

Corner Transfer Matrix Renormalization Group

an Algorithm for 2D Classical Lattice Models

a Literature Digestion of T. Nishino and K. Okunishi's article (1996).

Chengzhi Ye

ENS de Lyon

January 5, 2026

Outline

- 1 From 1D to 2D: The Contraction Issue
- 2 Going to 2D: A Fundamentally Harder Problem
- 3 CTMRG: Extending to 2D
- 4 Summary

Review: 1D Transfer Matrix (Audience's Pre-knowledge)

1D Ising Model:

$$H = -J \sum_i \sigma_i \sigma_{i+1}$$

Partition function as contraction:

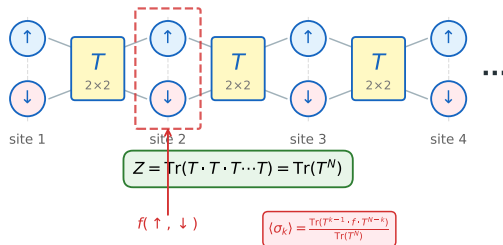
$$Z = \text{Tr}(T \cdot T \cdot T \cdots T) = \text{Tr}(T^N)$$

where $T \in \mathbb{R}^{2 \times 2}$ encodes local interactions.

Local observable:

$$\langle \sigma_k \rangle = \frac{\text{Tr}(T^k \cdot f(\uparrow, \downarrow) \cdot T^{N-k})}{\text{Tr}(T^N)}$$

All T identical \Rightarrow commute \Rightarrow diagonalize once!

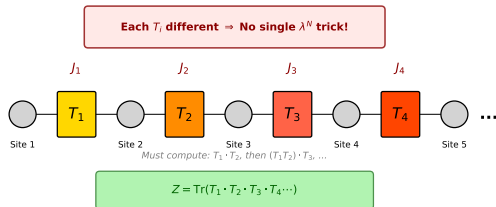


When Even 1D Classical Needs Renormalization Group

Recall: For uniform 1D Ising, $Z = \text{Tr}(T^N)$ has no contraction issues.

But what if...

- Couplings J_i are **site-dependent**?
- Random disorder: $J_i \sim \mathcal{N}(\bar{J}, \sigma)$?
- Open boundary conditions (no translation symmetry)?



Then: Cannot diagonalize T once!

$$Z = \text{Tr}(T_1 \cdot T_2 \cdot T_3 \cdots T_N)$$

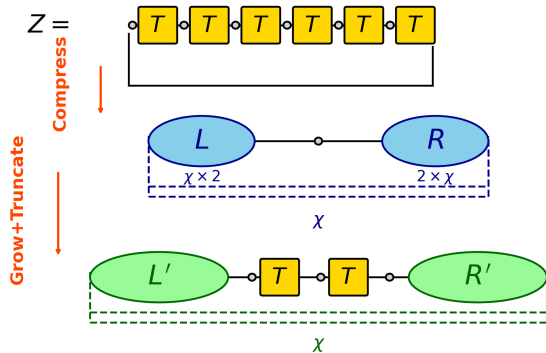
Each T_i is different \Rightarrow no simple λ^N formula.

The Problem

For N sites: need $O(N)$ matrix multiplications.
Not exponential, but **no closed-form solution**.

Solution: Transfer Matrix Renormalization Group (TMRG)

Insight: Not all configurations contribute equally to $Z \Rightarrow$ Keep only the **most relevant** ones!



χ = bond dimension (most relevant configurations in compressed basis)

TMRG Algorithm:

- 1 **Grow:** $L' = L \cdot T$, $R' = T \cdot R$
(add sites)
- 2 **Truncate:** SVD \rightarrow keep χ largest
(coarse-grain)

RG Fixed Point

Iterate until L^*, R^* converge \Rightarrow
Thermodynamic limit!

Kind-of like infinite-DMRG by adding new sites

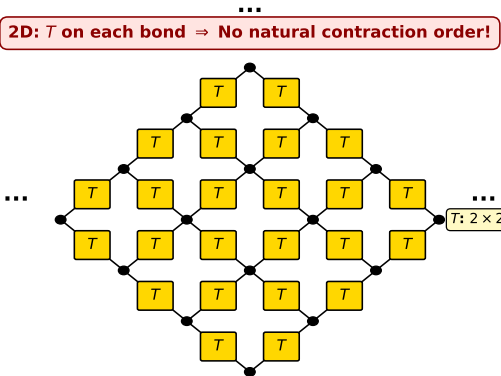
2D: The Transfer Matrix Becomes a **Network**

1D: Ordered product of matrices

$$Z = \text{Tr}(T_1 \cdot T_2 \cdots T_N)$$

Contract left-to-right: $O(N)$.

