

# Corner Transfer Matrix Renormalization Group

an Algorithm for 2D Classical Lattice Models  
a Literature Digestion of T. Nishino and K. Okunishi's article (1996).

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# Outline

- 1 From 1D to 2D: The Contraction Issue
- 2 Going to 2D: A Fundamentally Harder Problem
- 3 CTMRG: Extending to 2D
- 4 Summary

# Review: 1D Transfer Matrix (Audience's Pre-knowledge)

## 1D Ising Model:

$$H = -J \sum_i \sigma_i \sigma_{i+1}$$

All  $T$  identical  $\Rightarrow$  commute  $\Rightarrow$  diagonalize once!

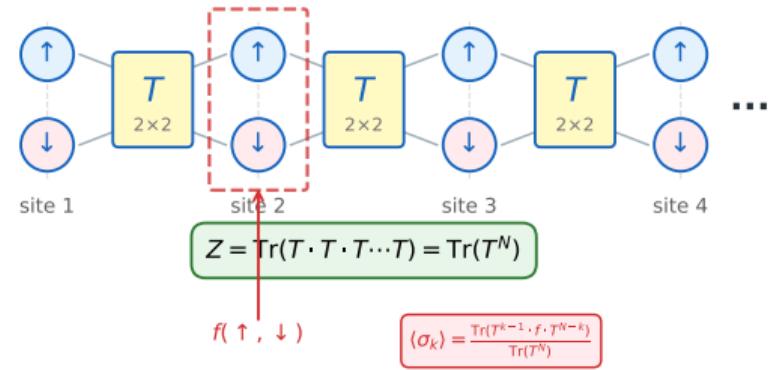
## Partition function as contraction:

$$Z = \text{Tr}(T \cdot T \cdot T \cdots T) = \text{Tr}(T^N)$$

where  $T \in \mathbb{R}^{2 \times 2}$  encodes local interactions.

## Local observable:

$$\langle \sigma_k \rangle = \frac{\text{Tr}(T^k \cdot f(\uparrow, \downarrow) \cdot T^{N-k})}{\text{Tr}(T^N)}$$



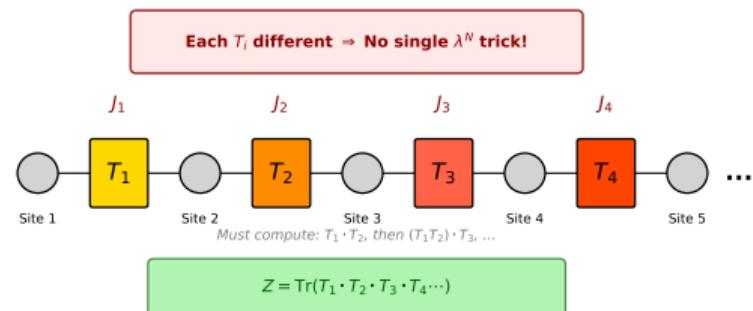
## Key Point

# When Even 1D Classical Needs Renormalization Group

**Recall:** For uniform 1D Ising,  $Z = \text{Tr}(T^N)$  has no contraction issues.

**But what if...**

- Couplings  $J_i$  are **site-dependent**?
- Random disorder:  $J_i \sim \mathcal{N}(\bar{J}, \sigma)$ ?
- Open boundary conditions (no translation symmetry)?



