

UNIVERSITY NAME (IN BLOCK CAPITALS)

Thesis Title

by

Author Name

A thesis submitted in partial fulfillment for the
degree of Doctor of Philosophy

in the

Faculty Name

Department or School Name

March 2018

Declaration of Authorship

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- Where any part of this thesis has previously been submitted for a degree or any other qualification at this University or any other institution, this has been clearly stated.
- Where I have consulted the published work of others, this is always clearly attributed.
- Where I have quoted from the work of others, the source is always given. With the exception of such quotations, this thesis is entirely my own work.
- I have acknowledged all main sources of help.
- Where the thesis is based on work done by myself jointly with others, I have made clear exactly what was done by others and what I have contributed myself.

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UNIVERSITY NAME (IN BLOCK CAPITALS)

Abstract

Faculty Name

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Doctor of Philosophy

by Author Name

The Thesis Abstract is written here (and usually kept to just this page). The page is kept centered vertically so can expand into the blank space above the title too...

Acknowledgements

The acknowledgements and the people to thank go here, don't forget to include your project advisor...

Contents

Declaration of Authorship	i
Abstract	iii
Acknowledgements	iv
List of Figures	vi
List of Tables	vii
Abbreviations	viii
Physical Constants	ix
Symbols	x
1 Results	1
1.1 Linear Case	1
1.2 Anistropic Case	3
A An Appendix	24

List of Figures

1.1	The Binder cumulant for different sizes with $\lambda_x = \lambda_y = 0$.	2
1.2	The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = \lambda_y = 0$.	2
1.3	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0$.	3
1.4	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0$. There is no value for 128.	3
1.5	The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.2$.	4
1.6	The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 0.2$.	4
1.7	The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.4$.	5
1.8	The Binder cumulant as a function for different sizes with $\lambda_x = -\lambda_y = 0.4$.	5
1.9	The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.6$.	6
1.10	The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 0.6$.	6
1.11	The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.8$.	7
1.12	The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 0.8$.	7
1.13	The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 1$.	8
1.14	The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 1$.	8
1.15	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 0.2$.	9
1.16	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 0.4$.	9
1.17	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 0.6$.	10
1.18	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 0.8$.	10
1.19	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 1$.	11

1.20	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.2$.	11
1.21	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.4$.	12
1.22	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.6$.	12
1.23	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.8$.	13
1.24	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 1$.	13
1.25	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.2$. There is no value for 128.	14
1.26	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = \lambda_y = 0.4$. There is no value for 128.	14
1.27	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.6$. There is no value for 128.	15
1.28	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.8$. There is no value for 128.	15
1.29	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = -\lambda_y = 1$. There is no value for 128.	16
1.30	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = \lambda_y = 0.2$. There is no value for 128.	21
1.31	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = \lambda_y = 0.4$. There is no value for 128.	21
1.32	The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = \lambda_y = 0.6$. There is no value for 128.	22

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- 1.33 The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0.8$. There is no value for 128. 22
- 1.34 The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 1$. There is no value for 128. 23

List of Tables

Abbreviations

LAH List Abbreviations **Here**

Physical Constants

$$\text{Speed of Light } c = 2.997\,924\,58 \times 10^8 \text{ ms}^{-\text{S}} \text{ (exact)}$$

Symbols

a	distance	m
P	power	W (Js^{-1})
ω	angular frequency	rads^{-1}

For/Dedicated to/To my...

Chapter 1

Results

Simulations were performed on system sizes $L = 40, 48, 64, 80, 104, 128$ with the λ s taken from zero to one in steps of 0.2 and also 1.5 where λ_x and λ_y had the same magnitude and the same or opposite signs.

1.1 Linear Case

Figure. 1.1 overlays the Binder cumulant evolution for different sizes in the linear ($\lambda_x = \lambda_y = 0$) case which corresponds to the XY model. Figure. 1.2 does the same but plotted with respect to $t/L^2 \log t$ where the collapse is expected. As is clear, this reobtains the result from XX while using Stochastic evolution rather than Monte Carlo and provides confidence in the results for the non-linear cases. The collapse is not perfect: in Figures 1.3 and ??, which shows the Binder cumulant as a function of the number of realisations for all the system sizes at two different points, shows how far apart the curves are. For the first case, the uncertainties of the points do not all cross, although none are isolated. In the second case, the collapse is weaker and $L = 40$ is consistently isolated, as can be seen from the graph of the collapse. This also occurs for $L = 64$ earlier on in the simulation. For $L = 48$ this could be an underestimation of the error, and there is a slight downward decrease in its value as the number of realisations is increased. Ideally one would perform further simulations to ensure this, but this was not possible in the time constraints of the project.

With those caveats in mind, we may be assured that the result has been obtained, and this is further corroborated by the plots of $\log n_v$, where n_v is the number of vortices, as a function of $\log(t/\log t)$, shown in XX to XX for each system size. For the linear case, it is expected that $\log n_v = -\log(t/\log t)$. This can be explained as follows: using the result

in XX that the distance between vortices R is approximately proportional to $(t/\log t)^{1/2}$, we have that the vortex density is then $1/R^2$, and the number of vortices $n_v \approx L^2/R^2$, and therefore $\log n_v \approx -\log R^2 = -\log(t/\log t)$ up to some additive constants. As the figures show, the gradient was obtained to good accuracy for all system sizes. Due to the finite size of the system, the approximation breaks down at large times once the majority of the vortices have annihilated, and therefore a time scale must be chosen with which to calculate the gradient. In practice, the time scale maximised but chosen to be prior to the finite size behaviour of the system. Therefore it is sensible not to read the values literally, but take them as indicative of the XY model behaviour for the linear case.

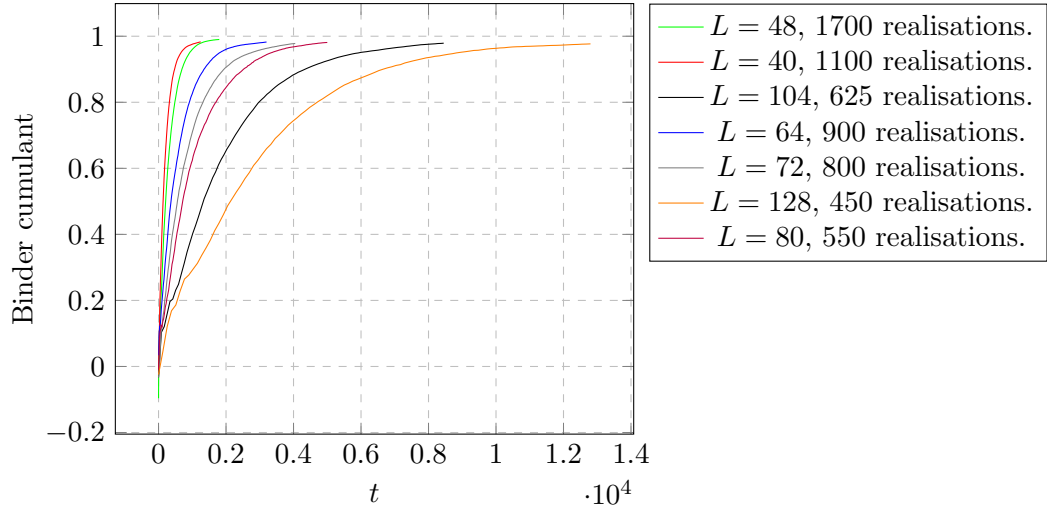


FIGURE 1.1: The Binder cumulant for different sizes with $\lambda_x = \lambda_y = 0$.

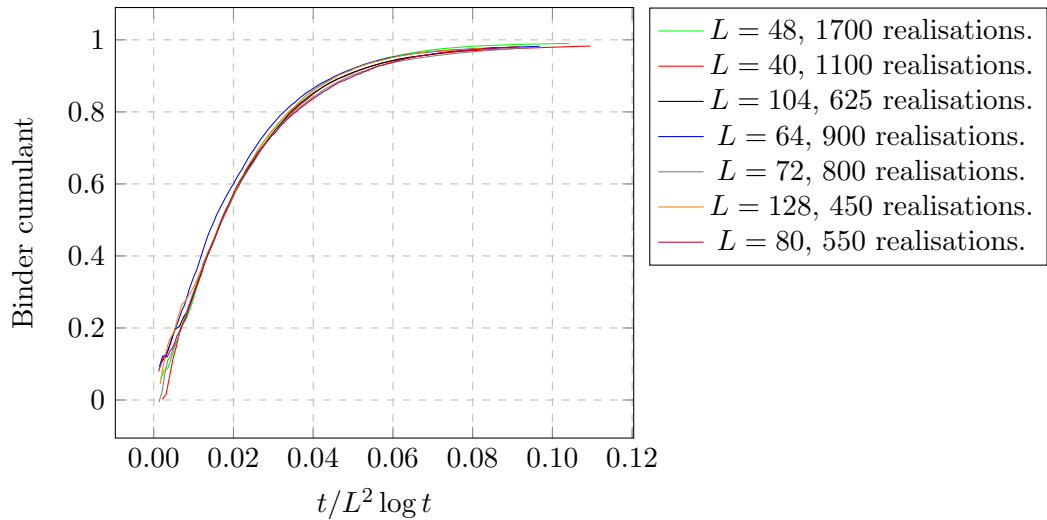


FIGURE 1.2: The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = \lambda_y = 0$.

