



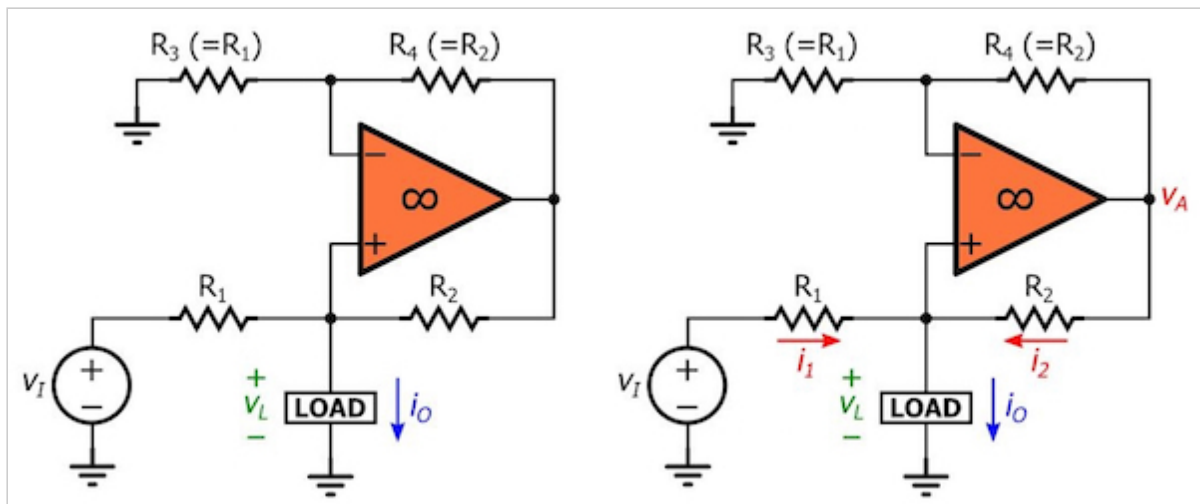
# The Howland Current Pump

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The Howland current pump, invented by MIT's Professor Bradford Howland in the early 1960s, consists of an op-amp and a balanced resistor bridge and puts out current in either direction.

The Howland current pump, shown in Figure 1a, is a circuit that accepts an input voltage  $v_I$ , converts it to an output current  $i_O = Av_I$ , with  $A$  as the transconductance gain, and pumps  $i_O$  to a load LD, regardless of the voltage  $v_L$  developed by the load itself. To see how it works, label it as in Figure 1b, and apply Kirchoff's Current Law (<https://www.allaboutcircuits.com/textbook/direct-current/chpt-6/kirchhoffs-current-law-kcl/>) and Ohm's Law (<https://www.allaboutcircuits.com/textbook/direct-current/chpt-2/voltage-current-resistance-relate/>).



**Figure 1.** (a) The Howland pump. (b) Properly labeling the circuit for its analysis.

$$i_O = i_1 + i_2 = \frac{v_I - v_L}{R_1} + \frac{v_A - v_L}{R_2}$$

Equation 1

The op-amp, together with  $R_3$  and  $R_4$ , forms a non-inverting amplifier with respect to  $v_L$ , thus giving

$$v_A = (1 + R_4 / R_3) v_L$$

Equation 2

Substituting  $v_A$  into Equation 1 and collecting, we put  $i_O$  into the insightful form

