

# Modelling the Hydrogen Molecule (H<sub>2</sub>) with Quantum Circuits

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## 1 - Background

Quantum vs Classical

Quantum Computer	Classical Computer
Uses quantum bits [qubits]	Uses bits
Has quantum properties [eg. superposition & entanglement]	Does not have quantum properties
Innately probabilistic	Deterministic

Table 1.1 - Differences between a Quantum and a Classical computer  
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Classical Quantum Correspondence

Quantum Gates	Matrices
Wavefunctions	Sum of complex vector states
Measurement	Distribution sampling

Hence, classical computers can simulate quantum circuits

Table 1.4 - Made by finalists with Canva

Operator	Gate(s)	Matrix
Pauli-X (X)		$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$
Phase (S, P)		$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$
$\pi/8$ (T)		$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/8} \end{bmatrix}$
Controlled Not (CNOT, CX)		$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$

Figure 1.5 [2]

Applications of Quantum Chemistry

- Determining fundamental molecular characteristics
- Drug discovery and biochemistry
- Innovations in material science (nanotech, superconductors, etc)
- Integrations in engineering processes and solutions
- Analyse catalytic and reaction mechanisms

Graphic 1.6 - Schematic showing how Quantum Computational Chemistry is carried out on Quantum Computers [1]

Table 1.7 - Made by finalists with Canva

## 2 - Objective

- Assess the **effectiveness and accuracy** of the Variational Quantum Eigensolver [VQE] in computing molecular ground state energies and molecular geometry
- Quantitatively characterise the **effects of noise**, in particular thermal relaxation noise, on simulation outcomes
- Develop a **well documented, streamlined framework** for future researchers that leverages the VQE approach

## 3 - Hamiltonian Representation

The Hamiltonian is the sum of potential and kinetic energy of particles in the system. The **Born-Oppenheimer Approximation** is used to assume the nucleus as stationary and hence the nuclear repulsion energy as a constant  $C_n$ . Atomic units [ $\hbar = m_e = e = 4\pi\epsilon_0 = 1$ ] are used to simplify the expression further.

Chemical Hamiltonian  $\hat{H} = \hat{T} + \hat{V}$

Second-Quantized  $\hat{H} = \sum_{pq} h_{pq} a_p^\dagger a_q + \frac{1}{2} \sum_{pqrs} h_{pqrs} a_p^\dagger a_q^\dagger a_s a_r$

Parity-Mapped Qubit Hamiltonian  $\hat{H}_{\text{qubit}} = \sum_{ij} t_{ij} \sigma_i^x \sigma_j^x + \sum_{ij} v_{ij} \sigma_i^z \sigma_j^z + \sum_i \epsilon_i \sigma_i^z$

Measured Expectation Value  $\langle \hat{H}_{\text{qubit}} \rangle = -0.32 + 0.18 X_0 X_1 + 0.40 Z_0 + 0.01 Z_0 Z_1 + 0.40 Z_1$

Diagram 3.1 - Schematic showing the transformations applied to the Hamiltonian  
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$$\hat{H} = \hat{V} + \hat{T} = \underbrace{\sum_{i,j}^{\text{electrons}} \frac{e^2}{4\pi\epsilon_0 |r_i - r_j|}}_{\text{electron-electron repulsion}} + \underbrace{\sum_{i,j}^{\text{nuclei}} \frac{Z_i Z_j e^2}{4\pi\epsilon_0 |R_i - R_j|}}_{\text{nuclear-nuclear repulsion}} - \underbrace{\sum_i^{\text{electrons nuclei}} \frac{Z_i e^2}{4\pi\epsilon_0 |r_i - R_j|}}_{\text{electron-nuclear attraction}} - \underbrace{\sum_i^{\text{nuclei}} \frac{\hbar^2 \nabla_i^2}{2m_i}}_{\text{nuclei kinetic energy}} - \underbrace{\sum_i^{\text{electrons}} \frac{\hbar^2 \nabla_i^2}{2m_e}}_{\text{electron kinetic energy}}$$
$$\approx - \sum_i^{\text{electrons}} \frac{1}{2} \nabla_i^2 + \sum_{i,j}^{\text{electrons}} \frac{1}{|r_i - r_j|} + \sum_i^{\text{electrons nuclei}} \frac{1}{|r_i - R_j|} + \sum_j^{\text{electrons nuclei}} \frac{Z_j}{|r_i - R_j|} + C_n' \quad (1)$$