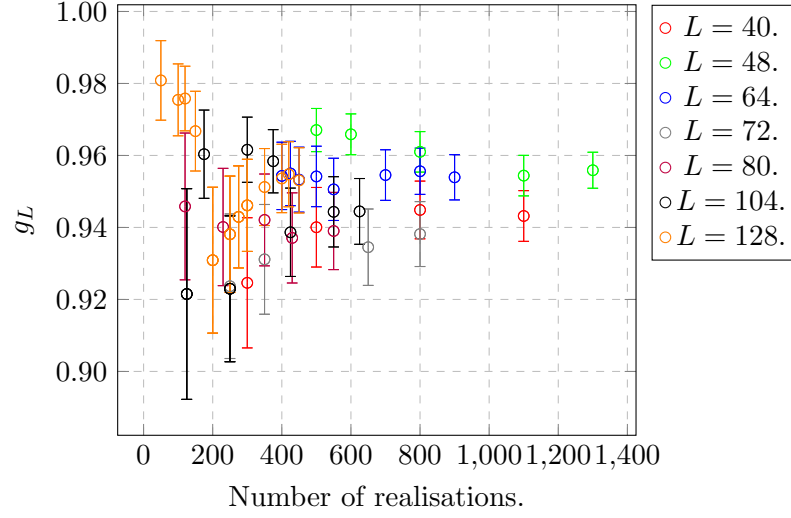


FIGURE 1.3: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0$.

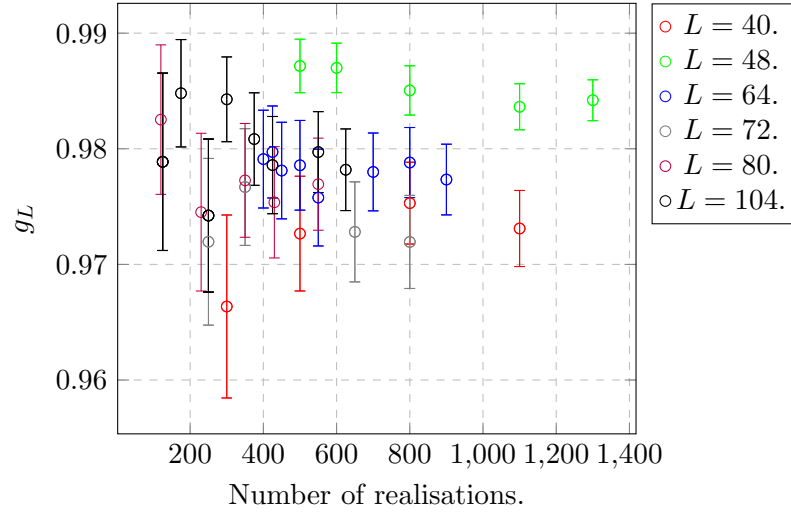


FIGURE 1.4: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40, \lambda_x = \lambda_y = 0$. There is no value for 128.

1.2 Anisotropic Case

For non-zero λ s, when $\lambda_x = -\lambda_y$, we expect to obtain the same result as the XY model. Figures ?? to ?? show the Binder cumulant plotted as a function of t and also $t/L^2 \log t$ for several values of increasing λ_x . Although the collapse is not as tight as for the linear case, it clearly still occurs to an extent, confirming the expectation. Due to the errors on the Binder cumulant, which are shown in Figures XX to XX, is not clear whether further realisations are necessary to tighten the collapse further, or if the dynamical exponent growth rate $(t/\log t)^{1/2}$ is simply not as good an approximation as in the linear case. There are clear differences here, however, as the convergence to near one occurs more

rapidly on increasing λ . This is corroborated by fig XX which shows that the gradient in $\log n_v$ where n_v is then number of vortex pairs as a function of $\log(t/\log t)$ decreases on increasing λ , indicating a faster transition to the ordered phase.

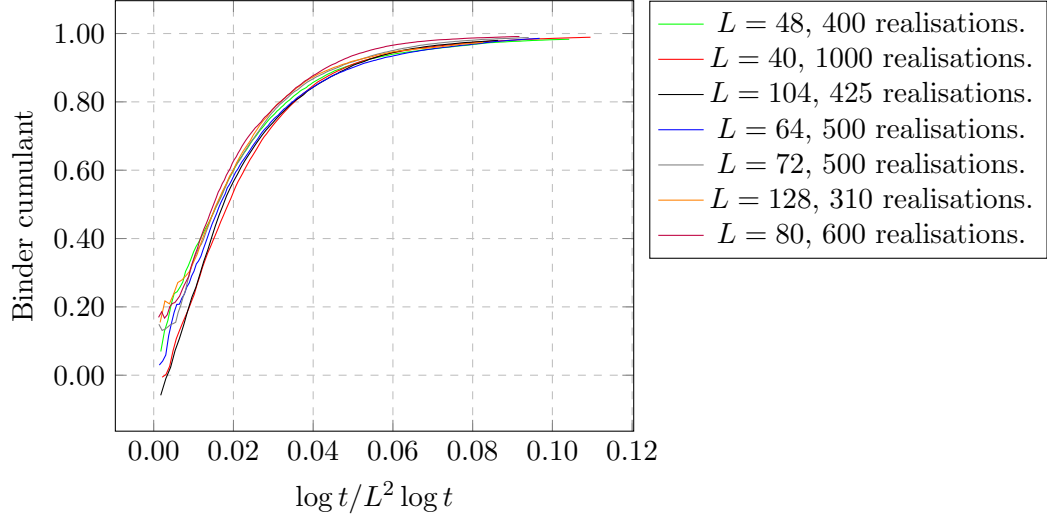


FIGURE 1.5: The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.2$.

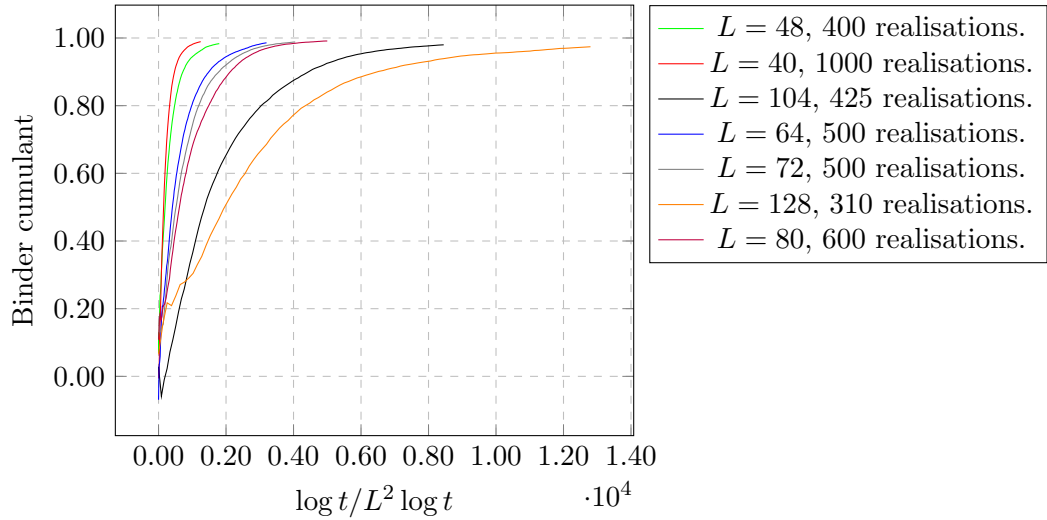


FIGURE 1.6: The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 0.2$.

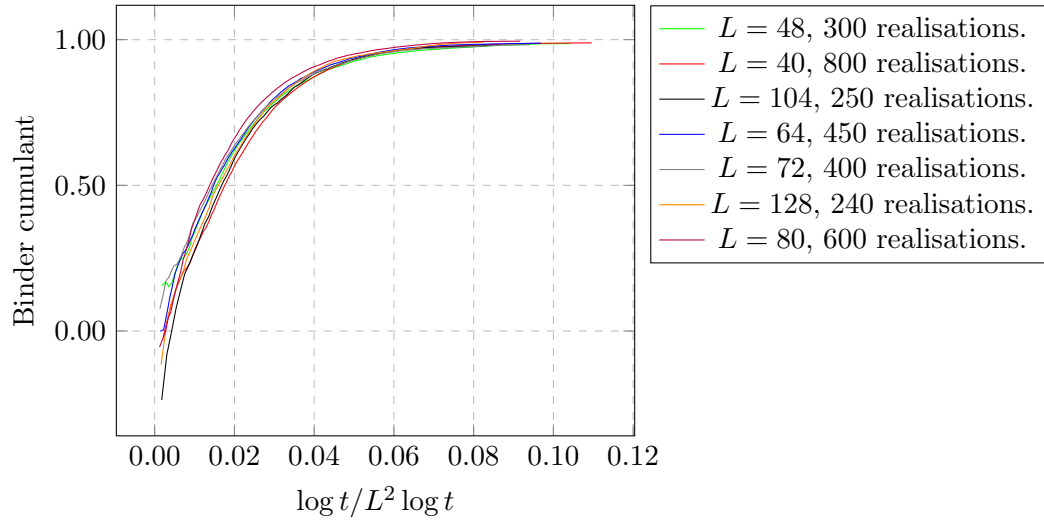


FIGURE 1.7: The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.4$.