$$i_O = A v_I - \frac{v_L}{R_o}$$

Equation 3

where  $A$  is the transconductance gain, in A/V,

$$A = \frac{1}{R_1}$$

## Equation 4

and where  $R_o$  is the output resistance presented by the circuit to the load,

$$R_o = \frac{R_2}{R_2/R_1 - R_4/R_3}$$

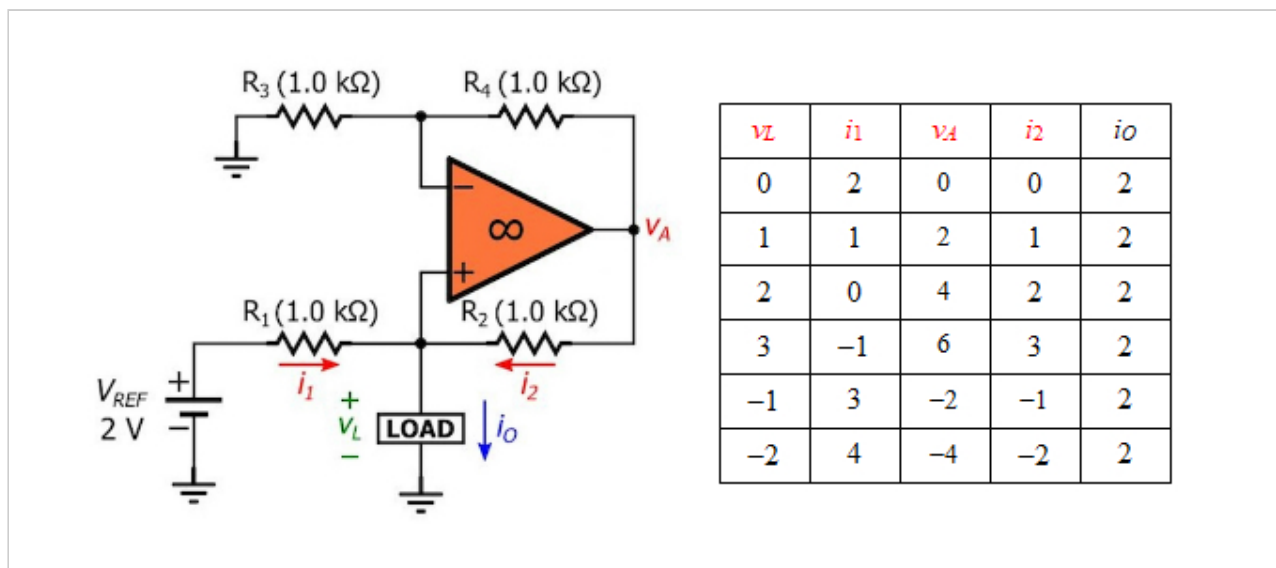
## Equation 5

To make  $i_o$  independent of  $v_L$  we must impose  $R_o \rightarrow \infty$ , or the balanced-bridge condition.

$$\frac{R_4}{R_3} = \frac{R_2}{R_1}$$

## Equation 6

Take a look at the example in Figure 2 and observe, row-by-row, how the op-amp adjusts  $i_2$ , via  $v_A$ , so as to ensure the same current  $i_o$  regardless of  $v_L$ .



**Figure 2.** (a) A 2 mA current source, and (b) its inner workings for different values of  $v_L$  (voltages in volts, currents in milliamps; a negative current value means that current flows in the direction opposite to the arrow).

With the polarity of  $V_{REF}$  as shown, the pump sources  $i_O$  to the load. Inverting the polarity of  $V_{REF}$  will cause the pump to sink  $i_O$  from the load. Note that for the pump to work properly  $v_A$  must always be confined within the linear range of op-amp operation. If the op-amp is driven into saturation, the pump will cease to operate properly.

## The Effect of Resistance Mismatches

A practical bridge is likely to be unbalanced because of resistance tolerances, so  $R_o$  is likely to be less than infinity. Denoting the tolerances of the resistances in use by  $p$ , we note that the denominator  $D$  of Equation 5 is maximized when  $R_2$  and  $R_3$  are maximized and  $R_1$  and  $R_4$  are minimized. For  $p \ll 1$ , we write

$$D_{\max} = \frac{R_2(1+p)}{R_1(1-p)} - \frac{R_4(1-p)}{R_3(1+p)} \cong \frac{R_2}{R_1}(1+p)^2 - \frac{R_4}{R_3}(1-p)^2 \cong \frac{R_2}{R_1}[(1+2p) - (1-2p)] \cong \frac{R_2}{R_1}4p$$

