;

# Return time points and solutions
return ts, solutions
end

```

[30]: solve_heat_equation

```

[33]:  $\alpha$  = 1.0e-4
L = 1.0
Nx = 101
Nt = 2000
 $\Delta t$  = 0.1
 $\sigma$  = 0.05
u_boundary = (300.0,300.0)

 $\Delta x$  = L/(Nx-1)
r =  $\alpha * \Delta t / (\Delta x * \Delta x)$ 

xs = range(0.0, L, length=Nx)

# Initial condition: Gaussian bump over 300 K

```

```

u_ini = @. 300 + 100*exp(-(xs - L/2)^2/2/σ^2)

# Pack trial parameters
params = α, L, Nx, 2, Δt, u_boundary, u_ini, r

# Compile
@time ts, solutions = solve_heat_equation(params);

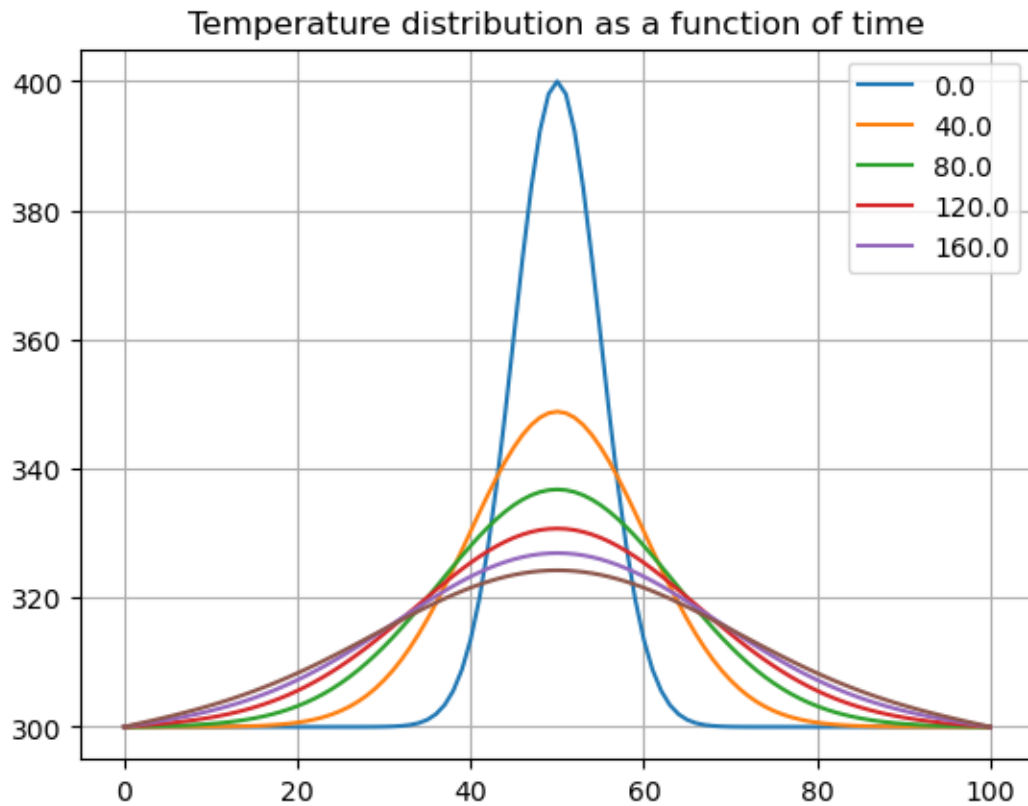
# Pack parameters
params = α, L, Nx, Nt, Δt, u_boundary, u_ini, r

# Final run
@time ts, solutions = solve_heat_equation(params);

# Plot
plot(solutions[:,1:400:end]); legend(ts[1:400:end])
title("Temperature distribution as a function of time")
grid()

```

0.000020 seconds (34 allocations: 25.656 KiB)
0.004092 seconds (26.01 k allocations: 23.759 MiB)



[]: