

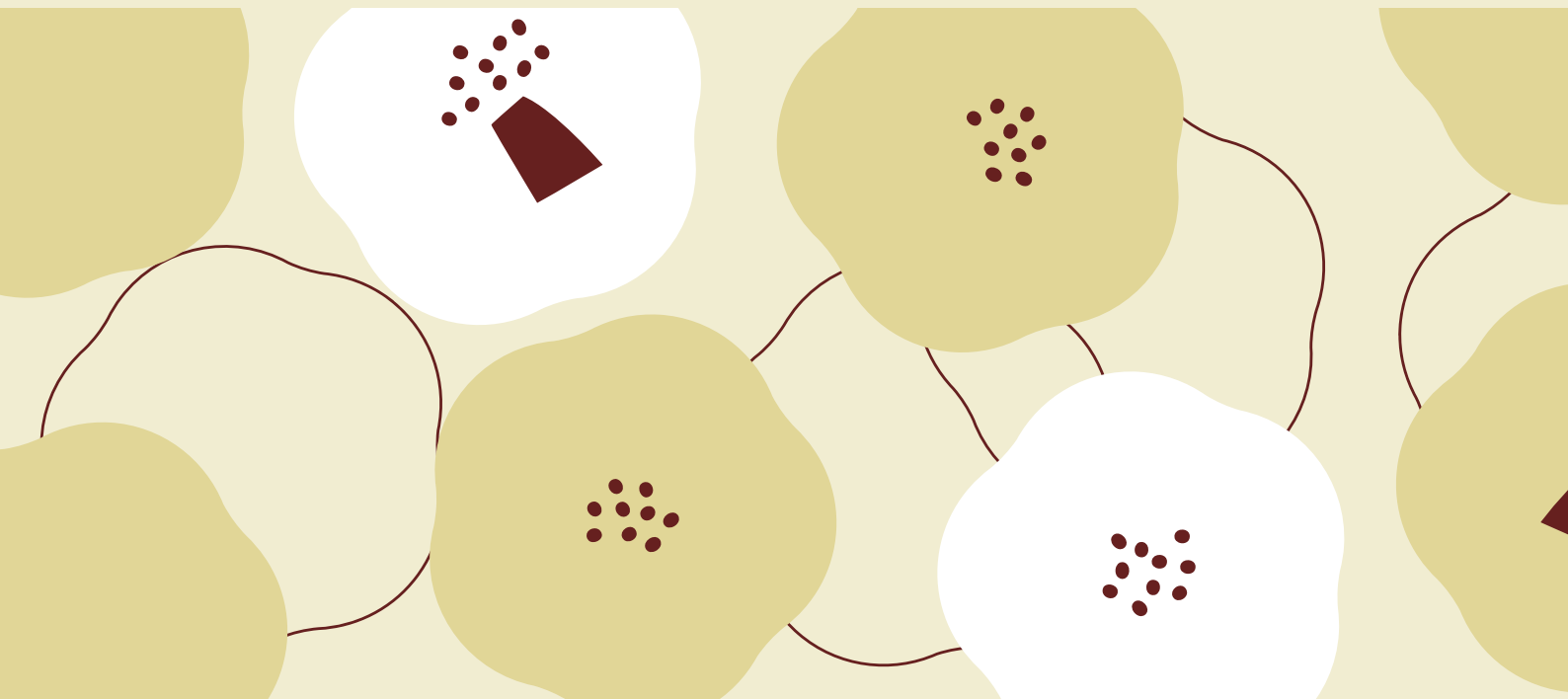


Example For PDE

with

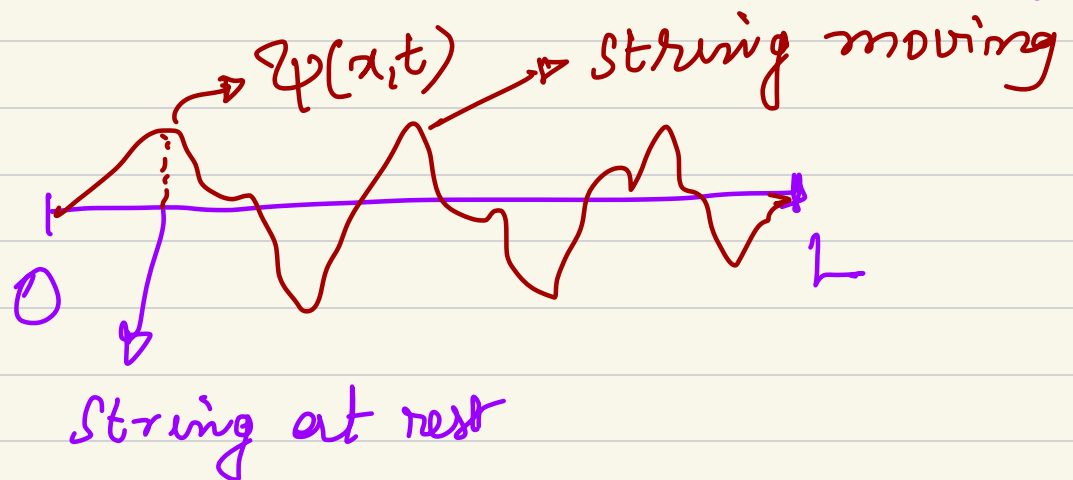
Boundary condition

MA2103, 2023



# A Vibrating String

A string of length ' $L$ ' is fixed at both end is left with an initial shape of  $\psi_0(x)$  find the motion of string as function of  $x, t$ .



