

Quantum Machine Learning

State of the Art and Future Directions



Abstract

Since the training of modern machine learning systems is a computation intensive endeavour, the idea of using quantum computing in the machine learning pipeline attracts increasing interest. Indeed, since several quantum machine learning algorithms have already been implemented on present day noisy intermediate-scale quantum devices, experts expect that machine learning on reliable, large scale quantum computers will soon become reality. However, next to potential benefits due to quantum speedup, quantum machine learning may also entail risks regarding reliability, trustworthiness, safety, and security. In preparation of an indepth assessment of these issues, this report therefore surveys the current state of the art in quantum machine learning. It reviews quantum inspired classical algorithms for classical data, genuine quantum algorithms for classical data, classical algorithms for quantum data, and quantum algorithms for quantum data. The particular focus of this review is on methods that are practically feasible on existing quantum computers.

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1 Introduction

One of the most striking and likely most disruptive technological developments in the early 21st century is the maturing of *Artificial Intelligence* (AI). After decades of worldwide research which began in the 1950s, (weakly) intelligent technical systems have now become reality and are increasingly deployed in practice. Prominent examples of recent accomplishments include cognitive systems for robust text- and speech recognition and generation [1, 2, 3], image- and video understanding and synthesis [4, 5, 6], game play [7, 8], recommendation [9, 10], (medical) diagnostics [11, 12], planning and decision making [13, 14, 15], and convincing conversational agents [16, 17, 18].

Achievements like these are mainly due to progress in *Machine Learning* (ML), a branch of AI concerned with adjusting the parameters of a software system in a training process such that it can develop cognitive capabilities. In other words, while there were many innovations in many branches of AI, the recent performance boost is due to systems which learn an intended input-output behavior from analyzing task specific example data [19, 20, 21, 22].

In and off itself, machine learning is not a new idea; it was initially conceptualized by Turing and had its first boom period in the 1980s after algorithms for training artificial neural networks became available. Its accelerated improvement over the past decade can be attributed to four technological trends:

- 1. exponentially growing amounts of (training) data available on the Internet, mainly due to social media platforms and mobile communication devices, but in no small parts also due to the digitization of industry and the life- and natural sciences
- 2. increasing capabilities and decreasing costs of high performance computing hardware, most notably the availability of GPU computing for serious applications
- 3. ever more rapidly developed and publicly disseminated open source software and frameworks for the design, training, and deployment of machine learning systems
- 4. last but not least, scientific progress in neurocomputing, most notably in the areas of deep neural networks and deep learning [23]

Indeed, a conclusion to be drawn from the past decade of machine learning research is that these trends had to come together to enable results such as mentioned above. Although there exist solutions for very capable learning systems of comparatively moderate sizes [24, 25], the capability of a learning system to exhibit human-like performance seems to be predicated on the fact that its size far exceeds traditional system sizes. In other words, the underlying mathematical models seem to require more flexibility and thus more adjustable parameters than have been considered historically.

For instance, the adjustable parameters of an artificial neural network are its synaptic weights and modern deep networks often come with hundreds of billions of such weights (see

e.g. [18]). In order for such flexible models to reliably learn a desired cognitive skill, classical results in learning theory suggest that they must be trained with large amounts of representative training data [26]. Indeed, state of the art systems are commonly trained with up to several billion data points (see e.g. [18]). Given numbers like these, it is clear that the training of modern learning systems is a resource intensive endeavor that easily translates to several hundred GPU years of computation time (see e.g. [18]). Modern machine learning has therefore reached a point where practical feasibility and success are conditioned on access to high performance computing hardware.

Against this backdrop, it is not surprising that machine learning researchers are beginning to look at another technology that has recently seen substantial progress, namely *quantum information processing* (QIP) or *quantum computing* (QC).

First ideas for quantum information processing were published in the 1970s [27, 28, 29, 30] and the theoretical foundations of quantum computing date back to the 1980s when Russian and American physicists first conceptualized quantum computers [31, 32, 33, 34]. The appeal of these early contributions was that they showed that quantum bits can carry more information than classical bits so that quantum computing could be expected to solve certain difficult problems much faster than classically possible.

It then took about another decade until first quantum algorithms were conceived which, when running on a quantum computer, would exhibit these hypothesized advantages [35, 36, 37]. However, at the time these algorithms were presented they were mere theoretical exercises in "pen and paper programming" since quantum computers on which they could have been implemented did not yet exist. Nevertheless, Grover's amplitude amplification algorithm convincingly demonstrated that quantum computers can exhaustively search unstructured lists quadratically faster than digital computers [37, 38] and Shor's celebrated quantum algorithm for large integer factorization even revealed exponential speedup over classical approaches [36, 39]. The latter sparked broader attention to the idea of quantum computing, because Shor's discovery implied that common data encryption schemes would no longer be secure if working quantum computers were to become technical reality. In short, these pioneering contributions therefore demonstrated that quantum algorithms running on ideal quantum computers would indeed exhibit quantum supremacy when dealing with certain difficult problems. They also made clear that this quantum supremacy would have considerable practical implications.

First prototypes of genuine quantum computers appeared in the late 1990s. They were genuine in that they were made up of quantum hardware for physical quantum information processing as opposed to digital hardware for simulated quantum information processing. For example, physical implementations of the Deutsch-Jozsa algorithm [35] were reported in 1998 [40, 41] and Shor's algorithm was first physically implemented by a team at IBM in 2001 [42]. Since then, technical developments progressed quickly and have meanwhile led to commercially accessible quantum computing platforms developed by companies such as IBM, Google, Microsoft, Rigetti, IonQ, or D-Wave.

This state of affairs and the ever intensifying efforts towards marketed solutions confirm that quantum computing really is a viable idea; proof of concepts exist, claims as to first practical observations of quantum supremacy have been made [43], and a growing number of practitioners see remaining challenges in the development of industrial strength quantum computers as mere engineering problems rather than as fundamental research problems.

It therefore seems likely that *quantum machine learning* (QML) will become practical, too, and that quantum information processing will become applicable at various stages of the machine learning pipeline.

Indeed, since many of the fundamental problems in machine learning are demanding optimization or search problems of the kind quantum computing can excel at, research on QML is noticeably intensifying. The corresponding literature has seen rapid growth and first textbooks and tutorials have become available [44, 45, 46, 47, 48]. Moreover, more and more quantum coding challenges are being held that specifically focus on aspects of machine learning and there are recognizable efforts by industrial players to get machine learners involved in quantum computing research.

Given these developments, experts predict that quantum machine learning on reliable, large scale quantum computers will soon facilitate demanding learning tasks and even put problems into reach which are intractable on digital computers. Should these predictions materialize, QML technology will further accelerate progress in AI which, in turn, will likely lead to novel marketable products and commercial solutions impacting economies and societies [49, 50].

However, next to its potential benefits, QML may also come with certain risks. Since learning systems now see commercialization and industrial deployment even in sensitive areas such as autonomous driving or financial services [51, 52, 53, 54], questions as to their ethics, reliability, trustworthiness, and safety are becoming more urgent than when machine learning was a mere academic endeavour [55, 56, 57, 58]. In fact, now that machine learning and quantum computing are beginning to coalesce, these questions are likely to become ever more important. In other words, given the expected impact of quantum machine learning on capabilities and utilizability of artificial cognitive systems, it seems appropriate to asses potential security issues related to quantum machine learning.

In preparation of such an assessment, this present report is intended to take stock and to give an overview of the current state of the art in quantum machine learning. In section 2, we first provide a short introduction to quantum computing, its basic principles, and major paradigms. Next, in section 3, we clarify the notion of machine learning and elaborate on common settings and methods.

Section 4 then surveys and reviews the scientific literature on quantum machine learning. There, we will adhere to the currently established QML taxonomy and consider quantum inspired classical algorithms for classical data, genuine quantum algorithms for classical data, classical algorithms for quantum data, and quantum algorithms for quantum data. Our pre-

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sentation with respect to genuine quantum algorithms will assume a point of view where we abstract from physical aspects of the hardware of quantum computers. Rather, we will follow common practice and focus on logical aspects or higher level abstractions that programmers and software developers are typically exposed to.

Having said this, it is nevertheless important to acknowledge that, as of this writing, quantum algorithm design is still very much a theoretic endeavor and that many theoretically valid ideas cannot be implemented on present day quantum computing devices. In other words, physical or hardware related aspects have to be kept in mind when assessing the validity of proposed quantum machine learning solutions. To this end, section 5 will review common assumptions as to data pre- and post-processing, robustness against measurement noise, numbers of available quantum bits, or required quantum circuit depths and will also asses how realistic such assumptions are and what they mean for the practical feasibility of currently proposed QML solutions.

Finally, in section 6, we will summarize key points of this report and provide an outlook to planned studies on quantum machine learning security.

2 Quantum Computing in a Nutshell

Quantum computing exploits quantum mechanical phenomena for information processing. This offers great computational power but also requires a different kind of thinking than in the domain of digital computing. The latter may explain why quantum computing has still not yet received widespread attention in mainstream computer science. As of this writing, major obstacles for a broader engagement with the topic likely are as follows: 1) quantum mechanical phenomena such as superposition, entanglement, tunneling, decoherence, or the uncertainty principle appear to be weird, abstract or unintuitive, and hard to accept for they cannot be observed in our daily macroscopic environments; 2) the mathematical tools used to model these phenomena are complex (no pun intended) and again rather abstract; and 3) to the uninitiated, the mathematical notation used in the quantum computing literature appears even more abstract and needs getting used to. Acknowledging these difficulties, this section provides a brief introduction to the basic terminology of- and fundamental concepts behind quantum computing.

To begin with, we recall that a *quantum mechanical system* (QM system) is characterized by a state vector $|\psi\rangle$ in a Hilbert space $\mathbb H$ over $\mathbb C$. The temporal evolution of such a system is governed by the Schrödinger equation and its observable physical properties (such as position, momentum, or energy) correspond to Hermitian operators. A measurement of an observable property of a QM system collapses its state to an eigenvector of the respective operator. This is to say that a measurement of a quantum variable results in an eigenvalue of the operator under consideration and the state "jumps" to the corresponding eigenvector. Which of the eigenvalues and corresponding eigenvectors this will be depends on certain probabilities encoded in the state vector. The crucial question as to why this kind of mathematics describes the behavior of nature on one of its most fundamental levels still awaits conclusive answers; for now, the empirically undeniable success of this linear algebraic framework has to be attributed to "the unreasonable effectiveness of mathematics in the natural sciences" [59].

Quantum computing (QC) deals with quantum mechanical systems that represent quantum bits or qubits. All throughout our following discussion, we will focus on logical qubits which are the abstract basic building blocks on which quantum algorithms operate. A physical qubit, on the other hand, is a physical realization of a logical qubit, namely a physical device that behaves like a logical qubit and forms a component of the hardware of a quantum computing system. We will neither discuss possible technologies for building such devices, nor will we discuss the practical advantages and disadvantages of different such technologies. An extensive overview of the state of the art in this area can be found in a recent study published by the Federal Office for Information Security [60].

From the point of view of quantum algorithms, our focus on logical qubits makes sense. It mimics common levels of abstraction in modern software development where programmers usually do not have to pay attention to hardware details. In fact, the distinction between hard- and software aspects of computing was essential for the success of digital computers

once they matured. Current quantum computers, however, have not yet reached this level of maturity. They are *noisy intermediate-scale quantum* (NISQ) devices [61] which contain less than a hundred physical qubits and still suffer from limited coherence times and low fault-tolerance. From the point of view of existing quantum hardware, logical qubits are therefore still an idealization since they abstract away technological shortcomings. This may render some of the mathematically valid ideas for quantum algorithms practically infeasible and, as of this writing, constitutes yet another difference between quantum- and digital computing. We will return to this issue in section 5.

On a classical computer, the basic units of information are bits and the mathematics that describes their behavior is Boolean algebra. On a quantum computer, the basic units of information are qubits and the mathematics that models their behavior is complex linear algebra.

While a classical bit is always in one and only one of two possible states (0 or 1), the basic tenet of quantum information processing is that a qubit exists in a *superposition* of two *basis states* and collapses to either one once measured. Examples of well known physical *two-state quantum systems* that exhibit this phenomenon include the polarization of a photon, the spin of an electron, or the ground- and first excited state of an atom. Yet, it is interesting to note that quantum mechanical superposition is not exclusive to the subatomic world but may just as well occur in completely isolated macroscopic systems. This is, for instance, used in quantum information processing devices made of superconducting circuits in which electrical currents flow in two directions simultaneously.

Superposition is crucial because it implies that the state space of a qubit is not confined to only two states. It rather consists of infinitely many states which can be represented as two-dimensional, complex-valued unit vectors that are linear combinations of two distinguished, linearly independent, orthonormal basis states. Using the Dirac notation, these basis states are commonly written as $|0\rangle$ and $|1\rangle$ and typically thought of as

$$|0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
 and $|1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ (1)

In terms of a mathematical equation, the above translates to the statement that (the state vector of) a qubit can be written as

$$|\psi\rangle = \alpha_0 |0\rangle + \alpha_1 |1\rangle \tag{2}$$

where the coefficients $\alpha_0, \alpha_1 \in \mathbb{C}$ are called the *amplitudes* of the basis states $|0\rangle$ and $|1\rangle$, respectively. Importantly, these amplitudes have to obey the normalization condition

$$|\alpha_0|^2 + |\alpha_1|^2 = 1 \tag{3}$$

and are interpreted as follows: if a measurement is performed on qubit $|\psi\rangle$, the probability of finding it in basis state $|0\rangle$ is $|\alpha_0|^2$ whereas the probability of finding it in basis state $|1\rangle$ corresponds to $|\alpha_1|^2$. In other words, the probability of measuring the qubit in, for instance, state $|0\rangle$ is given by $|\langle 0|\psi\rangle|^2$, because the inner product of $|0\rangle$ and $|\psi\rangle$ evaluates to

$$\langle 0|\psi\rangle = \alpha_0 \langle 0|0\rangle + \alpha_1 \langle 0|1\rangle = \alpha_0 \cdot 1 + \alpha_1 \cdot 0 = \alpha_0 \tag{4}$$

Note that measurements performed on $|\psi\rangle$ are irreversible operations because they constitute interactions with the outside world and therefore lead to loss of superposition or quantum *decoherence*. In other words, once a qubit has collapsed to either one of its basis states, it henceforth behaves like a classical bit.

Operations on qubits which preserve their quantum mechanical nature are called *reversible*. Mathematically, these correspond to unitary linear operators $U \in SU_2(\mathbb{C})$ for which $UU^\dagger = U^\dagger U = I$. Reversible operators can also be written as $U = e^{-iHt/\hbar}$ where H is yet another operator called the Hamiltonian. It corresponds to the total energy of a quantum system in the sense that its spectrum is the set of possible outcomes of measurements of the system's total energy.

Another tenet of quantum information processing is that qubits can be combined to form qubit registers. While a single qubit $|\psi\rangle$ exists in a superposition of 2 states, a quantum register $|\psi\rangle$ of n qubits exists in a superposition of 2^n states. Mathematically, this is to say that

$$\ket{\psi} = \sum_{i=0}^{2^n-1} \gamma_i \ket{\psi_i}$$
 (5)

Here, the γ_i once again obey $\sum_i |\gamma_i|^2 = 1$ and the $|\psi_i\rangle$ are 2^n -dimensional tensor products of single qubit states. Mathematically, these tensor products are Kronecker products and behave as follows: If we consider two single qubits represented as

$$|\psi\rangle = \alpha_0 |0\rangle + \alpha_1 |1\rangle = \begin{bmatrix} \alpha_0 \\ \alpha_1 \end{bmatrix}$$
 and $|\phi\rangle = \beta_0 |0\rangle + \beta_1 |1\rangle = \begin{bmatrix} \beta_0 \\ \beta_1 \end{bmatrix}$ (6)

then the state of a system composed of these two qubits is given by the four-dimensional vector

$$|\psi\rangle\otimes|\phi\rangle = \begin{bmatrix} \alpha_{0} \\ \alpha_{1} \end{bmatrix} \otimes \begin{bmatrix} \beta_{0} \\ \beta_{1} \end{bmatrix} = \begin{bmatrix} \alpha_{0}\beta_{0} \\ \alpha_{0}\beta_{1} \\ \alpha_{1}\beta_{0} \\ \alpha_{1}\beta_{1} \end{bmatrix} = \alpha_{0}\beta_{0}|0\rangle\otimes|0\rangle + \alpha_{0}\beta_{1}|0\rangle\otimes|1\rangle + \alpha_{1}\beta_{0}|1\rangle\otimes|0\rangle + \alpha_{1}\beta_{1}|1\rangle\otimes|1\rangle$$
(7)

This example illustrates that the 2^n dimensional state space of an n qubit system is spanned by a *computational basis* that consists of 2^n basis vectors which are tensor products of individual basis states. Since such tensor products over basis states occur in many quantum computing equations, they are typically written more succinctly. For instance, for a quantum register where n=3, we would write the $2^3=8$ basis states as

$$|\psi_0\rangle = |0\rangle \otimes |0\rangle \otimes |0\rangle \equiv |000\rangle \tag{8}$$

$$|\psi_1\rangle = |0\rangle \otimes |0\rangle \otimes |1\rangle \equiv |001\rangle$$
 (9)

:

$$|\psi_7\rangle = |1\rangle \otimes |1\rangle \otimes |1\rangle \equiv |111\rangle$$
 (10)

or, even shorter as $|i\rangle$ where $i\in\{0,\dots,7\}$. Since the state space of an n qubit system is 2^n -dimensional, operators acting on such states are (products of) $2^n\times 2^n$ dimensional unitary matrices. Yet another tenet of quantum information processing is that qubits can be entangled. The physical phenomenon of entanglement is indeed an essential aspect of quantum information that has no immediate classical counterpart. To understand its characteristics, we note that the 2 qubit system in (7) has a complete state that is a tensor product of individual qubit states. However, systems of qubits can also be in states that can not be expressed in terms of tensor products. These are called *entangled states* and a canonical example of such a state is

$$|\psi\rangle = \frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle \tag{11}$$

The following simple computations will illustrate why entangled states are special: Consider two systems $|\psi_1\rangle$ and $|\psi_2\rangle$ of 2 qubits where $|\psi_1\rangle$ is in a tensor product state

$$|\psi_1\rangle = \gamma_0|00\rangle + \gamma_1|01\rangle + \gamma_2|10\rangle + \gamma_3|11\rangle \tag{12}$$

and $|\psi_2\rangle$ is in the entangled state

$$|\psi_2\rangle = \frac{1}{\sqrt{2}}|00\rangle + \frac{1}{\sqrt{2}}|11\rangle \tag{13}$$

If we now perform a partial measurement of both systems where only the first qubit of either system is measured and the second one is left intact, the probability of finding the first qubit of system $|\psi_1\rangle$ in, say, state $|0\rangle$ is given by

$$\sum_{x \in \{0,1\}} |\langle 0x | \psi_1 \rangle|^2 = |\gamma_0|^2 + |\gamma_1|^2 \tag{14}$$

Moreover, once this partial measurement has been performed, the system will be in the new (re-normalized) state

$$|\hat{\psi}_{1}\rangle = \frac{\sum_{x \in \{0,1\}} \langle 0x | \psi_{1} \rangle | 0x \rangle}{\sqrt{\sum_{x \in \{0,1\}} |\langle 0x | \psi_{1} \rangle|^{2}}} = \frac{\gamma_{0} |00\rangle + \gamma_{1} |01\rangle}{\sqrt{|\gamma_{0}|^{2} + |\gamma_{1}|^{2}}}$$
(15)

In other words, measuring the first qubit of the first system will cause this system to still be in a superposition of two computational basis states. In our example, this means that the state of the second qubit remains undetermined since it could still be found either in state $|0\rangle$ or in state $|1\rangle$.

If, on the other hand, the same kind of partial measurement is performed on the second system $|\psi_2\rangle$, the probability of finding its first qubit in state $|0\rangle$ is given by

$$\sum_{x \in \{0,1\}} |\langle 0x | \psi_2 \rangle|^2 = |\frac{1}{\sqrt{2}}|^2 = \frac{1}{2}$$
 (16)

and, upon this measurement, the second system will be in state

$$|\hat{\psi}_2\rangle = \frac{\sum_{x \in \{0,1\}} \langle 0x | \psi_2 \rangle |0x\rangle}{\sqrt{\sum_{x \in \{0,1\}} |\langle 0x | \psi_2 \rangle|^2}} = |00\rangle$$
(17)

By the same token, if we had found the first qubit in state $|1\rangle$, the final state of the system would have been $|\hat{\psi}_2\rangle = |11\rangle$ and similar outcomes would have been observed if we had measured the second qubit of the entangled system. That is, whenever two qubits are entangled, their individual states cannot be measured separately; a measurement of either one of two entangled qubits also determines the state of the other.

Entanglement is of vital importance for quantum information processing because it is this phenomenon which makes it possible to consider only n (physical) qubits to perform computation in a 2^n dimensional state space. In other words, whereas doubling the "processing power" of a digital computer would require doubling the number of bits, doubling the "processing power" of a quantum computer only requires adding one extra qubit. Without entanglement, this exponential increase in performance is provably impossible [62].

Given these prerequisites, we can now discuss the two major paradigms for practical quantum computing, namely quantum gate computing and adiabatic quantum computing. The former approaches problem solving by sequencing unitary operators, U_1, U_2, \ldots into quantum circuits and well known examples include Grover's search algorithm [38] or Shor's integer factorization algorithm [39]. The latter considers problem solving as an energy minimization task and is concerned with finding ground states of problem specific Hamiltonian operators H. Adiabatic quantum computing is sometimes falsely said to be less general than quantum gate computing but both paradigms are provably equivalent with respect to computational power or expressiveness [63, 64]). Indeed, translating a quantum algorithm from one formalism to the other requires only polynomial efforts (for instance, in the number of additionally required qubits or quantum gates) and is therefore possible in principle. Nevertheless, from the programmer's point of view, either approach requires its own distinct form of thinking and of formalizing problems and we next elaborate on the underlying principles.

2.1 Adiabatic Quantum Computing

Adiabatic quantum computing (AQC), often also called adiabatic quantum optimization (AQO), relies on yet another quantum mechanical phenomenon. This phenomenon is formalized in the adiabatic theorem [65] which states that, if a quantum system starts in the ground state of a Hamiltonian operator which then gradually changes over a period of time, the system will end up in the ground state of the resulting Hamiltonian. Since Hamiltonians are energy operators, their ground states correspond to the lowest energy state of the quantum system under consideration. The adiabatic theorem therefore characterizes a form of quantum tunneling which refers to the fact that quantum systems can tunnel through energy barriers.

To harness AQC for problem solving, one prepares a system of qubits in the ground state of a simple, problem independent Hamiltonian and then adiabatically evolves it towards a Hamiltonian whose ground state corresponds to a solution to the problem at hand [66]. This can be done on D-Wave computers [67, 68, 69] which are particularly tailored towards solving quadratic unconstrained binary optimization problems (QUBOs) of the form

$$s^* = \underset{s \in \{-1, +1\}^n}{\operatorname{argmin}} s^{\mathsf{T}} Q s + q^{\mathsf{T}} s. \tag{18}$$

Here, the 2^n different bipolar vectors s over which to minimize represent possible global states of a system of n entities each of which might be in one of two local states (+1 or -1). The coupling matrix $\mathbf{Q} \in \mathbb{R}^{n \times n}$ models internal interactions and the vector $\mathbf{q} \in \mathbb{R}^n$ models additional constraints or external influences.

In physics, QUBOs are known as Ising energy minimization problems [70] and, in machine learning, they are fundamental to the theory of Hopfield networks [71]. QUBOs pose discrete- or combinatorial optimization problems which are NP-hard and therefore notoriously difficult to solve. They also are surprisingly universal and of considerable practical importance since they often arise in the context of subset selection- or bipartition problems.

Given a set \mathcal{X} of n elements x_i , subset selection problems ask for a subset $\mathcal{X}' \subset \mathcal{X}$ such that the elements of \mathcal{X}' meet certain criteria and bipartition problems ask for a partition $\mathcal{X} = \mathcal{X}_1 \cup \mathcal{X}_2$ such that the elements of either of the disjoint subsets \mathcal{X}_1 and \mathcal{X}_2 meet certain criteria. Now, if a subset selection- or bipartition problem can be (re)written as a QUBO over bipolar vectors s, the entries of the solution s^* can be understood as indicator variables. For a subset selection problem, they define the sought after subset as $\mathcal{X}' = \{x_i \in \mathcal{X} \mid s_i^* = +1\}$ and, for a bipartition problem, we would have $\mathcal{X}_1 = \{x_i \in \mathcal{X} \mid s_i^* = +1\}$ and $\mathcal{X}_2 = \{x_i \in \mathcal{X} \mid s_i^* = -1\}$, respectively.

Subset selection- or bipartition problems of this kind abound in data mining, pattern recognition, machine learning, or artificial intelligence. General use cases can be found in data base search, the computation of medoids in k-medoids clustering, or the identification of support vectors in support vector machine training. Prominent practical applications where QUBOs are central include verification-, planning-, or assignment problems in areas such as logistics or finance [72, 73, 74, 75].

What renders quantum computing an attractive approach for dealing with difficult QU-BOs is that seminal work by Farhi et al. [76] provided a simple general recipe for how to solve them via adiabatic evolution.

Assuming the parameters Q and q of a QUBO to be given, the basic idea is to consider a time dependent system of n qubits

$$|\psi(t)\rangle = \sum_{i=0}^{2^n-1} \alpha_i(t) |\psi_i\rangle.$$
 (19)

which evolves under a time-dependent Hamiltonian H(t) so that its behavior is governed by the Schrödinger equation

$$i\hbar \frac{d}{dt} |\psi(t)\rangle = H(t) |\psi(t)\rangle$$
 (20)

The qubits are prepared in the ground state of a beginning Hamiltonian H_B which gradually changes towards the problem Hamiltonian H_P . For the latter, [76] proposes

$$H_P = \sum_{i=1}^{n} \sum_{j=1}^{n} Q_{ij} \, \sigma_z^i \, \sigma_z^j + \sum_{i=1}^{n} q_i \, \sigma_z^i$$
 (21)

where $\sigma_z^i = I \otimes \ldots \otimes I \otimes \sigma_z \otimes I \otimes \ldots \otimes I$ denotes the Pauli spin matrix σ_z acting on the *i*-th qubit. For the beginning Hamiltonian, [76] suggests

$$H_B = -\sum_{i=1}^n \sigma_x^i \tag{22}$$

where σ_x^i is the Pauli spin matrix σ_x acting on the i-th qubit. Considering an evolution from t=0 to t=T, the Hamiltonian in (20) is assumed to be a convex combination

$$H(t) = \left(1 - \frac{t}{T}\right) \cdot H_B + \frac{t}{T} \cdot H_P \tag{23}$$

and can be used to evolve $|\psi(t)\rangle$ from $|\psi(0)\rangle$ to $|\psi(T)\rangle$ where $|\psi(0)\rangle$ is the ground state of H_B .

Finally, at time T, a measurement is performed on the qubits. This causes the whole system to collapse to one of its 2^n basis states and the probability for this state to be $|\psi_i\rangle$ is given by the amplitude $|\alpha_i(T)|^2$. However, since the adiabatic evolution was steered towards the problem Hamiltonian H_P , basis states $|\psi_i\rangle$ that correspond to ground states of H_P are more likely to be found.