**Then:** Cannot diagonalize  $T$  once!

$$Z = \text{Tr}(T_1 \cdot T_2 \cdot T_3 \cdots T_N)$$

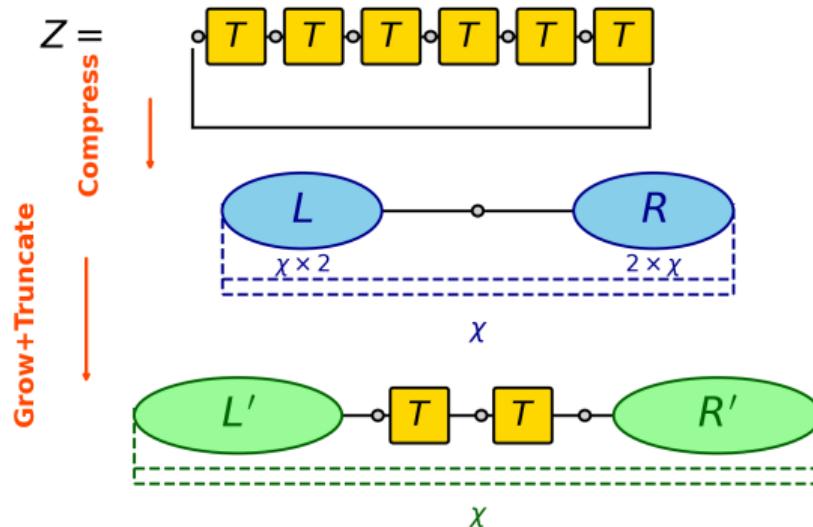
Each  $T_i$  is different  $\Rightarrow$  no simple  $\lambda^N$  formula.

## The Problem

For  $N$  sites: need  $O(N)$  matrix multiplications.  
Not exponential, but **no closed-form solution**.

# Solution: Transfer Matrix Renormalization Group (TMRG)

**Insight:** Not all configurations contribute equally to  $Z \Rightarrow$  Keep only the **most relevant** ones!



$\chi =$  bond dimension (most relevant configurations in compressed basis)

## TMRG Algorithm:

- ① **Grow:**  $L' = L \cdot T$ ,  $R' = T \cdot R$   
(add sites)
- ② **Truncate:** SVD  $\rightarrow$  keep  $\chi$  largest  
(coarse-grain)

## RG Fixed Point

Iterate until  $L^*$ ,  $R^*$  converge  $\Rightarrow$   
Thermodynamic limit!

Kind-of like infinite-DMRG by adding new  
of sites

## 2D: The Transfer Matrix Becomes a **Network**

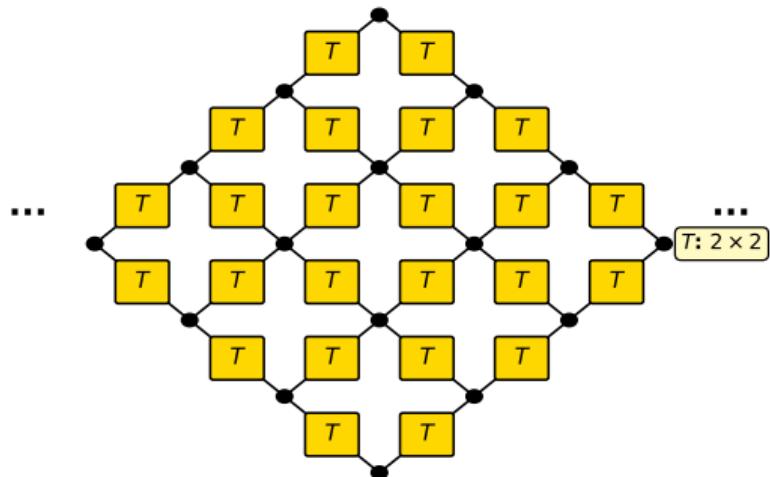
**1D:** Ordered product of matrices

$$Z = \text{Tr}(T_1 \cdot T_2 \cdots T_N)$$

Contract left-to-right:  $O(N)$ .

**2D:**  $T$  still  $2 \times 2$  on each **bond**,  
but bonds form a **2D network!**  
 $\Rightarrow$  No natural ordering.

2D:  $T$  on each bond  $\Rightarrow$  No natural contraction order!



