Time Travelling to Avoid Trump
A Mathematical Analysis

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Abstract

We solve the twin paradox without invoking general relativity. The purpose of this is to show that it is within the laws of physics to build a one-way time machine to the future, allowing us to avoid having to experience Donald Trump's presidency. This part of the paper, part I was written by Vincent Macri. We will mainly rely on an analysis of the relativistic Doppler effect.

We then explore possible way to actually make a one-way time machine, ignoring the economic costs of actually creating one, as they will certainly be prohibitive to doing this in practice anyway. This part of the paper was written by David White and Aviv Haber.

Contents

Part I

The Twin Paradox

Chapter 1

Stating the Twin Paradox

1.1 The problem

Imagine two people, we will call them Alice[1] and Donald[2].

Donald is currently the president of the United States, and he is doing a bad job. Alice doesn't like this situation, so she is going to to time travel to the future when Donald is no longer president.

1.2 Solution to Trump

Special relativity tells us that time will pass slower for an object in motion than an object at rest (Bruni, Dick, Speijer, & Stewart, 2012; Einstein, 1916). This means that if Alice moves very fast, she will experience less time pass for her than on earth, and will effectively time travel to the future.

1.3 The paradox

The paradox here is that special relativity states that the laws of physics are the same in all inertial frames of reference (Bruni et al., 2012; Einstein, 1916). This means that we could also argue that Alice is not moving, and instead the earth is moving, which would result in Alice aging faster than Donald, which is the opposite of what we want and the opposite of what actually happens (Bruni et al., 2012, pp. 593–594).

This is called the twin paradox.

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[1]Of computer security white paper fame.
[2]Of presidential infamy.

1.4 Approach to resolution of the paradox

It is commonly thought that general relativity is needed to resolve this paradox because Alice is in an accelerating frame of reference and special relativity cannot handle accelerating frames of reference. That is not true. Special relativity can indeed handle accelerating frames of reference, but it is more difficult (Gibbs, 1996; Weiss, 1992–2016). However, it is much easier to use special relativity to solve this problem than it is to use general relativity. And in fact, we will later see that the acceleration of Alice is irrelevant to the resolution of the paradox.

1.5 Notation and assumptions

1.5.1 Notation

Our notation

In math, Alice will be referred to as A, and Donald will be referred to as D.

Other physics notation

Following is some common non-trivial physics notation that we will be using:

Speed of light
$$c_0 = 299\,792\,458 \, \text{m/s}$$

We will use c_0 as the speed of light in a vacuum. We are using c_0 instead of c because c_0 is the recommended SI notation (BIPM, 2006).

β notation
$$\beta = \frac{v \, \text{m/s}}{c_0 \, \text{m/s}}$$

Where v is velocity and c_0 is the speed of light.

1.5.2   Assumptions

Values

We will say that Alice travels a distance of 4 light years at a speed of $0.8\,c_0$ since these are nice numbers to work with (Kogut, 2012, p. 35). She then turns around and comes back to earth.

Turnaround

We will assume that Alice makes an instantaneous turnaround. Later we will show that this assumption has no effect on the resolution of the paradox.

Chapter 2

Resolving the Twin Paradox

2.1   A naive analysis

We have:

$$v = 0.8\,c_0 \implies \beta = 0.8$$
$$d = 4\,\text{ly}$$

And we know this formula from Bruni et al., 2012, p. 583 (modified to use our notation):

$$\Delta t_D = \frac{\Delta t_A}{\sqrt{1 - \beta}}$$

Which can be rearranged into:

$$\Delta t_A = \Delta t_D \sqrt{1 - \beta^2} \tag{2.1}$$

We can trivially calculate how much time should pass for Donald:

$$\Delta t_D = \frac{2d}{v} = \frac{2 \times 4\,\text{ly}}{0.8\,c_0} = \frac{8\,\text{ly}}{0.8\,c_0} = 10\,\text{y}$$

And how much time should pass for Alice follows by simply plugging this into (2.1):

$$\Delta t_A = \Delta 10\sqrt{1 - 0.8^2}$$
$$\Delta t_A = \Delta 10\sqrt{\frac{9}{25}}$$
$$\Delta t_A = \Delta 10\frac{3}{5}$$
$$\Delta t_A = 6$$

So, Donald ages by 10 years, and Alice ages by 6 years.