## 4 - Methodology [VQE - Varitional Quantum Eigensolver]

### 1. Second Quantisation

- Accounts for fermionic properties [Pauli Exclusion Principle]
- Easier to solve for **pairwise** electron interactions.
- Transforms continuous problem to **discrete, matrix-like** form
- Formulated with a **finite set of orbitals** [i.e. STO-3G], each being a solution to the Schrödinger equation

One-electron integrals:

$$h_{pq} = \int \phi_p^*(r) \left( -\frac{1}{2} \nabla^2 - \sum_j^{\text{nuclei}} \frac{Z_j}{|r_j - R_j|} \right) \phi_q(r) d^3r \quad (2)$$

Two-electron integrals:

$$h_{pqrs} = \iint \frac{\phi_p^*(r_1) \phi_q^*(r_2) \phi_r(r_2) \phi_s(r_1)}{|r_1 - r_2|} d^3r_1 d^3r_2 \quad (3)$$

Computing all integrals and expressing in terms of fermionic operators:

$$\hat{H} = \sum_{p,q} h_{pq} a_p^\dagger a_q + \frac{1}{2} \sum_{p,q,r,s} h_{pqrs} a_p^\dagger a_q^\dagger a_r a_s \quad (4)$$

### 2. Fermionic Mappings

The **Jordan-Wigner Transformation [JWT]**:

- converts a **fermionic Hamiltonian** to a **qubit Hamiltonian**
- maps creation and annihilation operators to **Pauli operators**

$$\begin{aligned} \tilde{a}_p |z_0, \dots, z_{p-1}, 1, z_{p+1}, \dots, z_{N-1}\rangle &= (-1)^{\sum_{q=0}^{p-1} z_q} |z_0, \dots, z_{p-1}, 0, z_{p+1}, \dots, z_{N-1}\rangle \\ \tilde{a}_p |z_0, \dots, z_{p-1}, 0, z_{p+1}, \dots, z_{N-1}\rangle &= 0 \quad (5) \\ \tilde{a}_p^\dagger |z_0, \dots, z_{p-1}, 0, z_{p+1}, \dots, z_{N-1}\rangle &= (-1)^{\sum_{q=0}^{p-1} z_q} |z_0, \dots, z_{p-1}, 1, z_{p+1}, \dots, z_{N-1}\rangle \\ \tilde{a}_p^\dagger |z_0, \dots, z_{p-1}, 1, z_{p+1}, \dots, z_{N-1}\rangle &= 0 \quad (6) \end{aligned}$$

Expressing the above in terms of Pauli Matrices:

$$a_j^\dagger = \frac{1}{2}(X_j - iY_j) \prod_{k=0}^{j-1} Z_k \quad a_j = \frac{1}{2}(X_j + iY_j) \prod_{k=0}^{j-1} Z_k \quad (7)$$

### Parity Mapper

Some states are impossible and qubits can be tapered off to improve efficiency.

- Spin must be conserved
- Total number of electrons is constant

## 3. Ansatz Preparation [UCC - Unitary Coupled Cluster]

UCC approximates a molecule's wavefunction by applying a **unitary exponential operator** to a reference state, typically the Hartree Fock [HF] state, which is an initial approximation of the molecule's ground state.

**H<sub>2</sub> HF state and excitation parameters**

Spin-Orbital Mapping  $|\Psi_{\text{HF}}\rangle = |1100\rangle$ :  
 $|0\rangle_3 \rightarrow 2\beta$  (virtual spin-orbital for  $\beta$ )  
 $|0\rangle_2 \rightarrow 2\alpha$  (virtual spin-orbital for  $\alpha$ )  
 $|1\rangle_1 \rightarrow 1\beta$  (occupied spin-orbital for  $\beta$ )  
 $|1\rangle_0 \rightarrow 1\alpha$  (occupied spin-orbital for  $\alpha$ )

UCCSD Excitations:

$\theta_\alpha$ :  $1\alpha \rightarrow 2\alpha$   
 $\theta_\beta$ :  $1\beta \rightarrow 2\beta$   
 $\theta_{\alpha\beta}$ :  $1\alpha 1\beta \rightarrow 2\alpha 2\beta$

$$|\Psi_{\text{UCC}}\rangle = e^{T-T^\dagger} |\Phi_{\text{HF}}\rangle \quad (9)$$