### The Fundamental Problem

**No natural contraction order!** The 2D tensor network cannot be reduced to a simple trace of matrix products. **Exact contraction is way too slow!**

# One Approach: Row-to-Row Transfer (DMRG-style)

**Idea:** Group one row of  $L$  spins  $\rightarrow$  treat as a “super-spin” with  $2^L$  states.

**Apply DMRG/TMRG ideas:**

- Row config space  $\rightarrow$  MPS
- Row-to-row transfer  $\rightarrow$  MPO
- Truncate via SVD

[Figure: Row  $\rightarrow$  super-spin,  $T_{\text{row}} \in \mathbb{R}^{2^L \times 2^L}$ ]

## Pros & Cons

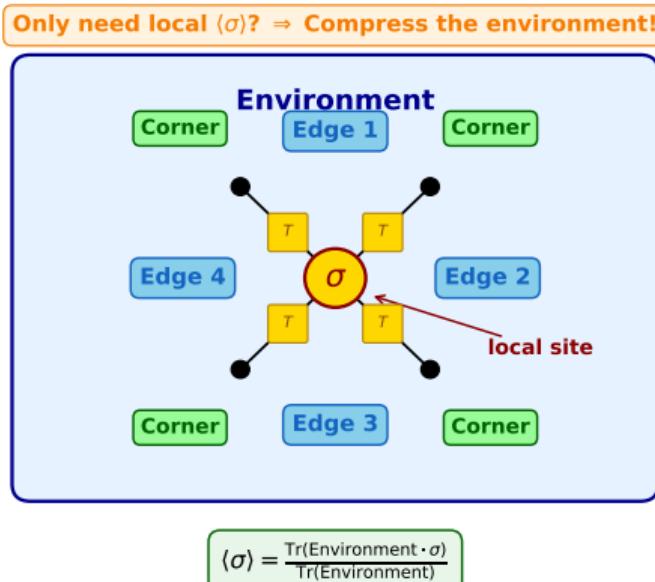
- + Systematic, well-understood
- Breaks 2D symmetry
- Hard to generalize to other lattices

Then  $Z = \text{Tr}(T_{\text{row}}^M)$  looks like 1D!

(Details in Appendix)

# A More Natural Approach: Use Environment just like in 1D TMRG

**Question:** If we only want **local observables**  $\langle \sigma_{ij} \rangle$ , do we really need to contract the *entire* infinite lattice?



## Key Insight:

Decompose the infinite 2D environment into **geometrically natural** pieces:

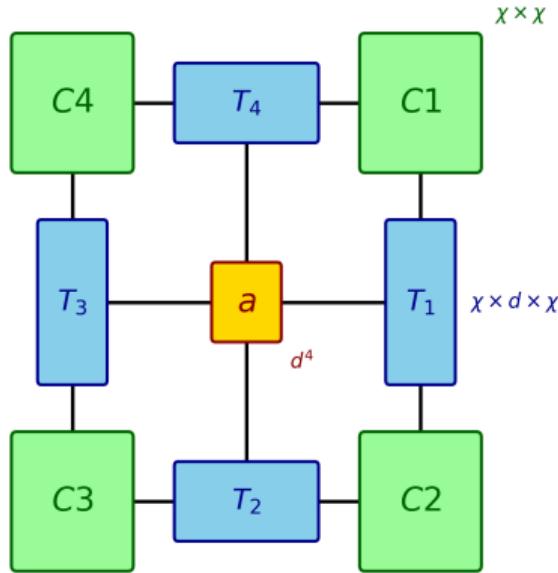
- **4 Corners** (quarter-planes)
- **4 Edges** (half-infinite strips)

## Baxter's Insight (1968)

The **Corner Transfer Matrix** encodes a quarter of the infinite plane!

# 2D Analog: Four Corners + Four Edges

CTMRG: Decompose 2D environment into 4 Corners + 4 Edges



**Corners  $C$ :**

- Quadrant ( $\chi \times \chi$ )

**Edges  $T$ :**

- Half-strip ( $\chi \times d \times \chi$ )