This answer is right, but the problem is that we started by assuming that Donald is stationary and Alice is moving. However, we could have said that Alice is stationary and Donald is moving, and then we would calculate $\Delta t_A = 10$ and $\Delta t_D = 6$, which is wrong. So doing the analysis this way leads to ambiguity. We must develop a more rigorous way to analyze this problem.

2.2 Deriving the relativistic Doppler effect

The simplest way to solve the twin paradox is to use the relativistic Doppler effect in order to analyze what Donald sees and what Alice sees.

We will have both Alice and Donald flash a light at the other once per second, according to their own proper time. Their lights are infinitely powerful, and can be seen from light years away (once the light has travelled there of course). We will assume the path of the lights are entirely through a perfect vacuum.

2.2.1 More notation

First, let's define some notation specific to this section:

f_s The frequency the emitter (source) flashes their light at. We will make Alice the source.

f_o The frequency the other person (observer) sees the light flashing at. We will make Donald the observer of Alice's flashing light.

t_s The proper time to the next wavefront as seen by the emitting source, Alice.

t_o The proper time to the next wavefront as seen by person observing the flashes, Donald.

λ The distance to the next wave front of the approaching light wave. λ is calculated as:

$$\lambda = \frac{c_0 \, \mathrm{m/s}}{f_s \, \mathrm{s}^{-1}} = \frac{c_0}{f_s} \, \mathrm{m}$$

2.2.2 The derivation

We start by relating Δt_s to λ and v when Alice is moving away from Donald:

$$\Delta t_s = \frac{\lambda}{c_0} + \frac{v \times \Delta t_s}{c_0}$$

$$c_0 \Delta t_s = \lambda + v \Delta t_s$$

$$c_0 \Delta t_s - v \Delta t_s = \lambda$$

$$\Delta t_s (c_0 - v) = \lambda$$

$$\Delta t_s = \frac{\lambda}{c_0 - v}$$

$$\Delta t_s = \frac{1}{c_0 - v} \times \lambda$$

Now substitute in $\lambda = \frac{c_0}{f_s}$:

$$\Delta t_s = \frac{1}{c_0 - v} \times \frac{c_0}{f_s}$$

$$\Delta t_s = \frac{c_0}{c_0 - v} \times \frac{1}{f_s}$$

$$\Delta t_s = \frac{1}{1 - \frac{v}{c_0}} \times \frac{1}{f_s}$$

$$\Delta t_s = \frac{1}{f_s(1 - \beta)} \tag{2.2}$$

Next, we will perform a unit analysis to verify that (2.2) gives us a value in seconds:

$$\Delta t_s = \frac{1}{\text{s}^{-1}}$$

$$\Delta t_s = \text{s}$$

So, we have derived (2.2) and verified that it gives us a value in seconds. Now we need to use this formula.

The next step is to develop the actual Doppler effect formula. We will work off of the special relativity time dilation formula given in Bruni et al., 2012, p. 593, modified to use our notation, and rearrange it into a form more useful for our purposes:

$$\Delta t_s = \frac{\Delta t_o}{\sqrt{1 - \beta^2}}$$

$$\Delta t_o = \Delta t_s \sqrt{1 - \beta^2} \tag{2.3}$$

We will now substitute (2.2) into (2.3) to combine our two equations in order to develop a third formula:

$$\Delta t_o = \frac{\sqrt{1 - \beta^2}}{f_s(1 - \beta)} \tag{2.4}$$

Next, we will finish developing the relativistic Doppler shift formula.

Note that, by definition:

$$f_o = \frac{1}{\Delta t_o} \tag{2.5}$$

We will now substitute (2.4) into (2.5):

$$f_o = \frac{1}{\frac{\sqrt{1-\beta^2}}{f_s(1-\beta)}}$$

$$f_o = \frac{f_s(1 - \beta)}{\sqrt{1 - \beta^2}}$$

$$f_o = f_s \left(\frac{1 - \beta}{\sqrt{1 - \beta^2}} \right)$$

$$f_o = f_s \left(\frac{\sqrt{(1 - \beta)^2}}{\sqrt{1^2 - \beta^2}} \right)$$

$$f_o = f_s \left(\frac{\sqrt{(1 - \beta)(1 - \beta)}}{\sqrt{(1 - \beta)(1 + \beta)}} \right)$$

$$f_o = f_s \sqrt{\frac{(1 - \beta)(1 - \beta)}{(1 - \beta)(1 + \beta)}}$$

$$f_o = f_s \times \sqrt{\frac{1 - \beta}{1 + \beta}} \tag{2.6}$$

This is the formula for the relativistic Doppler effect when Alice and Donald are moving away from each other. If Alice is moving towards Donald, then we simply change the sign on β to get:

$$f_o = f_s \times \sqrt{\frac{1 + \beta}{1 - \beta}} \tag{2.7}$$

2.3 Minkowski spacetime diagrams and world lines

We must now introduce a new way of analyzing problems in special relativity: the Minkowski spacetime diagrams and world lines.

We will draw the world lines of our situation onto Minkowski spacetime diagrams. This is simply a diagram with time on the vertical axis, and distance on the horizontal axis.

We will start with a spacetime diagram of what we would expect to see, disregarding special relativity, using Donald as the frame of reference. See Figure 2.1.
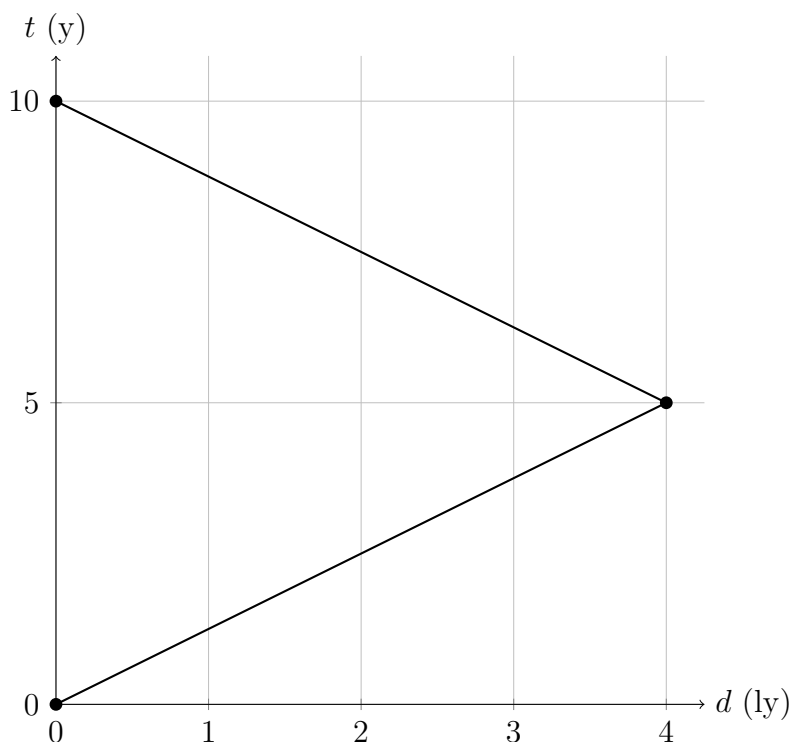

Figure 2.1: Minkowski spacetime diagram of what we would expect to see according to classical physics with Donald as the frame of reference.

Of course, figure 2.1 does not apply, since $\beta = 0.8$ (which is quite large), however, I find it helpful to visualize the problem and introduce the notion of spacetime diagrams.

2.4 Minkowski spacetime, the relativistic Doppler effect, and the twin paradox

Recall the assumptions and values we decided to use in subsection 1.5.2:

Distance Alice travels a distance of 4 light years, so $d = 4\,\mathrm{ly}$.

Velocity Alice travels at a speed of $0.8\,c_0$, so $v = 0.8\,c_0$ and $\beta = 0.8$.

The actual values we use do not matter for resolving the paradox. These values were chosen because they give nice numbers when we perform the calculations, which makes the analysis easier to follow.

We also assumed an instantaneous turnaround. This too makes the math easier, but we will show that it does not change the resolution to the paradox.

Analyzing the twin paradox with the relativistic Doppler effect is helpful because it allows us to calculate what each person sees, and show that they are seeing different things, which solves the ambiguity stated in section 2.1.

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