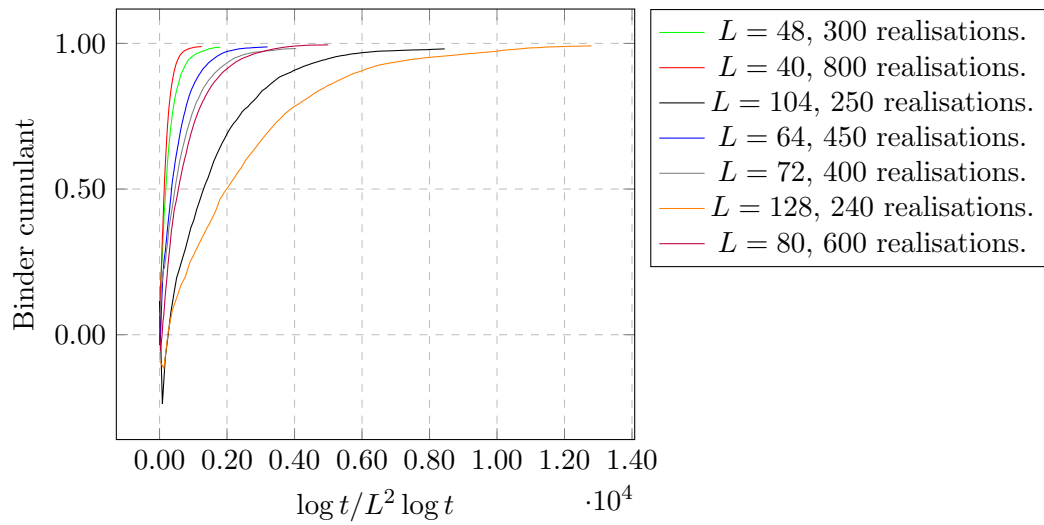


FIGURE 1.8: The Binder cumulant as a function for different sizes with $\lambda_x = -\lambda_y = 0.4$.

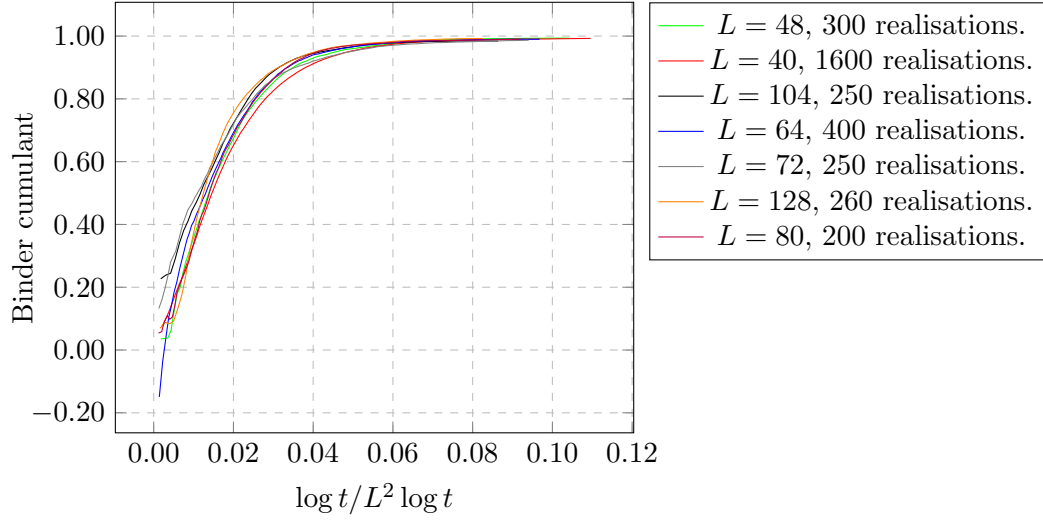


FIGURE 1.9: The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.6$.

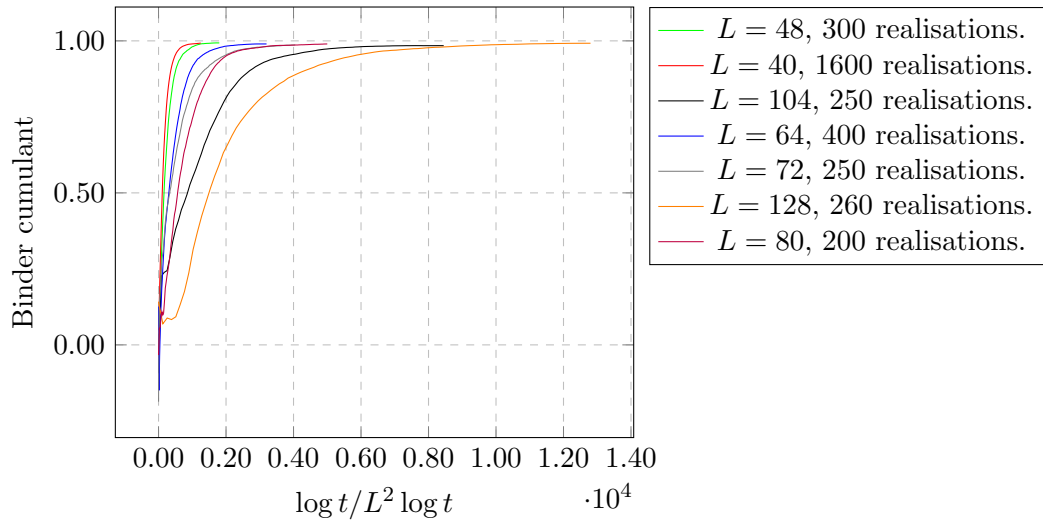


FIGURE 1.10: The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 0.6$.

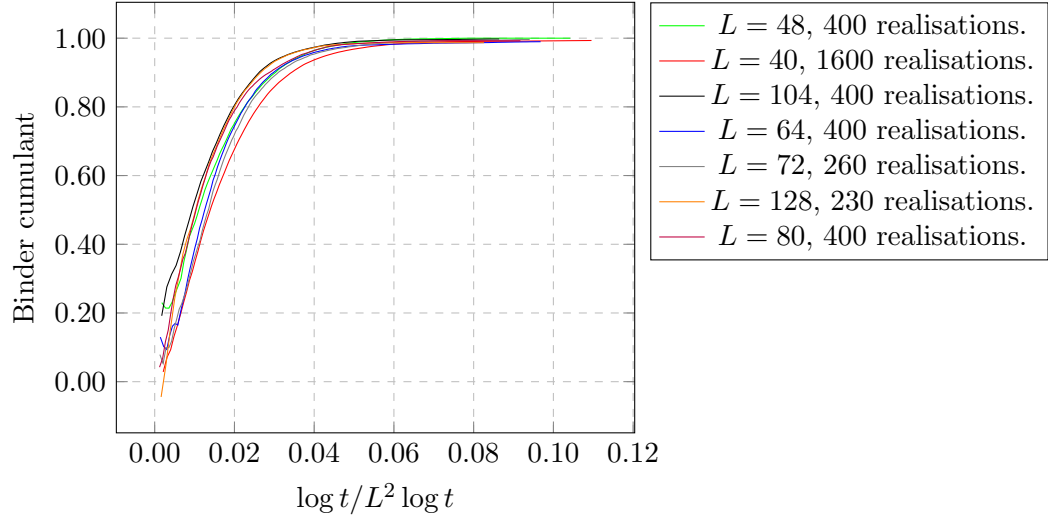


FIGURE 1.11: The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 0.8$.

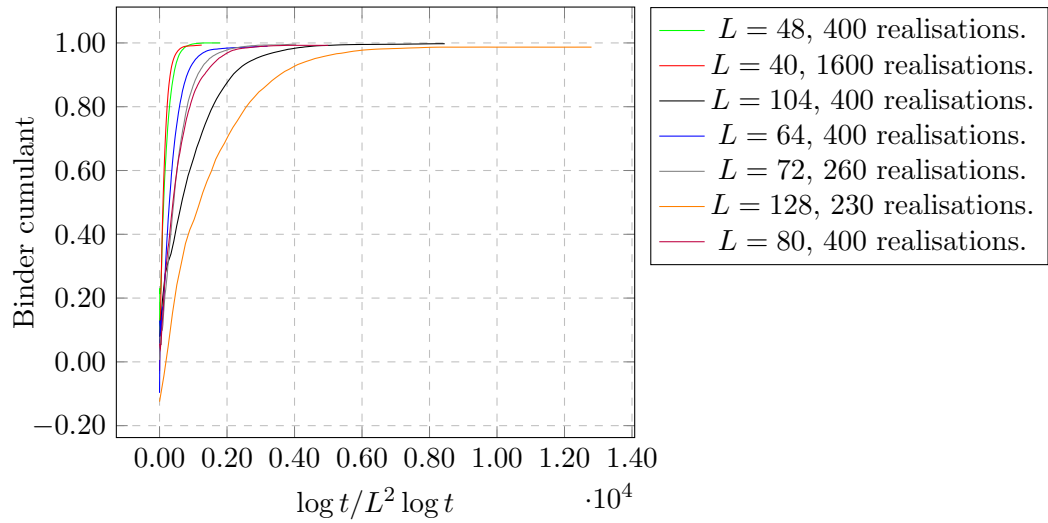


FIGURE 1.12: The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 0.8$.

