PH3102: QM Assignment 08

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1 A system of two Interacting Spins

Answer 1.1

There are 2 independent spins $\vec{S}_1$, $\vec{S}_2$ which have two independent Hilbert spaces assosciated with them. Thus this systems having 2 spins will be described by the tensor product of both Hilbert spaces giving us the labelling $|S_1, m_{S_1}\rangle$ and $|S_2, m_{S_2}\rangle$. Thus to label all the states of the system we need 4 eigenstates.

Answer 1.2

The hamiltonian is given as

$$\begin{aligned} \widehat{H} &= J\hat{\boldsymbol{S}}_1 \cdot \hat{\boldsymbol{S}}_2 \\ &= \frac{J}{2}\left[\hat{\boldsymbol{S}}^2_{tot} - \hat{\boldsymbol{S}}^2_1 - \hat{\boldsymbol{S}}^2_2\right] \end{aligned} \tag{1}$$

Let us write our eigenstates $|S_1, m_{S_1}\rangle \otimes |S_2, m_{S_2}\rangle$ as the tensor product of the eigenstates with the good quantum numbers.

Answer 1.3

Since the Spin operators $\hat{\boldsymbol{S}}_1$ and $\hat{\boldsymbol{S}}_2$ acts on independent Hilbert spaces, the operators are independent of each other. We thus have $\left[S_{1i}, S_{2j}\right] = 0$ where $i, j \in \{x, y, z\}$. This gives us the commutation relations $\left[\hat{\boldsymbol{S}}_{tot}, \hat{\boldsymbol{S}}_i\right] = 0, i \in \{1, 2\}$. Thus we see that $\hat{S}_{tot}, \hat{S}_1$ and $\hat{S}_2$ all commute with the Hamiltonian, giving us

$$\left[H, \hat{S}_{tot}\right] = \left[H, \hat{S}_1\right] = \left[H, \hat{S}_2\right] = 0 \tag{2}$$

. We know that we need 4 operators to form a Complete Set of Commuting Observables(CSCO). Let us define $\hat{S}_z = \hat{S}_{1z} + \hat{S}_{2z}$. We know that Spin angular momentum also follows total angular momentum commutation relations. From the addition of angular momenta relations, we get

$$\begin{aligned} &\left[\hat{\boldsymbol{S}}^2_{tot}, S_z\right] = 0 \\ &\Rightarrow [H, S_z] = 0 \end{aligned} \tag{3}$$

Let us denote our eigenstates as $|\ s, m_s; s_1, s_1\rangle$ where we have

$$
\begin{aligned}
\hat{\boldsymbol{S}}^2_{tot}|s, m_s; s_1, s_2\rangle &= s(s+1)\hbar^2|s, m_s; s_1, s_2\rangle \\
\hat{\boldsymbol{S}}^2_1|s, m_s; s_1, s_2\rangle &= s_1(s_1+1)\hbar^2|s, m_s; s_1, s_2\rangle = \frac{3}{4}\hbar^2|s, m_s; s_1, s_2\rangle \\
\hat{\boldsymbol{S}}^2_2|s, m_s; s_1, s_2\rangle &= s_2(s_2+1)\hbar^2|s, m_s; s_1, s_2\rangle = \frac{3}{4}\hbar^2|s, m_s; s_1, s_2\rangle \\
S_z|s, m_s; s_1, s_2\rangle &= m_s\hbar\ |s, m_s; s_1, s_2\rangle
\end{aligned} \tag{4}
$$

We can then write our reformed hamiltonian as

$$
\widehat{H} = \frac{J}{2}\left(\hat{\boldsymbol{S}}^2_{tot} - \frac{3}{2}\hbar^2\mathbb{I}\right) \tag{5}
$$

We can see that the energy eigenvalue is independent of m_s, thus the levels with same m_s are degenerate. We have already seen before that this is the addition of two spin 1/2 momenta we thus get the eigenstates

<u>Singlet State</u>

$$
|0,0\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \quad \left[E_{S=0} = -\frac{3}{4}J\hbar^2\right] \tag{6}
$$

<u>Triplet State</u>

$$
\left.\begin{aligned}
|1,1\rangle &= |\uparrow\uparrow\rangle \\
|1,0\rangle &= \tfrac{1}{\sqrt{2}}[|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle] \\
|1,-1\rangle &= |\downarrow\downarrow\rangle
\end{aligned}\right\} = \left[E_{S=1} = \frac{1}{4}J\hbar^2\right] \tag{7}
$$

