#### LECTURE

# 7

*C*Y11001 Spring 2018

- Entropy and Spontaneity
- Carnot Engine



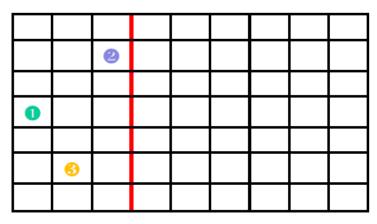
### **Statistical View of Entropy**

The equilibrium thermodynamic state of an isolated system is the most probable state.

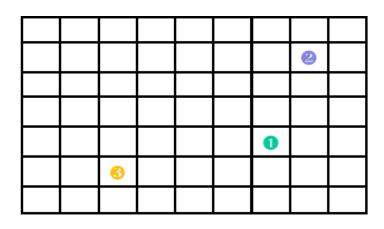
$$S = k_B \ln W$$

 $k_{\rm B}$  is Boltzmann's constant W is number of different ways in which the energy of the system can be arranged (number of microstates)

As  $T \rightarrow 0$ ,  $W \rightarrow 1$  and  $S \rightarrow 0$  (Nernst Heat Theorem/Third Law of TD)



$$W = 21 \times 20 \times 19 = 7,980$$



$$W = 63 \times 62 \times 61 = 2,38,266$$

### The Clausius Inequality:

$$dS = dq_{rev} / T$$
  
$$dS_{surr} = - dq / T_{surr}$$

$$\Delta S = \int_{1}^{2} \frac{dq_{rev}}{T}$$

$$\Delta S_{surr} = \frac{-q}{T_{surr}}$$

For expansion/compression work,  $dw - dw_{rev} \ge 0$ 

Since 
$$dU = dq + dw = dq_{rev} + dw_{rev}$$

$$dq_{\rm rev} \ge dq$$
  
 $dq_{\rm rev}/T \ge dq/T$ 

For a closed system,

$$dq = dU - dw$$

$$TdS \ge dU - dw$$

$$dU - TdS - dw \le 0$$

 $dS \ge dq/T$ 

 $TdS - dq \ge 0$ 

The Clausius Inequality

Entropy can not decrease when a spontaneous change occurs in an isolated system (dq = 0)

### **Entropy and Reversibility**

• Reversible process:

Reversible heat transfer between system and surrounding must occur with no finite temperature difference.

$$T_{syst} = T_{surr} \Rightarrow dS_{univ} = dS_{syst} + dS_{surr} = \frac{dq_{rev}}{T_{syst}} - \frac{dq}{T_{surr}} = 0$$

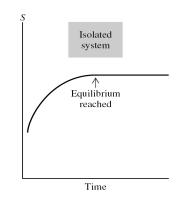
• Irreversible process (spontaneous):

The system + surrounding = universe, can be considered an isolated system. For any isolated system,  $dS_{\rm univ} = dS_{\rm syst} + dS_{\rm surr} > 0$ 

$$\Delta S_{univ} \ge 0$$

### **Entropy and Equilibrium**

• For an isolated system, spontaneous changes will occur until the entropy is maximized. This leads to equilibrium.