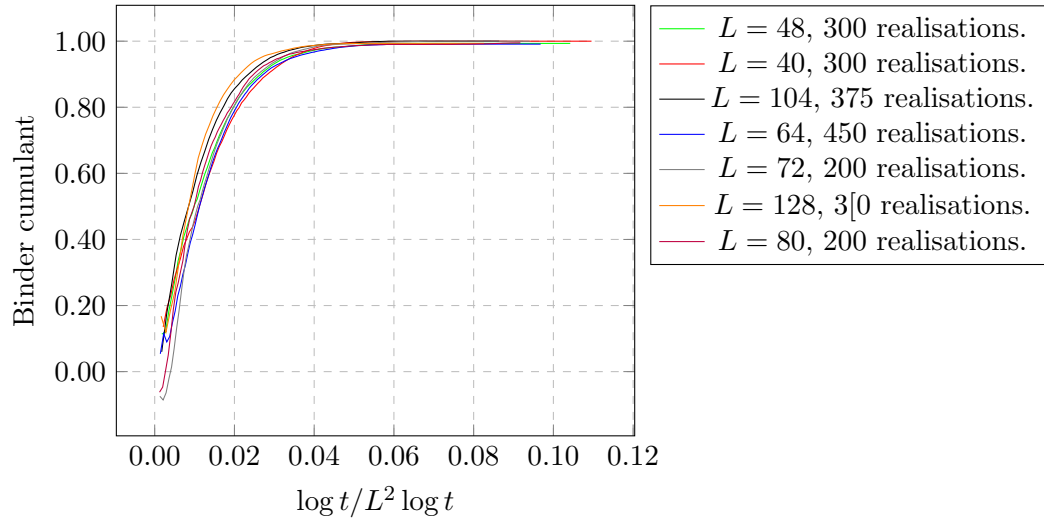


FIGURE 1.13: The Binder cumulant as a function of $t/L^2 \log t$ for different sizes with $\lambda_x = -\lambda_y = 1$.

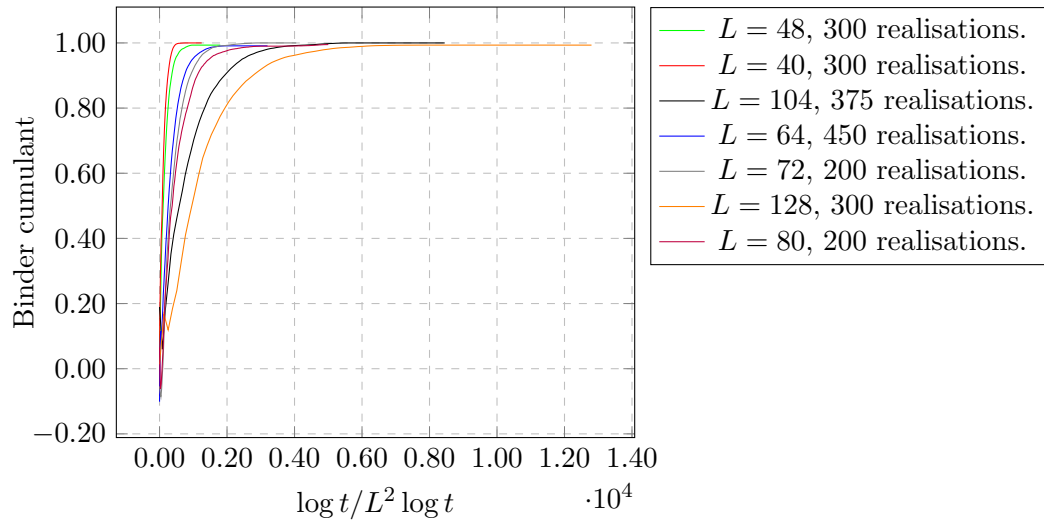


FIGURE 1.14: The Binder cumulant for different sizes with $\lambda_x = -\lambda_y = 1$.

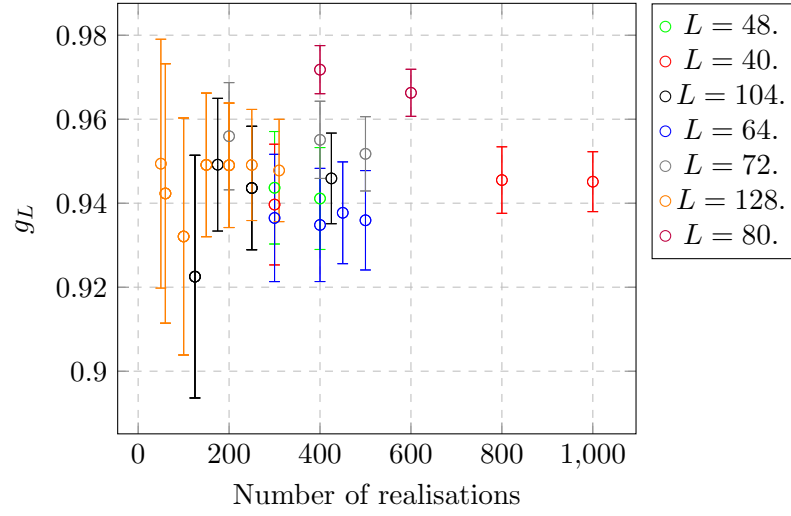


FIGURE 1.15: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.2$.

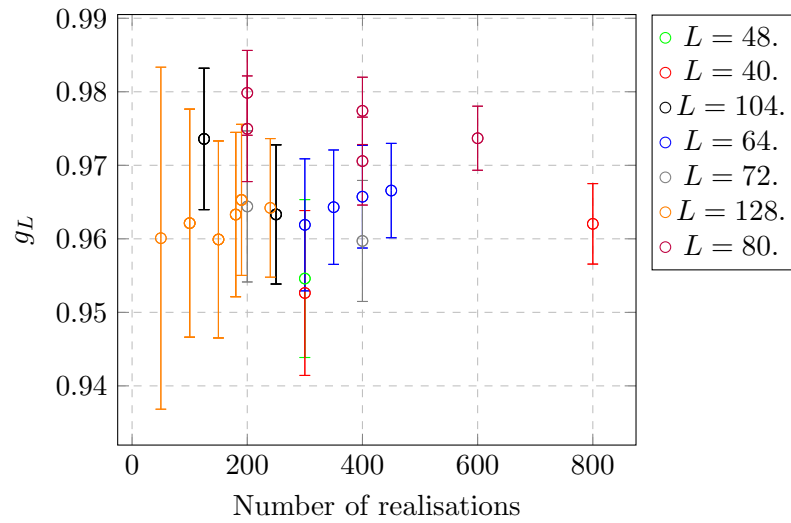


FIGURE 1.16: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.4$.

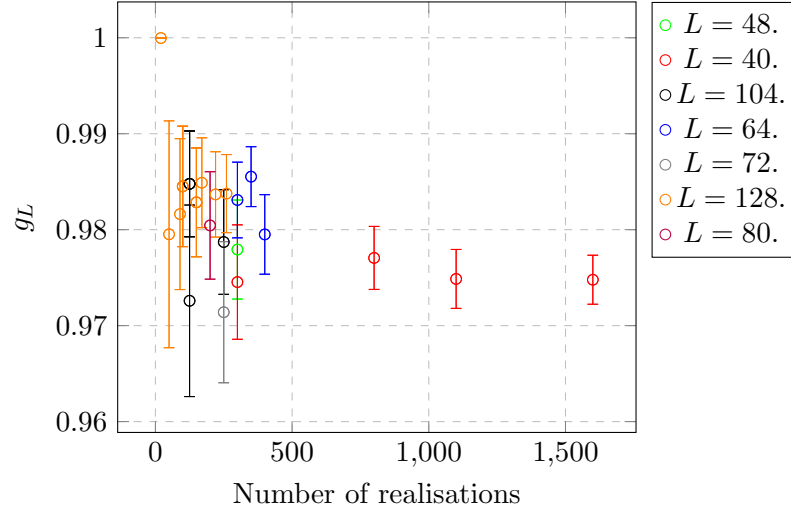


FIGURE 1.17: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.6$.

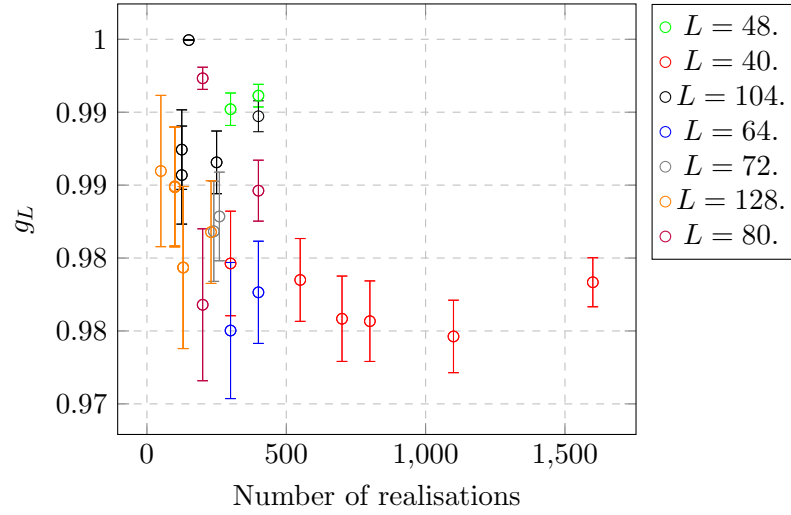


FIGURE 1.18: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = -\lambda_y = 0.8$.