This motion of string can be modelled by wave equation

$$\frac{\partial^2 \psi}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2} = 0$$

speed of the wave

the boundary condition given are

$$\psi(x=0, t) = 0 \text{ and } \psi(x=L, t) = 0$$

both end of string is at rest all the time

$$\text{Also } \psi(x, t=0) = \psi_0(x)$$

$$\text{we also take } \frac{\partial \psi}{\partial t} \bigg|_{t=0} = 0$$

we start with PDE

$$\frac{\partial^2 \psi}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 \psi}{\partial t^2} = 0$$


$\psi(x, t)$  is assumed to be separable

i.e.  $\psi(x, t) = \psi_x \psi_t$

Substituting in PDE and after some simplification, we get

$$\underbrace{\frac{1}{\psi_x} \frac{\partial^2 \psi_x}{\partial x^2}}_{(1)} - \underbrace{\frac{1}{c^2 \psi_t} \frac{\partial^2 \psi_t}{\partial t^2}}_{(2)} = 0$$

Since term (1) is purely function of  $x$ , and term (2) is purely a function of  $t$ , both must be a const.

Let  $-\alpha^2$  be constant of separability 

From (1) we get

$$\frac{1}{\psi_x} \frac{\partial^2 \psi_x}{\partial x^2} + \alpha^2 = 0$$

$$\frac{\partial^2 \psi_x}{\partial x^2} + \alpha^2 \psi_x = 0$$

from (2) we get

$$\frac{1}{c^2 \psi_t} \frac{\partial^2 \psi_t}{\partial t^2} + \alpha^2 = 0$$

Since  $\psi_x$  is a smooth sol<sup>n</sup> in the interval, it can be expressed in terms of Fourier series in interval  $[0, L]$  as

$$\psi_x = a_0 + \sum_{n=1}^{\infty} \left\{ a_n \cos\left(\frac{2\pi n x}{L}\right) + b_n \sin\left(\frac{2\pi n x}{L}\right) \right\}$$

$$\sin \psi_x(0) = \psi_x(L) = 0 \quad a_0 = 0, \text{ also } a_n = 0 \forall n$$

We left with

$$\psi_x = \sum_{n=1}^{\infty} b_n \sin\left(\frac{2\pi n x}{L}\right)$$

To fix the constant, we substitute  $\psi_x$  in the ODE

$$\frac{d^2 \psi_x}{dx^2} + \alpha^2 \psi_x = 0$$

$$\sum_{n=1}^{\infty} \left\{ -a_n \left( \frac{2\pi n}{L} \right)^2 \sin \left( \frac{2\pi n}{L} x \right) + \alpha^2 a_n \sin \left( \frac{2\pi n}{L} x \right) \right\} = 0$$

$$\sum_{n=1}^{\infty} \left[ \alpha^2 - \left( \frac{2\pi n}{L} \right)^2 \right] a_n \sin \left( \frac{2\pi n}{L} x \right) = 0$$

This can happen only if

$$\alpha^2 = \left( \frac{2\pi n}{L} \right)^2 \quad \text{i.e } \alpha \text{ depends on } n$$

$$\alpha_n = \frac{2\pi n}{L}$$

Now we look at time part

$$\frac{d^2 \psi_t}{dt^2} + \alpha^2 c^2 \psi_t = 0$$

and the sol<sup>n</sup> is

$$A_n \cos(\alpha_n c t) + B_n \sin(\alpha_n c t)$$

$$\text{At } t=0 \text{ we need } \left. \frac{\partial \psi}{\partial t} \right|_{t=0} \text{ should be '0'}$$

This condition gives us  $B_n = 0$

$$\psi_t = A_n \cos(\alpha_n c t)$$

The complete sol<sup>n</sup> is given by

$$\psi(x,t) = \sum_{n=1}^{\infty} a'_n \sin(\alpha_n x) \cos(\alpha_n ct)$$

$$\alpha_n = \frac{2\pi n}{L}$$

finally at  $t=0$

$$\psi(x,t=0) = \psi_0(x) = \sum_{n=1}^{\infty} a'_n \sin \alpha_n x$$

from this we get-

$$a'_n = \frac{2}{L} \int_0^L \sin(\alpha_n x) \psi_0(x) dx$$

This is the end !