**Local tensor  $a$ :**

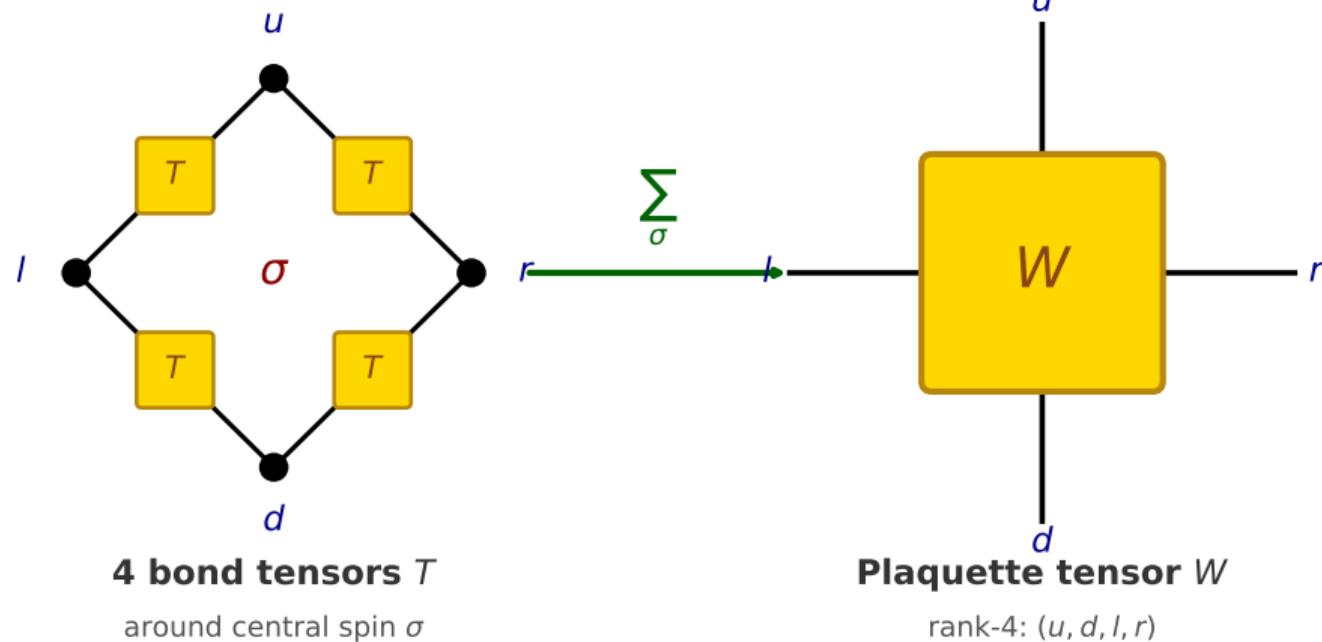
- Rank-4 ( $d^4$ )

Key: What is  $a$ ?