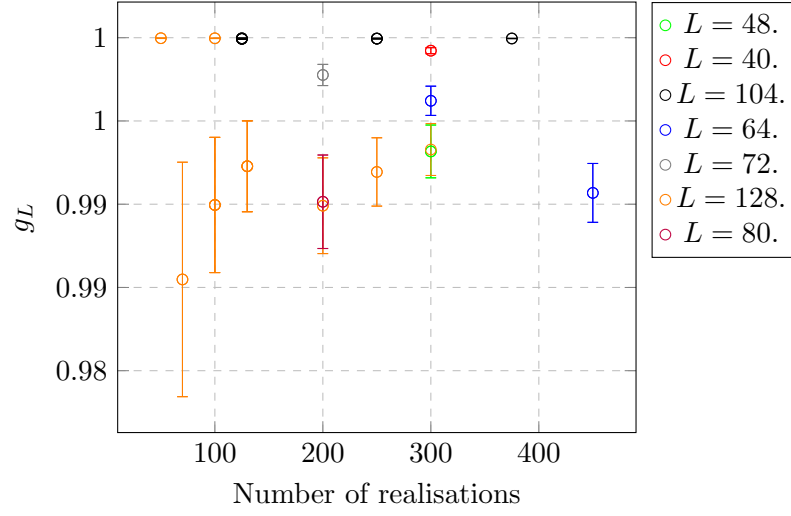


FIGURE 1.19: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 1$.

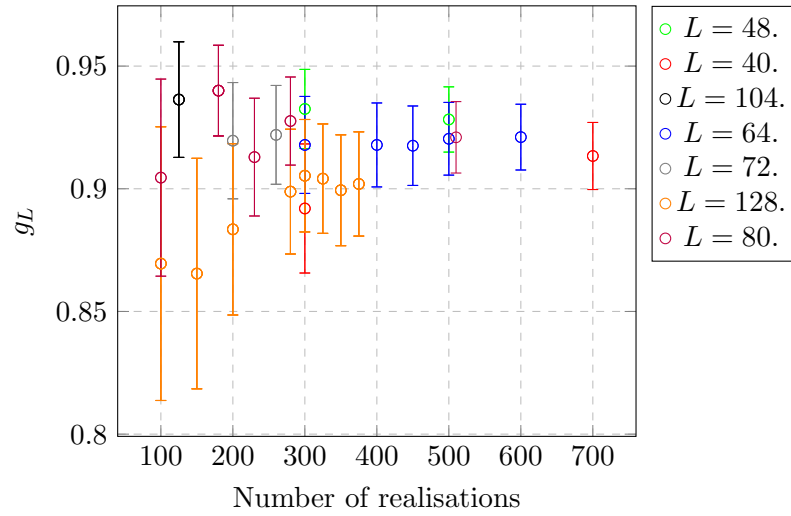


FIGURE 1.20: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0.2$.

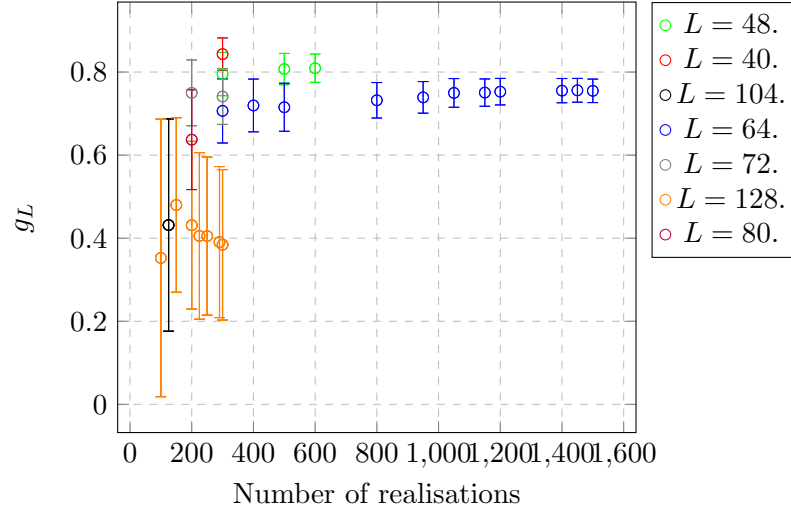


FIGURE 1.21: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.4$.

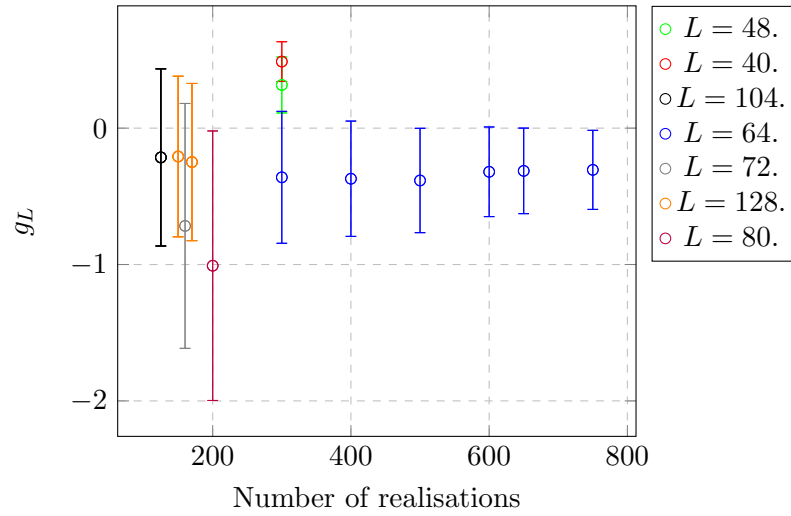


FIGURE 1.22: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.6$.

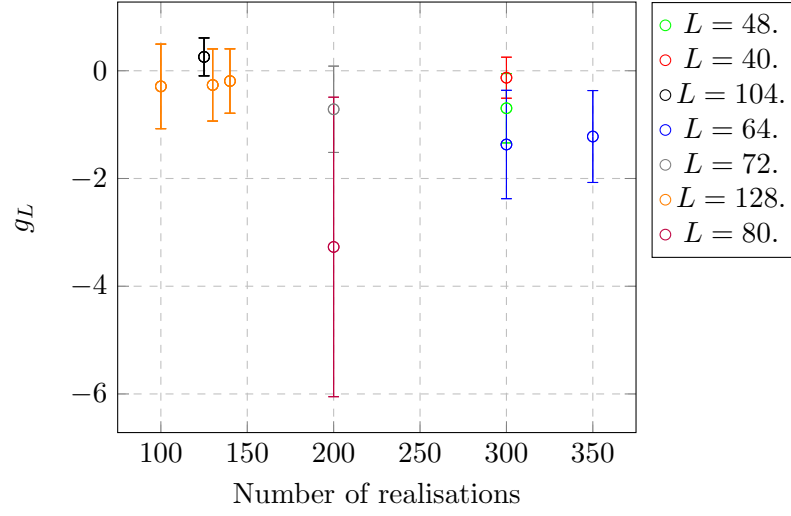


FIGURE 1.23: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 0.8$.

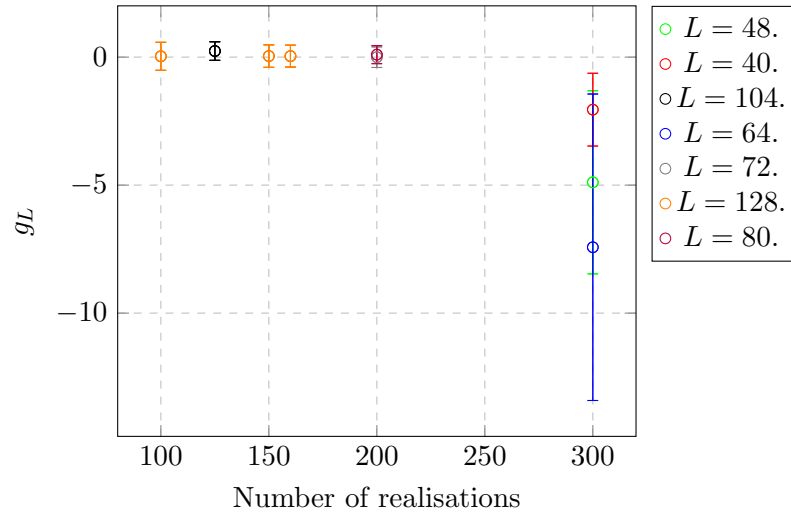


FIGURE 1.24: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 6500$ (the mid-point of the simulation) for $L = 40, \lambda_x = \lambda_y = 1$.