2D: T still 2×2 on each **bond**,
but bonds form a **2D network**!
 \Rightarrow No natural ordering.



The Fundamental Problem

No natural contraction order! The 2D tensor network cannot be reduced to a simple trace of matrix products. **Exact contraction is way too slow!**

One Approach: Row-to-Row Transfer (DMRG-style)

Idea: Group one row of L spins \rightarrow treat as a “super-spin” with 2^L states.

[Figure: Row \rightarrow super-spin, $T_{\text{row}} \in \mathbb{R}^{2^L \times 2^L}$]

Apply DMRG/TMRG ideas:

- Row config space \rightarrow MPS
- Row-to-row transfer \rightarrow MPO
- Truncate via SVD

Pros & Cons

- + Systematic, well-understood
- Breaks 2D symmetry
- Hard to generalize to other lattices

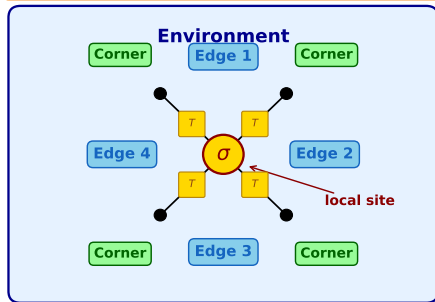
Then $Z = \text{Tr}(T_{\text{row}}^M)$ looks like 1D!

(Details in Appendix)

A More Natural Approach: Use Environment just like in 1D TMRG

Question: If we only want **local observables** $\langle \sigma_{i,j} \rangle$, do we really need to contract the *entire* infinite lattice?

Only need local $\langle \sigma \rangle$? \Rightarrow Compress the environment!



$$\langle \sigma \rangle = \frac{\text{Tr}(\text{Environment} \cdot \sigma)}{\text{Tr}(\text{Environment})}$$

Key Insight:

Decompose the infinite 2D environment into **geometrically natural** pieces:

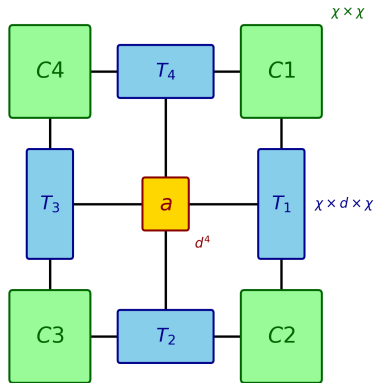
- 4 **Corners** (quarter-planes)
- 4 **Edges** (half-infinite strips)

Baxter's Insight (1968)

The **Corner Transfer Matrix** encodes a quarter of the infinite plane!

2D Analog: Four Corners + Four Edges

CTMRG: Decompose 2D environment into 4 Corners + 4 Edges



Corners C :

- Quadrant ($\chi \times \chi$)

Edges T :

- Half-strip ($\chi \times d \times \chi$)

Local tensor a :

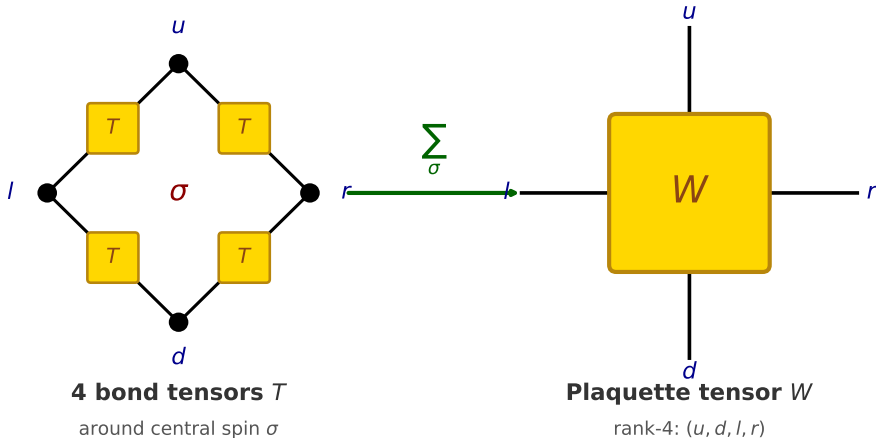
- Rank-4 (d^4)

Key: What is a ?