The efficiency of this computational paradigm depends on the choice of T which is known to depend on the minimum energy gap between the ground- and first excited state of H(t). While exponentially small gaps will render adiabatic quantum computing inefficient, problems with gaps that scale inverse polynomially can be solved efficiently. In fact, it is known that energy gaps are inversely proportional to the square root of the number of basis states that have energies close to global minimum [77]. For problems with a small number of valid solutions, an appropriate choice is $T \in O(\sqrt{2^n})$ [78]. This, in turn, is to say that adiabatic quantum computing can solve certain binary optimization problems quadratically faster than classically possible.

To illustrate the steps involved in setting up practical problems for solution on an adiabatic quantum computer and to better explain the process of adiabatic optimization, we briefly discuss an application example.

The *max-sum diversification problem* is of importance in location- or portfolio optimization and also arises in the context of information retrieval, recommendation, clustering, and other settings where one has to identify very distinct elements in a given data set. To be more specific, given a set $\mathcal{X} = \{x_1, \dots, x_n\}$ and an appropriate, problem specific distance measure $d(\cdot, \cdot)$, max-sum diversification asks for a subset $\mathcal{S}^* \subset \mathcal{X}$ of k < n elements of maximum dispersion and thus consists in solving

$$\mathcal{S}^* = \underset{\mathcal{S} \subset \mathcal{X}}{\operatorname{argmax}} \sum_{\boldsymbol{x}_i \in \mathcal{S}} \sum_{\boldsymbol{x}_j \in \mathcal{S}} d(\boldsymbol{x}_i, \boldsymbol{x}_j)$$
s.t. $|\mathcal{S}| = k$. (24)

Collecting the $d(\boldsymbol{x}_i, \boldsymbol{x}_j)$ in a distance matrix $\boldsymbol{D} \in \mathbb{R}^{n \times n}$ and considering a binary indicator vector $\boldsymbol{z} \in \{0,1\}^n$ where $z_i = 1$ if $\boldsymbol{x}_i \in \mathcal{S}$ and $z_i = 0$ otherwise and letting $\boldsymbol{1} \in \mathbb{R}^n$ denote the vector of all ones, the problem in (24) can also be written as

$$z^* = \underset{z \in \{0,1\}^n}{\operatorname{argmax}} z^{\mathsf{T}} D z$$
s.t.
$$1^{\mathsf{T}} z = k$$
(25)

Furthermore, when introducing a generic Lagrange multiplier λ and using the fact that z^TDz is quadratic in z, this NP-complete integer programming problem can be written as a QUBO over z, namely

$$z^* = \underset{z \in \{0,1\}^n}{\operatorname{argmin}} - z^{\mathsf{T}} D z + \lambda \left(\mathbf{1}^{\mathsf{T}} z - k \right)^2 \equiv \underset{z \in \{0,1\}^n}{\operatorname{argmin}} - z^{\mathsf{T}} P z - p^{\mathsf{T}} z$$
 (26)

However, the decision variables (18) are bipolar vectors $s \in \{-1, +1\}^n$ whereas (26) minimizes over binary vectors $z \in \{0, 1\}^n$. We therefore emphasize a fact often used when (re)writing problems for adiabatic quantum computing, namely that it is easy to convert between binary and bipolar vectors, because $z = (s + 1)/2 \Leftrightarrow s = 2 \cdot z - 1$.

Hence, when plugging z=(s+1)/2 into (26), it becomes an exercise in algebra [79] to show that max-sum diversification consists in solving the following energy minimization problem

$$s^* = \underset{s \in \{-1, +1\}^n}{\operatorname{argmin}} - s^{\mathsf{T}} Q \, s + q^{\mathsf{T}} s \tag{27}$$

whose parameters amount to $Q = -\frac{1}{4} \left(D - \lambda \, \mathbf{1} \mathbf{1}^{\mathsf{T}} \right)$ and $q = -\frac{1}{2} \left(D - \lambda \, \mathbf{1} \mathbf{1}^{\mathsf{T}} \right) \mathbf{1} - \lambda \, k \, \mathbf{1}$. With respect to the original formulation in (24), we note that, once the solution s^* to (27) has been found, entries $s_i^* = +1$ indicate which $x_i \in \mathcal{X}$ to select into \mathcal{S}^* . Figure 1 shows a specific, rather didactic instance of the problem of max-sum diversification and illustrates the inner workings of the AQC approach towards solving it. The data in Fig. 1(a) consists of a set \mathcal{X} of n=12 data points $x_i \in \mathbb{R}^2$, each of which represents monthly climate conditions (average temperature and precipitation) in the city of Hamburg. Given this data, the task is to determine k=4 months of diverse climatic characteristics.

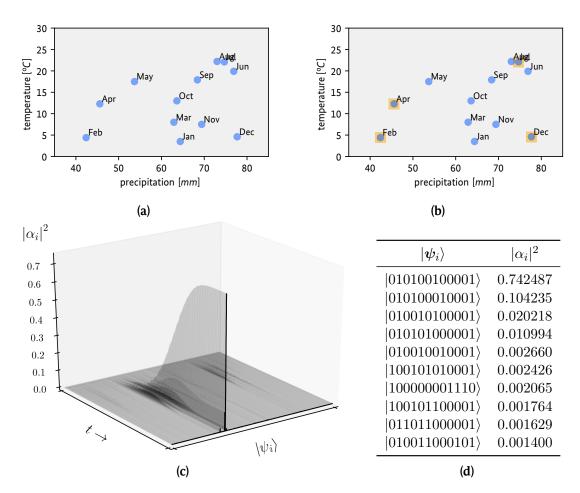


Figure 1: Adiabatic quantum computation for max-sum diversification. (a) A didactic data set of n=12 two-dimensional data points representing monthly climate conditions in the city of Hamburg. (b) Result obtained for k=4 and Euclidean distances between data points. (c) Visualization of the evolution of the amplitudes of a corresponding 12 qubit system $|\psi(t)\rangle$. During its evolution over time t, the system is in a superposition of $2^{12}=4096$ basis states $|\psi_i\rangle$ which represent potential solutions to the problem. Initially, each potential solution (reasonable or not) is equally likely. At the end of the process, one basis state has a noticeably higher amplitude $|\alpha_i|^2$ than the others and is thus more likely to be measured. (d) Ranking of the ten most likely states for the qubit system to be found in at the end of the adiabatic evolution; the likeliest final state indexes the months of February, April, July, and December and thus represents the solution shown in (b).

To set up the QUBO parameters Q and q, the data was normalized to zero mean and unit variance (a common step in machine learning). After this normalization, Euclidean distances $D_{ij} = \|x_i - x_j\|$ between far apart data points are of the order of $\mathcal{O}(2)$ so that setting $\lambda = 2n$ will cause neither of the terms in (27) to dominate the minimization problem.

Given ${\bf Q}$ and ${\bf q}$, one can prepare the beginning- and problem Hamiltonian according to the above recipe and then let an n=12 qubit system $|\psi(t)\rangle$ adiabatically evolve over $T\in \mathcal{O}(\sqrt{2^{12}})$ time steps. To visualize the dynamics of this adiabatic quantum search for a solution, the process was simulated on a digital computer and the quantum computing toolbox QuTiP [80] was used to solve the Schrödinger equation in (20).

Figure 1(c) shows how the qubit system evolves after it has been prepared in a superposition of $2^{12}=4096$ basis states $|0000000000000\rangle, |000000000001\rangle, \ldots, |11111111111111\rangle$ each of which represents a subset of the given set of data points. In particular, the figure visualizes the behavior of the amplitudes $|\alpha_i(t)|^2$ of the states the system can be measured in. At time t=0, all basis states are equally likely but over time their amplitudes begin to diverge; amplitudes of basis states that correspond to low energy states of the problem Hamiltonian increase while amplitudes of basis states that could hardly be considered a solution to the problem decrease. At t=T, certain basis states are therefore more likely to be measured than others and the table in Fig. 1(d) ranks the ten most likely ones.

The most likely final state in this example is $|010100100001\rangle$ which, when understood as an indicator vector, indexes the subset consisting of the months of February, April, July, and December which are indeed diverse with respect to their climatic conditions (see Fig. 1(b)). Interestingly, the next most likely solution $|010100010001\rangle$ would swap July for August and, looking at the data in Fig. 1, this behavior of the quantum computing algorithm appears reasonable.

2.2 Quantum Gate Computing

Beside AQC, *Quantum Gate Computing* (QGC) is the dominant paradigm for quantum information processing systems. A detailed introduction into that topic can be found in [81].

Whether we are using qubits or bits, we need to manipulate them in order to turn the inputs we have into the outputs we need. Thus, QGC borrows many of its defining properties from classical computers: On a low level, digital computers process information by manipu-

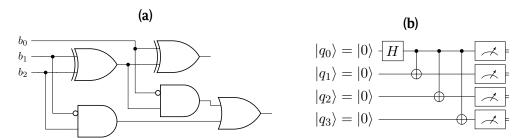


Figure 2: Left: Example of a digital subtractor circuit. Right: Example of a quantum gate circuit for generating a 4-qubit Bell state.

lating sets of bits via atomic logical operations between them. Input bits, being either in the state 0 or 1, are processed by so-called gates. Each gate performs a fundamental logical operation, e.g., AND, OR, XOR, NOT. For computations on a few bits, it is useful to represent this process in a diagram known as a circuit diagram or just circuit. These have input bits on the left, output bits on the right, and gates represented by specific symbols in between. An exemplary circuit that computes the difference of two bits, given a carry bit, is shown in Fig. 2 (a).

The evolution of a closed quantum system is described by a unitary transformation. Thus, all operations that we perform on qubits must be unitary too. In theory and practice, such operators typically act on only one or two qubits at a time but can be sequenced to form more complicated qubit transformations. Borrowing terminology from digital computing, quantum operators acting on qubits are also called *quantum gates*. An exemplary quantum gate circuit that creates a Bell state is shown in Fig. 2 (b).

A crucial difference to the classical setting is that, being unitary operators, quantum gates have to have the same number of input and outputs. Moreover, since the effect of any unitary gate U can be reversed by its Hermitian conjugate U^{\dagger} , a quantum gate's input can always be reconstructed from its output. For many classical gates such as the AND gate, this is not the case. At first sight, this reversibility constraint on quantum gates therefore seems to render them less expressive or powerful than classical ones. However, by introducing (several) ancillary bits, irreversible classical gates can be made reversible and there exist universal sets of reversible primitives [82]. Reversibility does therefore not restrict computability. Quantum gates can instead be seen as a generalization of classical reversible gates and the use of ancillary qubits is common practice in quantum computing.

Let us quickly recall the frequently appearing quantum gates and their relation to classical digital logic. The X-gate constitutes the counterpart of the logical NOT-gate. Applying X to a qubit switches the amplitudes of $|0\rangle$ and $|1\rangle$.

$$X|0\rangle = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} = |1\rangle$$

Closely related are the Y-gate

$$Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} = -i|0\rangle\langle 1| + i|1\rangle\langle 0|$$

and the Z-gate

$$Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = |0\rangle\langle 0| - |1\rangle\langle 1|$$
.

At a first glace, X, Y, and Z have not much in common. However, each of them can be interpreted as a rotation around the corresponding axis (x-axis, y-axis, z-axis) by π radians in a conceived 3-dimensional coordinate system. However, applying the Z-gate to one of the basis states $|0\rangle$, $|1\rangle$ seems to have no effect:

$$Z|0\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} = |0\rangle \qquad \qquad Z|1\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix} = -|1\rangle$$

Clearly, $|0\rangle$ and $|1\rangle$ are eigenstates of Z. They form the so-called Z-basis. There is, in fact, an infinite number of bases from which some have specific names in the context of quantum gate computing. One that appears directly in various quantum algorithms is the X-basis, constituted by $|+\rangle$ and $|-\rangle$.

$$|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\1 \end{bmatrix} \qquad \qquad |-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) = \frac{1}{\sqrt{2}} \begin{bmatrix} 1\\-1 \end{bmatrix}$$

The specific reason for why this basis appears frequently is the Hadamard gate, also known as H-gate. It allows us to create a superposition of $|0\rangle$ and $|1\rangle$, effectively realizing a uniform distribution over $\{0,1\}$.

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1\\ 1 & -1 \end{bmatrix}$$

Interpreted as a change of basis, it allows us to move from the Z-basis to the X-basis:

$$H|0\rangle = |+\rangle$$
 $H|1\rangle = |-\rangle$

An important insight for digital circuits is the universality of specific gate sets. There is no need to have hardware implementations of all possible logical 2-bit gates available—it can be shown that a single NAND-gate suffices to implement all Boolean operations. Similar relations can be found for quantum gates as well, e.g. the following equivalence:

$$HZH = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = X$$

Indeed, there is a finite number of possible logic gates. However, quantum circuits allow for *parametrized gates*. These gates require numeric parameters to define their actual function. They are of utmost importance for quantum machine learning in general, since they allow us to define families of functions via a single circuit. Such families constitute model classes for which classical machine learning techniques can be applied to select those models which fit best to some user specified data. We will revisit this topic in Section 4.

The P-Gate (also known as *phase gate*) generalizes the Z-gate. It performs a rotation by ϕ radians around the Z-axis.

$$P(\phi) = \begin{bmatrix} 1 & 0 \\ 0 & \exp(i\phi) \end{bmatrix}$$

Due to closedness, all quantum gates can in fact be absorbed into a single canonical representation. This U-gate is parameterized and subsumes all possible single qubit gates.

$$U(\theta,\phi,\lambda) = \begin{bmatrix} \cos\left(\frac{\theta}{2}\right) & -\exp(i\lambda)\sin\left(\frac{\theta}{2}\right) \\ \exp(i\phi)\sin\left(\frac{\theta}{2}\right) & \exp(i(\phi+\lambda))\cos\left(\frac{\theta}{2}\right) \end{bmatrix}$$

All gates discussed so far act only on a single qubit. However, multi-qubit operations are required to realize non-trivial algorithms and to facilitate quantum entanglement between qubits. The most common way to work with multiple qubits is via *controlled not* (CNOT) gates. They are the quantum counterpart to the classical XOR-gate. While single qubit gates can be fully described by a single matrix, CNOT has two representations,

$$CNOT = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \qquad CNOT = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

depending on which qubit is the controller and which is the target. The semantic of the CNOT is as follows: if the control-qubit is in state $|0\rangle$, the target remains unchanged. If the control-qubit is in state $|1\rangle$, an X-gate (NOT-gate) is applied to the target. This effect can be applied to create a Bell state over n qubits: One initializes all qubits to $|0\rangle$ and applies a H-gate to one of them, say, q_0 . q_0 is then in a uniform superposition of $|0\rangle$ and $|1\rangle$. Then, n-1 CNOT-gates are added, connecting q_0 with each of the remaining n-1 qubits. In all cases, q_0 is the controller. Now, measuring any of the n qubits in state $|0\rangle$ (or $|1\rangle$) implies that q_0 must be $|0\rangle$ (or $|1\rangle$) too, and hence all other qubits. In other words, measuring any of the n qubits determines the state of all remaining qubits. The resulting circuit is depicted in Fig. 2 (b) for n=4.

As mentioned earlier, the action of any quantum circuit can be written as a product of unitaries.

$$U = U_1 U_2 U_3 \dots U_d$$

Here, d is called the *depth* of the circuit. Formally, each 1-qubit and 2-qubit operator has to be extended to an 2^n -dimensional operator in order to fit into the above expression.

Measurements are special in that they enforce the quantum state to collapse into a classical one. The state will lose its probabilistic nature and all subsequent measurements of the outputs will yield the same result deterministically. A simple but important example of a measurement is that of a qubit in the computational basis. This is a measurement on a single qubit with two outcomes defined by the two measurement operators $M_0 = |0\rangle\langle 0|$ and $M_1 = |1\rangle\langle 1|$. We observe that each measurement operator is Hermitian, and that $M_0^2 = M_0$ as well as $M_1^2 = M_1$. Thus, $I = M_0^\dagger M_0 + M_1^\dagger M_1 = M_0 + M_1$. Suppose the state being measured is $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$. Then, the probability of obtaining the measurement outcome 0 is

$$\mathbb{P}(0) = |\psi\rangle M_0^{\dagger} M_0 \langle \psi| = |\psi\rangle M_0 \langle \psi| = |\alpha|^2$$

Similarly, the probability of obtaining the measurement outcome 1 is $\mathbb{P}(1) = |\beta|^2$.

2.3 Further Quantum Computing Paradigms

The above discussion might lead to a misconception, namely that adiabatic quantum computing and quantum gate computing are the only "paradigms" in the broad field of quantum information processing. This short section is therefore supposed to help avoid possible

confusion. The following briefly discusses the notions of topological quantum computing and variational quantum computing. The latter is of considerable interest in the context of quantum machine learning and will be discussed in detail in section 4.

2.3.1 Topological Quantum Computing

Another term ostensibly similar to adiabatic quantum computing or quantum gate computing is *topological quantum computing* [83]. However, whereas the former two refer to logical aspects of quantum computing and abstract away from how to physically implement logical qubits, the latter refers to physical aspects and a physical implementation of qubits.

We therefore only briefly mention that a topological quantum computer is still a mainly theoretical device that considers two-dimensional quasi-particles called anyons, their braids, and their characteristics in three-dimensional space-time. If such a quantum computer could be built, for instance utilizing quantum Hall effects, an expected advantage over currently common technical realizations based on, say, superconducting circuits or trapped ions, is that it would be more robust against decoherence.

2.3.2 Variational Quantum Computing

Variational quantum algorithms (VQA) and variational quantum approximate optimization algorithms (QAOA) combine quantum computations with classical computations in iterative feedback loops [84, 85]. This is supposed to address limitation of current NISQ devices (limited numbers of qubits, limited circuit depths, measurement noise) and aims at harnessing the best of both worlds (classic and quantum).

Either classical optimizers are used to design working quantum circuits, or quantum computation outcomes are measured and classical analysis is used to update Hamiltonian operators for the next round of quantum computations. By now, VQAs have been proposed for many applications and they appear to be the currently best way of exploiting quantum advantages in the NISQ era [86]. Indeed, one can show that variational quantum computing is universal and that efficient input/output strategies for looped classic-to-quantum optimization are possible [87]. Corresponding methods are also known as hybrid quantum-classical computing methods and we will return to these ideas in section 4.

2.4 Technical State of the Art and Limitations

As of this writing, most technologies used to physically implement qubits still face issues of stability, decoherence, error tolerance, and scalability. As a consequence, current NISQ devices therefore still need many physical qubits just for the purpose of error-correction. In

other words, present day logical qubits typically consist of many physical qubits in order to be able to perform useful computations. This can cause a gap between theory and practice of quantum algorithm design. For instance, theoretically valid algorithms which assume a large number of logical qubits to be available may not yet be practical.

Moreover, as of today, quantum computing is still bit level computing. That is, abstract data structures (such as linked lists or binary trees) or control structures (such as if-then-else-statements or for- or while-loops known from higher level programming languages or are not yet available to quantum programmers. Even common programming patterns such as variable assignments (x = y or x = x+1) are not possible on present day quantum computers because the *no-cloning theorem* and no-broadcast theorems state that it is impossible to create independent identical copies of arbitrary quantum states [88].

Although there are efforts towards identifying quantum programming patterns [89] and although (open source) software development kits for working with quantum computers such as IBM's QISKIT [90], Google's CIRQ [91], Rigetti's Forest [92] or PennyLane resulting from a collaboration of several smaller quantum computing companies [93] become increasingly available, it is important to note that these are tools for very low-level programming. For instance, the software development kits we just mentioned are Python libraries for the mathematical design of quantum circuits. They do allow for running quantum circuits on quantum computers or simulators but do not yet provide features equivalent to high-level data structures or control flow structures.

We will return to these issues and their impact on algorithm design for quantum machine learning in section 5.

3 Machine Learning in a Nutshell

Machine learning is the science of fitting parameterized mathematical models to data in order to realize intelligent systems that can make predictions or perform inference.

While the term *data* is often understood to mean known facts that can be recorded and have an implicit meaning, machine learning generally assumes a more technical point of view. In what follows, we, too, will assume this point of view and understand *data* to refer to collections of numeric values that can be obtained from measurements, surveys, interactions with technical devices, or comparable procedures and can be processed by computers. At first sight, this focus on numeric values may appear to restrict generality. Note, however, that everything that is held in the memory of a computer (a text, a piece of music, an image, ...) is encoded in terms of bit patterns and therefore in terms of numbers.

We will further assume that any individual data object can be encoded in terms of a tuple of m real numbers. Using a general notation, we will write such a tuple as $x \in \mathbb{R}^m$ and call it a data point; a collection $\mathcal{D} = \{x_1, \dots, x_n\}$ of n data points will be called a data set.

Data can be annotated by *metadata* which contain additional information. Common everyday examples are: a time stamp indicating when a picture was taken, a caption describing the content of a figure, the resource description framework (RDF) tags of a Website, or object identifiers such as the international standard book number (ISBN) of a book. Again resorting to an abstract notation, we will express metadata as $y \in \mathbb{R}^l$ and refer to an annotated data set $\mathcal{D} = \{(x_1, y_1), \dots, (x_n, y_n)\}$ as a labeled data set.

In machine learning, data sets are seen as a source of information about a specific scenario. The fundamental assumption is that data which were gathered in a certain context or with respect to a certain application domain are not arbitrary. On the contrary, machine learners always suppose that application specific data sets exhibit certain regularities which, in general thought, may be intricate or complex and therefore not immediately obvious to human analysts.

In other words, any machine learning solution for a real world problem implicitly posits that there exists some generally unknown process or mechanism which can explain variations in appearance within a given sample of example data. Mathematically, this translates to the assumption that there exists some function $f \in \mathcal{F}$ with parameters $\theta \in \Theta$ which models salient as well as latent characteristics of a set of data. For in stance, in the case of labeled data, this means

$$\mathbf{y}_{i} = f(\mathbf{x}_{i} \mid \boldsymbol{\theta}) + \boldsymbol{\varepsilon}_{i} \tag{28}$$

where the *noise term* ε_j accounts for possible inaccuracies which may, for instance, be due to corrupted measurements or faulty annotations. The family of functions \mathcal{F} to which the model f belongs to is called the hypothesis- or *model class* and Θ denotes the corresponding *parameter space*.

Other categorical terms which commonly occur in discussion about machine learning are supervised learning, unsupervised learning, and reinforcement learning. These refer to general machine learning philosophies which differ with respect to the richness of data available for training and testing.

Supervised learning deals with labeled data (x_j, y_j) and the goal is to train a model such that it can produce an appropriate output y for a given input x. A variant of this setting is semi-supervised learning which aims at assigning a label to unlabeled data using the knowledge contained in a small set of labeled data. If a semi-supervised learning system also incorporates its own earlier predictions into this assignment processes, it is sometimes referred to as a corrective learning system.

Unsupervised learning, on the other hand, works with unlabeled data x_j and aims at uncovering latent or inherent structures within a given data set. Examples of structures users might be interested in are clusters, higher order correlations, or relational dependencies.

Finally, reinforcement learning is a type of supervised learning that receives feedback in place of a label [94]. Based on such feedback, reinforcement learning methods tune models that are typically intended for decision making in feedback situations where a current output may impact the next input. Such a setting can often be formalized in terms of a partially observable Markov decision process (POMDP) [95]. Roughly speaking, learning in such settings happens in a (guided) trial and error procedure where the software agent is in some state, decides for an action which leads to a successor state, and only later finds out if a sequence of actions had a desired effect. This delayed feedback is then used to adjust the action selection mechanism via dynamic programming approaches and, over time, the agent learns which action to perform when in order to achieve certain (long term) goals.

In what follows, we give a more detailed account of the stages involved in training a machine learning system. We will largely focus on the case of unsupervised learning, yet, everything we discuss also applies to the other learning paradigms.

3.1 The Machine Learning Pipeline

Given the above premises, the process of developing a machine learning solution for a practical application can be broken down into several distinct phases. It begins with a data collection campaign in which representative examples for the intended application scenario are gathered and possibly annotated.

Here, it is of utmost importance to insist on representative data. This means the collected data points have to cover every foreseeable situation that may occur during the later deployment of the system. Indeed, a common and easily avoidable cause for insufficient reliability of machine learning solutions are biased training data which omit or neglect some of the inputs a system is expected to handle. In practice, this can have embarrassing consequences (for instance in incidents where software classified pictures of African Americans to depict goril-

las¹) to downright catastrophic ramifications (such as when a cars autopilot software failed to recognize a truck, crashed into it, and caused the death of its driver²).

The crucial point is that if a deployed system is confronted with a kind of input which has been overlooked during its development, its output can be unpredictable. Ideally, the collected data would be balanced and reflect all possible use case situations in about equal parts. This, however, may not always be possible. For example, in a predictive maintenance scenario where the task is to evaluate sensor data in order to predict whether a machine will continue to run smoothly or whether a failure is imminent, there may be much fewer examples of failure cases than of normal conditions.

Next, there may be a data pre-processing phase in which the collected data might be brought into a form more amenable for processing or where flawed data points might be filtered out. For instance, text data might be transformed into an appropriate numeric representation or measurement noise in sensor data might be smoothed.

After pre-processing, the collected data is split into two disjoint subsets, namely a set of training data and a set of test data and developers have to decide for a model class. This decision usually requires some domain expertise and will generally depend on the nature of the data as well as the intended use case. For instance, when dealing with a time series prediction scenario, one might opt for a polynomial function, a Markov chain, a Gaussian process, a decision tree, or a recurrent neural network. Each such general decision is followed by more specific decisions: Which degree is the polynomial supposed to have? How many states should the Markov chain contain? How to parameterize the kernel of the Gaussian process? How deep should the decision tree be? How many neurons should the neural network have and how should they be interconnected?

As these examples indicate, there often are numerous mathematical models that may apply to a given problem. Alas, clear cut criteria or simple suggestions for which model to use when are hard to come by. Rather, substantial research efforts are spent on understanding the usefulness of different models in different contexts and on developing new models for new contexts. However, a general recent trend is to avoid overly specialized models but to consider very flexible and thus widely applicable models such as deep neural networks instead.

In the training phase, the parameters of the chosen model are then adjusted such that the model matches the training data to the best extend possible. This adjustment happens automatically by means of running learning algorithms which are typically based on (statistical) optimization techniques. In order for this kind of mathematical learning to be possible at all, there has to be a criterion for how to measure how well the model and its current choice of pa-

¹see, for example, https://www.forbes.com/sites/mzhang/2015/07/01/google-photos-tags-twoafrican-americans-as-gorillas-through-facial-recognition-software/ Or https://www.nytimes. com/2021/09/03/technology/facebook-ai-race-primates.html

²see, for example, https://arstechnica.com/cars/2019/05/feds-autopilot-was-active-duringdeadly-march-tesla-crash/orhttps://www.forbes.com/sites/bradtempleton/2020/06/02/tesla-intaiwan-crashes-directly-into-overturned-truck-ignores-pedestrian-with-autopilot-on/

rameters matches the training data. In other words, there has to be an objective for the parameter optimization process and these objectives are often formalized in terms of an error- or *loss function*. A simple example of such a loss function often used when learning from labeled data is the mean squared error

$$E(\boldsymbol{\theta}) = \frac{1}{n} \sum_{j=1}^{n} \| \boldsymbol{y}_j - f(\boldsymbol{x}_j \mid \boldsymbol{\theta}) \|^2$$
(29)

and the corresponding training objective would be to solve

$$\boldsymbol{\theta}^* = \underset{\boldsymbol{\theta} \in \boldsymbol{\Theta}}{\operatorname{argmin}} \ E(\boldsymbol{\theta}) \tag{30}$$

Note, however, that potential loss functions are as numerous as potential model choices and that loss functions should, in fact, be chosen with respect to the data and model at hand. For example, when working with binary labels, one often considers the so called hinge loss or, when training probabilistic models, one might want to maximize a likelihood. Even other loss functions are based on divergence measures, entropy criteria, or problem- or model specific distances. In fact, the investigation of loss functions, their characteristics, and applicability is an important topic of ongoing machine learning research.