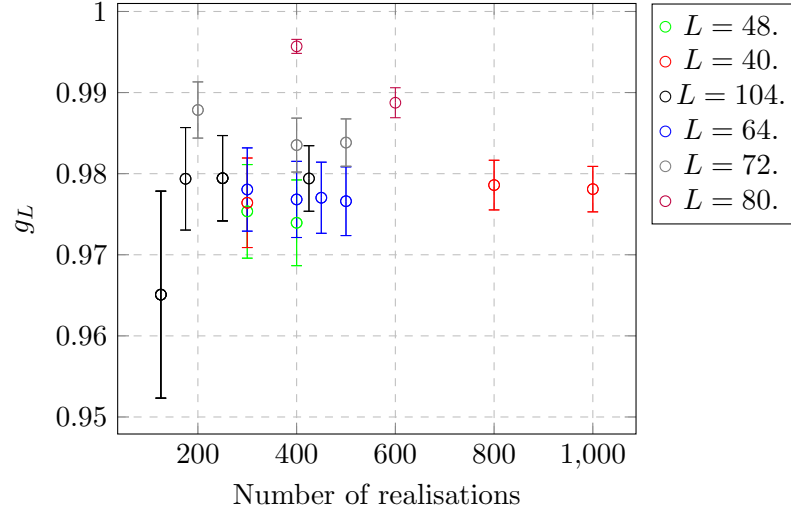


FIGURE 1.25: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 0.2$. There is no value for 128.

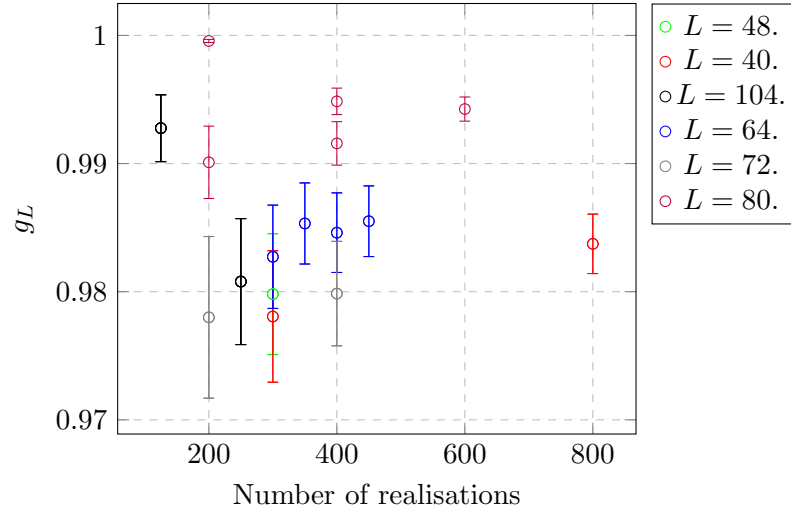


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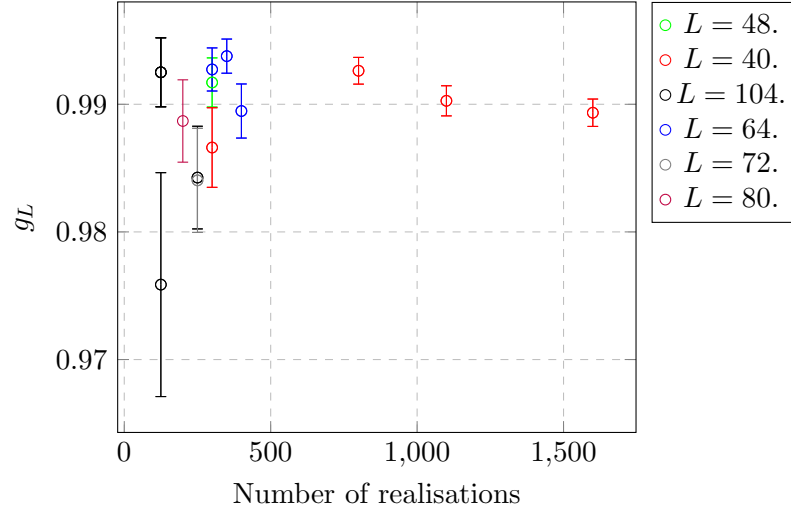


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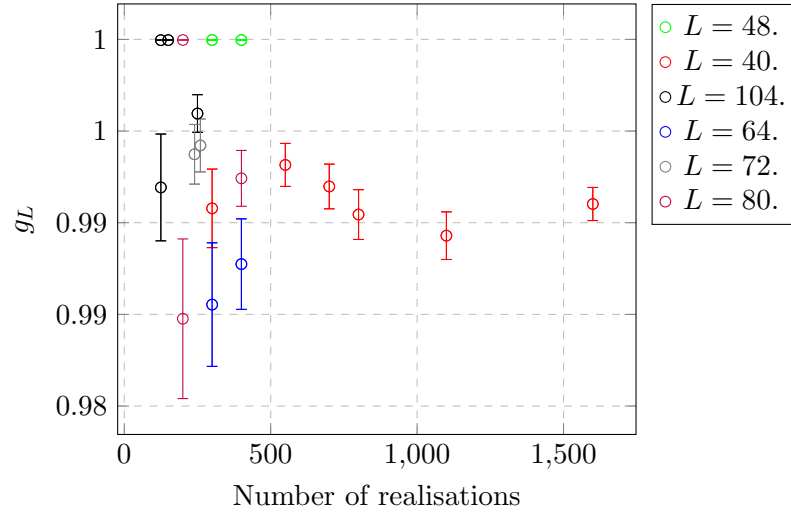