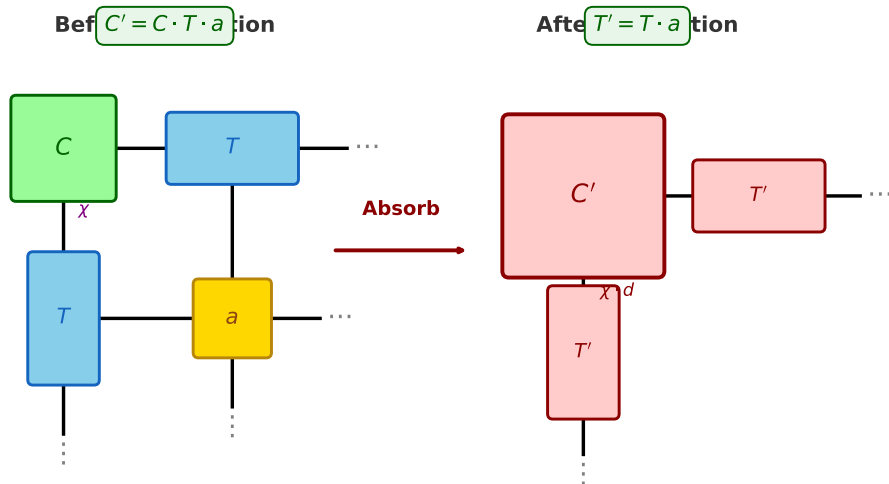
a is a **rank-4 tensor** — for computing Z : $a = W$ *plaquette tensor*; for $\langle \sigma \rangle$: $a = s_1 \delta_{s_1 s_2 s_3 s_4}$, etc.

The Plaquette Tensor: From Bonds to Faces

$$W_{u,d,l,r} = \sum_{\sigma} T_{\sigma,u} T_{\sigma,d} T_{\sigma,l} T_{\sigma,r}$$



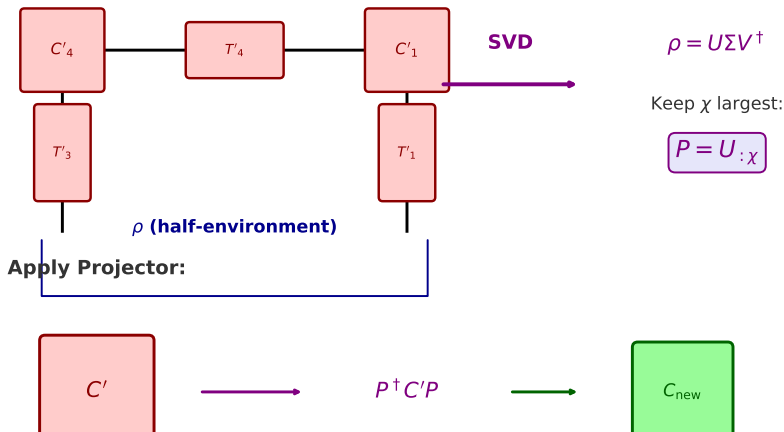
CTMRG Iteration: Step 1 — Grow (Absorb)



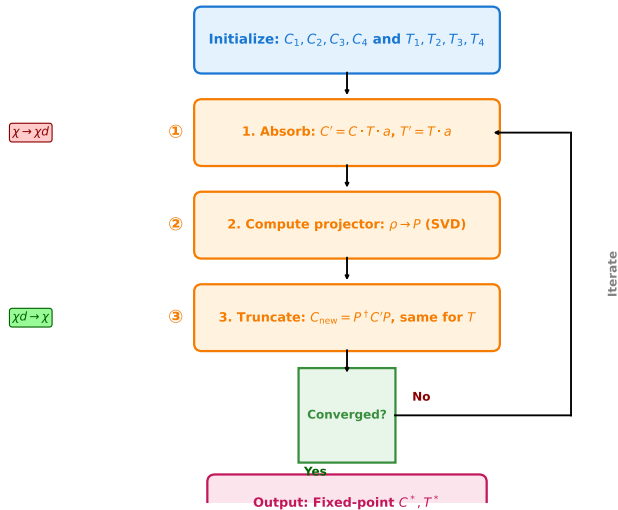
Bond dimension grows: $\chi \rightarrow \chi \cdot d$. Need truncation!