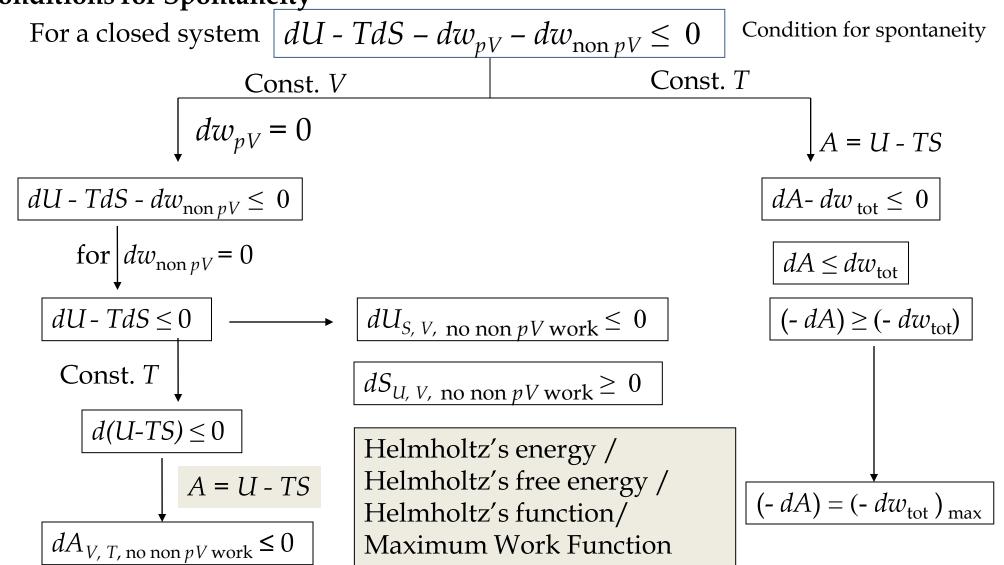


• For a closed system (heat and work exchange with surrounding is allowed), spontaneous changes will occur until the entropy of system plus surrounding, is maximized. This leads to equilibrium.

$$dS_{total} = dS_{sys} + dS_{surr} > 0$$
 Spontaneous process  $dS_{total} = dS_{sys} + dS_{surr} = 0$  Equilibrium process  $dS_{total} = dS_{sys} + dS_{surr} < 0$  Impossible process

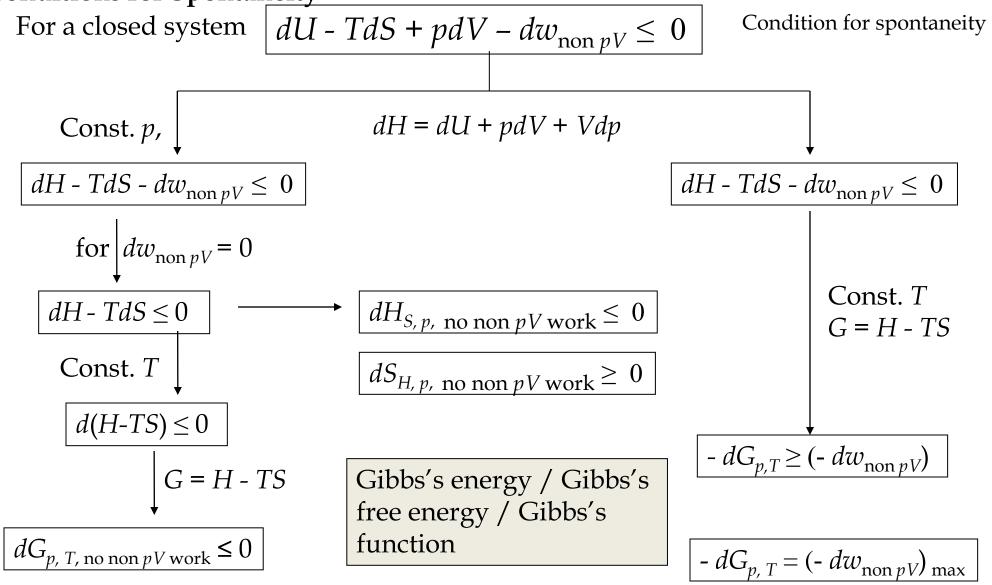
For any process 
$$dS_{sys} + dS_{surr} \ge 0$$
 "=" for reversible, equilibrium ">" spontaneous (irreversible) (real)





For a closed system at constant T and V, the state function A = U-TS decreases during the spontaneous irreversible process

**Conditions for Spontaneity** 



At constant temperature and pressure, chemical reactions are spontaneous in the direction of decreasing Gibbs energy.

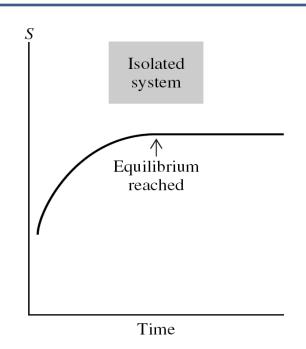
### Criteria for spontaneity

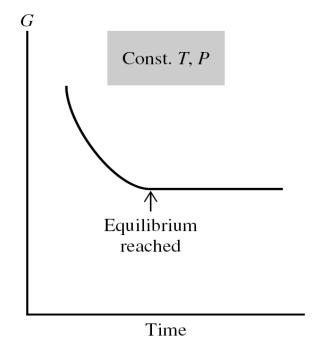
$$dS_{\text{sys}} + dS_{\text{surr}} > 0$$

For a closed system, no non p-V (additional) work

$$dU_{S,V} \le 0 \qquad dH_{S,p} \le 0 \qquad dA_{V,T} \le 0$$
$$dS_{U,V} \ge 0 \qquad dS_{H,p} \ge 0 \qquad dG_{p,T} \le 0$$

Reversible processes carry equal sign.