Answer 1.4

The operator $\hat{\boldsymbol{S}}_{tot}(\hat{n}) = n_x S_x + n_y S_y + n_z S_z$. Let us now calculate the expectation value of $\hat{\boldsymbol{S}}_{tot}(\hat{n})$ in the ground singlet state.

$$
\begin{aligned}
\left\langle \hat{\boldsymbol{S}}_{tot}(\hat{n}) \right\rangle &= n_x\langle 0,0|S_x|0,0\rangle + n_y\langle 0,0|S_y|0,0\rangle + n_z\langle 0,0|S_z|0,0\rangle \\
&= \frac{n_x}{2}\langle 0,0|S_+ + S_-|0,0\rangle + \frac{n_y}{2i}\langle 0,0|S_+ - S_-|0,0\rangle \\
&= 0
\end{aligned} \tag{8}
$$

We can see that the expectation value is independent of the value of the $\hat{n}$ and thus is exhibits SU(2) symmetry.

Answer 1.5

We now calculate the magnetisations of the three excited states

$$
\begin{aligned}
\left\langle \hat{\boldsymbol{S}}_{tot}(\hat{n}) \right\rangle &= n_x\langle 1,m|S_x|1,m\rangle + n_y\langle 1,m|S_y|1,m\rangle + n_z\langle 1,m|S_z|1,m\rangle \\
&= n_z\langle 1,m|S_z|1,m\rangle
\end{aligned} \tag{9}
$$

The total magnetisation for the triplet excited states are

$$\left\langle \hat{\boldsymbol{S}}_{tot} \right\rangle = \begin{cases} n_z & m = 1 \\ 0 & m = 0 \\ -n_z & m = -1 \end{cases} \tag{10}$$

Classically ferromagnets are magnets with all dipoles aligned in the same direction. Thus the expectation value in the $\hat{z}$ direction is non zero. Thus we have $|1,1\rangle$ and $|1,-1\rangle$ to be classical ferromagnets.

To adequately describe the singlet state we would require another operator called the staggered magnetism operator which gives us anti-ferromagnets. This results in the singlet states being antiferromagnetic.

2 Broken symmetry and a quantum "phase" transition in a toy model

Answer 2.1

We have already calculated for the above Hamiltonian. The new Hamiltonian has a small modification over the previous one. Using Equation 5, The new Hamiltonian simplifies to

$$\begin{aligned} \widehat{H} &= J\left(\hat{\boldsymbol{S}}_1 \cdot \hat{\boldsymbol{S}}_2\right) + B(S_1^z + S_2^z) \\ &= \frac{J}{2}\left(\hat{\boldsymbol{S}}_{tot}^2 - \frac{3}{2}\hbar^2\mathbb{I}\right) + BS_z \end{aligned} \tag{11}$$

The eigenstates are the same with different eigenvalues for energy. We can see that the introduction of the magnetic field dependence lifts the degeneracy we saw in the previous step. The Energy eigenvalues are given below

Singlet State $\quad |0,0\rangle \longrightarrow E_{0,0} = -\frac{3}{4}J\hbar^2$

Triplet States

$$s = 1 : \begin{cases} |1,1\rangle & \longrightarrow E_{1,1} = \frac{1}{4}J\hbar^2 + B\hbar \\ |1,1\rangle & \longrightarrow E_{1,1} = \frac{1}{4}J\hbar^2 \\ |1,-1\rangle & \longrightarrow E_{1,-1} = \frac{1}{4}J\hbar^2 - B\hbar \end{cases} \tag{12}$$

Answer 2.2

We have $J > 0, B > 0$. From the plot it be can be easily seen that the state $|1,-1\rangle$ crosses the state $|0,0\rangle$

$$\frac{1}{4}J\hbar^2 - B\hbar = -\frac{3}{4}J\hbar^2 \Rightarrow B = J\hbar \tag{13}$$

Before the crossing the ground state was $|0,0\rangle$, which after crossing become $|1,-1\rangle$. This causes symmetry breaking. As we saw in Q1 $|0,0\rangle$ has SU(2) symmetry, but $|1,-1\rangle$ does not have SU(2) symmetry. Thus the ground state no longer has SU(2) symmetry.

$\boxed{\textbf{Answer 2.3}}$

If we take $J < 0$, as shown in the graph, there is no "phase" transition taking place. $|1, -1\rangle$ is always the ground state. In this case however, the ground state doesnt exhibit SU(2) symmetry.

This can also be expected as again if we consider the staggered magnetism operator, $|0, 0\rangle$ is an antiferromagnetic state. That means that the state has equal spins in opposite directions. However the state $|1, -1\rangle$ is ferromagnetic and has a preferred direction. This causes symmetry breaking.