Here we have incorporated the relationship of Equation 6, applied approximation

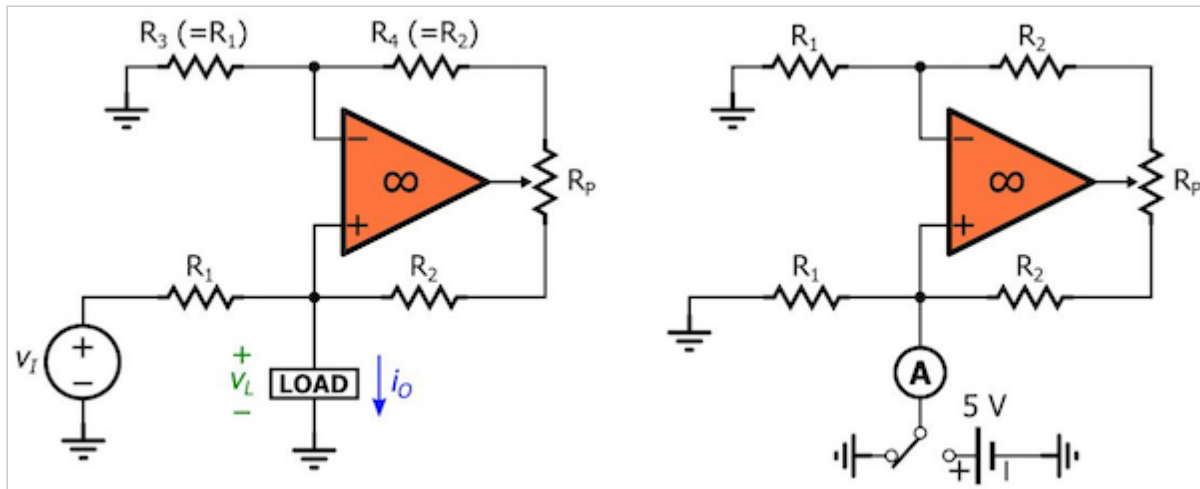
$$1/(1 \mp p) \cong 1 \pm p$$

and ignored quadratic terms in  $p$ . Substituting into Equation 5 gives

$$R_{o(\min)} = \frac{R_2}{D_{\max}} \cong \frac{R_1}{4p}$$

*Equation 7*

As an example, using 1% ( $p = 0.01$ ) resistances in Figure 2a can lower  $R_o$  from  $\infty$  to as little as  $1,000/(4 \times 0.01) = 25 \text{ k}\Omega$ , thus making  $i_O$  depend upon  $v_L$ , by Equation 3. If the bridge is unbalanced in the opposite direction of above, then the worst-case condition for  $R_o$  is  $-25 \text{ k}\Omega$ . So, depending on the mismatch,  $R_o$  may lie anywhere from  $+25 \text{ k}\Omega$  to  $\infty$  to  $-25 \text{ k}\Omega$ .



**Figure 3.** (a) Using a potentiometer  $R_p$  to balance the resistive bridge. (b) Calibration set up.

For improved performance, we must either use lower-tolerance resistances or balance the bridge using a potentiometer  $R_p$ , as in Figure 3a. To calibrate the circuit, ground the input as in Figure 3b and use an ammeter  $A$ . First, flip the switch to ground, and if necessary, zero the op-amp's input offset voltage until the ammeter reads zero. Then flip the switch to a known voltage, such as 5V, and adjust  $R_p$  until the ammeter reads again zero. By imposing that  $i_O$  with  $v_L = 5\text{ V}$  be equal to  $i_O$  with  $v_L = 0\text{ V}$ , we are making  $i_O$  independent of  $v_L$ , in effect driving  $R_O$  to infinity, by Equation 3.