Learning algorithms are numerous, too, and the question of which algorithm to consider in the training phase should again be decided with respect to the given data, model, and loss function. Indeed, depending on the chosen model and loss function, it might be trivial to neigh impossible to optimally solve the learning problem in (30). For instance, if $f(x \mid \theta)$ is linear in its parameters θ and $E(\theta)$ is convex, there will be a closed form solution θ^* . Depending on the numbers of training data points and model parameters, it may still require considerable efforts to practically compute this solution, but one knows that it exists. On the other hand, if the model $f(x \mid \theta)$ is highly non-linear, the error landscape $E(\theta)$ will have numerous local minima and an algorithm that is guaranteed to find the globally optimal solution θ^* might not even exist. In situations like these, one often applies optimization techniques such as gradient descent

$$\boldsymbol{\theta}_{t+1} = \boldsymbol{\theta}_t - \eta_t \cdot \nabla E(\boldsymbol{\theta}_t) \tag{31}$$

or other iterative methods. Note, however, that $f(x \mid \theta)$ may be a complicated, composite function so that the computation of $\nabla E(\theta_t)$ itself may already require a whole chain of complex operations (such as in case of the *backpropagation* algorithm for training neural networks). Moreover, the choice of meta-parameters of a training algorithm such as the step size η_t in (31) is typically a non trivial matter and might require careful tuning. All in all, it can thus be very burdensome to train expressive models with many degrees of freedom or parameters. In fact, modern models such as deep neural networks with billions of parameters, typically necessitate the use of high performance computing hardware for training. Improved training algorithms, their convergence rates, and performance guarantees are therefore yet another topic of ongoing machine learning research.

In the *test phase*, the trained model is then validated and evaluated on the test data set. This requires an appropriate performance measure for which there are again abundantly

many (application dependent) choices. For instance, in the context of a classification task, one often evaluates the accuracy, i.e. the percentage of correct class predictions, of a trained classifier. For an information retrieval system, one may be interested in its recall and precision which measure the percentage of relevant retrieved instances and the percentage of relevant instances among the retrieved instances. But there also exist highly specialized performance indicators such as the perplexity of a language model or the BLEU score of a language translation system.

What is of pivotal importance in this phase is that training data and test data must be independent. In other words, data points considered during training must not be reconsidered during testing. This is crucial because what really needs to be assessed are the generalization capabilities of the trained model. That is, one needs to evaluate how well the function $f(x \mid \theta^*)$ performs on novel data, i.e. on data it has not seen during training. For reasons we will explain below, it is actually dangerous to evaluate the trained model on its training data and, for reasons we explained above, it is again pivotal that the testing data are representative.

Once a trained model has been thoroughly tested and found to perform reliably as well as accurately, it can finally be deployed in practice and be released into its application phase.

Use cases for this general methodology are manifold but most commonly found in settings where systems need to make data-based predictions, generate data-based suggestions or decisions, or classify novel observations or measurements. We also note that terms such as prediction, classification, or decision making can have a rather broad meaning. For instance, a prediction could be an estimate of tomorrow's stock market closing price or the statement that this English sentence translates into that German sentence. A classification could be as simple as "this picture shows a cat" or as sophisticated as "this picture shows a little boy on a sunny beach playing with a red ball". What kind of capabilities a learning system can achieve largely depends on the nature of the available training data and on the nature of the chosen model. With respect to the latter, notable strides have been made using modern neural network models and architectures such as variational auto-encoders, generative adversarial networks, or transformer networks [96, 97, 2].

Lazy vs. Eager Machine Learning

Yet another categorization of machine learning techniques contrasts lazy learners with eager learners. It is worthwhile to briefly mention this distinction as it points to different kinds of levers one may consider when trying to incorporate quantum computing techniques into the machine learning pipeline.

A lazy learner is a machine learning system that simply stores data and delays modeling until asked to make predictions. Such approaches are particularly suitable when dealing with large and frequently changing databases since only the affected part of the model needs to be retrained for new data. A simple yet common example of a lazy learner for classification problems is the k-nearest neighbor method which classifies novel observations based on their distances to known, i.e. previously learned, prototypes. Given an incoming data point, it searches a database of prototypes for the k nearest ones, performs a (weighted) majority voting over their class labels, and assigns the result to the new observation. Here, training is rather simple and can usually be accomplished quickly as the training problem basically consists in automatically identifying suitable prototypes. Application of such a classifier is usually simple and not too time consuming as well. However, depending on the size of the prototype database and the number k of nearest neighbors to consider, the required computations be burdensome and may necessitate the use of specific data- or index structures to guarantee fast runtimes. A considerable advantage of methods such as a k-nearest neighbor classifier is that they can easily be (re)trained during the application phase of the system because their prototypes may be extended or modified.

An eager learner, on the other hand, is an algorithm that trains a model in a dedicated offline training phase and considers a usually large but static training data set. This can lead to a very tedious or computationally intensive learning phase, but the application of the trained model to new examples is very efficient. Most of the well known machine learning models such as, say, neural networks are trained eagerly. Indeed, since we just mentioned prototype techniques as a prime example for lazy learning, it is important to note that not all prototype-based models are lazy. A prominent example of a class of eager prototype-based models are support vector machines.

3.3 Sources of Uncertainty in Machine Learning

To conclude this short overview of machine learning, we need to point out that the process of fitting a parameterized mathematical model to data is inherently statistical. This statement applies to any machine learning paradigm and should be understood as follows: From an abstract point of view, a fitted model provides a summary or compressed representation of the information contained in a training sample and there are several uncertainties involved in its training that may lead to different results in different training runs.

First of all, data collection is an informed but random process since different machine learning practitioners may collect different data samples when tackling the same problem. Second of all, modelling is an informed but random process since different practitioners may opt for different models when tackling the same problem. Third of all, model training may start with random initializations of the model parameters and thus settle in different local minima of the chosen loss function.

These sources of uncertainty can lead to undesirable outcomes. For instance, models can be biased and inappropriate for the data at hand. This is mainly the case when models are too simple or not flexible enough to capture relevant input output relations. In this case, training will lead to *underfitting*. Models can also be too flexible and thus overly sensitive to small fluctuations in their training data. Training such models on (slightly) different data samples

will likely produce considerably different results and they are said to show high variance. High variance models tend to learn minute irrelevant details and may therefore suffer from *overfitting*.

Phenomena like these explain the need for training and testing on independent data sets because only if training and test data differ can under- or overfitting be identified.

The so called bias-variance dilemma poses a considerable challenge in machine learning as one typically strives for models that capture patterns in the training data and also generalize well to unseen data. Alas, it may not be possible to achieve both these goals simultaneously since bias and variance tend to be reciprocal. As a rule of thumb, bias can be reduced by focusing on local information as it is done in nearest neighbor models or in radial basis function models. Variance, on the other hand, can be reduced through averaging over multiple data points or larger regions in data space as it is done in most other common, typically more complex models.

A common fallacy in this context is to assume that model complexity (measured in terms of the number of model parameters) causes variance and overfitting. However, this need not be the case; instead, overfitting in complex models is mainly due to too much freedom in the choice of their parameters. If this freedom is restricted, for instance by constraining the value a parameter can assume, overfitting can often be avoided. Imposing restrictions on model parameters is known as *regularization* and there exists a host of methods for how to accomplish this.

Other methods for mitigating variance include data dimensionality reduction, feature selection, or increasing the size of training data sets. At the same time, adding features or increasing data dimensionality can decrease bias. Moreover, many models and algorithms come with specific parameters which allow for trading off bias and variance. For instance, choosing a higher value of k in a k-nearest neighbor model will increase its bias and decrease its variance whereas a lower value of k will decrease bias and increase variance.

Yet another way of addressing the bias-variance dilemma consist in using *ensemble learning* techniques such as boosting or bagging. While boosting algorithms combine many models of individually high bias into an ensemble of low bias, bagging methods combine individual models of high variance into an ensemble with low variance. In this context, it is interesting to observe that recent empirical results indicate that modern neural networks with very wide layers (which can be seen as ensembles of individually weak learners) do not seem to suffer from the reciprocal bias-variance characteristics of more traditional models [98, 99].

4 Quantum Machine Learning

Modern machine learning is undeniably successful but also a very resource intensive endeavor. By now, the field has reached a point where practical computational efforts for training state-of-the-art models can only be dealt with on high performance computing hardware. Since this trend is likely to continue, it is no surprise that a growing number of machine learning researchers are beginning to look at the potential benefits of quantum computing.

By now, the scientific literature on quantum machine learning is vast and taking stock of the state of the field is warranted. In the following survey, we will consider quantum inspired classical algorithms for classical data, genuine quantum algorithms for classical data, classical algorithms for quantum data, and quantum algorithms for quantum data.

4.1 Quantum Inspired Machine Learning

Quantum inspired models constitute a broad class of frameworks, methods, and algorithms for classical data processing on classical computers that involve quantum mechanical concepts, in particular, the mathematics of quantum mechanics and quantum information processing. Alas, the understanding of what kind of methods and techniques are inspired by quantum mechanical insights varies widely. Some authors adhere to a narrow interpretation and only consider methods or algorithms which clearly would not exist without quantum mechanics. Others assume a rather broad point of view and consider, say, general optimization techniques such as simulated- or mean field annealing to be quantum inspired just because they occur in- or may have originated from the study of certain quantum mechanical systems.

An example of an early contribution that is quantum inspired in the narrow sense can be found in a book by van Rijsbergen in which he develops a quantum theory of information retrieval [100]. Modern information retrieval deals with the problem of searching data sets of unstructured media objects (texts, images, videos, etc.) for items or content related to a user query. A well known instance of this setting is Web search and Web search engines have become the best known examples of information retrieval systems.

The information retrieval problem is often formalized in terms of linear algebraic vector space models. Here, n data objects are encoded in terms of m dimensional vectors which are then gathered in an $m \times n$ matrix. A decomposition of this matrix into factor matrices of lower rank can reveal latent structures in the data and, if queries are encoded in terms of m dimensional vectors, too, query matching simply becomes the evaluation of inner products.

Based on these ideas and borrowing inspiration from quantum logic, a venerable quantum theoretic approach towards reasoning [101], van Rijsbergen argues that information retrieval should ideally be formulated in terms of superpositions in (infinitely dimensional) Hilbert spaces. There, eigenvectors of quantum density operators would represent basic (latent) concepts to be searched for, and the corresponding eigenvalues would measure overlaps

between concepts and queries. His ideas are convincing and compelling but also of limited practical value. Suitable quantum hardware to implement them on does not yet exist and digital emulations that would go beyond toy examples are impossible, because they would require exponentially large amounts of digital memory to store adequate Hilbert space representations of media objects.

Examples of a more generous interpretation of quantum inspired algorithms can be found in a recent contribution by Arrazola et al. [102] who, among others, are concerned with recommendation systems. Recommendation systems are a topic closely related to information retrieval, and play an important role in e-commerce. Here, too, popular models are based on linear algebraic formulations. If user-item preferences are represented in terms of a typically sparse $m \times n$ matrix, the problem of recommending items to users becomes a problem of predicting missing matrix entries. Just as in the case of vector space information retrieval, this prediction problem can be tackled using matrix decomposition methods.

While matrix factorization problems frequently occur in computational intelligence applications, their exact solutions are often hard to come by because the amount of floating point operations required for modern large scale matrices exceeds what is reasonably possible on digital computers. Therefore, there exists a vast amount of literature on randomized or probabilistic algorithms for fast, approximate numerical linear algebra [103, 104, 105, 106, 107, 108].

Indeed, Arrazola and his colleagues consider the FKV sampling algorithm [103], replace remaining exact procedures (for the computation of eigenvalues) by expected value estimations, and deem this a quantum inspired approach because similar estimators occur in quantum mechanics. However, probabilistic approaches to matrix factorization (with proper renormalization) have a venerable history, among others in machine learning [104, 105, 109], where they have hardly been thought of as inspired by quantum mechanics. In fact, if one were to subscribe to the broad interpretation of quantum inspired techniques, any statistical inference involving, say, Monte Carlo sampling, would count as such just because its mathematical apparatus was first developed by the quantum physics community.

To be frank, the use of the term *quantum inspired* often appears to be a marketing instrument for making certain approaches look more interesting than they turn out to be upon inspection. Consider, for instance, methods surveyed by Zhang [110] who benchmarks more than a hundred quantum inspired evolutionary optimization algorithms. Evolutionary algorithms take inspiration from biology because they work on populations of genotypes (i.e. encoded possible solutions to a problem) which express phenotypes (i.e. correspondingly decoded solutions). Problem solving or optimization happens in an iterative manner, where current auspicious genotypes are recombined or mutated into new ones whose phenotypes then represent new possible solutions. Since better solutions are used to replace worse ones, the population as a whole becomes better over time. This exploration of the solution space continues until the population contains enough members which express acceptable solutions. Traditionally, genotypes are often just bit string encodings of more complex objects,

and recombination and mutation happen in a purely random fashion. Obvious improvements over this baseline include more expressive encodings and more informed or probabilistically guided updates. Granted, some of the methods surveyed in [110] apply genuine quantum concepts (such as encodings in terms of sets of two-dimensional complex vectors or qubits) but many others simply rely on probabilistic update mechanisms whose connection to quantum mechanics has to be considered loose at best.

In what follows, we will therefore attempt to focus on quantum inspired techniques that are relevant to machine learning and, at the same time, are recognizably rooted in the mathematics of quantum mechanics and quantum computing.

4.1.1 Tang's Quantum Inspired Algorithm for Recommendation Systems

A broad class of techniques which are genuinely quantum inspired consists of methods that "dequantify" quantum algorithms. A widely reported example of such a method is due to Tang [111] who, as a graduate student, famously discovered a classical analogue of a quantum recommendation system algorithm introduced by Kerenidis and Praksh [112].

The quantum algorithm in [112] repeatedly samples from a low-rank matrix approximation by means of running quantum phase estimation [113, 114, 115] and quantum projections (partial measurements). This allows for sampling matrix elements proportional to their magnitude and thus for sampling matrix elements most relevant to the recommendation task without having to access every element of the matrix. Kerenidis and Praksh thought this quantum algorithm to be exponentially faster than classically possible, because a classical implementation of the procedure seemed to necessitate iterations over all matrix elements. Crucially, this supposed quantum advantage hinges on the assumption that incoming classical data can efficiently be encoded in terms of quantum states. However, while the QML literature often assumes such state preparations to be given and posits the existence of quantum random access memories (QRAMs) in which these states are available, Kerenidis and Praksh actually provided a protocol and data structure for this purpose.

Tang's fundamental insight was that this protocol also allows for encoding classical data in a representation that satisfies norm constraints required in randomized linear algebra. There, it is known that norm-based sampling of, say, matrix columns minimizes variance among all unbiased estimators of factor matrices. Moreover, robust estimates can be obtained from samples smaller than the size of the matrix to be factored [103, 108]. Tang therefore swapped quantum state preparation for preparation of a classical sampling procedure and thus obtained a classical algorithm only polynomially slower than the original quantum method. In other words, Tang's work showed that, in this particular matrix completion setting, potential quantum speedup is not as substantial as it appeared to be at first sight. This led her to conclude that claims as to speedups of QML algorithms over classical ML algorithms should take into consideration any state preparation assumptions in a QML model and match them against sampling assumptions in a corresponding classical ML model.

Tang's work therefore provides new directions for classical algorithm research and helps better understanding for which kind of problems one can or cannot expect exponential quantum speedup. Indeed, in the wake of her discovery, efforts in "dequantifying" quantum algorithms have noticeably increased and combinations of her encoding scheme with randomized numerical linear algebra are becoming ever more popular. Examples include new methods for estimating pseudoinverses [116] based on quantum algorithms for the singular value decomposition [117], sub-linear algorithms for solving linear systems of equations [118] based on the HHL algorithm [119], and even faster algorithms for recommendation systems and regression problems [120].

4.1.2 **Tensor Networks**

Tensor networks are mathematical models originally developed for the study of manybody quantum systems in condensed matter physics [121, 122]. More recently, they were found to be helpful tools in quantum information processing and connections to machine learning models have been established [123, 124].

Tensor networks represent quantum states based on local entanglement structures. This is particularly useful whenever one is dealing with states that have a tensor product structure. Consider a system of n qubits whose description would usually require $\mathcal{O}(2^n)$ complex coefficients or, equivalently, a complex-valued tensor Γ with n indices

$$|\psi\rangle = \sum_{i_1, i_2, \dots, i_n} \gamma_{i_1 i_2 \dots i_n} |i_1\rangle \otimes |i_2\rangle \otimes \dots \otimes |i_n\rangle$$
 (32)

However, similar to the matrix factorization techniques discussed above, such tensors can be modeled in terms of contractions of products of lower order tensors, for example

$$\gamma_{i_1 i_2 i_3 i_4} = \sum_{j,k,l} \alpha_{i_1 j_k} \, \beta_{k i_2 l} \, \delta_{l i_3 i_4} \tag{33}$$

This allows for graphically describing the tensor as a network of interconnected tensors and, interestingly, when modeling physical systems, such networks tend to have common basic building blocks.

Just as the matrix factorization techniques sketched above, low rank tensor decompositions have a venerable history in machine learning and data science [125]. Among others, they allow for flexible recommendation, fast video analysis, or robust document understanding [126, 127, 128]. Since data scientists increasingly deal with multidimensional data of high volume and higher order latent structures, tensor decompositions and tensor networks together with corresponding inference- or learning algorithms are attracting more and more interest. There exist kernelized or non-Euclidean extensions, various cost- or loss functions to guide the estimation of meaningful factors, and distributed computing approaches. Tensor networks thus allow for addressing large scale problems and are at the heart of generalized regression

and classification techniques, support tensor machines, higher order canonical correlation analysis, higher order partial least squares, and generalized eigenvalue decomposition. Last but not least, they also allow for the optimization of deep neural networks [123, 124].

Indeed, studying the capabilities of deep learning, Lin, Tegmark, and Rolnick recently argued that tensor network models may provide an explanation as to the good performance of deep neural networks [129]. Although it has been known for long that mathematical theorems guarantee that very wide neural networks are universal approximators, i.e. are able to learn arbitrary functions arbitrarily well, they observe that functions of practical interest can often be learned more "cheaply" by hierarchical networks with comparatively narrow layers. They argue that this is due to characteristics such as symmetry, local structure, compositions, and polynomial combinations of log-probabilities which frequently occur in physics and that deep networks capture such aspects in a manner similar to tensor network models in quantum information processing. Indeed, they prove several "no-flattening theorems" which establish that certain deep neural networks cannot be approximated by shallow ones without loss of accuracy and efficiency. Of particular interest in the context of quantum machine learning is the authors' result that a shallow neural network cannot multiply n variables using fewer than 2^n neurons in its single hidden layer.

In an important practical contribution, Stoudenmire and Schwab assumed a tensor network point of view on deep learning [130]. They showed how algorithms for the optimization of tensor networks can be adapted to supervised learning tasks. In particular, they work with matrix product states to parameterize deep networks for image classification. On the MNIST data, a standard benchmark data set in machine learning, they observed less than 1% test error and thus reached state of the art performance.

Results like this spawned further research into tensor network inspired deep learning. For instance, Glasser, Pancotti, and Ciracet recently explored the connection to probabilistic graphical models, a venerable class of machine learning models [131]. They considered generalized tensor networks where information from a tensor might be copied and reused in other parts of the network, showed how to integrate this idea into common deep learning architecture, and derived an algorithm to train such networks in a supervised manner. This proved to overcome the limitations of regular tensor networks in higher dimensions while not loosing computation efficiency. In experimental evaluations with image- and sound data, they found their method to improve on previously proposed tensor network algorithms. However, an observation most interesting in the context of quantum machine learning is that their approach can, in principle, be implemented on quantum computers and may therefore impact future research on quantum assisted machine learning.

Tensor networks therefore constitute interesting models in the intersection of the disciplines of quantum computing and machine learning and there is a quickly growing scientific community dedicated to the topic. For instance, several workshops on tensor networks for machine learning have already been held at the NeurIPS conference, a primary venue of machine learning research. However, so far, corresponding methods have not yet widely caught

on and not yet found their way into mainstream machine learning technology. What is noticeable, though, is that researchers from Google, the Perimeter Institute, and the company X Development LLC recently released TensorNetwork, an open source library for implementing tensor network algorithms [132]. This library is supposed to support physicists and machine learners alike and provides functionalities for efficient high volume sparse data handling and tensor factorization. Whether or not initiatives like this will boost the use of tensor networks in theory and practice just as other libraries did in the case of deep learning technology remains to be seen.

4.1.3 Digital Annealing

Adiabatic quantum computers such as produced by D-Wave realize an energy minimization process called quantum annealing and they are specifically tailored towards solving quadratic unconstrained binary optimization problems (QUBOs) of the general form in equation (18) in section 2.1. Specific instances of QUBOs occur in the context of verification-, planning-, or assignment problems in areas such as, say, logistics or finance. While QUBOs therefore are of considerable practical importance, they are also difficult to solve in general because they pose combinatorial optimization problems which are NP-hard in general. It is therefore expected that (adiabatic) quantum computing will have economic impact as it provides novel approaches towards industrially relevant problems.

Alas, the technical effort required for running present day adiabatic quantum computers is substantial. While they have the potential to deliver unprecedented computing power, they must be maintained at temperatures near absolute zero and be protected against magnetic interference, thermal variation, and mechanical vibration in order for their physical realizations of logical qubits to remain in superposition. Just the amount of energy required to maintain a reliable cryogenic environment for a current D-Wave 2000Q machine is estimated to exceed 25 kWh [133]. Present day adiabatic quantum computing is thus an expensive endeavor and may not yet amortize its costs in industrial applications.

However, in machine learning, QUBOs as in (18) are known as Hopfield energy minimization problems and play a central role in the theory of Hopfield networks [71]. Hopfield networks are a special kind of neural networks inspired by physical representations of spin glass phenomena [70] and they are of considerable theoretical interest because they provide simple models for higher cognitive processes such as memory retrieval. Hopfield networks have a venerable history and are established textbook material [22]. Among others, there exist various classical algorithms for Hopfield energy minimization ranging from random updates, over greedy gradient descent methods to simulated- and mean field annealing.

Against this backdrop, Fujitsu has developed special purpose digital hardware for solving QUBOs. They refer to their technology as a *digital annealer* and market it as a solution that rivals the utility of quantum computers [134].

Fujitsu's system is based on conventional complementary metal-oxide-semiconductor (CMOS) technology and von Neumann architectures; their dedicated chip fits onto a single circuit board, works at room temperatures, and does not require a complex support infrastructure. The optimization algorithm implemented on this hardware is based on simulated annealing which is extended in several directions. For instance, the method employs a parallel trial Monte Carlo procedure which considers several switches of decision variables (simulated qubits) in parallel. This is supposed to overcome problems with respect to low acceptance probabilities in common annealing algorithms. The method also involves a so-called dynamic offset mechanism which raises acceptance probabilities after those iterations in which the Monte Carlo scheme did not make any progress. This is empirically observed to help the algorithm overcome narrow barriers in the energy landscape of the problem to be solved, and thus to make faster progress towards solutions. In its current version, the system can solve QUBOs with up to 1024 variables [135].

A similar digital annealing technology has been independently developed by a team of researchers at TU Dortmund and Fraunhofer IAIS [133, 136]. Their system is implemented on very affordable field programmable gate arrays (FPGAs). It can currently solve QUBOs of up to 2048 variables and thus exceeds the number of variables of the D-Wave 2000Q adiabatic quantum computer at a mere 0.006% of its power consumption. Moreover, while D-Wave machines can currently only solve fully connected problems with up to 119 variables, the FPGA-based solution supports dense parameter matrices for all 2048 variables.

The system has been shown to successfully solve machine learning problems such as k-means clustering, maximum a-posterior prediction, or binary support vector machine training. The optimization algorithm running on this hardware takes parallelism into account and is based on a customizable $(\mu + \lambda)$ evolutionary algorithm which allows for tuning the maximal problem dimension n, the number of parent solutions μ , the number of offspring solutions λ , and the number of bits per coefficient Q_{ij} . When working with low-budget FPGAs, this makes it possible to allocate more FPGA resources either for parallel computation (parameters μ and λ) or for the problem size (parameter n and bit depth of Q_{ij}).

4.1.4 Quantum Inspired Data Clustering

The term quantum clustering refers to a class of quantum mechanically inspired density-based clustering algorithms which assume that clusters are defined in terms of regions of more densely distributed data points. Specifically, quantum clustering algorithms model a given set of data in terms of a Gaussian mixture model consisting of one mixture component per data point. This Gaussian mixture is then considered as a quantum mechanical wave function for the data set and a quantum potential is constructed such that the data wave function becomes a stable solution to the time-independent Schrödinger equation. Gradient descent on this potential causes data points to move towards nearby local minima and data points that end up close to one another are assumed to belong to the same cluster [137].

More elaborate versions of this approach replace gradient descent by quantum evolution which can be understood as a kind of non-local gradient descent that, in turn, is capable of tunneling through potential barriers. While this has considerably higher computational costs, it allows for interesting data visualizations and the identification of substructures turning the method into a hierarchical clustering approach [138].

4.1.5 Quantum Inspired Gravitational Search

Gravitational search algorithms (GSAs) are among the newest in the class of swarm optimization algorithms which rely on the metaphor of gravitational interaction between data objects [139]. GSAs have the advantage that they are easy to implement and capable of escaping from local optima of an objective function. They have by now become established tools for effective and efficient global optimization in solving various kinds of continuous problems. In particular, the binary version, BGSA, applies to solving binary encoded problems and the discrete version, DGSA, allows for solving combinatorial problems [139].

Nezamabadi-pour [139] introduced a population based meta-heuristic search algorithm that is a binary quantum inspired gravitational search algorithm (BQIGSA) combining both gravitational search and quantum computing. His underlying idea is to solve binary encoded problems by a quantum bit superposition together with a modified rotation Q-gates strategy. Nezamabadi-pour evaluated the algorithm's effectiveness by performing experiments on combinatorial 0–1 knapsack problems or Max-ones functions. He found that BQIGSA can compete with classical BGSA, conventional genetic algorithms, binary particle swarm optimization including a modified version, a binary differential evolution, a quantum inspired particle swarm optimization, and three well known quantum inspired evolutionary algorithms.

Lou et al. [140] provide a concrete use case for quantum inspired binary gravitational search algorithms, namely the problem of predicting failure times of cloud services. They argue the importance of this case by pointing out that data centers coordinate several hundred thousand heterogeneous tasks to provide the services' high reliability. The authors motivate their research by the great challenge to acquire the optimal parameters for a relevance vector machine approach towards solving nonlinear predicting problems. In practical evaluations, they observed similar to better predicting performance of their IQBGSA-RVM algorithm compared to the baselines of chaotic genetic algorithms, binary gravitational search algorithms, binary particle swarm optimization, quantum inspired binary particle swarm optimization and standard BQIGSA (which all employ relevance vector machines).