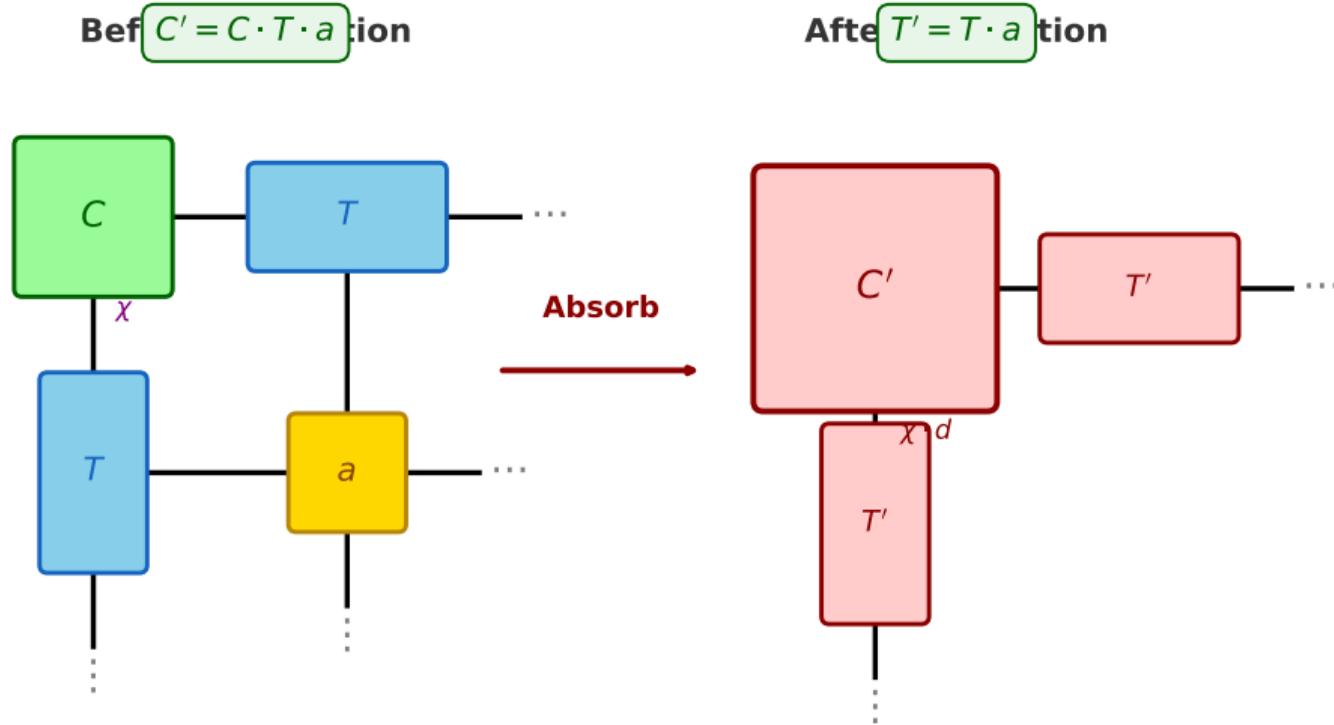
$a$  is a **rank-4 tensor** — for computing  $Z$ :  $a = W$  plaquette tensor; for  $\langle \sigma \rangle$ :  $a = s_1 \delta_{s_1 s_2 s_3 s_4}$ , etc.

# The Plaquette Tensor: From Bonds to Faces

$$W_{u,d,l,r} = \sum_{\sigma} T_{\sigma,u} T_{\sigma,d} T_{\sigma,l} T_{\sigma,r}$$



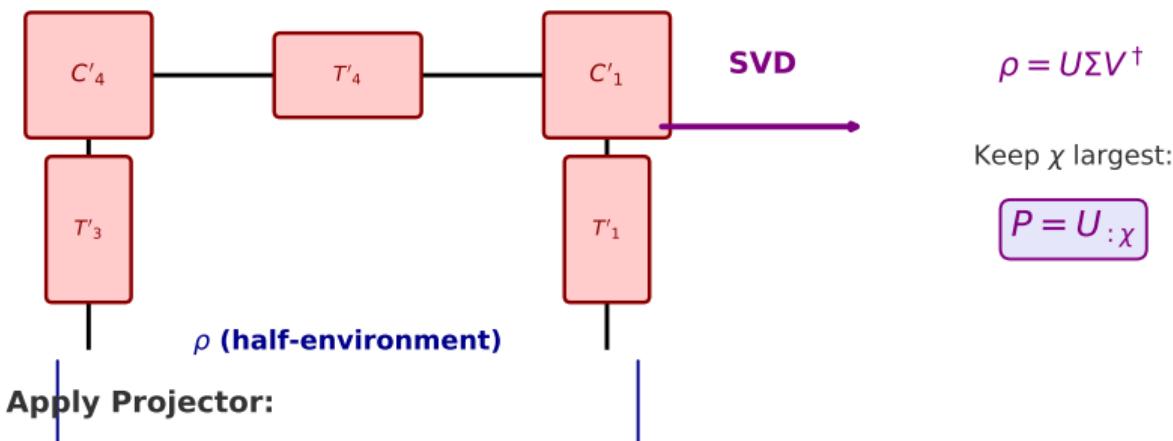
# CTMRG Iteration: Step 1 — Grow (Absorb)