The cluster operator  $T$  encodes electron excitations with weights  $\theta$  that modify the **electron density** in different orbitals.

$$\hat{T} = \hat{T}_1 + \hat{T}_2 = \sum_{i,a} \theta_{ia}^\dagger a_i^\dagger a_a + \frac{1}{4} \sum_{i,j,a,b} \theta_{ij}^{ab} a_i^\dagger a_j^\dagger a_b a_a \quad (8)$$

## 4. Measurement and Optimisation

Figure 4.4.1 - VQE iteration  
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Figure 4.4.2  
Made by finalists using LaTeX

The Simultaneous Perturbation Stochastic Approximation (SPSA) algorithm is typically used due to its **effectiveness in noisy environments**. Noisy effects tend to cancel out across multiple iterations due to its stochastic nature.

## 5 - Quantum Circuit Diagram

Diagram 5.1 - circuit diagram of our quantum circuit  
Rendered by finalists using Qiskit

## 6 - Thermal Relaxation Noise

Qubits face noise from many sources. Besides measurement and readout noise, 2 critical parameters in characterising such noise are T1 and T2 times.

T1: **lifetime** of a qubit's excited state

[ $\infty$  time to relax from  $|1\rangle$  to  $|0\rangle$ ]

T2: **coherence** of superposition

[ $\infty$  time to lose its relative phases]

### Superconducting Qubits Noise Levels

IBM Cloud Median T1 time: 200-300us  
Median T2 time: 100-250us

Gate	Gate Time
Single Qubit Gates	40ns
Swap	40ns
Controlled Not (CX)	150ns

Table 6.3 - Made by finalists with Canva

## 7 - Implementation and Graph Plotting

By varying the H-H distance, we used VQE in a noiseless simulation to map the **energy landscape** [Fig. 8.1]. To find the **true ground state energy** and **optimal geometry**, we concurrently optimized interatomic distance and electron excitation parameters using a custom algorithm [Fig. 8.2, noisy conditions].

We then explored noise effects—specifically **thermal relaxation**—with T<sub>1</sub> set to 2T<sub>2</sub>. Averaging 20 runs per point yielded a smooth curve described by the product of an **exponential decay** and **inverse power function** [Fig. 8.3].

For all graphs, we compare our results with Self Consistent Field [SCF] and Full Configuration Interaction [FCI], where FCI offers the most accurate energy. Closer agreement with FCI indicates a better result.

## 8 - Results

Figure 8.1 - Applying the VQE approach to varying interatomic distances of H<sub>2</sub>, the two horizontal lines representing the minimum ground state energy calculated via the 2 different methods [HF and FCI]

Figure 8.2 - The intermediate results of one execution from our algorithm that optimises interatomic distance [bond length] and electron excitation parameters concurrently.

Figure 8.3 - Noise graph with the fitted equation characterises the relationship between thermal relaxation noise and simulation outcome. Energy approaches the FCI Energy as the simulation decreases in noise.

## 9 - Conclusion

- VQE is accurate and effective, even under noisy environments.** Extremely close agreement (**0.0134%** error) of the true ground state energy to FCI energy given by the energy landscape simulation. **Concurrent optimisation** of molecular geometry and excitation parameters is feasible and yielded great results with **0.11%** error on energy and **1.9%** error on bond length.
- Quantitatively characterised the correlation between thermal relaxation and simulation outcome**, modelled as the product of an exponential decay and inverse power function.
- In conclusion, the project is **successful** in meeting its objectives. A comprehensive Jupyter Notebook with **code and detailed documentation** of the VQE process in quantum chemistry is **available on GitHub for future researchers**.

## 10 - Future Work

Investigation into its actual applicability on Quantum Computers

- Requires comparison with the data obtained from running the circuit on an actual Quantum Computer
- If applicable, investigate other parameters that affect the final graph shape/coefficients
- If not applicable, investigate possible reasons for this discrepancy and if it reflects improvement that could be made to current simulation methods
- Research on Quantum Error Correction Codes (QECC) and their effectiveness in quantum chemistry

Investigation the effectiveness of VQE in simulating more complex molecules

- Likely only achievable on Quantum Computers
- Assess its accuracy, computational resources, and simulation results

## References

- Figure 1.5 - Taken from Wikipedia before making secondary cropping [https://en.wikipedia.org/wiki/Quantum\_logic\_gate#/media/File:Quantum\_Logic\_Gates.png]
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