CTMRG Iteration: Step 2 — Truncate

Truncation: Use environment to find optimal projector P



CTMRG: Complete Algorithm



Summary: The Core of CTMRG

CTMRG in Three Steps

- ① **Decompose:** Infinite square lattice \rightarrow 4 corners + 4 edges (customized for other lattice);
- ② **Grow and Truncate (RG):** add row+column then absorb, SVD to keep χ most relevant;
- ③ $Z = W_{s_1 s_2 s_3 s_4} \text{Tr}(C_1 T_1^{s_1} C_2 T_2^{s_2} C_3 T_3^{s_3} C_4 T_4^{s_4})$; $\langle \mathcal{O} \rangle = \frac{\mathcal{O}_{s_1 s_2 s_3 s_4}}{Z} \text{Tr}(C_1 T_1^{s_1} C_2 T_2^{s_2} C_3 T_3^{s_3} C_4 T_4^{s_4})$

[Figure: Visual summary — Decompose \rightarrow Grow \rightarrow Truncate]

Thank You All for Your Attention!

And thanks to Prof. Tommaso Roscilde, Dr. Fabio Mezzacapo, Filippo Caleca and Saverio Bocini for your teaching and assistance!

Appendix: TMRG Setup in Detail

Goal: Compute $Z = \text{Tr}(T_1 \cdot T_2 \cdots T_N)$ for large N .

[Figure: Left environment L , local site T , right environment R]

Left Environment L :

$$L_\sigma = \sum T_1 \cdot T_2 \cdots T_{k-1}$$

Right Environment R :

$$R_\sigma = \sum T_{k+1} \cdots T_N$$

Appendix: TMRG Step 1 — Grow (Absorption)

Absorption: Add one more site to the environment.

[Figure: $L' = L \cdot T_k$, growing the left environment]

Mathematically:

$$L'_{\sigma_k} = \sum_{\sigma_{k-1}} L_{\sigma_{k-1}} \cdot (T_k)_{\sigma_{k-1}, \sigma_k}$$

Appendix: TMRG Step 2 — Truncate (SVD)

Idea: Compress L' back to dimension χ using SVD.

Form the “density matrix”:

$$\rho = L' \cdot R'^T$$

SVD:

$$\rho = U \Sigma V^\dagger$$

Truncate:

$$P = U_{:,1:\chi}$$

New environment:

$$L_{\text{new}} = P^\dagger L'$$

[Figure: SVD spectrum, keep χ largest]

Eckart-Young Theorem

SVD gives the **optimal** rank- χ approximation

Appendix: TMRG Convergence and Observables

Iterate: Grow \rightarrow Truncate \rightarrow Grow \rightarrow Truncate $\rightarrow \dots$

Convergence criterion:

Fixed point: L^*, R^* such that

$$L^* \xrightarrow{\text{grow+truncate}} L^*$$

Free energy:

$$\begin{aligned} f &= -k_B T \lim_{N \rightarrow \infty} \frac{1}{N} \ln Z \\ &= -k_B T \ln \sigma_1^* \end{aligned}$$

Local observables:

$$\langle \sigma_k \rangle = \frac{L^* \cdot \sigma_k \cdot R^*}{L^* \cdot R^*}$$

Correlation functions:

$$\langle \sigma_i \sigma_j \rangle = \frac{L^* \cdot \sigma_i \cdot T^{|i-j|} \cdot \sigma_j \cdot R^*}{Z}$$

Physical Meaning

L^*, R^* encode the **thermodynamic limit**.

Appendix: 2D Row-to-Row — Row as Super-Spin

Idea: Treat one row of L spins as a single “super-spin”.

Row configuration:

$$\vec{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_L)$$

Total states: 2^L .

Row-to-row transfer:

$$(T_{\text{row}})_{\vec{\sigma}, \vec{\sigma}'} = \prod_{i=1}^L e^{\beta J \sigma_i \sigma'_i} \prod_{i=1}^{L-1} e^{\beta J \sigma_i \sigma_{i+1}}$$

$$T_{\text{row}} \in \mathbb{R}^{2^L \times 2^L}.$$

[Figure: Row \rightarrow super-spin with 2^L states]

Key Point

Now $Z = \text{Tr}(T^M)$ looks like 1D problem with 2^L dimensional transfer matrix

Appendix: Row Transfer Matrix = MPO

Key insight: T_{row} has a **tensor network** structure!

Decompose T_{row} into local tensors:

$$T_{\vec{\sigma}, \vec{\sigma}'} = \sum_{\alpha_1, \dots} W_{\alpha_0 \alpha_1}^{\sigma_1 \sigma'_1} W_{\alpha_1 \alpha_2}^{\sigma_2 \sigma'_2} \dots$$

[Figure: MPO structure of T_{row}]

Local tensor W :

$W_{\alpha\beta}^{\sigma\sigma'}$ = local Boltzmann weight

- σ, σ' : spins in rows $n, n+1$
- α, β : auxiliary (horizontal bonds)

No Quantum!