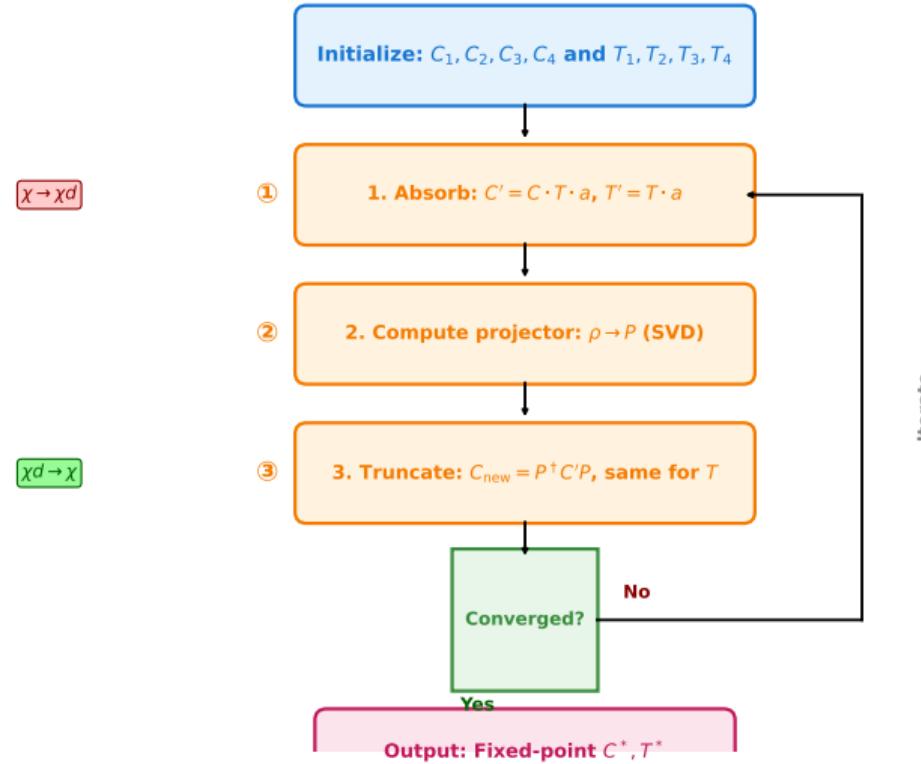
Bond dimension grows:  $x \rightarrow x^d$  Need truncation!

# CTMRG Iteration: Step 2 — Truncate

**Truncation: Use environment to find optimal projector  $P$**



# CTMRG: Complete Algorithm



# Summary: The Core of CTMRG

## CTMRG in Three Steps

- ① **Decompose:** Infinite square lattice → 4 corners + 4 edges (customed for other lattice);
- ② **Grow and Truncate (RG):** add row+column then absorb, SVD to keep  $\chi$  most relevant;
- ③  $Z = W_{s_1 s_2 s_3 s_4} \text{Tr}(C_1 T_1^{s_1} C_2 T_2^{s_2} C_3 T_3^{s_3} C_4 T_4^{s_4}) ; \langle \mathcal{O} \rangle = \frac{\mathcal{O}_{s_1 s_2 s_3 s_4}}{Z} \text{Tr}(C_1 T_1^{s_1} C_2 T_2^{s_2} C_3 T_3^{s_3} C_4 T_4^{s_4})$

[Figure: Visual summary — Decompose → Grow → Truncate]

# Thank You All for Your Attention!

And thanks to Prof. Tommaso Roscilde, Dr. Fabio Mezzacapo, Filippo Caleca and Saverio Bocini for your teaching and assistance!