4.2 Machine Learning for Quantum Computing

The notion of machine learning *for* quantum computing commonly refers either to the use of classical methods for preparing inputs for quantum processing units and processing

outputs obtained from quantum hardware, or to the use of classical methods for designing quantum circuits.

4.2.1 Quantum Circuit Design

Quantum circuit design for quantum gate computing requires considerable experience, a deep understanding of the mathematics of quantum information processing, as well as a good deal of creativity. Because quantum circuit design can thus be considered a difficult task [89], the idea of using machine learning or optimization algorithms for this purpose arose early on and can be traced back to the 1990s.

For instance, Rubinstein [141] is concerned with evolutionary algorithms for quantum circuit design. He discusses possible encoding schemes and fitness functions and presents evolved circuits that allow for the production of entangled states. Leier [142] also works with genetic algorithms for quantum circuit design, yet focuses on how to address exponentially large search spaces for operators on exponentially large quantum states. His almost twenty year old findings suggest that quantum circuits can be evolved to a certain extent or that human experts can manually infer suitable quantum circuits from evolved solutions for small problem instances. Notably, he observed that the use of evolutionary crossover operators seemed to deteriorate the quality of the solutions found.

Now that working prototypes of quantum computers have become available, design procedures similar to the above can be tested in practice. For instance, Franken et al. [143] are concerned with variational quantum eigensolvers (VQEs). These are hybrid algorithms that combine classical and quantum computing steps in order to determine the eigenvalues of large matrices. Solutions to this general problem are sought after in many areas of science and engineering. For instance, in quantum simulations, the matrix in question often is the Hamiltonian of a quantum system and its lowest eigenvalue is of interest as it characterizes the ground state of the simulated system.

Using the VQE approach, a quantum subroutine is run inside of a classical iterative optimizer. This quantum subroutine prepares a state based on a set of given parameters and performs a series of measurements in the appropriate basis. Measurement results are read into a classical memory and are then used to classically estimate expected values of the parameters for the next iteration. Franken et al. observe that the efforts for estimating gradients of the kind of cost function that occur in this process are considerable. As a remedy, they therefore work with a weight-agnostic evolutionary scheme. They test their approach on real quantum hardware in the IBM quantum experience environment and use the automatically determined circuits to solve benchmark problems such as the transverse field Ising Hamiltonian and the Sherrington-Kirkpatrick spin model [144].

Indeed, the idea of parameterized quantum circuits has lately attracted increasing attention. It allows for the use of reinforcement learning [145, 146, 147] and other learning algo-

rithms [148, 149] for quantum circuit design. Similarly, machine learning methods are increasingly considered as tools for solving quantum circuit mapping problems, i.e. problems of mapping a given general quantum circuit design onto a specific NISQ architecture [150, 151, 152, 153].

4.2.2 Quantum State Preparation

Machine learning approaches are more and more considered for other applications beside automatic circuit design. For example, they are used as a tool for preparing the initial state and for modelling the noise characteristics of a quantum computing system.

As will be detail below, the ability to load classical data efficiently into quantum states is the basis for the realization of many quantum algorithms. However, the best known general methods require $\mathcal{O}(2^n)$ gates to load an exact representation of a generic data structure into an n-qubit state. Therefore, scaling issues need to be taken into account because scaling can easily predominate the complexity of a quantum algorithm and thereby impair any potential quantum advantage.

Grover and Rudolph demonstrate how log-concave and other efficiently integrable probability distributions can be approximately encoded into an m-qubit register [154]. Even though the method in Grover and Rudolphs seminal note is presented for the univariate case, they explain that an extension to multiple dimensions, as it is usually the case in machine learning, is possible. Nevertheless, the described approach does require that the underlying distribution can be efficiently integrated by a classical algorithm. This precondition is frequently violated for most high-dimensional distributions.

Zoufal et al. [155] present a hybrid quantum-classical algorithm for efficient, approximate quantum state loading as an alternative to the Grover and Rudolph Method. For this purpose, they describe quantum generative adversarial networks (qGANs) to learn a generative model that prepares the desired probability distribution. For this a probability distribution is given implicitly via data samples. The qGAN is composed of two components, the generator, which is a parametrizable quantum circuit and the discriminator, a classical neural network. Measuring the result of the quantum circuit corresponds to a sample from the distribution that is induced by the circuit. The parameters of the classical network are tuned to distinguish these generated samples from the true data samples. Based on this classification, the parameters of the quantum circuit are tuned to output to improve the quality of the learned distribution. During training, the qGAN learns a low-dimensional representation of the probability distribution underlying the data samples and loads it into a quantum state. The loading requires $\mathcal{O}(\text{poly }n)$ gates and can thus enable the use of quantum algorithms, such as quantum amplitude estimation, which require to load a specific distribution. The idea of qGAN distribution learning and its loading method have been shown to work in simulations as well as in actual implementations on superconducting quantum processors.

4.2.3 Quantum Noise Modelling

Harper et al. [156] focuses on the output distribution that is observed through measurements and the noise that affects the accuracy of those measurements, instead of manipulating the input distribution.

When building large-scale quantum computers, Noise arising from various sources is the main obstacle. Quantum systems with sufficiently uncorrelated and weak noise are required for solving large real-world problem instances. Even though there has been substantial progress in designing hardware specific error correcting codes, as well as improved measurement and qubit hardware, continued progress depends on the ability to characterize quantum noise faithfully and efficiently with high precision[157].

In [156], the parameters of a classical probabilistic graphical model are tuned such that the model reproduces the noise distribution of a superconducting quantum processor. The model puts out an estimate of the effective noise and can detect correlations within arbitrary sets of qubits. It can also be applied to understand how the noise between pairs of qubits correlates. Visualization can be generated to discover long-range dependencies between the noise of specific qubits. In fact, the method revealed previously unknown error correlations within the device that was used for the experimental evaluation. The method is the first implementation of a provably rigorous and comprehensive diagnostic protocol capable of being run on real-world quantum processors, according to the authors. It builds the foundation for calibration in the presence of cross-talk, bespoke quantum error-correcting codes, and customized fault-tolerance protocols that can greatly reduce the overhead in a quantum computation.

4.3 Quantum Enhanced Machine Learning

From the point of view of machine learning practitioners, the idea of quantum enhanced machine learning is arguably among the most exciting aspects in the broad field of quantum machine learning. It deals with the use of quantum computing algorithms for solving computationally demanding learning tasks, i.e. with the processing of classical data on quantum devices to realize intelligent systems.

The expectation is that quantum speedup will make it possible to considerably accelerate learning processes or even tackle problems which are still beyond reach even for current super computers. Indeed many common learning tasks involve linear algebra routines on very large systems of equations or optimization or search problems for which it seems likely that quantum advantages can be realized.

Since worldwide research on quantum enhanced machine learning has noticeably intensified over the past couple of years, the literature has already become vast, and numerous quantum methods and algorithms have recently been reported (see, for instance, the introductory papers by Dunjiko et al. or Biamonte et al. [46, 47] and the references therein). The fol-

lowing survey will thus begin with a broad overview and then review quantum algorithms for several specific machine learning problems where quantum solutions appear to be auspicious.

Aïmeur et al. [158] were among the first to address quantum information processing (QIP) for machine learning. They review QIP concepts and investigate novel learning tasks that run within environments where information is fundamentally quantum mechanical. They illustrate their idea by using the case of quantum data set clustering and providing examples of possible quantum clustering algorithms.

Riste et al. [159] specifically treat problems for which there exists a proven quantum advantage. That is, they consider rather didactic learning problems whose solution is expensive on classical computers but can be efficiently obtained on quantum computers. According to the authors, most such problems involve the repeated use of an oracle.

At this point it seems warranted to clarify the notion of oracles because many quantum computing algorithms posit their existence. We therefore note that, while it is often difficult to determine a solution to a given problem, it is often also simple to verify if an alleged solution really solves the problem at hand. Consider for instance the problem of prime factorization. While it requires some effort to determine that 1, 2, 3 and 7 are prime factors of 42, it is comparatively easy to verify that they are. By the same token, it is also easy to verify that, say, 1,5 and 13 are not the prime factors of 42. A function that easily accomplishes such a verification for a given problem is called an oracle. Just for completeness we also note that for many if not most computational intelligence problems oracles do not exist. For instance, in chess it is generally difficult (if not impossible) to determine the single best move for a given game state. Contrary to prime factorization, it will generally be also at least as difficult to verify if an alleged best move is indeed optimal.

Riste et al. observe that the costs incurred by an oracle-base quantum algorithm can be determined by its querying complexity, i.e. by the number of oracle calls needed to find a solution with a given probability. While showcases of oracle-based quantum algorithms have already been verified experimentally on various platforms with different physical realizations of logical qubits, the authors note that these showcases typically involved problems that could be modeled in terms of very few qubits. Based on toy scenarios like these, however, it is hard to argue for quantum supremacy because classical algorithms, too, could solve problems like these with a few queries to an oracle. The authors therefore investigated the performance of classical and quantum approaches in solving a more demanding oracle-based problem known as learning parity with noise. Using a custom made five-qubit superconducting processor (consisting of superconducting qubits) as well as classical Bayesian computing, they observed a query complexity in favor of quantum processing which substantially increases depending on the acceptable error rate and the problem size. Their important major finding is thus that a significant quantum advantage can already arise even in present day noise-intensive systems.

While general results such as these are encouraging, one has to keep in mind that current NISQ devices can only realize few logical qubits and still suffer from limited coherence times. That is, even if the findings in [159] suggest that limited fault tolerance may not be the most

pressing technical challenge in quantum machine learning, certain quantum learning algorithms may still not be practical. O'Quinn and Mao [160] emphasize this and point out that existing platforms such as IBM Q, IonQ, Rigetti Computing, D Wave, Microsoft Quantum and Google Quantum do not yet allow for practical implementations of some of the more fanciful ideas discussed in the quantum machine learning literature.

4.3.1 Quantum Algorithms for Linear Algebra

One of the reasons for why machine learning models often assume input data to be encoded in terms of data points $x \in \mathbb{R}^m$ is rather prosaic and due to technological circumstances: as long as the dimension m of a given vector is not too large, conventional digital computers can easily perform linear algebraic operations. Indeed, there exist venerable software libraries such as BLAS and LAPACK which date back to the 1970s, compile on almost any operating system, and allow for highly efficient linear algebra computations. Moreover, modern graphics processing units (GPUs) allow for rapid, parallel memory access and are specifically designed to facilitate matrix vector operations. In a sense, numerical computing technology and machine learning co-evolved and it is therefore not surprising that many machine learning algorithms involve matrix vector products, matrix inversion, or the decomposition of matrices into factors of lower rank.

Against this backdrop it seems reasonable to attempt to harness potential quantum advantages for machine learning sub-routines since quantum computing is all about matrices of size $2^n \times 2^n$ acting on 2^n dimensional vectors.

Harrow, Hassidim, and Lloyd [119] developed a quantum algorithm for solving systems of linear equation which can be written in terms of sparse and well conditioned matrices. This algorithm is considered to be one of the more fundamental, presently known methods with provable quantum speedup over their classical counterparts. It thus falls within the same category as Shor's factoring algorithm (based on Coppersmith's quantum Fourier transformation [161]), Grover's search algorithm, or quantum random walk algorithms [162, 163] and is by now simply known as the HHL algorithm.

Harrow, Hassidim, and Lloyd are specifically concerned with solving Ax = b for x where the $n \times n$ matrix A is Hermitian and b is an n-dimensional unit vector. They then suppose an amplitude encoding $|b\rangle$ of b (for instance using the method in [164]) and a Hamiltonian operator $\exp(iAt)$ and perform quantum phase estimation [113, 114], ancilla bit rotation, inverse quantum phase estimation, and measurement to obtain an estimate of the eigenbasis of the coefficient matrix in which it is easy to invert it.

If the coefficient matrix of the considered linear system of n variables is sparse (not densely populated) and has a low condition number k and—crucially—if users are mainly interested in the result of a measurement on the solution vector rather than in the solution vector itself, then HHL has a runtime of $\mathcal{O}(\log n \cdot k^2)$. This constitutes an exponential

speedup over the fastest classical algorithms which generally run in $\mathcal{O}(n \cdot k)$. Importantly, improvements over the originally proposed approach are possible if phase estimation is replaced by the "linear-combination of unitary method" which approximates matrix inverses via Fourier- or Chebyshev-series expansions [165]. Moreover, since the HHL algorithms has been verified experimentally, it is known to work on existing quantum computers, at least for small problems which meet its prerequisites [166, 167, 168].

The singular value decomposition, also dubbed a "singularly valuable decomposition" [169], is one of the most important matrix factorization techniques with numerous practical applications in science and engineering. Assuming that a universal quantum computer with access to a QRAM was available, Rebentrost et al. [170] present a quantum singular value decomposition which is exponentially faster than its classical counterparts. Making less wide reaching assumptions, Gilyen et al. [117] propose a quantum singular value "transformation" algorithm that is based on the idea of qubitization [171] and computes a bounded polynomial approximation of the singular value decomposition. Note that the term qubitization refers to a technique for simulating the time evolution operator $\exp(iAt)$ that features prominently in the HHL algorithm. For this, Low et al. [171] prove that it can be done with low error using comparatively simple quantum random walk circuits.

Gilyen et al. observe that their model leads to rather simple quantum circuits which tend to involve only constantly many ancilla qubits. From a general application point of view, their method allows for estimating Moore-Penrose pseudo-inverses and thus applies to regression problems. Moreover, the authors also observe that their quantum singular value transformation leads to a unified framework for several quantum algorithms. Methods that can be seen as special cases of their approach include fixed-point amplitude amplification, robust oblivious amplitude amplification, and certain quantum walks. Yet, overall the work in [117] is of mostly theoretical nature; attempts of a practical implementation are not reported even though the authors provide a detailed discussion of how to implement quantum linear algebra methods by representing matrices in terms of unitary circuits and vectors in terms of quantum states.

Clader, Jacobs and Sprouse [172] generalize the HHL algorithm. They propose a state preparation routine which is capable of initializing generic states and they integrate a quantum compatible preconditioner, which enlarges the problem space for which solutions can be obtained with an exponential speedup over classical linear systems solvers. The authors verified their algorithm's applicability by letting it compute the electromagnetic scattering cross section of an arbitrary target. As a result, it was found to be exponentially faster than the best classical algorithm.

Huang, Bharti and Rebentrost [173] deal with near-term quantum algorithms for linear equation systems of the form Ax = b and examine the use of variational algorithms. As part of their research, the authors develop a near-term algorithms, which is founded on the classical combination of quantum states (CQS) method. They conducted experiments of solving large linear systems by simulating the quantum algorithm on a classical computer and report the approach to work well.

Subasi, Somma and Orsucci [174] introduce two quantum algorithms based on evolution randomization, a simple variant of adiabatic quantum computing, to prepare a quantum state $|x\rangle$, which is proportional to the solution of the linear equation system Ax = b. Both algorithms are easy to implement and designed using families of Hamiltonians that are linear combinations of products of A. A quantum oracle is supposed to be given that, for any row of A, outputs the nonzero matrix elements and their indices. Given this prerequisite, their algorithms exhibits an exponential quantum speedup.

To solve a system of linear equations, Lee, Joo and et al. [175] introduce a hybrid quantum algorithm built on the HHL algorithm. It reduces the circuit depth from the original algorithm without accuracy loss of the results. On the contrary, the authors' experiments (using 4 qubits by IBM Quantum Experience) produced better results, that is a higher accurate performance on specific systems of linear equations.

Bravo-Prieto et al. [176] take as a starting point the limitation that existing quantum solvers of linear equation systems are hardly implementable due to the required circuit depth. The authors introduce a variational quantum linear solver (VQLS), which is a hybrid algorithm designed for near-term quantum computers. This solver seeks to variationally prepare $|x\rangle$ such that $A|x\rangle \propto |b\rangle$ and Bravo-Prieto et al. then derive a termination condition guarantying the achievement to a desired precision. The authors verified their algorithm using Rigetti's quantum computer.

Xu et al. [177] take up the challenge of a quantum algorithm's high demands on the circuits depth (leading to a high demand on the quantum device's universal fault-tolerance) to solve linear algebra tasks. The authors developed variational algorithms, which are compatible with noisy intermediate-scale quantum devices. They demonstrate that the linear equation system's solutions and matrix-vector multiplications are translatable as the ground states of the constructed Hamiltonians. They implemented their algorithm using the IBM quantum cloud services and observed a high solution fidelity of 99.95%.

4.3.2 Quantum Algorithms for Regression

The term regression analysis refers to a broad spectrum of statistical methods for estimating relationships between a dependent variable and one or more independent variables. While there exist many different algorithms for fitting a regression model to data, the arguably most well known approach is least squares regression. At its heart, this is a simple linear technique which nevertheless allows for dealing with non-linear models. The basic setting is as follows: Given a data matrix $\boldsymbol{X} \in \mathbb{R}^{m \times n}$ whose columns represent (possibly non-linearly transformed) data points and a target vector $\boldsymbol{y} \in \mathbb{R}^n$, the task is to identify a weight vector $\boldsymbol{w} \in \mathbb{R}^n$ such that the loss $\|\boldsymbol{X}^{\mathsf{T}}\boldsymbol{w} - \boldsymbol{y}\|^2$ is minimal. This is can be understood as an almost trivial supervised learning problem since its unique closed form solution is known to be given by $\boldsymbol{w} = (\boldsymbol{X}\boldsymbol{X}^{\mathsf{T}})^{-1}\boldsymbol{X}\boldsymbol{y}$. Moreover, for the practical computation of this expression there exist fast and numerically stable gradient descent schemes. Nevertheless, because least squares

regression is such a fundamental technique in data science and machine learning, it has attracted the attention of quantum computing researchers.

Initially, work in this direction was mainly theoretical. For instance, assuming the existence of a universal quantum computer which when queried about an element in a given row of a sparse matrix X, produces a quantum state that encodes the column number and, moreover, is capable of producing copies of an input state $|y\rangle$, Wiebe, Braun, and Lloyd [178] present an HHL-based algorithm that estimates $|w\rangle$ up to scale.

Schuld, Sinayskiy, and Petruccione [179] address certain shortcomings of this approach. In particular, they extend it towards dense matrices and improve on the dependency on the condition number. Their idea is to create amplitude encoded quantum states that represent the singular value decomposition of the data matrix and then to run phase estimation to invert the unknown singular values. They acknowledge that the problem of preparing quantum states which encode classical data is non-trivial and assume that states are either given in a quantum random access memory (QRAM) from which the algorithm could access them or that there is an oracle that loads data into an entangled register.

Some of the computational steps in [179] apply ideas which were first brought forth by Wang [180] who presents a quantum algorithm for fitting a linear regression model and focuses on estimating the quality of the fit and on assessing whether the given data set is suitable for quantum regression at all. This allows his algorithm to output optimal parameters in classical form. However, Wang, too assumes availability of a QRAM which might not be technically feasible in the foreseeable future.

Yu et al [181] introduce an improved quantum algorithm for ridge regression, a regularized version of ordinary least squares regression. Ridge regression allows for more robust estimates but comes with the need of having to estimate additional hyper-parameters. The authors develop a quantum algorithm for ridge regression that utilizes the parallel Hamiltonian simulation technique for simulating Hermitian matrices in parallel and for implementing a quantum k-fold cross-validation that estimates the ridge regression's predictive performance. Yu et al. compose their algorithm in two stages: Stage 1 searches for a suitable hyperparameter α that optimizes the predictive performance of the regression model by running quantum cross-validation. Stage 2 creates a quantum state which encodes the optimal fitting parameters of ridge regression with such an α , that is later used to predict new data. Again, a major (practical) weakness of the proposed method is that the authors quantum encoded data to be available in a QRAM.

Hou et al. [182] make such an assumption, too. The authors propose a quantum partial least squares regression algorithm. Partial least squares regression is yet another variant of the ordinary least squares approach which allows for addressing multiple correlation problems. The authors develop a method for quantum eigenvector search in order to accelerate the regression parameter selection and propose a density matrix product to avoid multiple QRAM queries for constructing residual matrices. If a QRAM were available, this algorithm would run exponentially faster than classical partial least squares estimation techniques.

Liu and Zhang [183] further improve on techniques such as the above. To this end, their algorithm takes statistic leverage scores of a data matrix into account. They show how to generate a quantum state proportional to $|w\rangle$ in $\mathcal{O}(\log n)$ time for sparse and well-conditioned \boldsymbol{X} . Again, they assume efficient quantum access to classical data was possible.

Gilyen, Song, and Tang [184] rightfully point out that many QML algorithms for low-rank regression require input to be stored in a QRAM especially those that theoretically achieve runtimes that are poly-logarithmic in the dimension of the data under consideration. As a remedy, they describe a classical algorithm for linear regression that borrows from the HHL algorithm and improves on the method due to Chia et al. [116]. Their stochastic gradient algorithm exploits recent numerical sparsity techniques for fast eigenvector computation and regression analysis and though not designed as a quantum algorithm point towards possible genuine quantum solutions.

Finally, Date and Potok [185] recently showed how to perform linear regression on adiabatic quantum computers. In order to accomplish this, they cast the regression problem as a quadratic unconstrained binary optimization problem which they then solve on a D-Wave 2000Q adiabatic quantum computer. While this approach does not suffer from any (as of yet technically unrealistic) assumptions as to data encoding, it incurs a certain loss in numerical precision. However, the authors compare their results obtained on a D-Wave machine to those resulting from classical numerical solution implemented in Python running on desktop computer equipped with a current multi-core Intel processor. Their experiments reveal that the adiabatic quantum implementation achieves up to 2.8 fold speedup over the digital implementation but is on par with respect to the regression error metric.

4.3.3 Quantum Algorithms for Clustering

Prototype-based clustering is an unsupervised machine learning problem where minimization of a loss function allows for partitioning a set of n data points into $k \ll n$ clusters which are defined in terms of representative elements. The arguably most well known instance of a prototype-based clustering method is k-means clustering where the n data points $\boldsymbol{x}_j \in \mathbb{R}^m$ are Euclidean vectors and the k prototypes $\boldsymbol{\mu}_i \in \mathbb{R}^m$ are the centroids of their respective clusters. These are typically found through an iterative minimization of an variance-based loss; alas, while this loss is commonly considered to be intuitive, its minimization actually poses an NP-hard problem even for the simple case where k=2 [186]. Popular classical k-means clustering algorithms such as those due to Lloyd, MacQueen, or Hartigan are therefore but heuristic approaches to the problem and there is no guarantee that they will determine the optimal solution.

However, for k=2, k-means clustering becomes a bipartition problem that can be cast as a QUBO and thus be solved via adiabatic quantum computing [73]. The corresponding Hamiltonians are derived from a reformulation of the conventional k-means objective which is based on Fisher's analysis of variance and leads to an equivalent criterion that also allows for

kernel k-means clustering via adiabatic quantum computing [187]. In both cases, the adiabatic quantum optimization procedure performs an exhaustive search over the space of all possible solutions and does so quadratically faster than classically possible. While the validity of the ideas in [73, 187] was originally verified in simulation experiments, the approaches have meanwhile been implemented on D-Wave computers and were found to be practically feasible [188, 189, 190].

If the QUBO in [73] is modified appropriately, the approach also allows for the estimation of k-medoids, i.e. cluster prototypes that approximate cluster means [191] and the resulting algorithm has been successfully applied to solve cardinality-constrained index-tracking problems in the financial industry [192].

The work in [73] also shows how to extend adiabatic quantum bipartition clustering towards non-numeric data for which one can define similarity- or distance measures. Given such measures, relations among non-numeric data points (such as text strings) can be modeled in a graph and the bipartition clustering problem becomes a minimum graph cut problem. Hamiltonians for this general problem class are algebraically identical to those for k=2-means clustering and the approach is practically viable.

A remaining practical challenge with these kind of quantum algorithms for bipartition clustering is that they require as many logical qubits as there are data points to be clustered. While this is no different on digital computers, quantum computers that could manipulate as many logical qubits as modern, large scale data sets would require, do not yet exist.

Aïmeur, Brassard, and Gambs [193, 194] discuss how quantum gate computing can speed up unsupervised learning algorithms used in data clustering. Specifically, they propose quantum versions for minimum spanning tree clustering, divisive clustering, and k-medians clustering. The basic idea is to consider an oracle which provides knowledge as to distances between data points and to use this oracle in algorithms based on Grover's phase amplification.

Tomesh et al. [195] observe that quantum machine learning often assumes classical data to be encoded in term of quantum states in superposition which, due to the difficulty of known encoding schemes, can annihilate potential quantum speedup. They therefore investigate a coreset-based approach towards prototype-based clustering with reduced loading overhead. Coresets are an important concept in machine learning as they can allow for approximating large sets of data points in terms of small representative sets that can quickly be determined (often in linear time). Given a coreset for a problem under consideration, Tomesh et al. apply the quantum approximate optimization algorithm (QAOA) [84] and run numerical simulations to compare the their approach against classical k-means clustering. Their results indicate that there exist data sets where QAOA might outperform standard k-means clustering on a coreset. However, the authors point out that, for their method to show a quantum advantage over conventional k-means algorithms, problems have to be rather special in that they have to have appropriate coresets. This is, because their method crucially hinges on the existing of a characteristic coreset which may not exit in general.

4.3.4 Quantum Algorithms for Nearest Neighbor Search

In what follows, the term nearest neighbor search is used to refer to the following general setting: Given an unstructured set $\{x_i\}_{i=1}^n \subset \mathbb{R}^m$ of prototypes and an input data points $x \in \mathbb{R}^m$, determine which prototype is closest to the input, i.e. solve $x^* = \operatorname{argmin}_i \|x_i - x\|$. Given the above discussion, it is clear that this problem occurs in the context of assigning data to clusters. Moreover, if the prototypes are labeled, nearest neighbor search also allows for nearest neighbor classification in that input data points x can be automatically labeled using the label of their respective closest prototype x^* .

Classically, naïve nearest neighbor search requires efforts of $\mathcal{O}(n)$. However, there exist specific data structures such as k D-trees which allow for nearest neighbor search in $\mathcal{O}(\log n)$. This gain comes at the cost of having to pre-process the prototypes in order to structure them correspondingly. This typically requires efforts of $\mathcal{O}(n \cdot \log n)$ which amortize if the prototypes have to be searched repeatedly, for instance, in data base search scenarios. While nearest neighbor search is thus not too demanding a problem for classical computers, it is interesting to observe that quantum computing can, in principle, accomplish search in $\mathcal{O}(\sqrt{n})$ even without data pre-processing.

Wiebe et al. [196] stress the importance of nearest neighbor search in clustering and classification and present quantum nearest neighbor approaches for binary classification and k-means clustering. In essence, their approach allows for quantum computation of Euclidean distances between sparse data points and assumes the availability of quantum oracles that can efficiently look up the j-th value of prototype x_i and identify the l-th non-zero value in prototype x_i . They furthermore assume that is is possible to prepare a quantum state that encodes $\|x_i - x\|$ and present a quantum circuit for this purpose. They then apply a classical quantum minimum search algorithm due to Dürr and Høyer [197] which itself is a variant of Grover's amplitude amplification procedure for searching with oracles [37, 38]. They prove that the number of oracle queries of their method scales as $\mathcal{O}(n \cdot \log n)$ which is thus as fast as informed classical search. In numerical simulation experiments on common machine learning benchmark data they find that their technique is robust to noise arising from coherent amplitude estimation, performs well, and asymptotically outperforms classical sampling techniques for nearest neighbor search.