## The Effect of Op-Amp Nonidealities

### Common-Mode Rejection Ratio

A practical op-amp is sensitive to its common-mode input voltage, a feature that is modeled with a small internal offset voltage in series with the noninverting input. In the case of the Howland pump, this offset voltage can be expressed as  $v_L/\text{CMRR}$ , where CMRR is the common-mode rejection ratio as reported in the op-amp's datasheet. With reference to Figure 4a, we note that Equation 1 still holds, but Equation 2 changes to

$$v_A = \left(1 + \frac{R_4}{R_3}\right) \times \left(v_L + \frac{v_L}{\text{CMRR}}\right) = \left(1 + \frac{R_2}{R_1}\right) \times v_L \times \left(1 + \frac{1}{\text{CMRR}}\right)$$

Substituting into Equation 1, solving for  $i_O$ , and putting  $i_O$  in the form of Equation 3 gives

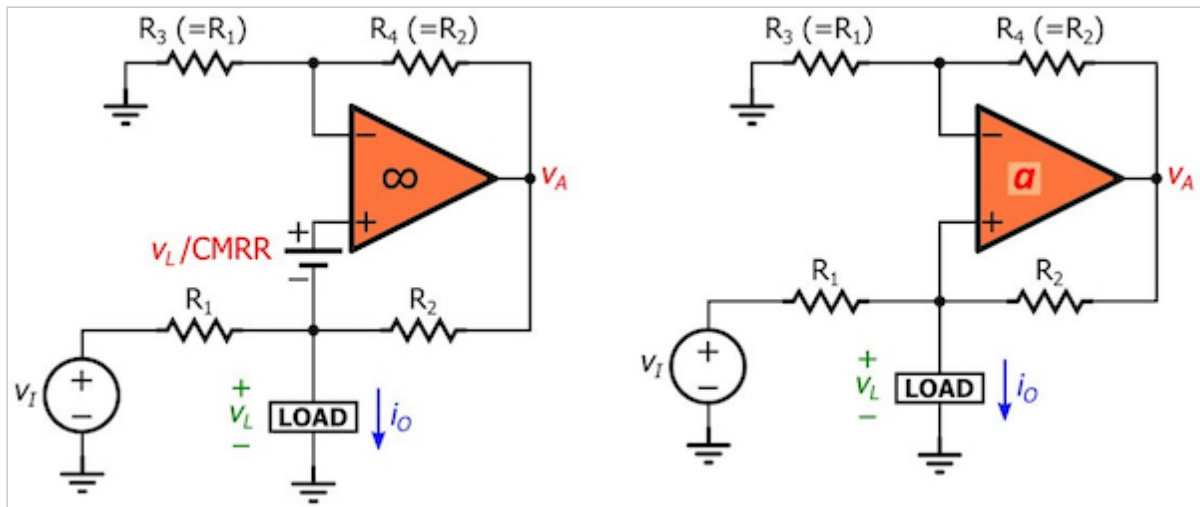
$$R_o = (R_1 \parallel R_2) \times \text{CMRR}$$

Equation 8

As an example, using an op-amp with CMRR = 60 dB (=1000) in the above example will lower  $R_o$  from  $\infty$  to  $(10^3 \parallel 10^3) \times 1000 = 500 \text{ k}\Omega$ . With an arrangement of the type of Figure 3b, we can use the potentiometer to compensate for the cumulative effect of bridge imbalance as well as non-infinite CMRR.

## Open-Loop Gain

So far we have assumed the op-amp to have infinite open-loop gain. The gain  $a$  of a practical op-amp is finite, so let us now see how this affects circuit behavior.



**Figure 4.** Circuits to investigate the effect of (a) non-infinite common-mode rejection ratio and (b) non-infinite open-loop gain.