## Appendix: TMRG Setup in Detail

**Goal:** Compute  $Z = \text{Tr}(T_1 \cdot T_2 \cdots T_N)$  for large  $N$ .

[Figure: Left environment  $L$ , local site  $T$ , right environment  $R$ ]

**Left Environment  $L$ :**

$$L_\sigma = \sum T_1 \cdot T_2 \cdots T_{k-1}$$

**Right Environment  $R$ :**

$$R_\sigma = \sum T_{k+1} \cdots T_N$$

## Appendix: TMRG Step 1 — Grow (Absorption)

**Absorption:** Add one more site to the environment.

[Figure:  $L' = L \cdot T_k$ , growing the left environment]

**Mathematically:**

$$L'_{\sigma_k} = \sum_{\sigma_{k-1}} L_{\sigma_{k-1}} \cdot (T_k)_{\sigma_{k-1}, \sigma_k}$$

## Appendix: TMRG Step 2 — Truncate (SVD)

**Idea:** Compress  $L'$  back to dimension  $\chi$  using SVD.

**Form the “density matrix”:**

$$\rho = L' \cdot R'^T$$

**SVD:**

$$\rho = U \Sigma V^\dagger$$

**Truncate:**

$$P = U_{:,1:\chi}$$

**New environment:**

$$L_{\text{new}} = P^\dagger L'$$

[Figure: SVD spectrum, keep  $\chi$  largest]

Eckart-Young Theorem

SVD gives the **optimal** rank- $\chi$  approximation

# Appendix: TMRG Convergence and Observables

**Iterate:** Grow → Truncate → Grow → Truncate → ...

**Convergence criterion:**

Fixed point:  $L^*, R^*$  such that

$$L^* \xrightarrow{\text{grow+truncate}} L^*$$

**Free energy:**

$$f = -k_B T \lim_{N \rightarrow \infty} \frac{1}{N} \ln Z$$

$$= -k_B T \ln \sigma_1^*$$

**Local observables:**

$$\langle \sigma_k \rangle = \frac{L^* \cdot \sigma_k \cdot R^*}{L^* \cdot R^*}$$

**Correlation functions:**

$$\langle \sigma_i \sigma_j \rangle = \frac{L^* \cdot \sigma_i \cdot T^{|i-j|} \cdot \sigma_j \cdot R^*}{Z}$$

**Physical Meaning**

$L^*, R^*$  encode the **thermodynamic limit**.

## Appendix: 2D Row-to-Row — Row as Super-Spin

**Idea:** Treat one row of  $L$  spins as a single “super-spin”.

**Row configuration:**

$$\vec{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_L)$$

Total states:  $2^L$ .

[Figure: Row  $\rightarrow$  super-spin with  $2^L$  states]

**Row-to-row transfer:**

$$(T_{\text{row}})_{\vec{\sigma}, \vec{\sigma}'} = \prod_{i=1}^L e^{\beta J \sigma_i \sigma'_i} \prod_{i=1}^{L-1} e^{\beta J \sigma_i \sigma_{i+1}}$$

$$T_{\text{row}} \in \mathbb{R}^{2^L \times 2^L}.$$

### Key Point

Now  $Z = \text{Tr}(\tau^M)$  looks like 1D problem with  $2^L$ -dimensional transfer matrix!

# Appendix: Row Transfer Matrix = MPO

**Key insight:**  $T_{\text{row}}$  has a **tensor network** structure!

**Decompose**  $T_{\text{row}}$  into local tensors:

$$T_{\vec{\sigma}, \vec{\sigma}'} = \sum_{\alpha_1, \dots} W_{\alpha_0 \alpha_1}^{\sigma_1 \sigma'_1} W_{\alpha_1 \alpha_2}^{\sigma_2 \sigma'_2} \dots$$

**Local tensor  $W$ :**

$W_{\alpha \beta}^{\sigma \sigma'}$  = local Boltzmann weight

- $\sigma, \sigma'$ : spins in rows  $n, n+1$
- $\alpha, \beta$ : auxiliary (horizontal bonds)

[Figure: MPO structure of  $T_{\text{row}}$ ]

No Quantum!

“MPO” is just a factorization of the classical transfer matrix.

## Appendix: Boundary = MPS

**MPS:** Efficient representation of row configuration space.

**Full vector:**

$$|R\rangle = \sum_{\vec{\sigma}} R_{\vec{\sigma}} |\vec{\sigma}\rangle$$

has  $2^L$  components.