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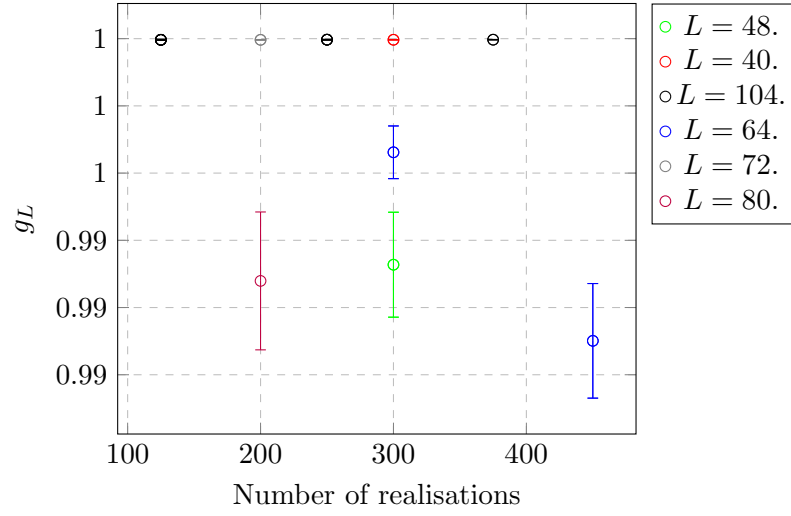
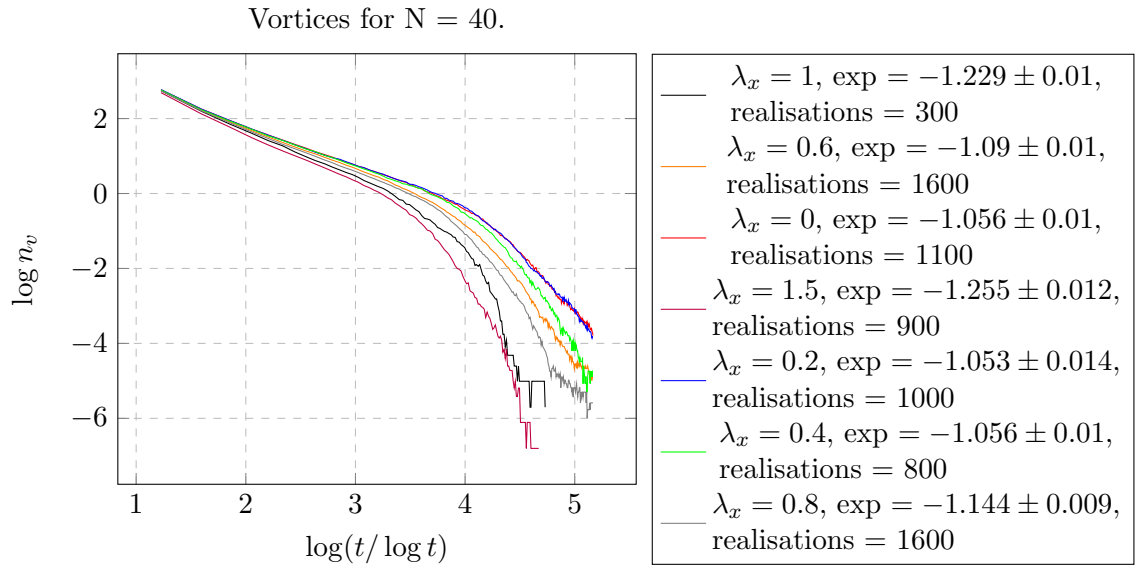
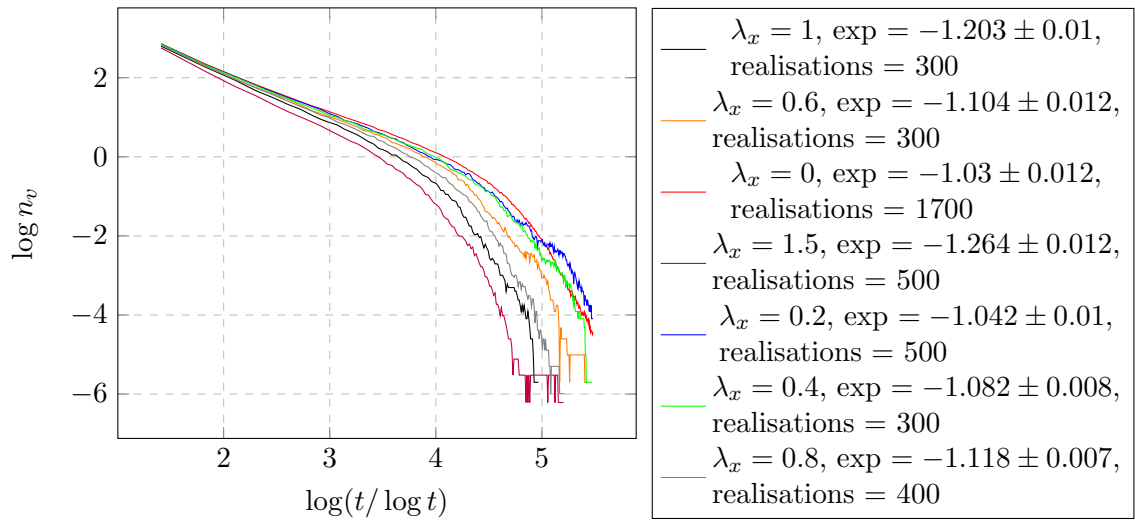
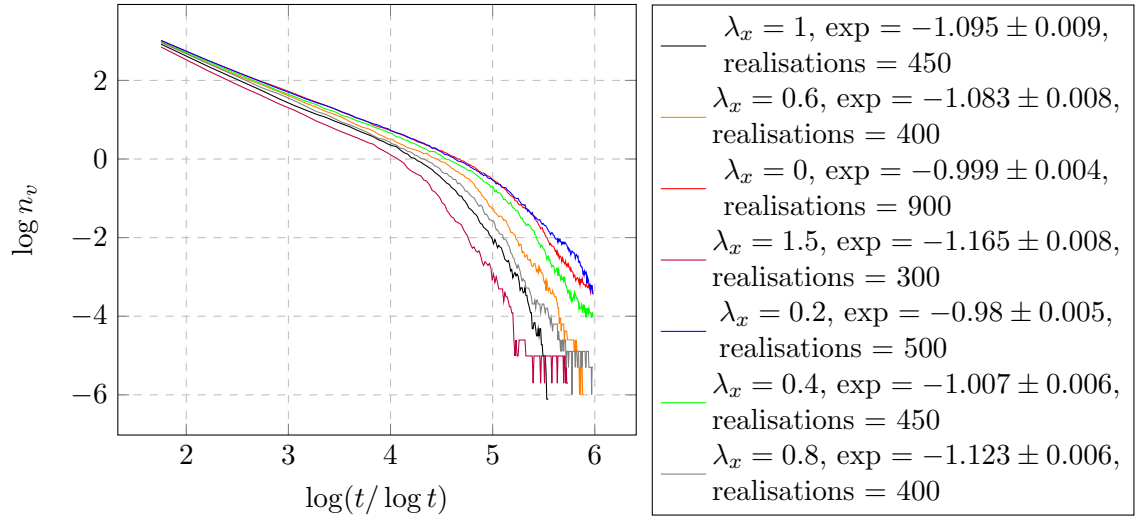
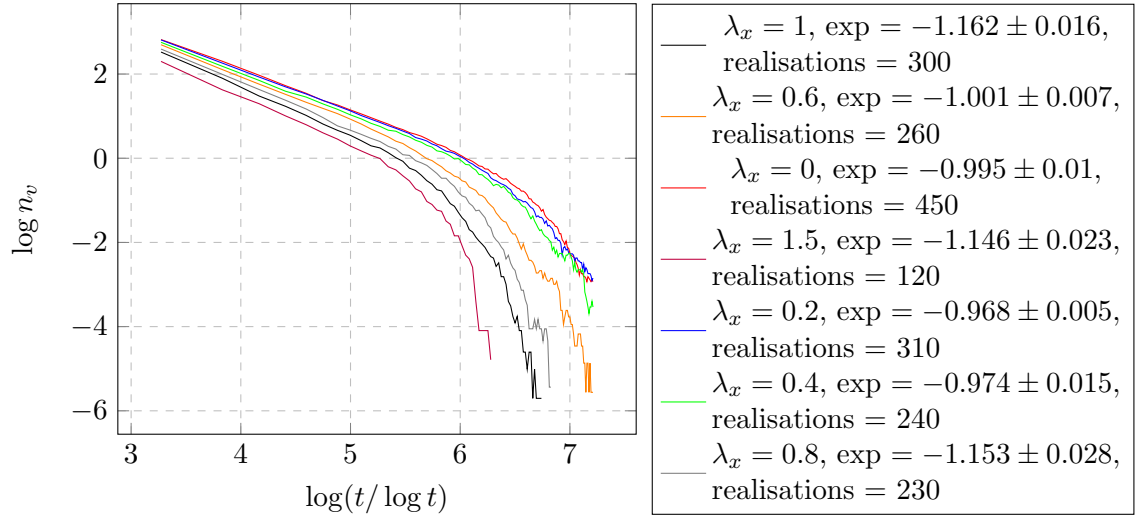
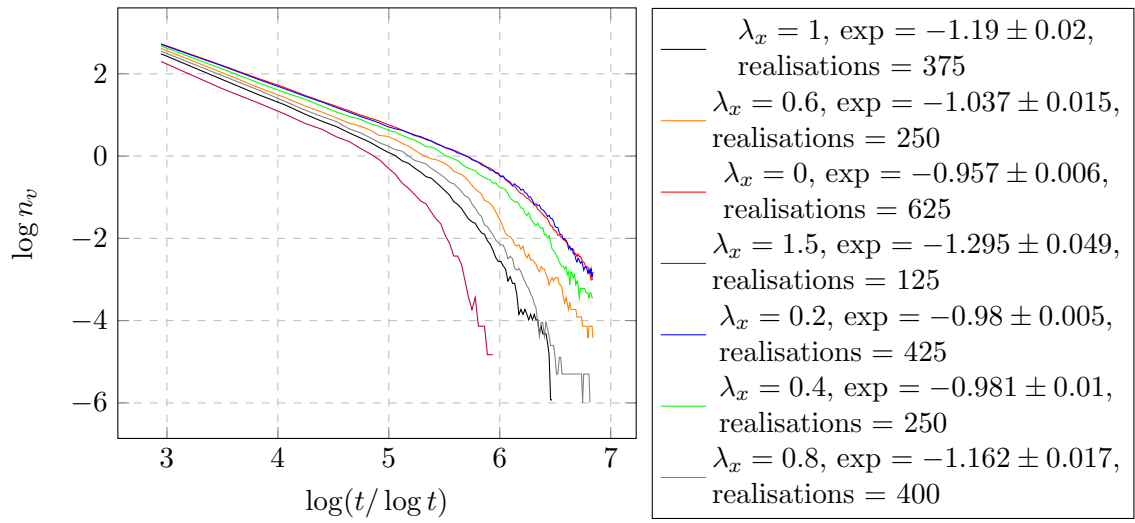


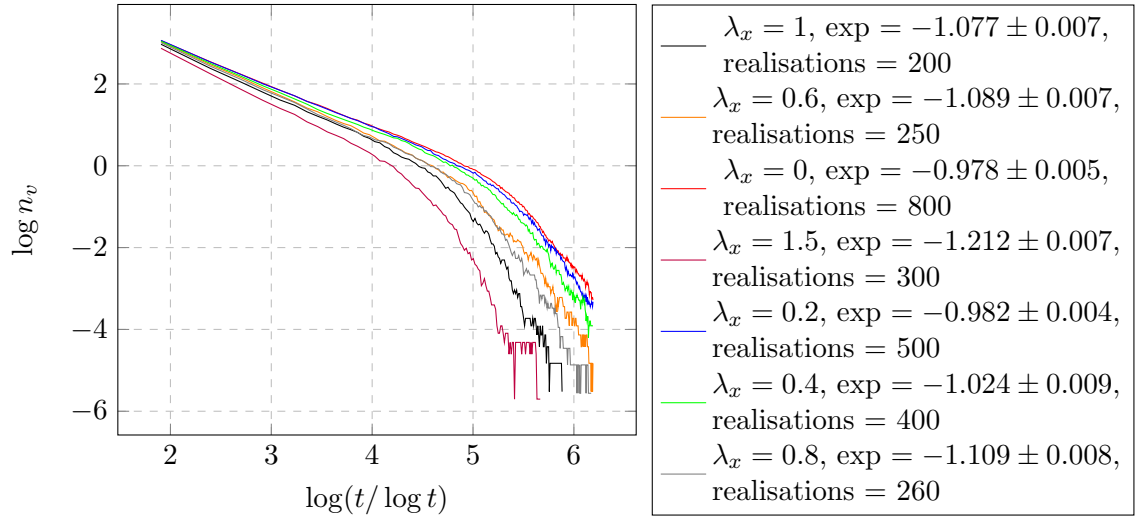
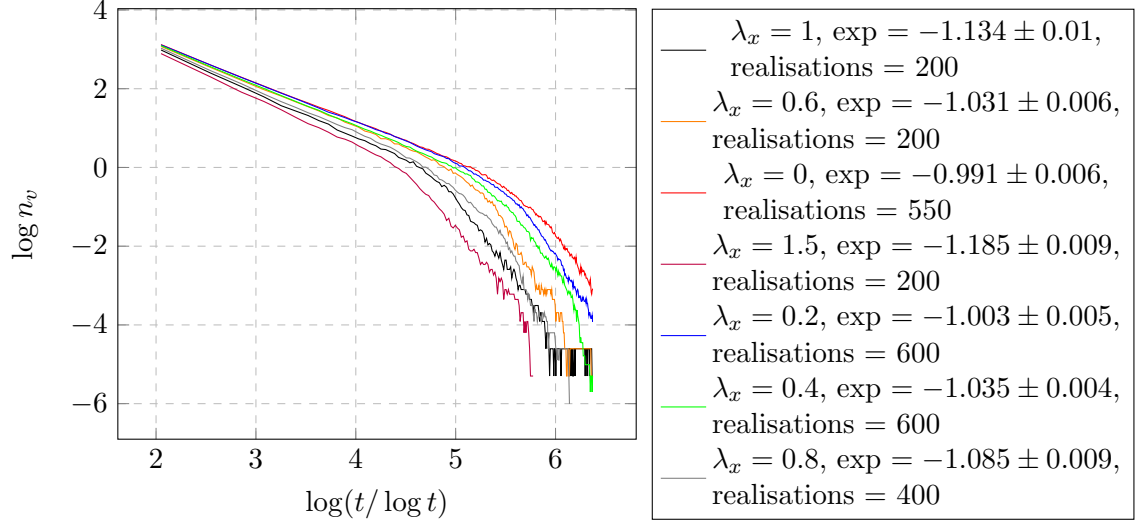
FIGURE 1.29: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = -\lambda_y = 1$. There is no value for 128.