Llyod et al. [198] are concerned with the problem of k-nearest neighbor search. Whereas naïve classical solutions to this problem would require efforts of $\mathcal{O}(kn)$, the quantum algorithm proposed in [198] has a runtime complexity of only $\mathcal{O}(\log kn)$ and thus achieves exponential speedup. However, the authors once again assume that classical data has efficiently been loaded into a QRAM so that their algorithm can access it in a quantum parallel manner. They then proceed to prepare states that represent sub-norms of the input data on which they run adiabatic quantum optimization to identify and select the k nearest prototypes to a given query.

In a recent contribution, Basheer et al. [199] integrate ideas from [196] and [198] with re-

cent techniques for preparing data into amplitudes of quantum states [200] and present specific quantum circuits for k-nearest neighbor search in classification. Their algorithm has a runtime complexity of $\mathcal{O}(\sqrt{kn})$ but, crucially, does not make any assumptions as to the availability of QRAMs. In numerical experiments, in which the authors simulate classification problems with few classes using only few qubits, they observe the approach to work well.

They also point out that their methods seamlessly applies to downstream processing in quantum computing settings. That is, the proposed algorithm can be directly used on quantum data and thus circumvent the need of reconstructing quantum states using measurements on an ensemble of identical states in physics applications.

Concerned with a much simplified nearest neighbor setting, Schuld et al. [201] present a distance-based classifier which they realize in terms of a simple quantum interference circuit. Excluding a simple state preparation procedure, the circuit consist of only one Hadamard gate two single-qubit measurements units. It computes distances between input data points and two prototypes in quantum parallel and has been implemented in the IBM quantum experience environment.

Quantum Algorithms for Classification

Pattern recognition or classification is the problem of assigning labels (signifying classes or categories) to observations of objects (represented in terms of data points) and a function or an algorithm that accomplishes classification, especially in a specific scenario, is called a classifier.

If a classifier can classify objects from two classes (say pictures of cats and dogs), it is said to be a binary classifier. A classifier that can classify multiple classes (say pictures of cats, dogs, mice, horses, elephants, ...) is called a multi-class classifier. Since any multi-class classification problem can be addressed by several one-versus-all classifiers, binary classification is the more fundamental task and the problem of training a robust and reliable binary classifier using labeled training data is one the most important problems in machine learning.

There exist numerous possible models that allow for binary classification. Some of these are probabilistic (naïve Bayes classifiers or Bayesian networks), some operate on categorical data (decision trees), and some are geometric or linear algebraic in nature (least squares classifiers, linear discriminant classifiers, support vector machines, etc.).

A simple binary linear classifier for input $x \in \mathbb{R}^m$ is a threshold function of the following form

$$y = f(\boldsymbol{x} \mid \boldsymbol{\theta}) = \operatorname{sign}(\boldsymbol{x}^{\mathsf{T}} \boldsymbol{w} - b) \tag{34}$$

whose parameters $\theta = \{w, b\}$ are a projection- or weight vector $w \in \mathbb{R}^m$ and a thresholdor bias value $b \in \mathbb{R}$. While the internal computations $x^{\intercal}w - b$ of this model are indeed linear, the function sign is said to be a non-linear activation function which produces outputs in $\{-1,+1\}$ and thus allows for binary decision making. It is also common to consider activation functions such as the hyperbolic tangent which produces outputs in (-1,+1), the logistic function with outputs in (0,1), or rectified linear units with outputs in $[0,\infty)$. The choice of activation can impact the ease of training but the fundamental problem in each case is to determine suitable \boldsymbol{w} and \boldsymbol{b} . Depending on the choice of learning algorithm (least squares training, linear discriminant training, support vector training, etc.) this can be more or less challenging but is well understood in general and computationally not too demanding.

Binary linear classifiers learn linear decision boundaries which divide the underlying data space into two disjoint half-spaces, namely $\{x \mid f(x) \geq 0\}$ and $\{x \mid f(x) < 0\}$. However, it is interesting to note that many linear classification models allow for invoking the *kernel trick* [202]. This way, data can implicitly be processed in very high to infinitely dimensional kernel Hilbert spaces. This then makes it possible to use simple, linear models in order to learn highly non-linear decision boundaries.

Since linear classifiers (kernelized or not) rely on linear algebraic computations, they often allow for corresponding quantum computing algorithms. Early attempts in this direction were made by Schuld et al. [203, 204]. In [203] they argue that quantum computing can outperform classical techniques in the case of ambiguous input patterns. They develop a quantum nearest neighbor classifier that involves Trugenberger's quantum algorithm for measuring Hamming distances in quantum associative memories [205, 206]. In simulation experiments on standard benchmark data, they find their method to work well. Its runtime behavior is similar to classical implementations but the authors remark that, if there was an efficient $\mathcal{O}(n)$ approach for constructing the required quantum superposition states or if a QRAM would exist, their quantum algorithm would be independent of the number of training data, a feat that seems impossible for the classical counterpart.

In [204], Schuld et al. present a quantum perceptron, i.e. a quantum circuit that realizes the binary linear classification function in (34). They propose quantum phase encoding of $\varphi = x^\intercal w$ as $|\varphi\rangle = \frac{1}{\sqrt{2}} \left(|0\rangle + e^{i\,\varphi}|1\rangle \right)$ and then use the phase estimation algorithm [113, 114] to determine the corresponding output $|y\rangle$. In terms of runtime or ease of implementation this does not seem to offer any benefits over the classical approach; however, the authors point out that their quantum perceptron yields outputs which contain quantum information. This could be used for superposition-based learning procedures in which training data enter the algorithm as a superposition of data points which would then allow for quantum parallel processing.

Wiebe et al. [207] ask if quantum computation could lead to non-trivial improvements in the computational and statistical complexity of perceptrons and present algorithms which answer both these questions affirmatively. To this end, they consider the notion of the version space of a set of training data for a binary classification problem. This term refers to the set of all hyper-planes that perfectly separate the two classes under consideration. Seen from the version space perspective, the problem of classifier training can be cast as a search problem which, in turn, can be solved using Grover's quantum search algorithm. Based on these

ideas, Wiebe et al. develop two algorithms with different speedups for perceptron training. The first trains in $\mathcal{O}(\sqrt{n})$ where n denotes the number of training data points; the second trains in $\mathcal{O}(\rho^{-1/2})$ where ρ denotes the optimal margin between the two classes under consideration. As of this writing, it is difficult to imagine how this could be achieved classically. The work in [207] therefore establishes that quantum computing can speed up perceptron training and thus improve on one of the most fundamental tasks in machine learning. The authors thus conclude that the quantum computing point of view can lead to a deeper general understanding of machine learning and may lead to models for which there is no classical analogue.

Tacchino et al. [208] have recently shown how to practically implement a simple binary-valued perceptron on an existing quantum device, namely a 5-qubits IBM quantum computer based on superconducting technology. In practical experiments, they demonstrate successful classification of 4 bits strings using 2 qubits and of 16 bits strings using 4 qubits. This hints at an exponential storage advantages over classical methods but the authors concede that generic, i.e. non binary patterns, may need quantum states or unitary transformations whose preparation requires exponentially many quantum gates. This, in turn, would eliminate the quantum advantage of their method. They also acknowledge that current NISQ era devices do not allow for arbitrary controlled operations but require them to be broken down into operations involving only single- and two-qubit gates which significantly increases circuit depth and thus the risk of decohenrence. As a possible future generalization away from binary inputs, the authors point to phase encoding similar to the approach in [204].

In another recent contribution, Schuld et al. [209] observe that quantum computing bears certain similarities to the idea of working with kernel methods in machine learning in that both, quantum computers and kernel machines, process information in (possibly infinitely high-dimensional) Hilbert spaces. They argue that this can lead to new ideas for the design of quantum machine learning algorithms and, in particular, interpret the problem of encoding classical data in a quantum state as the problem of computing a non-linear kernel function that maps the data into quantum Hilbert space.

To this end, they associate a quantum Hilbert space with a machine learning feature space and derive a kernel that corresponds to the inner product of quantum states. They discuss how this allows for the encoding of computational basis states, amplitude encoding of data points, and representations of product states and use these insights to devise a parameterized circuit model of a quantum support vector machine. The circuit is composed of displacement- and phase gates whose parameters are determined via classical stochastic gradient optimization. Schuld et al. thus consider a variational quantum training algorithm that combines quantum computations with classical computations in an iterative feedback loop. Experimental verification with a simple mini benchmark data set of two-dimensional data points from two classes suggests that the idea works well in practice.

Similarly, Grant et al. [210] are concerned with the capabilities of present day NISQ devices and point out that hybrid quantum-classical algorithms where quantum computers run model circuits and classical computers perform statistical loss minimization to train the

model circuit are currently the best strategy for quantum (assisted) machine learning. Building on concepts first proposed in [211], they therefore propose hierarchical quantum circuits for binary classification. In particular, they consider parameterized circuits of tensor network topology. In experiments with tree tensor networks and multi-scale entanglement renormalization (MERA) networks, they determine their parameters classically and then deploy them on an ibmqx4 machine available within the IBM quantum experience environment. The data they consider are small subsets of standard machine learning benchmarks and their results suggest that the method is robust to noise and can achieve high accuracy. Importantly, the authors stress that their method allows for classifying classical- as well as quantum data.

Finally, Date et al. [190] recently discussed how to cast the problem of training linear support vector machines for binary classification as a quadratic unconstrained binary optimization problem. While naïve classical algorithms would require efforts of $\mathcal{O}(n^3)$ for support vector machine training with n data points, setting up the QUBO requires only $\mathcal{O}(n^2)$ computations and its solution could subsequently happen on an adiabatic quantum computer. This is not demonstrated practically but, in line with their linear regression results in [185], should be entirely possible. The authors' mainly theoretical work thus suggests that well established machine learning baseline models can be trained on adiabatic quantum computers from where the resulting model parameters can easily be read into digital memories for further processing.

4.3.6 Quantum Boosting

Boosting is a machine learning technique tailored towards classifier training that is rather easy to implement and known to yield well working systems. There are several variants and flavors, but the original idea of adaptive boosting (AdaBoost) [212] is arguably still the most popular approach. Assuming a set of individually weak classifiers, AdaBoost evaluates an exponential loss to determine which (re-)weighted sum of their outputs provide a strong classifier. The resulting boosted classifier is also known as an ensemble classifier or, if the individual weak learners are non-linear threshold units such as in (34), a shallow neural network. Shortly after its inception in the 1990s, classifier boosting led to breakthroughs in rapid and reliable computer vision [213] and correspondingly trained systems were, for instance, implemented in consumer cameras to detect smiling faces. Alas, the good performance of boosted classifiers comes at the price of typically extensive training times, as it requires numerous rounds of training to select a high-performance ensemble among a very large number of individual learners. Yet, since boosting is essentially a subset selection problem, it has been identified as a setting for potential quantum speedup early on.

Neven et al. [214] propose adiabatic quantum optimization for boosting. They transform the originally continuous weight optimization problem into an optimization problem over discrete variables of low bit depth and consider an adapted quadratic loss function. This allows them to express boosted classifier training as a QUBO. In experiments with heuristic surrogates for quantum hardware as well as with an early D-Wave machine, their QBoost approach

compares favorably to classical AdaBoost. It reduces training efforts and leads to better generalization and faster runtimes because the resulting ensembles are typically found to be small.

However, the authors concede that these advantages are mainly due to their bit-constrained learning model and not a quantum effect per se. The notable result is thus that a restricted (and hence implicitly regularized) model that was specifically designed for implementation on quantum hardware can outperform unrestricted conventional models. Yet, since solving the underlying discrete optimization problem might be too demanding for conventional computers if the problem size is large, the work in [214] establishes that quantum computing enables approaches towards machine learning that would be classically infeasible.

Pudenz and Lidar [215], too, consider quantum adiabatic evolution for classifier training in a manner similar to that of Neven et al. However, they also apply adiabatic quantum computing in the application phase of their system and adiabatically evolve strong classifiers identified during training on a superposition of inputs in order to be able to identify likely anomalous elements in their data space in a quantum parallel manner. The practical application they consider is the problem of verification and validation of classical software where programming errors (bugs) are the anomalies to be detected. Extensive simulation experiments indicate that their methods works well.

Schuld and Petruccione [216] consider the problem of creating a classifier ensemble to be a problem of quantum state preparation. A quantum parallel evaluation of individual quantum classifiers on such states allows for estimating their combined or aggregated decision in terms of a single qubit measurement. The authors argue that their framework allows for exponentially large ensembles and thus for advantages over classical approaches.

As a quantum gate model of a weak classifier, they consider their earlier proposal of a quantum perceptron [204] and, based thereupon, describe a protocol for preparing states that represent classifier ensembles. Crucially, this requires estimates of the accuracy of the individual classifiers which can then be used in a weighting scheme. While the authors are more concerned with the feasibility of quantum ensembles rather than with their potential advantages, they do point to several possible applications in quantum physics and remark that their ideas suggest approaches towards optimization-free quantum machine learning. In other words, a potential benefit of the proposed approach is that measuring large enough ensembles of individually weak classifiers can lead to accurate decisions without that the ensemble would have to be tuned in a training process.

4.3.7 Quantum Support Vector Machines

Kernel machines are an important class of classical machine learning models. The key idea observation is that many (linear) machine learning methods require access to the data only through inner products between feature representations of two data points, i.e.

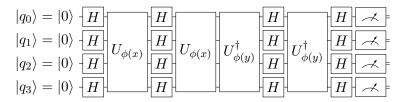


Figure 3: Variational 4-qubit circuit for quantum kernel machines. Data points x and y are both mapped into the 2^4 -dimensional Hilbert space via parameters of unitary gates $U_{\phi(x)}$ and $U_{\phi(y)}$.

 $\langle \phi(x) | \phi(y) \rangle$. The mapping $\phi: \mathcal{X} \to \mathbb{R}^d$ to the d-dimensional feature space has to be specified by the user and is crucial for the quality of the resulting model. When the dimension d is large, the explicit computation of ϕ can be circumvented via a so-called kernel function $k(x,y) = \langle \phi(x) | \phi(y) \rangle$.

The most prominent incarnation of this type of method is the support vector machine (SVM) [217]. Classically, the choice of the feature map $\phi(\cdot)$ or, more commonly, the kernel function $k(\cdot,\cdot)$ are either domain specific (for instance, string kernels, graph kernels, tree kernels, image kernels, document embedding kernels, time series kernels, ...) or generic (for instance, radial basis function kernels, polynomial kernels, Fisher kernels, ...).

In contrast, the feature map of quantum kernel machines is hardware specific. Instead of relying on problem specific knowledge to construct the feature map or the kernel, the intrinsically 2^n -dimensional Hilbert space of an n-qubit register is utilized to realize the feature map [218].

A schematic representation of a corresponding quantum gate circuit is shown in Fig. 3. There, a maximal superposition is prepared and then passed through n-qubit unitaries which create the feature space transformation of the data. It is important to understand that data does not enter the circuit in the discrete qubit state space. Instead, data is passed in form of parameters of universal unitary gates $U_{\phi(x)}$ and $U_{\phi(y)}$ representing (parts of) the high-dimensional feature map. Each classical n-bit binary string is interpreted as one feature and the corresponding probability amplitudes of the qubit state as feature values. The actual kernel value k(x,y) is then given by estimating the transition amplitude $|\langle \phi(x), \phi(y) \rangle|^2 = |\langle 0^n | U_{\phi(y)}^\dagger U_{\phi(x)} | 0^n \rangle|^2$. Clearly, the specific choice of $U_{\phi(\cdot)}$ is not fixed and can be tuned for the application at hand. Obtaining the full kernel matrix for n data points requires n(n+1)/2 runs of that circuit. The resulting quantum kernel matrix is then ready to be used in any kind of kernel machine, for instance, in support vector machines, kernel regression, or kernelized k-means clustering.

In addition to quantum k-nearest neighbors algorithms, Wittek [219] introduces the concepts of quantum support vector machine algorithms based on a slightly extended formulation of traditional SVMs using least-squares optimization. He indicates that quantum SVMs have a high generalization performance, even though the extension to general multi-class

problems is expensive.

Wang et al. [220] take advantage of the fact that the *least squares support vector machine* algorithm shows high efficiency in processing a small quantity of data for the problem of trend recognition. To optimize the parameters of the least squares support vector machine and to establish the curve fitting model, the authors use quantum particle swarm optimization to realize a faster global parameter search. To corroborate the practicality of their ideas, the authors perform on a mathematical error analysis.

Rebentrost, Mohseni, and Lloyd [221] demonstrate that support vector machines can be implemented on a quantum computer such that they achieve logarithmic complexity in the data dimension and the number of training examples. This constitutes an exponential speedup over classical sampling algorithms which run in polynomial time. The main component of the quantum algorithm developed by the authors is a non-sparse matrix exponentiation technique which inverts the training data inner-product (kernel) matrix. However, the latter is based on quantum regression methods such as those discussed above. Since such methods assume that quantum states can either be prepared efficiently or are available for convenient lookup in a quantum random access memory, the ideas in [221] may be of limited practical feasibility on current NISQ era devices.

Havenstein, Thomas and Chandrasekaran [222] compared the runtime and accuracy of a classical SVM against two QSVMs, a kernel-based QSVM and a variational QSVM. The authors have determined that, for binary classification problems, a QSVM does not provide any substantial improvement over a classical SVM. However, in some multiclass classification cases the variational QSVMs exhibit a higher accuracy than classical SVMs. The QSVM multiclass classifiers are therefore promising as the number of available qubits increases. A similar conclusion has been reached by Zahorodko et al. [223] who expect that quantum-enhanced machine learning is suitable for classifying high-dimensional data sets in cases where especially binary classification can be processed by single-qubit systems in an efficient way.

4.3.8 Quantum Neural Networks

Artificial neural networks are machine learning models that mimic the way information is processed in biological brains. They are composed of small interconnected computational units called neurons which receive weighted input from other neurons. The weighted input received by a neuron is then typically summed up and subjected to a non-linear activation function. In other words, individual neurons can typically be thought of as a kind of binary linear classifier similar to the one in (34).

An important result due to Hornik [224] is that large enough neural networks are universal approximators. This mathematical statement is to be understood in the sense that they can approximate any Borel measurable function arbitrarily well. Their general universal approximation characteristic makes neural networks powerful tools for a wide range of artifi-

cial intelligence problems. Because of this, neural networks had their first boom period in the 1980s when the backpropagation algorithm for their training became available [225]. They then faded out of popularity because, back then, computers were not yet powerful enough to perform the extensive computations required in supervised neural network training. Over the past decade they returned to the limelight and are currently the most common tool in machine learning.

As the task of training large neural networks is computationally burdensome, quantum neural networks have become a popular topic among quantum computing researchers. Indeed, parameterized quantum circuits, i.e. quantum circuits whose gates or unitary operators come with tunable phase parameters, bear a certain resemblance to layered neural networks. Moreover, similar to the universal approximation theorem for neural networks, large or complex enough quantum circuits are known to be able to represent any target function with arbitrary precision [129]. Lin, Tegmark, and Rolnick also argue that data sets arising from measurements of physical systems will exhibit symmetry and locality so that there should exist less than exponentially large models that lead to useful results. Against this backdrop, variational quantum circuit models are intended to approximate solutions to a task at hand while also restricting the number of quantum operators and quantum circuit depth.

The proven universal approximation capability of classical neural networks crucially depends on the fact that their synaptic summations are subjected to non-linear activation functions. However, this poses a challenge for quantum neural networks. For instance, when data is encoded in the amplitudes of a quantum state, one cannot apply an arbitrary non-linear function without distorting it. Allcock et al. [226] therefore suggest algorithms for training and evaluating feed forward neural networks which are based on canonical classical feed forward and backpropagation procedures. Their algorithms rely on a quantum subroutine for approximating inner products between vectors which would yield training times quadratically faster in the size of the network than the classical counterparts. In addition, the authors assume that their approach is intrinsically resilient to overfitting because it quantum mimics classical regularization techniques. However, although the hybrid quantum-classical training procedure in [226] closely follows classical neural network training, it may not yet be practical. This is because Allcock et al. assume all sequential steps during their training procedure to be computed classically while quantum operations are only used for estimating the inner products in these steps. This requires their in- and outputs to be read from- and written to a QRAM which is impossible on current quantum computers.

Other works approach the aforementioned challenge with measurements. Beer et al. [227] propose a quantum analogue of a classical neuron from which they construct quantum feed forward neural networks capable of universal quantum computation. The basic idea is to consider a quantum perceptron to be an arbitrary unitary operator that maps m input qubits to n output qubits. Inputs are initialised in a possibly unknown mixed state and outputs in a fiducial product state, that is, in quantum state one can reliably reproduce with low variability. Using fidelity as a cost function, the authors propose a training procedure where perceptron unitaries are updated using phase shifts proportional to the loss incurred for a given pair

of network in- and outputs. Regarding this procedure, they note that, if their quantum neurons are organized in layered architecture, updates can be accomplished layer by layer without having to access the full quantum circuit. They argue that this first of all reduces memory requirements and second of all decouples memory requirements from network depth. In other words, training efforts scale with the width of the proposed quantum neural network.

Running simulation experiments with quantum neural networks of moderate sizes, Beer et al. find their approach to generalize well to previously unseen testing data and also to be robust against noisy training data. A drawback with respect to present day quantum computers, however, is the assumption that quantum operators acting on arbitrarily many qubits were available. Given the present state of the art, this is technically not yet possible on existing devices.

Concerned with implementations on NISQ era quantum computers, Mitarai et al. [228] propose hybrid quantum-classical algorithm for quantum circuit learning. Their framework realizes task specific learning via iterative parameter tuning in circuits of a fixed depth. A theoretical analysis backed by simulation experiments reveals that their scheme can approximate nonlinear functions and the authors expect that the approach should be implementable on existing quantum computing devices.

Verdon et al. [229], too, focus on quantum neural network realizations for near-term quantum computers. In particular, they point out that it generally appears to be difficult to determine the initial parameters of a quantum circuit such that subsequent hybrid quantumclassical optimization will quickly converge towards local minima of the loss function under consideration. They therefore consider the use of classical neural networks which they train to generate initial parameters for the quantum learning process. They empirically find that this appears to allow for considerably faster variational quantum optimization in that the total number of optimization iterations required to reach a given accuracy is reduced considerably. The authors experimentally verify this for the quantum approximate optimization algorithm (QAOA) and for the variational quantum eigensolver (VQE) which they apply to problems such solving a graph max-cut task, the Sherrington-Kirkpatrick Ising model, or the Hubbard model. In fact, they observe that optimization strategies learned by the classical neural network seem to generalize across problem sizes. This leads them to expect that it may be possible to train their system on small problems that can be classically simulated and then transfer it to larger problems which are classically intractable. This way, the number of costly iterations in variational quantum-classical optimization for such settings could be reduced.

Indeed, the issue of how to initialize parameterized quantum circuits for quantum neural network training with hybrid quantum-classical gradient descent algorithms is critical. Mc-Clean et al. [230] observe that random initializations are popular and commonly considered by many researchers as they are simple as well as hardware efficient. However, they point to apparent inherent limitations of this approach when dealing with more than a few qubits. The authors demonstrate analytically as well as numerically that descent based training of even reasonably parameterized quantum circuits suffers from vanishing gradients and that there

exist "barren plateaus in quantum neural network training landscapes". This is to say that, during gradient based parameter optimization, expected values of observable required for the process tend to concentrate which causes their gradients to tend to zero. In other words, Mc-Clean et al. show that there is a high probability for gradient-based variational quantum neural network training to get stuck in local minima and thus to yield sub-optimal parameterizations. They argue that this is largely due to technical limitations of current quantum circuits and has to be taken into consideration when designing larger quantum neural networks.

Concerned with a practical implementation of parameterized quantum circuits on a trapped ion quantum computer, Zhu et al. [231] consider parameter optimization by means of particle swarm optimization and Bayesian optimization. Their work contains several noteworthy innovations. First of all they consider the problem of training generative model circuits in the spirit of generative neural networks. These are neural networks which learn to produce new, plausible data points. Often, this happens in encoder-decoder architectures where training data are first compressed into low-dimensional representations and then decompressed into their original form. Known applications range from automatic text generation, over deep faked images or videos to the synthesis of plausible particle trajectories in high-energy physics simulations.

Zhu et al. argue that generative models can be expected to benefit considerably from potential quantum speedup and, second of all, describe a quantum circuit that is able to learn to reproduce patterns in the bars-and-stripes data set, a very simple, low-dimensional, yet common benchmark for experiments with novel techniques that is attributed to MacKay [20]. Tailored towards this data, they propose a quantum circuit design consisting of parameterized rotation and entangling operators which they adjust in a hybrid quantum-classical manner. In the classical parameter optimization steps, they work with particle swarm- and Bayesian optimization methods both of which are well established tools in classical machine learning. They implement their approach on a custom made quantum computing platform presented in [232] and observe that the convergence of their quantum circuit towards an optimized model critically depends on the optimization strategy. Their practical results suggest that it is practically possible to successfully train high-dimensional universal quantum circuits and that quantum neural networks may thus play a role in future generative modeling applications.

Leyton-Ortega et al. [233], too, are concerned with generative modeling for the canonical bars-and-stripes data and further investigate the observation that the performance of hybrid quantum-classical algorithms seems to depend on the choice of classical optimizer and on the circuit model. They argue that any conclusive (empirical) statement to resolve this issue requires to conduct extensive experiments with different optimization algorithms and circuit designs. To accomplish this in practice, they work with Rigetti's quantum cloud service and investigate the practical performance of data-driven quantum circuit training for different classical solvers and different circuit designs. For the latter, they consider different entangling qubit connectivity graphs and varying circuit depths. Their extensive experiments reveal that gradient-free optimization algorithms lead to much better results than gradient-based solvers. In particular, the former seem to be much less affected by the noise characteristics of existing

quantum computing systems.

Another example of a real world implementation of quantum neurons can be found in work by Tacchino et al. [208]. They introduce a quantum neuron design where $m=2^n$ dimensional data- and weight vectors are encoded using only n qubits which makes use of the exponential storage advantage of quantum computers. To generate entangled states they consider hyper-graph states [234] and prepare them using several controlled Z gates. Non-linear outputs are realized by means of quantum measurements of ancilla qubits. The authors demonstrate the practical feasibility of their ideas through an implementation of the n=2 case on an IBM Q5 quantum computer and find that their quantum neuron can successfully learn a simple bars-and-stripes pattern completion task.