**MPS compression:**

$$R_{\vec{\sigma}} = A^{\sigma_1} A^{\sigma_2} \cdots A^{\sigma_L}$$

Only  $L \cdot \chi^2 \cdot 2$  parameters!

[Figure: MPS tensor network]

### Physical Meaning

MPS compresses  $2^L$  row configs into  $\chi$  effective states with **limited entanglement**.

# Appendix: Row-to-Row DMRG Algorithm

**Combine:** MPS (boundary) + MPO (transfer) + SVD (truncation).

- ① Initialize MPS  $|L\rangle$ ,  $|R\rangle$  for boundaries
- ② **Grow:** Apply MPO  $T_{\text{row}}$  to boundaries

$$|L'\rangle = T_{\text{row}}|L\rangle$$

- ③ **Truncate:** SVD to compress bond dimension back to  $\chi$
- ④ Iterate until convergence to fixed point  $|L^*\rangle$ ,  $|R^*\rangle$

## Pros

- Systematic, well-understood
- Controlled approximation

## Cons

- Breaks 2D rotational symmetry
- Hard to generalize to other lattices

## Appendix: No Monte Carlo Sign Problem

**Classical models:** All Boltzmann weights are **positive!**

$$W = e^{-\beta H} > 0 \quad \text{always}$$

[Figure: Classical (positive) vs Quantum (sign problem)]

## Appendix: Why is 3D Difficult?

In 2D: Environment tensors are **1D objects** (edges).

In 3D: Environment would be **2D surfaces** — back to exponential!

*[Figure: 2D boundary problem in 3D systems]*

# Appendix: Extensions and Generalizations

## Beyond square lattice Ising:

### Different lattices:

- Honeycomb
- Triangular
- Kagome

### Different models:

- Potts model
- Clock model
- Vertex models

### Quantum systems (via iPEPS):

- 2D Heisenberg model
- Frustrated magnets
- Topological phases

### Improvements:

- Directional CTMRG
- Full-update vs simple-update
- Gradient optimization

# Comparison: iDMRG vs CTMRG

Aspect	iDMRG (1D)	CTMRG (2D)
Dimension	1D chain	2D square lattice
Environment	Left + Right	4 Corners + 4 Edges
Grow step	Add site pair	Add row + column
Truncation	SVD on center bond	SVD on corner boundaries
Bond dimension	$\chi$ (MPS)	$\chi$ (environment)
Fixed point	$L^*, R^*$	$C_i^*, T_i^*$
Computational cost	$O(\chi^3)$	$O(\chi^6)$ or $O(\chi^5)$

# Appendix: Computational Complexity

## CTMRG scaling:

Operation	Cost	Bottleneck?
Corner absorption	$O(\chi^4 d^2)$	
Edge absorption	$O(\chi^3 d^2)$	
Build density matrix	$O(\chi^4)$	
SVD for projector	$O(\chi^3 d^3)$	✓
Apply truncation	$O(\chi^3 d)$	
<b>Total per iteration</b>	$O(\chi^3 d^3)$ to $O(\chi^6)$	

## Comparison

- iDMRG (1D):  $O(\chi^3)$  per sweep
- CTMRG (2D):  $O(\chi^5) - O(\chi^6)$  per iteration
- More expensive, but still **polynomial** in  $\chi$ !

# Appendix: Key References

## Original works:

- R. J. Baxter, *J. Math. Phys.* **9**, 650 (1968) — Corner transfer matrices
- R. J. Baxter, *J. Stat. Phys.* **19**, 461 (1978) — CTM method

## Modern CTMRG:

- T. Nishino & K. Okunishi, *J. Phys. Soc. Jpn.* **65**, 891 (1996)
- T. Nishino & K. Okunishi, *J. Phys. Soc. Jpn.* **66**, 3040 (1997)

## CTMRG for iPEPS:

- R. Orús & G. Vidal, *Phys. Rev. B* **80**, 094403 (2009)
- P. Corboz et al., *Phys. Rev. B* **84**, 041108(R) (2011)

## Reviews:

- R. Orús, *Ann. Phys.* **349**, 117 (2014) — Tensor networks review