Vortices for $N = 48$.

Vortices for $N = 64$.Vortices for $N = 128$.

Vortices for $N = 104$.

Vortices for $N = 72$.Vortices for $N = 80$.

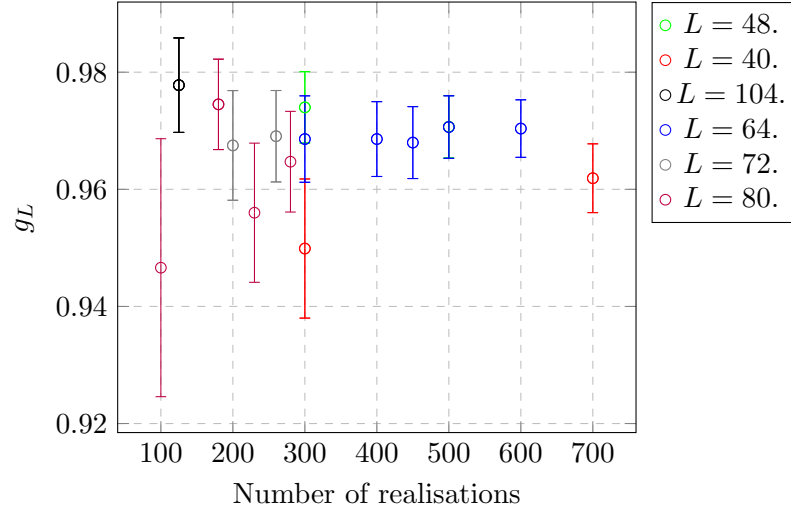


FIGURE 1.30: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0.2$. There is no value for 128.

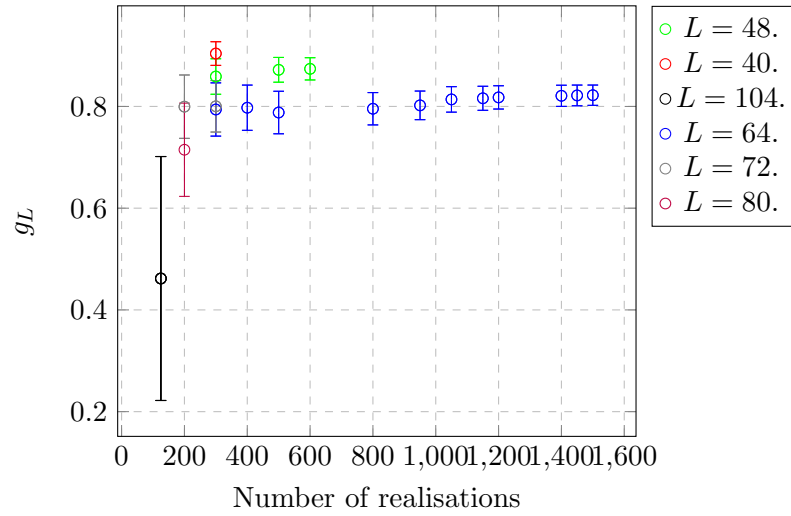


FIGURE 1.31: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0.4$. There is no value for 128.

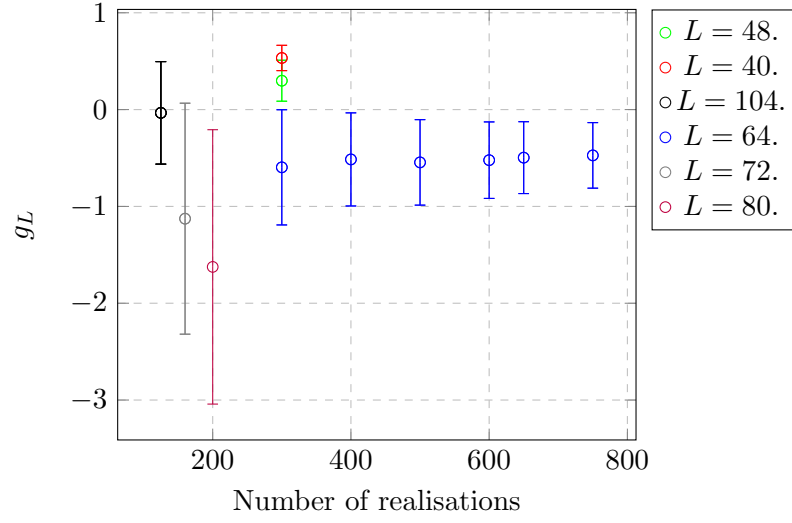


FIGURE 1.32: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0.6$. There is no value for 128.

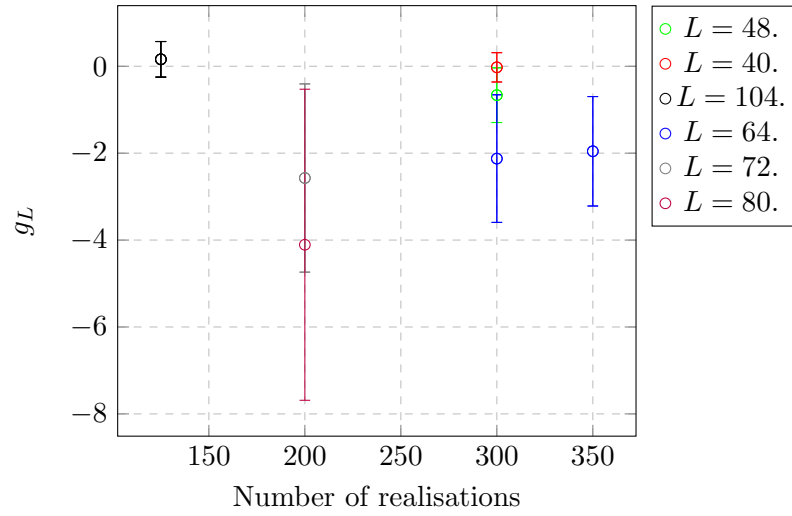


FIGURE 1.33: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 0.8$. There is no value for 128.

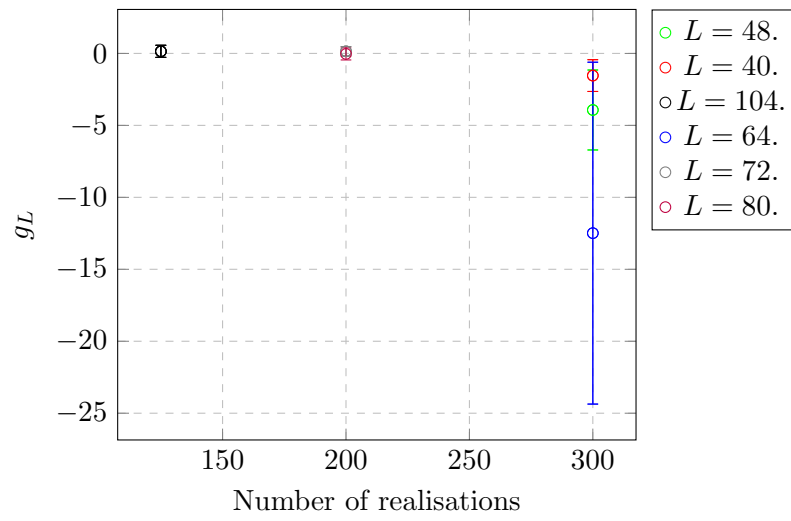


FIGURE 1.34: The uncertainty in the Binder cumulant as a function of the number of realisations at a point closest to $t/L^2 \log t$ for $t = 937.5$ (three quarters through the simulation) for $L = 40$. $\lambda_x = \lambda_y = 1$. There is no value for 128.

Appendix A

An Appendix

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