Zhao et al. [168] observe that Bayesian methods such as Gaussian processes regression provide proven and successful machine learning models for reasoning under uncertainty and that they have recently been adapted to deep learning problems. They also observe that there exist proposals for quantum Gaussian process regression based on the HHL algorithm [235] and discuss how this allows for deep network training without backpropagation. Having established this connection, the authors propose a hybrid quantum-classical algorithm for Bayesian deep learning. They show how non-linear kernel matrices can be approximated in terms of polynomial series and then easily lend themselves to be used as quantum density operators. Using the HHL algorithm, whose very specific prerequisites are met by Gaussian process kernel matrices, their quantum subroutine achieves polynomially faster matrix inversion than classical methods and the authors simulate their algorithm on a Rigetti quantum virtual machine. For a considerably reduced setting with 2×2 Gaussian process kernel matrices, they also present implementations on the Rigetti 8Q-Agave and IBM Q5 quantum computers. Interestingly, they find that the probability of successful training is much higher on the IBM architecture which they attribute to its longer coherence times. In particular, the authors observe the probability of their protocol to train successfully to amount to 89% which is an encouraging result with respect to future efforts involving larger problem sizes.

Helmholtz machines are a class of neural networks that specifically allow for generative modelling [236]. They consist of two sub-networks, a bottom-up recognition network that maps input data to a distribution over hidden variables and a top-down generative network that generates novel data from an instantiation of the values of the hidden variables. The training of a Helmholtz machine usually happens in an unsupervised manner using an algorithm known as the "wake-sleep algorithm" [237].

Van Dam et al. [238] present a hybrid quantum-classical approach for Helmholtz machines. Their corresponding parameterized shallow quantum circuit models can be trained in a gradient-free manner using an optimization scheme based on the wake-sleep algorithm. The authors implement their system on the Quantum Inspire simulator and evaluate its practical performance on a reduced bars-and-stripes data set consisting of binary images of size 2×2 pixels. For this data they consider Helmholtz machines with four visible neurons and three hidden ones. They observe that their hybrid quantum-classical algorithm learns better

network parameters than a corresponding purely classical implementation. They emphasize that their method can efficiently approximate the underlying probability distributions with only polynomially many evaluations of the quantum circuit and also expect their approach to be implementable on real quantum computing devices.

Boltzmann machines are yet another type of neural networks, in particular, a probabilistic extension of Hopfield networks. As such it is rather straightforward to conceive of quantum Boltzmann machines; this basically requires the definition of a suitable energy function whose minimizers constitute appropriate network parameters. Again due to their close connection to Hopfield networks, training can be realized via adiabatic quantum optimization [239] which, in turn, can be implemented on D-Wave computers [240].

Chen et al. [241] point out the the internal probabilistic states of a Boltzmann machine can be translated into quantum states which then leads to a purely quantum mechanical interpretation of Boltzmann machines called Born machines. These Born machines are yet another kind of generative model and can be used to represent distributions of classical data in terms of quantum states in superposition. In this regard, Liu and Wang [242] remark that quantum sampling should be of lower computational complexity than sampling in corresponding classical implementations. They realize this by means of projective measurements of qubits and propose a gradient-based quantum-classical algorithm for optimization of a quantum circuit such that the likelihood of generated samples can be estimated. In simulation experiments with bars-and-stripes data they observe their approach can learn data distributions, especially when deeper circuits are being used.

Coyle et al. [243] consider a restricted subset of Born machines, namely those whose Hamiltonians correspond to those encountered in Ising models. For these, they show that there is no efficient classical sampling scheme for the underlying quantum circuits and therefor propose a quantum-classical algorithm based on gradient descent. An innovative aspect of this approach is that the authors consider loss functions such as the Sinkhorn divergence or the Stein discrepancy. In experiments with practical implementation on Rigetti's forest platform (using a simulator as well as the Aspen quantum processing unit), the authors find that their procedure works well and reliably. As a potential future practical application of their generative modeling method, the authors point to the problem of quantum circuit compilation.

4.3.9 Federated and Distributed Quantum Machine Learning

In today's world, where cloud computing has established itself as a highly salable commodity tool for the of massive data sets, and where a great variety of data is available too, new concepts such as *federated learning* [244] promise further accelerated development. In federated learning, a central node holds the global model, it receives the trained parameters from client devices and then aggregates them to generate an updated and improved global model that is shared to all client nodes [244].

Chen and Yoo [244] deal with this concept in the context of quantum machine learning and discuss distributed training across several quantum computers. The authors expect a substantial reduction of training time and an advantage from the data security and data privacy perspective since the training is performed where the data is located. They demonstrate this federated training approach in experiments with hybrid quantum-classical machine learning models (which could potentially be translated to pure quantum machine learning models). In particular, in simulation experiments, they consider a quantum neural network combined with a classical pre-trained convolutional model on an image analysis problem. They observe that their federated training protocol does not sacrifice performance for accuracy (at least in the considered setting) and conclude that federated quantum machine learning might help to preserve privacy and to distribute computational efforts across arrays of NISQ devices.

Sheng and Zhou [245] introduced a protocol for distributed secure quantum machine learning. Their overall goal is to augment classical clients with basic added quantum computing technology to delegate a remote quantum machine learning tasks to a quantum server in a manner such that privacy data is preserved preserved. They present a distributed quantum machine learning protocol for the specific scenario where clients together with a remote server can assign two-dimensional data points to different clusters. Their protocol is secure in that it does not leak information relevant to this task and, similar to quantum information transmission, would immediately realize eavesdropping attempts. The authors suggest that this protocol could be extended to higher dimensional data and to bid data settings, but, for now, their work is rather conceptual and an interesting first step into a new direction.

4.3.10 Variational Quantum Algorithms

Above, we already saw several hybrid- or variational quantum computing approaches towards quantum machine learning. Examples included variational approaches towards quantum circuit design [143], quantum state preparation and loading [155], quantum linear system solving [175, 176], quantum support vector machine training [209], and ideas for the realization of quantum neural networks [210, 228, 229, 230, 231, 238]. Indeed, given the technical capabilities of present day NISQ devices, hybrid quantum-classical solutions, where quantum computers run model quantum circuits and classical computers repeatedly perform (statistical) optimization to adapt those models to a given task the task, currently appear to be the best strategy for quantum (assisted) machine learning in a manner that allows for harnessing quantum supremacy [246].

What all present such approaches have in common is that they involve parameterized quantum circuits. While these circuits may or may not have been specifically designed for a task at hand, they typically consist of tunable gates in sequence or in parallel. Most commonly, parameterized gates perform qubit rotations and the problem is to adjust their parameters such that the circuit as a whole performs the desired computation with high probability. In order to convey a better understanding of this general idea, we next discuss two prominent

examples of variational quantum computing algorithms.

A well known and fundamental example of a variational quantum computing algorithm is the *variational quantum eigensolver* (VQE) introduced by Peruzzo et al. [85]. Motivated by the problem of having to estimate the ground state energy of certain chemical molecules, they consider the general problem of estimating the bottom eigenvalues of Hamiltonian operators. They observe that traditional quantum algorithms for eigenvalue computation relied on the phase estimation methods [115]. While these methods promise exponential speedup, they also generally require exponentially large numbers of quantum gates so that real world implementations on existing quantum computers are only possible for problems of small size [247].

Peruzzo et al. therefore propose a different approach and begin by observing the classical result that, for any Hermitian operator H, the Rayleigh quotient

$$r = \frac{\langle \psi | H | \psi \rangle}{\langle \psi | \psi \rangle} \tag{35}$$

is minimal exactly if $|\psi\rangle$ is the bottom eigenvector of H. Even more, if $|\psi\rangle$ is the bottom eigenvector of H, then it is known that r corresponds to the sought after minimum eigenvalue.

Based on this observation, the authors propose to consider a random, parameterized state $|\psi(\theta)\rangle$ and to systematically vary θ to determine the minimizer of (35). In order to efficiently evaluate r on a quantum computer, they represent H as a linear combination of tensor products of Pauli operators and note that expected values of such operators can be estimated by means of local qubit measurements only. A considerable advantage their approach has over phase estimation techniques is that it requires only shallow quantum circuits as opposed to deep ones. It will therefore be less affected by short coherence times on existing quantum devices. A drawback is that the variational approach require polynomially many executions of the circuits as opposed to just in the case of phase estimation approaches. For each of these calls, $|\psi(\theta)\rangle$ needs to be prepared and, in order to do this efficiently, the authors propose to let $|\psi(\theta)\rangle = U(\theta)|\phi\rangle$. Here, $|\phi\rangle$ is an appropriate, easily prepared reference state (for their molecular energy application, the authors consider the Hartree-Fock ground state) and $U(\theta) = \exp(-iR(\theta))$ where R is a rotation that acts on each qubit individually and can be decomposed into products of two simple Pauli operators. A quantum circuit which implements these computations can then be called in each iteration $t = 1, 2, \dots, T$ of an outer loop executed on a classical computer which updates $heta_{t+1} = heta_t - d heta_t$ where, for instance, $d\theta_t = r(\theta_t + \Delta\theta) - r(\theta_t - \Delta\theta)$ is an approximation of the gradient $\nabla r(\theta_t)$ of the Rayleigh quotient r seen as a function of θ . Note, however, that gradient-free optimization schemes are possible as well. In fact, the latter do not seem to suffer from the phenomenon of "barren plataeus" [230] but often lead to better solutions than gradient-based schemes [143].

Peruzzo et al. demonstrate the practical feasibility of their method in experiments with an implementation an a custom made photonic quantum computer. In particular, they successfully show that ground state energies of the 4×4 Hamiltonian of a helium hydride ion can be calculated reliably.

The quantum approximate optimization algorithm (QAOA) due to Farhi et al. [84] is another variational method that was specifically developed to determine approximate solutions to combinatorial optimization problems. In the above discussion of adiabatic quantum computing, we already saw that potential solutions to such problems can often be encoded in terms of n binary variables z_j which have to meet m conditions $c_k(z) \in \{0,1\}$ and thus maximize $c(z) = \sum_k c_k(z)$. Translating this setting into a formulation involving an n-qubit system $|s\rangle$ in a 2^n dimensional state space, Farhi et al. propose to consider an operator

$$U(C,\gamma) = \exp(-i\gamma C) = \prod_{k=1}^{m} \exp(-i\gamma C_k)$$
 (36)

where C denotes the overall problem Hamiltonian. They also define an operator

$$U(B,\beta) = \exp(-i\beta B) = \prod_{j=1}^{n} \exp(-i\beta \sigma_{j}^{x})$$
 (37)

where $B=\sum_j \sigma_j^x$ and σ_j^x denotes the Pauli spin matrix σ^x acting on the jth qubit. The next crucial idea is to introduce a total of 2p angles $\gamma_1,\ldots,\gamma_p,\beta_1,\ldots,\beta_p$ and to consider the system

$$|\gamma, \beta\rangle = U(B, \beta_p) U(C, \gamma_p) \cdots U(B, \beta_1) U(C, \gamma_1) |s\rangle$$
 (38)

which, as the authors emphasize, can be computed by a quantum circuit of depth $(m+1) \cdot p$. The expected value of the Hamiltonian C under this "angle" state is $E_p(\gamma, \beta) = \langle \gamma, \beta | C | \gamma, \beta \rangle$ and the maximum value it can attain is called

$$M_p = \max_{\gamma, \beta} E_p(\gamma, \beta) \tag{39}$$

Crucially, the author prove that, if $p \to \infty$, then $M_p \to \max_{\boldsymbol{z}} c(\boldsymbol{z})$. This then suggests a variational algorithm for approximately solving the given combinatorial problem: Choose p, initialize $\gamma_1, \ldots, \gamma_p, \beta_1, \ldots, \beta_p$, run a quantum circuit to compute $|\boldsymbol{\gamma}, \boldsymbol{\beta}\rangle$ or to measure $E_p(\boldsymbol{\gamma}, \boldsymbol{\beta})$, use this observation to classically adjust the current angle parameters, and repeat until $E_p(\boldsymbol{\gamma}, \boldsymbol{\beta})$ does not improve anymore.

The reason why this idea works is that (38) can be seen as an approximation of an adiabatic quantum computing procedure with Hamiltonian $H(t) = (1 - t/T) \cdot B + t/T \cdot C$ [84].

Farhi et al. also consider several special cases where the problem Hamiltonian is of special structure and thus allows for particularly simple optimization routines. But, in general, the method is very versatile and can, in principle, be used to approximately solve any optimization problem that can be cast as a QUBO. The quality of the solution will depend on the choice of the parameter p. With respect to implementations on current NISQ devices, this allows for trading off circuit depth against approximation quality. Better approximations require larger values of p which lead to deeper quantum circuits which, in turn, may suffer from decohenrence. Smaller values of p entail shallower circuits which may be run in practice but also produce sub-optimal solutions.

Yet, with respect to quantum advantages that can be achieved by running QAOA, the jury is still out. According to a calculation by Dalzell et al. [248], QAOA would outperform any existing classical supercomputer for problems involving 420 qubits and 500 constraints. As practically relevant combinatorial optimization problems go, this is a rather moderate size, however, quantum computers that could work on problem sizes like this do not yet exist. Moreover, Akshay et al. [249] identified certain inherent limitations of the procedure. In particular, they found that its capabilities depends on the ratio of the number m of problem constraints and the number n of problem variables. In combinatorial optimization, it is known that problems where $\alpha=m/n<1$ tend to have several to many solutions whereas problems with $\alpha=m/n>1$ have only few solutions if at all. For the latter kind of problems, Akshay et al. found that, even for large p, QAOA does not seem to work well. The authors refer to this as a (solution) reachability deficiency of QAOA and concede that further research is required to better understand the phenomenon.

Overall, however, variational- or hybrid quantum-classical approaches can considerably reduce quantum computing resources (circuit depth, coherence time, qubit counts) required to run quantum machine learning methods in a stable manner. Importantly, while the state spaces considered in this context are exponentially large, parameterized circuits typically have fixed structures (mainly laid out by human experts) and the number of their parameters only scales polynomially with the number of qubits. In fact, a common current design principle for parameterized circuits is to consider NISQ hardware efficient layouts, i.e. circuits of low qubit connectivity and with comparatively simple gates. Here, tensor network designs are of increasing interest but another strategy is to apply variational algorithms to design simple circuits on which to run variational algorithms.

The appeal of variational quantum algorithms of parameterized quantum circuits to those who develop quantum neural networks can likely be attributed to the fact that both, deep neural networks and parameterized quantum circuits, consist of basic computational units arranged in layered structures. Moreover, both involve the use of classical optimization algorithms to adjust their parameters to a given, specific problem. Some authors therefore even refer to parameterized quantum circuits as quantum neural networks per se. However, there are differences which make this terminology questionable.

First of all, the basic computational units in a quantum circuit are unitary operators rather than non-linear function as in the case of neural networks. Above, we already discussed that the capabilities of neural networks crucially depend on the fact that the activation functions of neurons are non-linear. Non-linear computations in a quantum circuit can be realized in a manner that preserves coherence, e.g. through entanglement, or in a manner that destroys coherence, e.g. through qubit measurements. In particular in combination with ancialla qubits, these operations allow for modelling the behavior of artificial neurons.

Second of all, internal states of a quantum circuit cannot be read out in a manner that would preserve their quantum nature. Contrary to neural networks, it is therefore impossible to compute local gradients within a quantum circuit. Consequently, it is impossible to ap-

ply methods like the backpropagation algorithm to train quantum circuits. As of this writing, there thus do not exist any workable proposals for how to adjust the parameters of a parameterized circuit without truly external supervision through an independent (classical) optimization algorithm.

4.4 Quantum Machine Learning for Quantum Data

If the prospects of quantum enhanced machine learning for classical data are most exciting for machine learning researchers, then the prospects of quantum machine learning and quantum analytics for quantum data are likely most exciting for physicists and chemists. This is because the simulation of physical or chemical systems or processes on quantum computers is widely seen as another promising application of quantum computing [250]. Such quantum computing simulations produce quantum data whose subsequent analysis would either require to measure and read them into classical computers or, preferably, could itself involve quantum algorithms. Quantum machine learning for quantum data thus refers to the idea of using quantum machine learning methods to process quantum mechanical data on quantum devices.

For example, Grant et al. [210] note that their hierarchical quantum circuits for binary classification can be applied to the problem of quantum state classification. Basheer et al. [199], too, point out that their quantum circuits for k-nearest neighbor classification can be directly used on quantum data. Chen et al. [251] particularly focus on the question of how to classify quantum data and consider the specific problem of distinguishing between pure and mixed states. To accomplish this, they train a parameterized quantum circuit and report that their system performs close to optimal on the training data and, more importantly, also generalizes well to previously unseen inputs. Since they are dealing with the classification of genuine quantum mechanical data, they conclude that their approach is among the first to successfully solve a learning problem for which there is no classical counterpart.

The term quantum state tomography refers to the general problem of predicting (probabilities of) outcomes of measurements of hidden quantum states. Traditional approaches to this estimation problem would require exponentially many measurements but Aaronson [252] observes that the problem could also be thought of a a statistical learning problem. He then proves that estimates can be obtained from only linearly many measurements. Together with a team of coworkers he meanwhile verified this theoretical expectation in practice [253]. Working with a custom made photonic quantum computer, they experimentally demonstrate this linear scaling behavior for their system in optical systems and conclude that computational learning theory provides a useful tool for the analysis of information that is of genuinely quantum mechanical nature.

Legeza and Solyom [254] investigate quantum data compression for finite quantum systems for cases with dependent density matrices. The authors demonstrated a connection between the entropy of the left or right block and dimension of the Hilbert space of that block.

Romero et al. [255] are concerned with the problem of autoencoding quantum data. They propose a simple parameterized quantum circuit that can act as an autoencoder neural network, i.e. can compress and decompress its input data. In simulation experiments, they train their model in a quantum variational manner so that it learns to compress a set of quantum state data for which classical compression algorithms supposedly do not work. In particular, they consider the problem of compressing Fermionic wave functions, i.e. eigenstates of a number operator, of a system of n quantum particles, compare their results to analytically computed predictions, and find it to work well. Ding et al. [256], too, are concerned with autoencoding quantum states and propose a corresponding quantum circuit composed of quantum adders. In an experimental implementation in the Rigetti cloud, they consider states of three superposed qubits and find their method to produce results that agree well with theoretical expectations.

In classical machine learning, autoencoders are often used as generative models and have been extended towards more capable architectures such as generative adversarial networks (GANs). In this spirit, Benedetti et al. [257] propose a system composed of two parameterized quantum circuits for generative modeling. They then use their system to sample unknown pure quantum states. Numerical simulation experiments suggest that the approach works well and the authors foresee applications in quantum state tomography. Similarly, Hu et al. [258] demonstrate that quantum state synthesis based on generative adversarial learning is practically possible. They implement their approach on a superconducting quantum circuit and find that it can be trained to generate high fidelity quantum output from quantum input

In their article "quantum data processing and error correction", Schumacher and Nielsen [259] deal with the behavior of noisy quantum information channels. The authors introduce the concept of coherent information which specifies a (per quantum information processing non-increasing) quantity measuring the amount of quantum information conveyed in the noisy channel. This quantity indicates the necessary and sufficient condition for the existence of perfect quantum error correction. Quantum error correction, in turn, is of importance in the NISQ era since existing devices suffer from noise, unreliable operations, and decoherence so that methods for generating fault-tolerant result are sought after.

5 Characteristics of Quantum Machine Learning Methods

The ever faster growing literature on quantum machine learning and the many success stories reported therein seem to suggest that quantum computing for computational intelligence is just about to break into the mainstream. However, as of this writing, the content of many reports, including many of those surveyed above, has still to be taken with a grain of salt. The two major caveats are that reported results might be exaggerated or that reported results abstract away from technical reality and assume the existence of universal quantum computers whose hypothetical capabilities far exceed what is possible with current NISQ era devices.

As a testimony to the former, we refer to Aaronson's article "Read the Fine Print" [260] in which he critiques what he calls the "quantum machine learning mini-revolution". In particular, he presents a critical analysis of what the HHL algorithm [119] for linear system solving is capable of and contrasts that with how it is commonly perceived by enthusiasts in the public and the scientific community.

Aaronson observes that HHL is often seen as a quantum algorithm for solving general linear systems Ax = b in a time exponentially faster than classically possible. That is, if $A \in \mathbb{C}^{n \times n}$ and $b \in \mathbb{C}^n$, then HHL is often seen as a method for finding $x \in \mathbb{C}^n$ in a time proportional to $\log n$ rather than to n^2 . He then points to the fine print and emphasizes that HHL is actually a method for creating a quantum state $|x\rangle$ which approximately represents the sought after solution x. He further emphasizes that the quantum speedup of the method hinges on the assumption that a state $|b\rangle$ representing b can be prepared efficiently. This $encoding\ problem$ may be easy in certain special cases, for example, if b is a unit vector whose entries have about the same magnitude. In general, however, state encoding will require efforts at least polynomial in a. Moreover, once $|b\rangle$ has been prepared, it must be coherently subjected to the operator $\exp(-iAt)$ for a period of time proportional to $ks/\epsilon \cdot \log n$. Here, k and k measure condition number and sparseness of matrix k and k indicates the approximation error users are willing to tolerate. Aaronson finally stresses that not every matrix is well conditioned and sparse enough to meet the implicit requirements under which the HHL algorithm is guaranteed to run efficiently.

Indeed, the general utility of HHL has been scrutinized before. For instance, Childs [261] observes that producing a quantum state $|x\rangle = A^{-1}|b\rangle$ is not tantamount to solving the original linear system. Rather, he points out the *decoding problem* and notes that obtaining x from $|x\rangle$ requires $\mathcal{O}(n)$ quantum measurements. This leads him to conclude that, when it comes to really solving a system of linear equations, present quantum algorithms do not yet have an exponential edge over classical approaches.

Both Aaronson and Childs note that the inventors of the HHL algorithm acknowledge all of this. Indeed, the above shortcomings and restrictions might be the reason why Harrow, Hassidim, and Lloyd [119] remark that there are situations in which we may not be interested in the actual solution \boldsymbol{x} but only in some summary statistic of it, such as, say, the sum of its entries. They also stress that, in situations like these, quantum algorithms may work much

faster than their classical counterparts and move on to argue that any possible quantum computation could be understood as an instance of linear equation solving. This argument, in turn, would either imply that classical computers should be able to efficiently simulate quantum computations (which they are not) or that quantum computers have an inherent linear equation solving advantage over classical computers. As the latter has not yet been conclusively demonstrated in general, Childs is critical about hardness results such as this one, and questions their role as a means of arguing for quantum supremacy in applied settings.

Overall, Aaronson concedes that the HHL algorithm and its decedents constitute theoretical as well as practical progress in quantum computing. At the same time, he emphasizes the need to carefully examine conditions or assumptions required for a quantum algorithm to achieve exponential speedup. Especially if effort required for in- and output or pre- and post-processing of data are factored in, many quantum computing solutions turn out to be only polynomially faster than classical algorithms. As Tang's former supervisor, Aaronson certainly knows what he is talking about. We therefore recall that Tang famously found a novel classical algorithm which—after appropriate pre-processing steps—works only polynomially slower than its quantum counterpart. From this, she concluded that state preparation assumptions in a quantum computing solution need to be matched against classical pre-processing assumptions before any claims can be made about significant quantum speedup.

To add further credence to this conclusion, we mention another example from our own experience. Recall that reports such as [209] or [226] describe classifiers which involve quantum subroutines for computing approximate inner products between quantum state vectors. However, if "approximate" inner products between high dimensional vectors are all that is asked for, we note that there exist classical methods which—after appropriate pre-processing steps—can compute them in $\mathcal{O}(1)$, i.e. in a time independent of the data dimensionality [262]. Especially in cases where inner products need to be evaluated repeatedly, e.g. in nearest neighbor searches, the costs for pre-processing quickly amortize and potential quantum advantages may become less significant. Granted, techniques such as in [262] appear not to be widely known but they exist and should be taken into account by quantum machine learning researchers.

The second caveat we mentioned in the introductory paragraph has to do with the common assumption that universal quantum computers are taken as a given when developing quantum computing algorithms. Yet, given the current state of technical developments, this is a far reaching assumption as many of the expected capabilities of universal quantum computers can not be realized by current NISQ era devices. Hence, to asses the actual technical feasibility of proposed quantum machine learning approaches, we have to check them against a catalogue of criteria. Important such criteria will be discussed in the following.

5.1 Hardware Requirements

When quantum algorithms are designed, properties of the hardware at hand are usually not taken into account. This is not surprising, since abstracting the actual hardware away is at the core of computer science. Fundamental results about complexity, data structures, and algorithms rely on theoretical models of computation, e.g., Turing machines and other types of automata. This is clearly sufficient when the correctness of an algorithm, its asymptotic runtime, or bounds on its memory consumption are being analyzed. In these scenarios, it is not necessary to bother about the clock rate of a processor, the amount of available main memory, or the data width of some bus architecture. In fact, theoretical machines possess an infinite amount of memory. Nevertheless, when algorithms have to run on real computers, all these constraints materialize and have to be considered for implementation. Algorithms can be equivalent under a theoretical model while their actual number of consumed clock cycles on a specific processor can differ by a large amount. Hence, one of the algorithms would consume far more energy and they stop being equivalent in appliances where a steady energy supply cannot be guaranteed, e.g., in the context of autonomous systems. Moreover, without specific tuning of an algorithm, state of the art GPUs cannot unfold their potential and many of their computational capabilities will lie idle [263]. Some of these practical considerations have found their way back into the theoretical community. One striking example are cacheoblivious algorithms [264], where the existence of a memory hierarchy is introduced to the theoretical analysis. Moreover, the explicit consideration of hardware properties also found its way to machine learning, e.g., machine learning models which are designed to rely only on integer arithmetic [265].

While the explicit incorporation of computational architectures has led to various successes in classical computer science, similar ideas are rather unknown in the quantum machine learning literature: When QML methods are derived, there is no treatment of particularities of the underlying quantum compute architecture. Authors do neither discuss whether, e.g., the usage of a Toffoli gate in a quantum circuit is problematic, nor is it discussed if using $\mathcal{O}(\log n)$ ancilla qubits leads to a practical relevant method. Especially because practical quantum computing hardware is still in its infancy and resources are scarce, real world constraints must be considered in order to understand if a specific QML algorithm is viable.

5.1.1 Quantum Memory

Some low-end classical processors like microcontrollers are not equipped with a floating point unit or units for processing wide single instruction multiple data vector operations. When an algorithm designer assumes the availability of such functional units, the algorithms are unlikely to run satisfactorily on other processors. The same happens in the design of quantum algorithms: Various QML methods require the availability of some sort of quantum memory. Given the current technical state of the art, the existence of an exact quantum memory is at least questionable due to the no-cloning theorem [88]. The theorem asserts that there is no

unitary U for which two non-orthogonal quantum states $|\psi\rangle$ and $|\phi\rangle$ exist, such that

$$U(|\psi\rangle\otimes|x\rangle) = |\psi\rangle\otimes|\psi\rangle$$
 and $U(|\phi\rangle\otimes|x\rangle) = |\phi\rangle\otimes|\phi\rangle$

whereas $|x\rangle$ is some arbitrary state that we want to overwrite with the content from $|\psi\rangle$ or $|\phi\rangle$. In short, we will not be able to construct a unitary operator that allows us to copy the content of arbitrary qubit registers $|\psi\rangle$ and $|\phi\rangle$ to a target register $|x\rangle$. It is important to understand that the no-cloning theorem prohibits the existence of a quantum computing system in which $|x\rangle$ is a working register and $|\phi\rangle$ and $|\psi\rangle$ are states which are stored in some type of quantum memory—making it effectively impossible to load data from the memory or to write data into the memory.