### **Heat Engine**

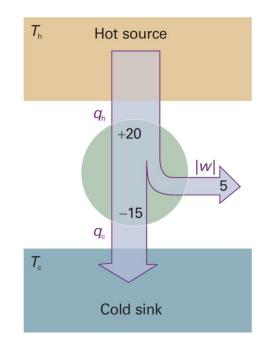
An engine is a device (system) that converts energy to work.

A heat engine draws heat from a hot reservoir, converts some heat to work, and releases some heat to a cold reservoir.

The engine itself is a system that undergoes a cyclical process

Experiments suggested:

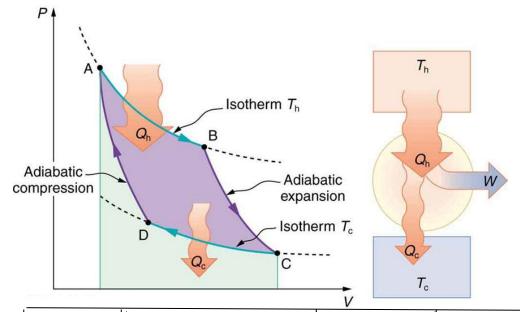
(Clausius) An engine does not exist whose sole effect is to transfer heat from a cold body to a hot body. (Kelvin) An engine does not exist that operates in a cycle and performs work by exchanging heat with only one reservoir!



**Fig. 3.7** Suppose an energy  $q_h$  (for example, 20 kJ) is supplied to the engine and  $q_c$  is lost from the engine (for example,  $q_c = -15$  kJ) and discarded into the cold reservoir. The work done by the engine is equal to  $q_h + q_c$  (for example, 20 kJ + (-15 kJ) = 5 kJ). The efficiency is the work done divided by the energy supplied as heat from the hot source.

#### Carnot Engine (1825, Sadi Carnot)

Ideal gas is used and all steps are reversible



$$TV^{\gamma-1} = \text{const.};$$

$$T_h V_B^{\gamma-1} = T_C V_C^{\gamma-1}$$

$$T_h V_A^{\gamma - 1} = T_C V_D^{\gamma - 1}$$

$$\left(\frac{V_B}{V_A}\right)^{\gamma-1} = \left(\frac{V_C}{V_D}\right)^{\gamma-1} \quad \left(\frac{V_B}{V_A}\right) = \left(\frac{V_C}{V_D}\right)$$

$$\ln \left( V_D / V_C \right) = -\ln \left( V_B / V_A \right)$$

			0 (0 : 5)	4 (D to A)	
	1 (A to B)	2 (B to C)	3 (C to D)	4 (D to A)	Total (A to A)
W	$-nRT_{\rm h} \ln(V_{\rm B}/V_{\rm A})$	$C_V(T_c-T_h)$	$-nRT_{\rm c} \ln(V_{\rm D}/V_{\rm C})$	$C_V(T_h-T_c)$	$-nR(T_h-T_c) \ln(V_B/V_A)$
q	$nRT_{\rm h} \ln(V_{\rm B}/V_{\rm A})$	0	$nRT_{\rm c} \ln(V_{\rm D}/V_{\rm C})$	0	$nR(T_{\rm h}-T_{\rm c}) \ln(V_{\rm B}/V_{\rm A})$
Δυ	0	$C_V(T_c-T_h)$	0	$C_V(T_h-T_c)$	0
(q <sub>rev</sub> / T)	nR In(V <sub>B</sub> /V <sub>A</sub> )	0	$nR \ln(V_{\rm D}/V_{\rm C})$	0	0

## **Efficiency of Carnot cycle**

$$\eta = \frac{\text{work performed}}{\text{heat absorbed from source}}$$

$$\eta = \frac{-w_{\text{tot}}}{q} = \frac{nR(T_h - T_c) \ln(V_B / V_A)}{nRT_h \ln(V_B / V_A)}$$

$$= \frac{T_h - T_c}{T_h} = 1 - \frac{T_c}{T_h} < 1$$

$$\eta = 1 - \frac{T_c}{T_h} = 1 + \frac{q_c}{q_h}$$

All reversible engines have same efficiency regardless of their construction

#### The total change in $q_{rev}/T$ during a Carnot cycle:

$$\frac{q_{rev}}{T} = \frac{q_h}{T_h} + \frac{q_c}{T_c} = 0$$

Change in  $q_{rev}/T$  around any closed path is **0**.

It is a state function.

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