“MPO” is just a factorization of the classical transfer matrix.

Appendix: Boundary = MPS

MPS: Efficient representation of row configuration space.

Full vector:

$$|R\rangle = \sum_{\vec{\sigma}} R_{\vec{\sigma}} |\vec{\sigma}\rangle$$

has 2^L components.

MPS compression:

$$R_{\vec{\sigma}} = A^{\sigma_1} A^{\sigma_2} \dots A^{\sigma_L}$$

Only $L \cdot \chi^2 \cdot 2$ parameters!

[Figure: MPS tensor network]

Physical Meaning

MPS compresses 2^L row configs into χ effective states with **limited entanglement**.

Appendix: Row-to-Row DMRG Algorithm

Combine: MPS (boundary) + MPO (transfer) + SVD (truncation).

- 1 Initialize MPS $|L\rangle$, $|R\rangle$ for boundaries
- 2 **Grow:** Apply MPO T_{row} to boundaries

$$|L'\rangle = T_{\text{row}}|L\rangle$$

- 3 **Truncate:** SVD to compress bond dimension back to χ
- 4 Iterate until convergence to fixed point $|L^*\rangle$, $|R^*\rangle$

Pros

- Systematic, well-understood
- Controlled approximation

Cons

- Breaks 2D rotational symmetry
- Hard to generalize to other lattices

Appendix: No Monte Carlo Sign Problem

Classical models: All Boltzmann weights are **positive**!

$$W = e^{-\beta H} > 0 \quad \text{always}$$

[Figure: Classical (positive) vs Quantum (sign problem)]

Appendix: Why is 3D Difficult?

In 2D: Environment tensors are **1D objects** (edges).

In 3D: Environment would be **2D surfaces** — back to exponential!

[Figure: 2D boundary problem in 3D systems]

Beyond square lattice Ising:

Different lattices:

- Honeycomb
- Triangular
- Kagome

Different models:

- Potts model
- Clock model
- Vertex models

Quantum systems (via iPEPS):

- 2D Heisenberg model
- Frustrated magnets
- Topological phases

Improvements:

- Directional CTMRG
- Full-update vs simple-update
- Gradient optimization

Comparison: iDMRG vs CTMRG

Aspect	iDMRG (1D)	CTMRG (2D)
Dimension	1D chain	2D square lattice
Environment	Left + Right	4 Corners + 4 Edges
Grow step	Add site pair	Add row + column
Truncation	SVD on center bond	SVD on corner boundaries
Bond dimension	χ (MPS)	χ (environment)
Fixed point	L^*, R^*	C_i^*, T_i^*
Computational cost	$O(\chi^3)$	$O(\chi^6)$ or $O(\chi^5)$

Appendix: Computational Complexity

CTMRG scaling:

Operation	Cost	Bottleneck?
Corner absorption	$O(\chi^4 d^2)$	✓
Edge absorption	$O(\chi^3 d^2)$	
Build density matrix	$O(\chi^4)$	
SVD for projector	$O(\chi^3 d^3)$	
Apply truncation	$O(\chi^3 d)$	
Total per iteration	$O(\chi^3 d^3)$ to $O(\chi^6)$	

Comparison

- iDMRG (1D): $O(\chi^3)$ per sweep
- CTMRG (2D): $O(\chi^5) - O(\chi^6)$ per iteration
- More expensive, but still **polynomial** in χ !

Appendix: Key References

Original works:

- R. J. Baxter, *J. Math. Phys.* **9**, 650 (1968) — Corner transfer matrices
- R. J. Baxter, *J. Stat. Phys.* **19**, 461 (1978) — CTM method

Modern CTMRG:

- T. Nishino & K. Okunishi, *J. Phys. Soc. Jpn.* **65**, 891 (1996)
- T. Nishino & K. Okunishi, *J. Phys. Soc. Jpn.* **66**, 3040 (1997)

CTMRG for iPEPS:

- R. Orús & G. Vidal, *Phys. Rev. B* **80**, 094403 (2009)
- P. Corboz et al., *Phys. Rev. B* **84**, 041108(R) (2011)

Reviews:

- R. Orús, *Ann. Phys.* **349**, 117 (2014) — Tensor networks review