Nevertheless, a whole line of QML research relies on this concept, including quantum recommendation systems [112], quantum deep convolutional neural networks [266], quantum gradient descent, methods for solving linear systems, and methods for ordinary least squares regression [267]. The primary building block of these approaches is the QRAM [268]. While being frequently cited, fundamental issues like the no-cloning theorem are not discussed in the paper that introduces QRAM. Instead, the authors refer to "quantum memory elements" without describing how they can be implemented. Moreover, a three-level quantum system is required to realize the routing in the so-called bucket-brigade memory architecture. While this is not a problem per se, commercial quantum computers with hardware for ternary quantum states are not available. Hence, rendering algorithms which rely on these concepts as not viable. Finally, 2^n memory slots are assumed to be available, where n is the number of bits in the address register. Clearly, in the light of currently available hardware and under consideration of vendor roadmaps, the availability of a large number of qubits is not given at the time of writing.

Physical implementations of superconducting qubits reside on the chip at fixed locations and are connected via a well-defined pattern, the so-called *connectivity structure*. Any specific structure originates mainly from practical considerations. In IBM Q Systems, the predominant connectivity is known as *Heavy Hexagon* structure, shown in Fig. 4 (a). These low-degree graphs are designed to minimize the possibility of frequency collisions and optimize the hardware performance within superconducting qubit architectures. The larger the number of neighbors of a qubit, the more frequencies are required to realize two qubit gates using cross-resonance interaction. A specialized family of error correcting codes for this structure exhibits outstanding properties with respect to decoding and frequency collisions [269].

Beside its favorable properties for error correction, the connectivity of a quantum gate processor has direct impact on the depth of actual quantum circuits. To see this, one has to consider the transpilation of user specified circuits. During transpilation an input circuit is compiled to a sequence of native gates such that all operations agree with the connectivity structure and noise properties of a specific quantum processor. For any non-trivial circuit, a series of transformations must be conducted to make it compatible with a given target device, and optimize them to reduce the effects of noise on the resulting outcomes. Due to high-dimensional noise distributions and limited gate sets, rewriting quantum circuits to match

hardware constraints and optimizing for performance can be far from trivial. Quantum circuit compilation can have a non-linear control flow and exhibits complex branching patterns. Most importantly, it encompasses the decomposition of gates involving three or more qubits into 2-qubit gates. Clearly, the heavy hexagon structure is pairwise and hence contains no connection between three or more qubits. As a direct result, an apparently "shallow" quantum gate circuit, consisting of a single unitary operation among 10 qubits, can thus eventually exhibit a high depth. On the other hand, a high depth requires long decoherence and dissipation times of the system, and thus, might not be viable.

5.1.2 Qubit Connectivity

To understand the impact on QML methods, let us consider quantum kernel machines from Sec. 4.3.7. The power of the quantum kernel comes from the fact that data is mapped into a feature space whose dimension is exponential in the number of data dimensions. It is important to understand that the dimensions of the feature map shall not be independent of each other. In other words, all qubits that underlie the feature space representation shall be entangled with each other. For quantum kernel machines [218], this is realized by considering the following unitary:

$$U_{\phi(oldsymbol{x})} = \expigg(-i\sum_{S\subset [n]}\phi_s(oldsymbol{x})\prod_{v\in S}\sigma_z^vigg)$$

Here, $\phi_s(x)$ denotes a data dependent angle. The expression $S \subset [n]$ denotes all possible subsets of n qubits. We hence conjecture that $U_{\phi(x)}$ creates an entanglement between all qubits. However, we know that the process of transpilation will break this down into 2-qubit gates, and thus, a very deep and most likely not viable circuit. The authors of [218] are, of course, aware of this particularity and propose to restrict the subsets S to those of size two, hence, taking the actual hardware structure into account. While this solution is perfectly reasonable, it greatly limits the expressiveness of the resulting kernel, and thus the accuracy of any QML method that relies on this kernel construction.

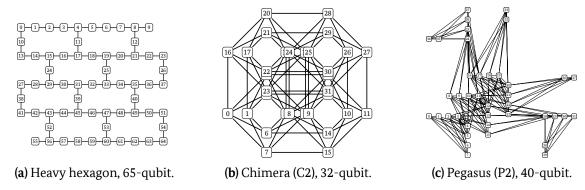


Figure 4: Qubit connectivity overview.

In case of the D-Wave quantum annealer, the *Chimera* and *Pegasus* structures define the qubit neighborhoods. They arise from considerations about a high inter-qubit connectivity. Formal definitions in terms of edge generating functions of the graphs shown in Figs. 4 (b) and (c) together with a brief discussion of their graph theoretic properties are provided in [270]. While connectivity seems to be a rather technical property of a quantum annealer, it has an immediate impact on the problems that can actually be solved on that chip. More precisely, the structure imposes constraints on the sparsity pattern of the QUBO matrix that encodes the problem that has to be solved. In Chimera, the qubits can couple to up to $\delta=6$ other qubits. In Pegasus, each qubits can couple to up to $\delta=15$ other qubits. This implies that each row of the QUBO matrix may only contain δ non-zero entries. Moreover, those entries must reside at fixed positions within each row. Revisiting the insights from the previous section, many QML methods that are based on quantum annealing require a completely dense QUBO matrix, e.g. the clustering method presented in [73]. There, the dimension of the QUBO matrix itself stems from the kernel matrix between all pairs of data points, and hence is dense.

Limited connectivity can be mitigated by introducing so-called chains into the problem's Hamiltonian. That is, the optimization problem is rephrased such that multiple qubits are enforced to take on the same value in the system's ground state, thus effectively increasing the fan-out of that qubit while reducing the total number of available qubits. To understand the consequences, consider the D-Wave Advantage 4.1 Quantum Annealer with 5627 qubits. When a complete dense QUBO matrix is required, chains of 17 qubits are introduced which results in an effective dimension of 177. The maximum dimension of the problem is hence reduced by orders of magnitude, severely reducing the number of data points that can be processed by the aforementioned kernel clustering approach.

5.1.3 Qubit Count

The number of qubits in real world systems is ever increasing. At the time of writing, at most 127 qubits are available for quantum gate devices, e.g., the IBM Eagle processor, and at most 5627 qubits for quantum annealing, e.g., D-Wave Advantage 4.1.

As mentioned in the previous section, the problem structure has a direct impact on the number of qubits that are actually available to a quantum annealer. When the problem dimension exceeds the number of available qubits, the problem is too large to be embedded in its entirety on the quantum processor. Instead, the problem is sequentially decomposed into smaller problems, each of which can fit on the quantum processor, with variables selected by highest impact on the objective function. Clearly, variables cannot be selected solely by maximal impact on the objective function, since those qubits may not be connected at all. In order to represent a local structure of the problem, traversal techniques such as breadth-first (BFS) or priority-first selection (PFS) can capture features that represent local structures within a problem. Instead of selecting the working set of qubits via their impact, *cutset conditioning*

5

can be applied. That is, the values of a subset B of variables (the "cutset") are fixed, effectively partitioning the set of variables $[n] = A \cup B$ with $A \cap B = \emptyset$. Thus, the QUBO problem $\operatorname{argmin}_{s \in \{-1,+1\}^n} Q(s) = s^\intercal Q s + q^\intercal s$ becomes

$$\operatorname*{argmin}_{\boldsymbol{s}_A \in \{-1,+1\}^n} Q(\boldsymbol{s}_A, \boldsymbol{s}_B)$$

Finding the optimal setting for s_A while values in s_B are fixed breaks the problem into separate components that can be solved independently. When A is chosen carefully, each independent problem is small enough to be solved on the quantum annealer. This procedure has to be repeated multiple times, choosing different subsets A. The general procedure is outlined in [271].

QUBO problems can indeed be solved on quantum gate computers, too. Based on the Variational Quantum Eigensolver [85], the Quantum Approximate Optimization Algorithm [84] provides specialized variational circuits, designed for addressing combinatorial optimization problems. The underlying variational forms have later been improved, yielding the Quantum Alternating Operator Ansatz [272]. As opposed to the quantum annealers with fixed connectivity, QAOAs dimension does not degrade as a function of the QUBO density. However, the number of qubits in contemporary quantum gate systems is lower than in quantum annealers. Thus, the same splitting techniques can be applied. Hence, larger optimization problems can be addressed even in case of a relatively low number of qubits.

For non-iterative or non-variational algorithms, it is in general not possible to split any quantum gate computation into smaller parts—mostly because entanglement cannot be maintained among the sub-problems. A reasonable direction can be the minimization of required qubits. However, the number of required "data" qubits or "input" qubits is problem specific. Methods for embedding high-dimensional points into a lower-dimensional manifold are available aplenty, but due to their highly approximate nature, it is unclear if the resulting embedded points still contain enough information. Assuming that the data qubits are fixed, the other major source of qubit utilization are so-called *ancilla qubits*. In a quantum computation, there is no way to deterministically construct a prescribed state unless one is given access to qubits whose original state is known in advance.

As an example of this concept, consider a matrix M whose singular values are upper bounded by 1. Let us denote the singular value decomposition of M by RDV. Then, the matrix

$$U = \begin{bmatrix} M & R\sqrt{I - D^2}V \\ R\sqrt{I - D^2}V & -M \end{bmatrix}$$

is unitary. This can be verified by observing that $UU^\dagger=I$. Note, however, that the dimension of the matrix has been doubled. Interpreting M as an operator that shall be applied on some state $|\psi\rangle$, we have to extend the state by one additional qubit $|0\rangle_a$, such that $U(|\psi\rangle\otimes|0\rangle_a)$ applies M to $|\psi\rangle$. A generalization of this idea is the linear combination of unitaries [273]. Unitaries are not closed under summation. However, a sum of m unitary matrices can be extended to an unitary operator by using $\log m$ ancilla qubits.

Ancilla qubits appear frequently in quantum machine learning [113, 114, 234]. However, care must be taken when the number of ancillas is provided in an asymptotic manner, e.g., [117, 274]. Given n input qubits, $\log n$ ancilla qubits can be reasonable. Nevertheless, $\mathcal{O}(\log n)$ may contain large constants, such that the number of required ancillas actually exceeds the number of available qubits by a large margin.

5.1.4 Quantum Circuit Depth

With the currently available NISQ technology, the result of a quantum computation on a physical device may deviate significantly from the expected outcome. A fundamental reason is that the quantum computer, despite all technical efforts, is not perfectly isolated and interacts (weakly) with its environment. In particular, there are two major effects of the environment that can contribute to computational errors: dissipation and decoherence in the sense of dephasing [275, 276]. Dissipation describes the decay of qubit states of higher energy due to an energy exchange with the environment. Decoherence, on the other hand, represents a loss of quantum superpositions as a consequence of environmental interactions. Typically, decoherence is more dominating than dissipation. In current quantum gate devices, these effects are captured by the quantities T1 and T2. Both can be interpreted as decay constants. That is, the probability that a qubit will stay in state $|1\rangle$ after time t is proportional to $\exp(-t/T1)$. Similarly, T2 indicates the decay constant for which an initial state $|+\rangle$ will evolve into an equal classical probabilistic mixture of the $|+\rangle$ and $|-\rangle$ states, so that one can no longer confidently predict the state. That is, it's the autocorrelation time after which the initial and final states become uncorrelated.

Decoherence effectively limits the length of the longest computation and thus the allowed depth of a quantum circuit. In other words, the more operations have to be performed on a qubit register, i.e. the deeper a quantum circuit has to be, the likelier it is that decoherence will happen. As explained in Section 5.1.2, transpilation will replace any unitary operator which acts on 3 or more qubits into a series of one-qubit and two-qubit gates. This implies that a single high-order operator might result in a circuit that is actually too deep to be executed by a system. Despite decoherence and dephasing constants T1 and T2 being available for most quantum processors, they are not utilized to adapt quantum machine learning methods to the underlying hardware and to account for the fact that the maximum computation time is limited.

Quantum devices based on adiabatic evolution are affected by decoherence as well [277]. The direct effect is a shorter usable annealing time and thus the time until a solution to the user specified optimization problem must be delivered. To see why this is significant, one has to recall that simple optimization problems can already be solved by classical computers. Thus, we are mostly interested in quantum solvers for hard optimization problems. However, hard problems, e.g., measured via the Hamiltonian's spectral gap, require a longer computation time—which is prohibited by the upper bound on the annealing time.

5.2 Algorithmic Requirements

We have discussed a series of pitfalls that arise in the design of quantum machine learning algorithms due to particularities of the quantum hardware. Of course, issues can also manifest solely from algorithmic aspects. In classical computing, efficient data structures and data types with a sufficient precision must be used. Components like pseudo random number generators and algebra sub-routines must be reliable and numerically stable. Parallel code has to be checked for race conditions and other problems that might emerge due to concurrency.

In quantum machine learning, similar issues can show up. In what follows, we discuss pre- and post-processing of inputs and outputs, stochasticity of the quantum computation, and noise sensitivity.

5.2.1 Data Pre- and Post-Processing

The question of how to encode data for training or inference with quantum machine learning methods so as to benefit from quantum advantages still awaits a final answer—if possible at all. Expected advantages, i.e. quantum speedup, may not emerge in practice if data preprocessing would require exponential efforts.

The root of the pre-processing problem lies in the preparation of states. To see this, consider the n-qubit state

$$|\psi\rangle = \sum_{j=0}^{2^n-1} \sqrt{p_j} |j\rangle$$

where $|j\rangle$ denotes the computational basis state that corresponds to the binary encoding of the integer j, and p_i is an arbitrary but fixed, classical probability mass function over $\{0,1\}^n$. Clearly, various types of data can be encoded into that representation: e.g., p_i could represent

- 1. (normalized) intensity values of pixels in an $2^{n/2} \times 2^{n/2}$ image,
- 2. term frequencies over some text corpus with a vocabulary of size 2^n ,
- 3. or the probability of observing a traffic jam on a specific set of street segments.

Related encodings indeed appear in the literature [278, 279, 280]. In that representation, a greyscale image in 4K resolution (3840 x 2160 pixel) would occupy $\lceil \log_2(3840 \times 2160) \rceil = 23$ qubits. It is, however, important to understand that preparing the state $|\psi\rangle$ requires resources that are proportional to the dimension of the Hilbert space of the underlying qubit register. In the above example n=23 was small enough such that processing 2^n values is feasible. If n is large, we cannot simply prepare $|\psi\rangle$. Even if a quantum algorithm would yield an exponential improvement over the best known classical algorithm given input state $|\psi\rangle$, it would not be possible to gain any advantage due to the complexity of preparing the initial state.

By assuming efficient integrability of an input probability distribution p_i , Grover and Rudolph explain how to efficiently generate the quantum state in a coherent fashion such that subsequent processing can be conducted on a quantum device [154]. However, their assumptions are not satisfied in most setups.

Recently, a classical model of computation was introduced which assumes that we can efficiently sample instances from a data set, where the probability to draw a specific data point is proportional to its ℓ_2 -norm. It can be shown that this is a natural analog to quantum algorithms that assume efficient state preparation of classical data. Based on this reasoning, classical versions of quantum algorithms for principal component analysis [281] and nearest-centroid clustering [198] can be devised. The runtimes of these classical algorithms are only polynomially slower which suggest that the exponential speedups of their quantum counterparts are an artifact of state preparation assumptions [282].

The same issues indeed arise when we want to read the result of a quantum computation. When the full Hilbert space representation of the state encodes our desired outcome, an exponential number of results has to be measured. Thus, even if the quantum computation itself delivers an exponential speedup, this advantage will vanish when all probability amplitudes of the final state have to be estimated.

5.2.2 Statistical Error

Statistical learning theory (SLT) is an integral part of machine learning. It may be the core distinguishing property between machine learning and plain numerical optimization. The theory itself is universal and independent of the underlying compute architecture. Through the eyes of statistical learning theory, there is nothing special about quantum machine learning. QML has to obey the laws and limitations predicted by SLT. At a first glance, this does not sound harmful. Nevertheless, some QML algorithms claim speed-ups over classical methods which can easily be falsified by basic insights from SLT. In what follows, we explain the basic reasoning behind this.

Assume that we try to find a function f such that the expected loss ℓ , also known as risk, is minimized:

$$R(f) = \mathbb{E}_{\mathbb{P}}[\ell(Y, f(X))]$$

Here, X is a random variable that represents the observed data or features, and Y is a random variable representing the class label (or regression target). The pairs (X,Y) follow the distribution \mathbb{P} . In practice, \mathbb{P} is indeed unknown, but a data set \mathcal{D} that contains samples from \mathbb{P} is available—allowing for the computation of the *empirical risk*:

$$\tilde{R}(f) = \frac{1}{|\mathcal{D}|} \sum_{(x,y) \in \mathcal{D}} \ell(y, f(x))$$

Let now \mathcal{H} denote a machine learning method, e.g., \mathcal{H} may contain all decision trees, all support vector machines with the same fixed kernel, or all neural networks with the same fixed

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network structure. Under rather mild assumptions, $\mathcal H$ contains a minimizer of $\tilde R$ [283]. Let us denote the empirical risk minimizer by $\tilde f=\arg\min_{f\in\mathcal H}\tilde R(f)$. The generalization error (also known as estimation error) can be expressed as distance between both minimizers in terms of R:

$$R(\tilde{f}) - \inf_{f \in \mathcal{H}} R(f) \leq 2 \sup_{f \in \mathcal{H}} |\tilde{R}(f) - R(f)| \leq \Theta\bigg(\sqrt{\frac{c(\mathcal{H}) - \log \delta}{|\mathcal{D}|}}\bigg)$$

where Ω denotes asymptotic equivalence and $\delta \in (0,1)$ is a probability. The last inequality holds with probability $1-\delta$ over the choice of the data set. Moreover, (\mathcal{H}) denotes a measure of complexity or expressivness of \mathcal{H} (e.g., VC dimension or Rademacher complexity [283]). When \mathcal{H} and δ are treated as fixed constants, the estimation error is upper bounded by $\Theta(\sqrt{1/|\mathcal{D}|})$. Hence, in order make use of all available data, any additional error that is introduced by a QML method shall not be larger than $\mathcal{O}(\sqrt{1/|\mathcal{D}|})$ —no matter if it is generated by an approximate optimization, measurement noise, or any other type of algorithmic error. When we denote the learned function that incurs an additional algorithmic error by \tilde{f}_{ε} , we have

$$R(\tilde{f}_{\varepsilon}) - R(\tilde{f}) = \Theta\big(\sqrt{1/|\mathcal{D}|}\big) + \underbrace{\big(\tilde{R}(\tilde{f}_{\varepsilon}) - \tilde{R}(\tilde{f})\big)}_{\text{Algorithmic error}} \; .$$

While this insight might not sound significant, it can be used to prove that quantum learning algorithms must have at least polynomial runtime in the dimension of the training data, and therefore cannot achieve exponential speedups over classical polynomial time machine learning algorithms—even when special data storage like QRAM is assumed [284].

To understand the impact, consider quantum least squares (QLS) regression [267]. The QLS algorithm is approximate. Its output is a state $|\hat{w}\rangle$ while the state that contains the true regression weights are $|w\rangle$. For the error, we have $||\hat{w}\rangle - |w\rangle| \leq \gamma$. According to [179], the runtime is $\mathcal{O}(\kappa^c \gamma^{-\beta} \operatorname{polylog}(|\mathcal{D}|))$ for some $c, \beta > 0$ and data condition κ . According to the reasoning above, the algorithmic error γ should not exceed the statistical error, since otherwise we would waste parts of the available data. We thus set $\gamma = |\mathcal{D}|^{-1/2}$. In this case, the runtime becomes $\mathcal{O}(\kappa^c |\mathcal{D}|^{\beta/2} \operatorname{polylog}(|\mathcal{D}|))$, and thus, polynomial in the number of data points—an exponential speedup over classical least squares cannot be justified.

Statistical effects also affect the accuracy of the readout. A specific number of readouts is required to obtain the result of the quantum computation with a desired accuracy. Using techniques developed within the field of quantum metrology, it is often possible to achieve a precision that scales as 1/m, where m is the number of readouts [285].

In the quantum regression setting, this implies that from state $|\hat{w}\rangle$, we can obtain \hat{w} by sampling the circuit m times such that $\||\hat{w}\rangle - \hat{w}\| \leq \Omega(1/m)$. By the same reasoning as above, this means that we have to obtain $m = \sqrt{|\mathcal{D}|}$ shots for the readout in order to not discard information from our data set \mathcal{D} . No matter how fast our regression algorithm is, the number of readouts will be a polynomial in the number of data points.

5.2.3 Noise Robustness

NISQ devices suffer from various types of noise, e.g., measurement noise or gate errors. These uncertainties arise in addition to the general probabilistic nature of quantum computation. One might wonder why we should even care about additional stochastic effects, instead of treating them as part of the probabilistic computation itself. Correction for measurement noise exists such that there is no direct practical problem with it. However, the error induced by noisy quantum gates is notoriously harder to characterize. One reason is crosstalk in superconducting quantum processors, e.g., performing operations on two pairs of qubits simultaneously can induce an error if the qubits are physically close on the chip. While crosstalk and other sources of noise are hardware related problems, it is up to the algorithm whether slight variations during the computation will lead to a completely different result.

As an example, consider adiabatic quantum computation. AQC is a stochastic procedure there is no guarantee that the outcome, measured after a finite amount of time, is the ground state of the target Hamiltonian. It can well be that the adiabatic evolution ends up in an excited (sub-optimal) state. The conditions under which the adiabatic process stays in the ground state until the end of the evolution are well understood [63, 286]. The key property is the Hamiltonian's spectral gap, that is, the difference between the ground state energy and the energy of the first excited state. Moreover, there is sufficient theoretical evidence that spectral gaps of random problem instances are likrly to become super-exponentially small [287, 288]. Thus, the time required for adiabatic evolution to remain in the ground state is longer than the time required for a classical brute force search through the full state space. It turns out that exponentially small gaps appear close to the end of the adiabatic evolution for large random instances of NP-complete problems. This implies that, unfortunately, adiabatic quantum optimization fails: The system gets trapped in one of the numerous local minima. Since the adiabatic evolution of NISQ devices is subject to minor perturbations, variations can affect the spectral gap and hence alter the success probability of the computation. Thus, QUBO formulations should take this problem into account. Nevertheless, experimental results show that the QAOA algorithm seems to be more robust against exponentially small spectral gaps [289].

Another line of research addresses the formal robustness verification of quantum machine learning algorithms against unknown quantum noise. As mentioned above, noise may arise from various sources, rendering a full characterization of the noise distribution almost infeasible. In [290], the authors relate noise robustness to adversarial examples and present an analytical robustness bound to assess the robustness of quantum classification algorithms. Their results are, however, not evaluated on real quantum computing devices. At the time of writing, evaluating the noise robustness of quantum machine learning algorithms is still an open problem.

6 Summary and Outlook

After decades of worldwide research, artificial intelligence has finally made considerable strides with respect to capabilities and applicability. As a consequence, weakly intelligent cognitive systems are now increasingly deployed in practice and appear in more and more areas of our daily- and professional lives. An interesting general observation in this context is that accelerated progress and noteworthy recent accomplishments in this area have mainly been driven by modern large scale machine learning.

Machine learning is a branch of artificial intelligence that deals with adjusting the parameters of software agents in a training process such that they can develop cognitive capabilities and problem solving skills. Put differently, the recent performance boost in artificial intelligence is mainly due to systems which analyze large amounts of task specific training data in order to learn an intended input-output behavior. The tacit assumption behind this idea is that any real world data must have been produced by some generally unknown process or mechanism and that this process or mechanism can be modeled mathematically. This translates to the assumption that there exists a suitable parameterized function whose parameters can be automatically adjusted such that the input-output behavior of that function reflects or mimics the characteristics of the given training data.

There exist numerous possible machine learning models (mathematical formalizations of a given application scenario) as well as numerous learning algorithms (mechanisms to fit a given model to a given set of training data). Traditionally, models and algorithms for their training were chosen with respect to the task at hand, often specifically tailored towards a specific setting. However, the dominating trend over the past decade and arguably the main reason behind the success of modern machine learning has been to work with domain agnostic models and training algorithms. In particular, deep learning has proven to be tremendously successful in a wide range of computational intelligence problems.

Deep learning involves deep (and wide) artificial neural networks which are composed of millions of artificial neurons which interact over even more artificial synapses. Since such networks thus come with billions of adjustable parameters (synaptic weights and activation function bias values), they provide very flexible general models which can be trained with (variants of) the backpropagation algorithm to learn a desired cognitive skill. However, in order for this to happen reliably, deep neural networks need to be trained with vast amounts of representative training data. Due to the number of data to be processed and the number of parameters to be adjusted, the training of modern deep networks is a formidable task that requires considerable computational resources in order to happen within reasonable time. Indeed, breaking down the efforts involved in training state of the art systems such as, say, OpenAI's GPT3 for text analysis- and synthesis reveals training times of several hundred GPU years which, to be possible within a matter of days, necessitates the use of dedicated compute clusters. In other words, state of the art machine learning has reached a point where its practical feasibility and success are conditioned on access to high performance computing hardware.

Given this state of affairs regarding the computational needs of current machine learning systems, it is not surprising to find that more and more researchers are beginning to look at quantum computing as a tool to be deployed at different stages of the machine learning pipeline. The basic observation is that quantum computing, too, has made considerable strides over the past decade and is now becoming practical. While quantum computing will not eliminate the dependency on special purpose hardware, it promises significantly faster computations across a wide range of scientific or industrial applications. With respect to machine learning, the ambition is to harness potential quantum speedup especially for the training of ever more complex systems.

The advantages of quantum computing over classical digital computing are rooted in the fact that it exploits quantum mechanical phenomena for information processing.

While digital computers operate on bits which are in one and only one of two possible states, quantum computers operate on qubits. These are logical interpretations of physical two-state quantum systems which exist in a superposition of two basis states. The state space of a qubit thus consists of infinitely many states so that a qubit can carry more information than a classical bit. While the mathematics that describes the behavior of classical bits is Boolean algebra, the mathematics that describes the behavior of qubits is complex linear algebra. Qubits can be represented as two-dimensional, complex-valued unit vectors that are formed as a linear combination of two distinguished, linearly independent, orthonormal basis vectors. The squares of the norm of the coefficients of a qubit state vector are called amplitudes and have an important probabilistic interpretation. If a measurement is performed on a qubit, it will decohere, i.e. loose the property of superposition, and collapse to either one of its basis states. The probability of the measured state to be the first basis state corresponds to the squared norm of the first coefficient and the probability of measuring the qubit in the second state corresponds to the squared norm of the second coefficient. Computations involving operations on- and subsequent measurements of qubits are therefore probabilistic rather than deterministic as in the case of classical bits. However, just as classical bits can be combined into bit registers, quantum bits can also be combined into qubit registers. Adding a single qubit to a register increases the dimension of the register's state space by a factor of two and so a quantum register of n qubits exists in a superposition of 2^n basis states. Another crucial difference to classical computing is that qubits can be entangled. Whenever two or more qubits are entangled, their individual states cannot be measured separately but a measurement of an entangled qubit will also determine the (combined) state of the others.

The interplay of these phenomena gives rise to the potential supremacy of quantum computing over digital computing: on a quantum computer it is possible to work with only n qubits in order to perform computations in an 2^n dimensional space. This is of considerable interest for combinatorial optimization which deals with exponential search spaces and plays a crucial role in a branch of artificial intelligence known as problem solving, as well as in general parameter selection problems in certain machine learning techniques.