With reference to Figure 4b, we now have

$$v_A = a \left( v_L - \frac{R_3}{R_3 + R_4} v_A \right)$$

Solving for  $v_A$ , substituting into Equation 1, solving for  $i_O$ , and putting  $i_O$  in the form of Equation 3 gives

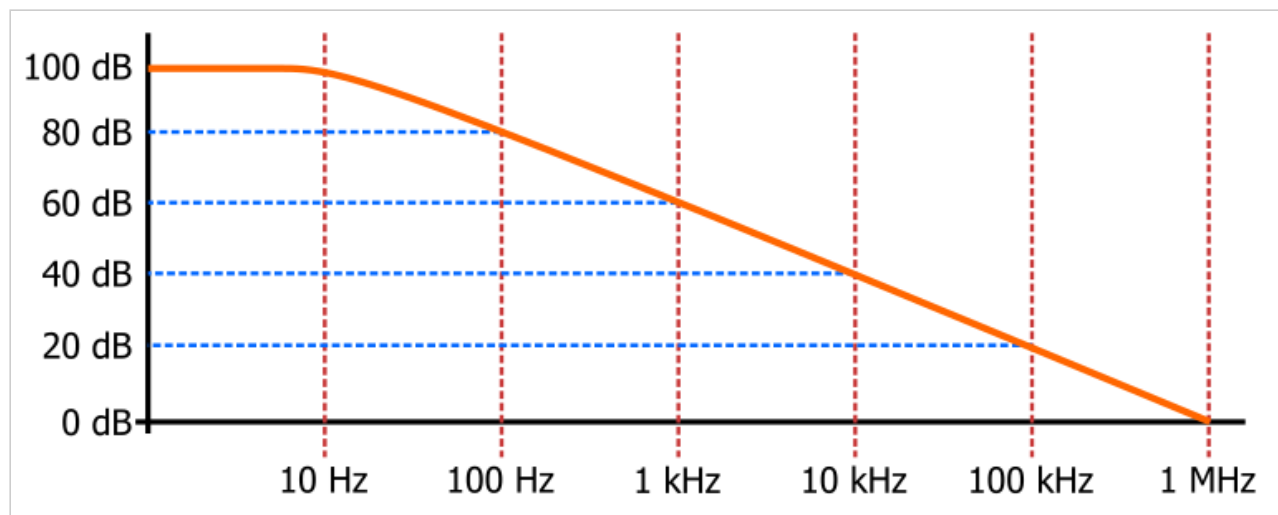
$$R_o = (R_1 \parallel R_2) \times \left( 1 + \frac{a}{1 + R_2 / R_1} \right)$$

*Equation 9*

As an example, using an op-amp with a DC gain of 100 dB (=100,000 V/V) will lower  $R_o$  from  $\infty$  to  $(10^3 \parallel 10^3) \times (1 + 100,000/2) \cong 25 \text{ M}\Omega$ . With an arrangement of the type of Figure 3b, we can use the potentiometer to compensate for the cumulative effect of bridge imbalance, non-infinite CMRR, and non-infinite open-loop DC gain, and raise  $R_o$  as close as possible to  $\infty$ .

However, as we increase the frequency of operation, the gain  $a$  rolls off with frequency, leading to a progressive deterioration of  $R_o$ . For example, if an op-amp with a DC gain of 100 dB has a gain-bandwidth product (<https://www.allaboutcircuits.com/technical-articles/negative-feedback-part-2-improving-gain-sensitivity-and-bandwidth/>) of 1 MHz, its open-loop gain vs. frequency (assuming a single-pole response) will look like this:





**Figure 5.** Single-pole frequency response of a 1 MHz op-amp with a DC open-loop gain of 100 dB.

Thus, the gain  $a$  drops to 60 dB ( $=1000$  V/V) at 1 kHz, and the value of  $R_o$  will drop to  $500 \times (1 + 1000/2) \cong 250$  k $\Omega$ . At 10 kHz  $R_o$  drops to  $500 \times (1 + 100/2) \cong 25$  k $\Omega$ , and so on.

## Further Reading

### A Comprehensive Study of the Howland Current Pump

(<https://www.ti.com/lit/an/snoa474a/snoa474a.pdf>) (PDF): an application note published by Texas Instruments.