Adiabatic quantum computers such as produced by D-Wave are especially tailored to-

wards solving an important class of combinatorial optimization problems, namely QUBOs, which occur in machine learning contexts such as data clustering, classifier boosting, or support vector machine training. The basic idea in adiabatic quantum computing is to express a given problem as an energy minimization problem and to devise an energy function which is known to attain its minima at the unknown solutions to the problem. The energy of a quantum system is characterized in terms of a Hamiltonian operator whose eigenvalue spectrum is the set of all possible outcomes when measuring the system's energy; its eigenstates reflect to which states of the system these energies correspond to. The lowest energy state of the quantum system under consideration is also called the ground state of the corresponding Hamiltonian and, if the energy function for a given problem can be translated into a Hamiltonian operator, its ground state can be found via adiabatic quantum computing. Adiabatic quantum computing exploits a phenomenon summarized in the adiabatic theorem. It states that, if a quantum system starts in the ground state of a Hamiltonian which then gradually changes towards another Hamiltonian, the system will end up in the ground state of the resulting Hamiltonian. On an adiabatic quantum computer one thus operates with two Hamiltonians, a problem independent one and the problem Hamiltonian and gradually changes the former into the latter. This general idea therefore exploits a form of quantum tunneling, i.e. a quantum mechanical principle which states that quantum systems can tunnel through energy barriers. The latter constitutes an advantage over classical energy-based optimization where optimizers may get trapped in local minima and thus may fail to find optimal solutions to a problem. Moreover, for many problems of practical importance, adiabatic quantum computers can search the solution state space polynomially faster than classically possible. While this may not sound like much, for large problems it can make the difference between tractable and intractable.

At first sight, the way of thinking required for adiabatic quantum computing may look abstract and unusual. However, problem solving by energy minimization is a well established general paradigm in classical machine learning. In fact, the well understood Hopfield networks, a venerable type of neural networks, can be seen as a classical or digital analogue to adiabatic quantum computing. This is to say that anything that can be accomplished by running a classical Hopfield network on a digital computer can, in principle, also be accomplished on an adiabatic quantum computer. In this sense, the conceptual gap between the classical paradigm of Hopfield neural networks and adiabatic quantum computing is therefore rather narrow.

Quantum gate computing manipulates qubits in a manner that more closely resembles classical digital computing. On the hardware level, digital computers process information by manipulating sets of bits via atomic logical operations which are implemented in terms of so-called gates. In quantum gate computing, corresponding operations on qubits are realized via the application of quantum mechanical operators. Mathematically, these are unitary transformations so that any operation on qubits must be unitary too. Such operators typically act on one or two qubits at a time but can be sequenced or executed in parallel to form more complicated operations on sets of qubits. In analogy to classical computing, complex computational

units that are composed of individual quantum gates are called quantum circuits. From a logical point of view, the fundamental problem in quantum gate computing is therefore to devise quantum circuits that show an intended input/output behavior. This is generally a non-trivial problem and it is interesting to note that classical machine learning is increasingly seen as tool for quantum circuit design.

What makes quantum gate computing attractive from the point of view of machine learning is that, mathematically, quantum gate computing is nothing but applied complex linear algebra. Due to mechanisms such as superposition or entanglement, however, the linear algebraic operations that occur in quantum gate computing implicitly act on state spaces that are exponentially larger than those considered in classical computing. Since many common machine learning tasks involve linear algebraic operations on large amounts of high dimensional data vectors, it seems auspicious to try to encode such data in qubit state spaces and to try to leverage quantum advantages when computing with such states. Here, the general expectation is that quantum speedup will considerably accelerate corresponding learning processes or even allow to tackle hitherto intractable problems.

Indeed, in the words of Aaronson, there has recently been a "quantum machine learning mini-revolution" and the number of scientific reports on quantum circuits for certain general problems in machine learning has grown considerably over the past decades. Problems which have been considered in this context include basic linear algebra routines, regression, or classification. With respect to the latter, there exist different paradigms so that it is no surprise to find numerous proposals of quantum circuits for nearest neighbor classifiers, basic linear classifiers (perceptrons), ensemble classifiers, support vector machines, and neural networks.

Present day ideas for how to realize quantum machine learning routines on quantum gate computers commonly involve variational quantum computing algorithms or hybrid quantum-classical methods. These consider parameterized quantum circuits which may be designed manually or automatically but in either case consist of tunable quantum gates. The problem solved by hybrid quantum-classical methods therefore is the problem of adjusting the parameters of individual gates such that the circuit as a whole performs the desired computation with high probability. The basic structure of variational quantum algorithms consist of an outer loop run on a digital computer which manages the current estimates of the sought after parameters. In each iteration, these parameters are used to setup computations on a quantum computer whose outcomes are then measured. Given these measurement results, classical optimization techniques are then used to estimate improved parameters and the whole processes is iterated until the quantum circuit shows the desired input/output behavior within a given tolerance.

Given the technical capabilities of existing, present day quantum computers, variationalor hybrid quantum-classical algorithms are appealing because they have been found to considerably reduce the quantum computing resources (circuit depth, coherence time, qubit counts) that are required to successfully run a quantum machine learning method. Moreover, researchers interested in developing quantum neural networks often even consider

parameterized quantum circuits as a quantum computing analogue of classical deep neural networks. This is likely because deep neural networks as well as parameterized quantum circuits consist of basic computational units arranged in layered structures and because both involve classical optimization algorithms for parameter adjustment. However, one has to be careful with such analogies since there also are crucial differences. First of all, quantum gates in a quantum circuit realize unitary operators rather than non-linear functions as in the case of neural networks. Non-linear computations are vital for the general problem solving capabilities of neural networks but they can not be realized through unitary operators alone. Non-linear quantum computations either involve the generous use of ancilla qubits or state measurements. Second of all, internal states of a quantum circuit cannot be read out without destroying their quantum coherence, i.e. without loosing quantum aspects such as superposition. Classical neural network training based on error backpropagation does therefore not apply to parameterized quantum circuits and the use of variational- or hybrid quantum-classical algorithms for circuit optimization currently constitutes the only viable approach to this problem.

The ever faster growing literature on quantum machine learning and the many success stories reported therein seem to suggest that the application of quantum computing algorithms to computational intelligence problems could soon break into the mainstream. However, in the present era of noisy intermediate-scale quantum (NISQ) computing, overly enthusiastic expectation still need to be reigned in. This has mainly to do with the technical capabilities of present day quantum computers which often differ from the assumptions made by quantum algorithm developers.

First of all, quantum algorithm design considers properties of logical qubits rather than those of physical qubits. The former are the basic building blocks on which quantum algorithms operate, the latter are physical devices inside of a quantum computer which behave like two-state quantum systems. From the point of view of algorithm design, the focus on logical qubits is reasonable and mimics levels of abstraction in modern software development where most programmers need not worry about hardware details. Current quantum computers, however, have not yet reached the same level of maturity as digital computers but still exhibit certain limitations. As of now, pure logical qubits are therefore an idealization since they abstract away shortcomings of physical qubits in present day NISQ devices. These typically contain less than a hundred qubits, often suffer from limited coherence times, and are susceptible to low fault-tolerance due to internal fluctuations or measurement noise. Indeed, it still is challenging to technically create and manipulate quantum states and to maintain their quantum mechanical properties over longer periods of time. It is rather marvelous that this has become possible at all and experts expect that continuing technological progress will lead to better and more powerful devices. Lack of fault tolerance is a critical issue and a technological milestone will be the implementation of quantum error correction mechanisms similar to those used in digital computing. Here, it is interesting to note that errors often result from the application of less robust quantum gates. As of this writing, only rather small quantum circuits work really stable. Since adiabatic quantum computers do not involve quantum

gate computations, they can more reliably manipulate larger (physical) qubit systems. At the same time, their use is currently restricted to rather specific energy minimization problems so that present day devices are not as universal as quantum gate computers. Even though both paradigms are theoretically equivalent, the emulation of quantum circuits on an adiabatic quantum computer would require qubit connectivity structures that have not yet been realized.

While presently realizable qubit counts, qubit connectivity, dimensionality of quantum gates, and quantum circuit depths impose practical restrictions on the kind of quantum algorithms that can run successfully on current NISQ devices, there are further limitations regarding the feasibility of certain logical quantum computing concepts. We therefore emphasize the following additional points.

Second of all, it is often not immediately clear how to encode classical data such that it can be processed on quantum computers; neither may it be obvious how to decode measured qubit states into classical representations that would allow for meaningful downstream processing. This input-output problem is often abstracted away by quantum algorithm designers which, in turn, may have dire consequences with respect to claims about quantum speedup. For instance, a quantum algorithm operating on a quantum encoding of classical data might be exponentially faster than its classical counterpart operating on the classical data. However, if the effort for preparing quantum states that encode the classical data is itself exponential, then the apparent quantum advantage vanishes. Here, it is noticeable that designers of quantum algorithms often simply hypothesize universal quantum computers. Among others, these machines are supposed to come with quantum random access memories (QRAMs) which, similar to their digital counterparts, can hold data for processing. However, given present day technologies, such QRAMs are not yet possible and it is even questionable if they will ever exist in their truest sense. This is due to another quantum mechanical principle, the no-cloning theorem, which states that it is impossible to create independent identical copies of arbitrary quantum states. To the best of our current knowledge, and unless reliable quantum error correction becomes available, a QRAM would thus only allow for approximate repeated access to quantum states for processing. Quantum state decoding, too, can diminish quantum advantages. Even if a quantum algorithm can produce a quantum state representation for a sought after solution much faster than classically possible, the effort for measuring and reading the resulting state into classical memory could still be so substantial that any advantage disappears. Before claims as to the superiority of quantum algorithms can be made, it is therefore pivotal to factor in efforts for state preparation or measurements and to match these against pre- and post-processing efforts of efficient classical algorithms.

Third of all, as of now quantum computing is essentially bit level computing. That is, present day quantum algorithm design deals with the design of problem specific quantum circuits or problem specific energy functions. Contrary to classical computing, there are no abstract data structures, such as linked lists, binary trees, or heaps. Neither are there or control structures such as if-then-else-statements or for- or while-loops known from higher level programming languages and common classical programming patterns such as the use of

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variables, too, are not immediately possible on present day quantum computers. This is challenging in so far as it suggests that certain existing machine learning algorithms build around such constructs may, for the foreseeable future, not be realized on quantum computers. While there are increased efforts towards quantum compilers which automatically translate higher level programs into quantum circuits, these are still in their infancy. Also, while high level application programmings interfaces (APIs) for quantum computing are increasingly available (e.g. QISKIT, CIRQ, Forest, PennyLane), it is important to note that these are mainly tools for setting up quantum computing processes. That is, users of these tools still have to think on the linear algebraic or "state and operator" level of quantum computing and use the API only to implement vectors, matrices and their interplay. Another caveat with respect to these APIs is that they often allow for efficient digital simulations of quantum information processing. While this is undoubtedly helpful in the design stage of algorithm design, it can also be misleading as the simulated quantum processors typically are universal quantum processors. This, in turn, means that a quantum machine learning algorithm that works on a simulated quantum computer may not work on a currently existing physical quantum computer. A final caveat with regard to quantum algorithms for machine learning is that quantum computations which involve measurement steps are inherently probabilistic. This is of course well known to quantum computing practitioners but may be overlooked by novices. Any quantum computation must therefore be repeated several times and any result obtained this way has to be interpreted in terms of expectations rather than in terms of deterministic outcomes.

Fourth of all, while machine learning is a comparatively new scientific discipline, quantum machine learning is even newer. This has consequences for the best practices and standards in the field. To clarify what this may mean, we note that, in its first couple of decades, classical machine learning has gone through a verifiability and reproducibility crisis. Practical results tended to be reported, without disclosing implementation details, data collections or processing protocols, or experimental procedures in detail. Regarding the validity of claimed capabilities of a reported method, such omissions can make a considerable difference. It is, for instance, pivotal that training and testing of a machine learning system happen on independent data sets, because a low error on the training data does not imply that the trained system can generalize well. On the contrary, an exceedingly good performance on the training data is often a symptom of overfitting. Over time, these issues have been recognized by the machine learning community and, as a consequence, present practice for scientific publication is to require authors to provide their code, data, and experimental protocols. In the new field of quantum machine learning, this is not yet the case. Indeed, reading the corresponding scientific literature, it is noticeable that crucial details as to how a practical result has been obtained are often missing. In many reports on practical quantum machine learning, it often not even recognizable if, say, rigorous evaluation practices have been followed so that one may argue that present day quantum machine learning is experiencing a reproducibility crisis period. For now, this may be acceptable as the field is in its nascent phase; it is, however, important to keep in mind that the good performance of some of the currently reported methods may not scale or generalize to large or different application settings.

While all these caveats may dampen enthusiasm for the prospect of quantum machine learning, recent technological progress has been substantial enough to merit serious engagement with the topic. In other words, although quantum computers and quantum machine learning algorithms are not yet mature enough to impact the way machine learning happens in practice, it currently seems reasonable to expect that—due to considerable investments by institutional and industrial stakeholders—the underlying technology will continue to develop and improve ever more quickly. This, in turn, may soon lead to practically viable solutions and cause hitherto unexpected developments and disruptions. A foresight process which fathoms potential benefits and risks of quantum machine learning therefore seems appropriate.

In particular, potential risks related to the use of quantum machine learning have not yet received the same amount of attention as its benefits and corresponding reports are scarce. In other words, while ethics, reliability, trustworthiness, and safety of classical machine learning have by now been recognized as important topics, quantum machine learning has not yet been scrutinized in these regards. However, given the expected impact of quantum machine learning on capabilities and utilizability of artificial cognitive systems, it seems appropriate to asses potential security issues related to quantum machine learning. In upcoming reports, we will therefore investigate questions pertaining to the reliability or vulnerability of quantum machine learning systems. Further considerations include whether or not quantum machine learning allows for new kinds of attacks on- or new defense mechanisms for critical digital infrastructures. The general goal will be to assess quantum machine learning from the point of view of cybersecurity and to determine what kind of measures, if any, will be required in this regard.

Glossary

Adiabatic quantum computing (AQC)

A quantum computing paradigm that is based on the adiabatic theorem and particularly tailored towards solving QUBOs; in adiabatic quantum computing, a qubit system gradually evolves from the ground state of a problem independent beginning Hamiltonian to the ground state of a Hamiltonian which models a given problem; the method thus uses quantum annealing or quantum tunneling for problem solving.

Adiabatic quantum optimization (AQO)

A synonym for adiabatic quantum computing.

Adiabatic theorem

A mathematical formulation of the quantum mechanical principle that, if a quantum system starts in the ground state of a Hamiltonian operator which then gradually changes for a period of time, the system will end up in the ground state of the resulting Hamiltonian.

Application phase (of an ML system)

The final stage in the practical development of a machine learning system; once the system has been trained and tested and was found to perform robustly and reliably, it can be deployed in practical applications.

Artificial Intelligence (AI)

A branch of computer science concerned with the design, development and deployment of technical systems or software agents that have cognitive capabilities such as text- or image understanding, planning, and decision making.

Backpropagation

A very widely used algorithm for the training of (feed-forward) neural networks; given an appropriate loss function, the algorithm essentially applies the chain rule of differentiation to adjust the parameters of all neurons in all layers of a neural net; put differently, the method performs gradient descend in a highly non-linear error landscape; there exist numerous variants of the original version of the algorithm, for instance, to automatically determine optimal step sizes for the descent procedure.

Basis states

States of quantum mechanical systems are described as linear combinations over basis states; for instance, the state of a qubit is represented as a two-dimensional, complex-valued unit vector that is a linear combination of two distinguished, linearly independent, orthonormal basis states.

Beginning Hamiltonian

A problem independent Hamiltonian operator used in adiabatic quantum computing; the beginning Hamiltonian is usually chosen such that its ground state can be easily prepared and used as the initial state for the adiabatic problem solving process.

Computational basis

A basis system used to express the state of a qubit system; the state of a system of qubits corresponds to a vector in a complex Hilbert space and the basis used to (numerically) express the entries of this vector can be chosen arbitrarily; the computational basis is often the most "natural" choice, i.e. a basis that suits the problem at hand or in which state vectors have a simple form.

Decoding problem

The problem of measuring or reading a quantum state which results from a quantum computation; quantum computing results usually need to be transferred to digital computers in order to enable further processing or analysis; if this transfer cannot be done efficiently, advantages due to quantum speedup can be lost.

Decoherence

A term referring to the loss of quantum coherence; if a quantum systems decoheres, it loses its quantum properties; decoherence of qubit happens through measurement or (unwanted) interactions with the outside world; qubit systems have thus to be kept as isolated as possible when performing quantum computations; if a qubit decoheres, it collapses to one of its basis states and henceforth behaves like a classical bit.

Digital annealer

A specialized digital computing device tailored towards solving QUBOs; digital annealers classically emulate the workings of adiabatic quantum computers; this requires efficient (hardware) implementations of classical optimization algorithms; compared to present day adiabatic quantum computers, digital annealers are much cheaper and much more energy efficient.

Eager learner

A machine learning system that trains a model in an offline training phase and considers a complete existing training data set for re-training and rather than only newly added data.

Encoding problem The problem of preparing quantum states that represent classi-

cal data; if classical data cannot efficiently be encoded in terms of quantum states, computational advantages due to quantum

speedup can be lost.

Ensemble learning Ensemble Learning methods combine different, typically sim-

ple (machine learning) models into a single stronger model that achieves better results than the individual member of the ensemble. Ensemble Learning allows, for example, for an average value

from the results of these different models.

Entangled state In quantum computing, this term refers to a state in which a sys-

tem of qubits can exists such that their individual states cannot be measured separately; rather, a measurement of an entangled qubit will also determine the (combined) state of the others; entanglement is a quantum mechanical phenomenon for which there is no classical analogue and which is essential for the inner

workings of quantum computers.

Federated learning In federated learning, each participant retains their own data and

contributes to a common learning process in which the system as

a whole benefits from the capabilities of each contributor.

Generalization The capability of a trained machine learning system to generalize

the predictive performance it acquired from analyzing training

data to previously unseen test data.

Ground state A term used to refer to the lowest energy state a quantum me-

chanical system can be in; mathematically the lowest energy of a quantum system corresponds to the bottom eigenvalue of its Hamiltonian operator and the system's ground state is given by

the corresponding eigenvector.

Hamiltonian operator A quantum mechanical operator acting on the state vectors of

quantum mechanical systems; Hamiltonians feature prominently in Schrödinger equations; their eigenvalues represent possible energy levels the corresponding quantum system can be measured

in.

Kernel machine A machine learning term used to refer to a wide class of machine

learning models; kernel machines are machine learning models

that make use of the kernel trick when processing data.

Kernel trick

A computational method used in kernel machines which allows for applying inherently linear methods to non-linear machine learning problems. Invoking the kernel trick involves two steps: 1) to rewrite a given machine learning method such that any occurrence of data vectors is in form of inner products and 2) to replace any computation of an inner product by the computation of an appropriate kernel function.

Lazy learner

A machine learning system that simply stores data and delays modeling until asked to make predictions; a simple example for this paradigm is a nearest neighbor classifier.

Logical qubit

A mathematical idealization of a physical two-state quantum system; logical qubits exist in a superposition of two basis states and can therefore represent more information than classical bits; they form the basic units in quantum computing.

Loss function

A criterion to measure how well a machine learning model and its current choice of parameters represent a given set of data; quality measurement is done by assigning to each prediction of the model the loss that occurs when a prediction deviates from the correct prediction; the more predictions deviate from the ground truth training data, the higher the output of the loss function.

Machine Learning (ML)

A branch of artificial intelligence that deals with adjusting the parameters of technical systems or software agents in a training process such that they can develop cognitive capabilities and problem solving skills.

Model class

A machine learning term used to refer to a parameterized family of functions; once a model class has been chosen for a machine learning problem, the task is to determine that model within the class that provides the best fit to a given set of training data; adjusting the model to the data at hand happens through machine learning algorithms which tune its parameters such that the model is in good agreement with the data; often, this happens by means of minimizing a loss function.

No-cloning theorem

A quantum mechanical principle that states that it is impossible to create independent identical copies of arbitrary quantum states.

Noisy intermediatescale quantum device (NISQ) A term used to describe existing, present day quantum computers which contain less than a hundred physical qubits and still suffer from limited coherence times and low fault-tolerance.

Oracle

A kind of function posited or required in many quantum computing algorithms; oracle functions can quickly verify if a supposed solution to a problem really is a solution; for instance, while the problem of finding the prime factors of 42 is somewhat difficult, it is easy to verify that 2, 3, and 7 are the solution; by the same token, it is also easy to verify that, say, 5 and 13 are not.

Overfitting

A type of machine learning modelling error that occurs when a model corresponds too closely to a given data set and is therefore likely to produce very different results for (slightly) different data samples.

Parameterized quantum gate

A kind of quantum gate that requires numeric parameters to define its actual function; parameterized quantum gates are of pivotal importance for variational quantum algorithms and for quantum machine learning; they allow for representing machine learning models via a single quantum circuit and for the use of optimization techniques to adjust the input/output behavior of the circuit such is agrees with given task specific data.

Phase gate

A single qubit quantum gate that modifies or shifts the phase of the quantum state of a qubit.

Physical qubit

A physical realization of a logical qubit, namely a device or apparatus that behaves like a two-state quantum system and forms a component of the hardware of a quantum computing system.

Problem Hamiltonian

A Hamiltonian operator used in adiabatic quantum computing; the problem Hamiltonian is designed in such a manner that its ground states represent the sought after solution to a given (combinatorial optimization) problem; in adiabatic quantum computing, a qubits system is prepared in the ground state of a problem independent beginning Hamiltonian which then gradually changes towards the problem Hamiltonian; this will cause the qubit system to end up in a ground state of the latter and thus solve the problem at hand. The design of problem Hamiltonians requires experience with (re)formulating QUBOs.

Quadratic unconstrained binary optimization (QUBO)

A kind of combinatorial optimization problem where the decision variables can only assume two values (typically either 0 and 1 or -1 and +1); many practically important subset selection or set bi-partition problem can be written as QUBOs but their solution is classically difficult as they are generally NP-hard.

Quantum approximate optimization algorithm (QAOA)

A variational quantum algorithm in which parameterized quantum gates are optimized to approximate adiabatic quantum computations.

Quantum bit (qubit)

A two-state quantum system; a qubit exists in a superposition of two basis states and can therefore represent more information than a classical bit; qubits are the basic unit of information in quantum computing or quantum information processing and are modeled in terms of vectors in a complex Hilbert space; one often distinguishes between logical qubits (which represent the mathematical essence of two-state quantum systems) and physical qubits (which are physical instances or realizations of two-state quantum systems); in other words, while there exists different physical manifestations of qubits, their essential (mathematical) properties are the same.

Quantum computing

A computational paradigm that harnesses the principles of quantum mechanics for information processing; exploiting quantum mechanical phenomena such as superposition or entanglement offers great computational power but also requires a different kind of algorithmic thinking than in classical digital computing.

Quantum gate computing

A quantum computing paradigm sometimes also referred to as the quantum circuit model of computation; in quantum gate computing, operations on qubits are realized via the application of quantum mechanical operators which are physically implemented as quantum gates; in analogy to classical computing, complex computational units which are composed of individual quantum gates are called quantum circuits.

Quantum gate

A basic computational unit within a quantum circuit; quantum gates are used to perform quantum mechanical operations on individual qubits or on systems of qubits.

Quantum information processing

A scientific discipline that deals with theory and practice of systems which process quantum information; an important aspect of quantum information processing is quantum computing; however, quantum information processing is sometimes seen to be more comprehensive than quantum computing because it also involves topics such as quantum communication or quantum sensing.

Quantum inspired computing

A term used mainly to describe classical algorithms whose design or operating principles are inspired by- or try to emulate quantum mechanical phenomena; for instance, a digital annealer digitally emulates the inner workings of an adiabatic quantum computer.

Quantum machine learning

A term mainly used to refer to the idea of incorporating quantum computing routines at different stages of the machine learning pipeline; in a much broader sense, the term is also used to refer to quantum inspired classical algorithms for classical data analysis, genuine quantum algorithms for classical data analysis, classical algorithms for quantum data analysis, and quantum algorithms for quantum data analysis.

Quantum supremacy

A term used to describe the fact that there exist problems which a quantum computer can solve but a classical computer can not in feasible time; in order to exhibit quantum supremacy with respect to a given problem, a quantum algorithm has to provide superpolynomial speedup over the best possible classical algorithm.

Quantum tunneling

A quantum mechanical principle that describes the fact that quantum systems can tunnel through energy barriers; quantum annealing or adiabatic quantum computing exploit this phenomenon.

Regularization

A mathematical technique or concept designed to prevent machine learning algorithms from overfitting by imposing constraints on model parameters that reduce the variance of a model.

Reinforcement learning

A machine learning paradigm tailored to the problem of learning "what to do when" in order to achieve a long term goal; reinforcement learning typically considers states an agent can be in, actions an agent can perform to transit to another state, and delayed rewards an agent has to maximize.

Superposition

A quantum mechanical principle which states that the sum (superposition) of two or more quantum states is another valid quantum state; superposition is often interpreted in the sense that a quantum mechanical system can be in many different states simultaneously. Upon measurement, the system will collapse to one of the superposed states.

Supervised learning

A machine learning paradigm where machine learning models are trained on annotated training data; the goal is to adjust the model parameters such that the model is able to map the given input data to the given output data with high fidelity; importantly, models trained in a supervised manner have to be evaluated on independent test data in order to verify that they can generalize to situations not contained in the training data; a common example of a supervised learning problem is classifier training.

Test phase (of an ML system)

A stage in the practical development of a machine learning system; given problem specific test data and a trained model, the model is evaluated on the test data using an appropriate performance measure; it is pivotal that the test data are independent of the data used to train the model, otherwise issues such as overfitting cannot be detected.

Topological quantum computing (TQC)

A still mainly theoretical concept for how to realize physical qubits based on the idea of working with quasi-particles called anyons.

Training phase (of an ML system)

A stage in the practical development of a machine learning system; given problem specific training data and a parameterized mathematical model, the model parameters are adjusted automatically such that the model matches the training data to the best extend possible; often, optimal model parameters are estimated by means of minimizing a loss function.

Two-state quantum system

A quantum mechanical system whose state space consists of infinitely many states which can be represented as two-dimensional, complex-valued unit vectors that are linear combinations of two distinguished, linearly independent, orthonormal basis states. Well known real world examples of two-state quantum systems include the polarization of a photon (with basis states vertical or horizontal) or the spin of an electron (with basis states up or down).

Unsupervised learning

A machine learning paradigm that aims to detect and uncover latent or inherent structures within a given data set of unlabeled data; a common example of an unsupervised learning problem is data clustering.

Variational quantum algorithm

An algorithm that combines quantum computations with classical computations in an iterative feedback loop; variational quantum algorithms or hybrid quantum-classical algorithms are used to adjust the parameters of a parameterized quantum circuit such that it shows an intended input/output behavior within a given tolerance; a variational quantum algorithm involves an outer loop run on a digital computer which manages the current parameter estimates, in each iteration, these parameters are used to setup computations on a quantum computer whose outcomes are measured, measurements are then used in classical optimization techniques to estimate improved parameters and the whole processes iterates until convergence.

Variational quantum eigensolver (VQE)

A variational quantum algorithm for estimating the bottom eigenvalues of Hamiltonian operators.

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