

# THESIS

BY

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## Emission kernel of parton shower

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statement of originality

I hereby confirm that I have written the accompanying thesis by myself, without contributions from any sources other than those cited in the text and acknowledgements. This applies also to all graphics, drawings, maps and images included in the thesis.

Karlsruhe, 22. Januar 2019

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Tigran Saidnia



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## 0.1 parametrisation

$$\left. \begin{aligned}
 q_i^\mu &= zp_i^\mu + y(1-z)p_j^\mu + \sqrt{zy(1-z)}m_\perp \\
 q^\mu &= (1-z)p_i^\mu + yzp_j^\mu - \sqrt{zy(1-z)}m_\perp \\
 q_j^\mu &= (1-y)p_j^\mu \\
 y &= \frac{q_i q}{p_i p_j} \\
 q_i + q &= p_i + yp_j \\
 q_j + q &= (1-z)p_i^\mu + (1+yz-y)p_j^\mu - \sqrt{zy(1-z)}m_\perp \\
 q_i \cdot q &= y(1-2z+2z^2)(p_i \cdot p_j) \\
 q_i \cdot q_j &= z(1-y)(p_i \cdot p_j) \\
 q_j \cdot q &= (1-z)(1-y)(p_i \cdot p_j)
 \end{aligned} \right\} \text{parametrisation} \quad (1)$$

## 0.2 Altarelli-Parisi splitting functions

$$\left. \begin{aligned} \langle \hat{P}_{qq} \rangle &= C_F \left[ \frac{1+z^2}{1-z} - \varepsilon(1-z) \right] \\ \langle \hat{P}_{gq} \rangle &= T_R \left[ 1 - \frac{2z(1-z)}{1-\varepsilon} \right] \\ \langle \hat{P}_{qg} \rangle &= C_F \left[ \frac{1+(1-z)^2}{z} - \varepsilon z \right] \\ \langle \hat{P}_{gg} \rangle &= 2C_A \left[ \frac{z}{1-z} + \frac{1-z}{z} + z(1-z) \right] \end{aligned} \right\} \text{splitting functions} \quad (2)$$

### 0.3 Colour factor calculation

fundamental representation in  $SU(2)$  and  $SU(3)$

$$\begin{aligned} T^a &= \tau^a \equiv \frac{\sigma^a}{2} \quad \text{with Pauli matrices } \sigma^a \\ T^a &= \vartheta^a \equiv \frac{\lambda^a}{2} \quad \text{with Gell - Mann matrices } \lambda^a \end{aligned} \quad (3)$$

$$\begin{aligned} \lambda^1 &= \begin{pmatrix} 0 & 1 & \\ 1 & 0 & \\ & & 0 \end{pmatrix}, \quad \lambda^2 = \begin{pmatrix} 0 & -i & \\ i & 0 & \\ & & 0 \end{pmatrix}, \quad \lambda^3 = \begin{pmatrix} 1 & & \\ & -1 & \\ & & 0 \end{pmatrix}, \quad \lambda^4 = \begin{pmatrix} & & 1 \\ & 0 & \\ 1 & & \end{pmatrix} \\ \lambda^5 &= \begin{pmatrix} & -i & \\ & 0 & \\ i & & \end{pmatrix}, \quad \lambda^6 = \begin{pmatrix} 0 & & \\ & 0 & 1 \\ & 1 & 0 \end{pmatrix}, \quad \lambda^7 = \begin{pmatrix} 0 & & \\ & 0 & -i \\ & i & 0 \end{pmatrix}, \quad \lambda^8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & & \\ & 1 & \\ & & -2 \end{pmatrix} \end{aligned} \quad (4)$$

As we can see,  $\lambda^3$  and  $\lambda^8$  are diagonal. These generators satisfy:

$$[T^a, T^b] = i\epsilon^{abc}T^c \quad (5)$$

The most common convention for the normalization of the generators in physics is:

$$\sum_{c,d} f^{acd} f^{bcd} = N\delta^{ab} \quad (6)$$

The main relation we will use later for  $SU(N)$ :

$$\text{tr}(T^a T^b) = T_{ij}^a T_{ji}^b = T_F \delta^{ab} \quad (7)$$

$$\sum_a (T^a T^a) = C_F \delta^{ij} \quad (8)$$

$$f^{acd} f^{bcd} = C_A \delta^{ab} \quad (9)$$

With  $T_F = \frac{1}{2}$ ,  $C_A = N$  and  $C_F = \frac{N^2-1}{2N}$ .

$$f^{abc} = -2i \text{tr}(T^a [T^b, T^c]) \quad (10)$$

$$d^{abc} = 2 \text{tr}(T^a T^b, T^c) \quad (11)$$

$$T^a T^b = \frac{1}{2} \left( \frac{1}{N} \delta_{ab} + (d^{abc} + i f^{abc}) T^c \right) \quad (12)$$

$$\text{tr}(T^a T^b T^c) = \frac{1}{4} (d^{abc} + i f^{abc}) \quad (13)$$

$$\text{tr}(T^a T^b T^a T^c) = \frac{-1}{4N} \delta_{bc} \quad (14)$$

$$f^{acd} f^{bcd} = N \delta^{ab} \quad (15)$$

$$f^{acd} d^{bcd} = 0 \quad (16)$$

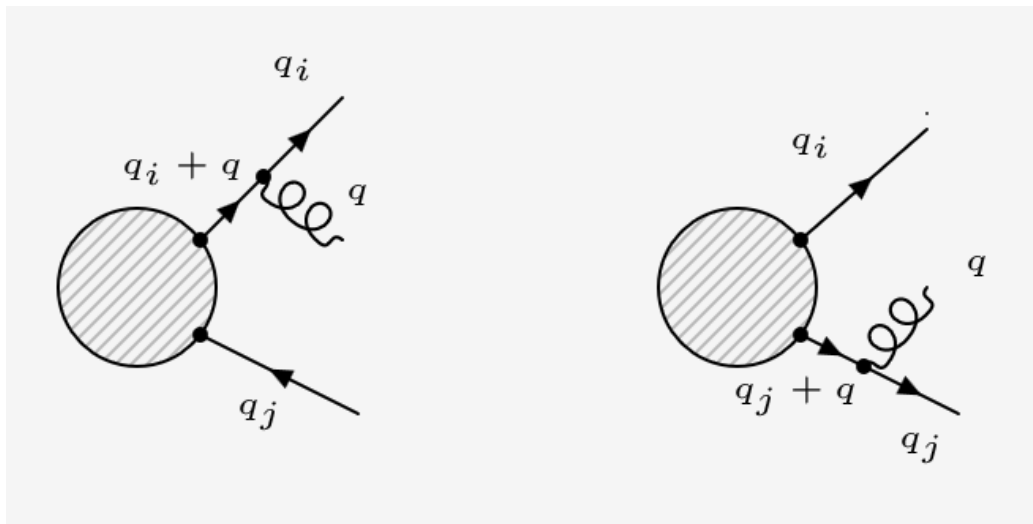
$$f^{ade} f^{bef} f^{cfd} = \frac{N}{2} f^{abc} \quad (17)$$

Fierz identity:

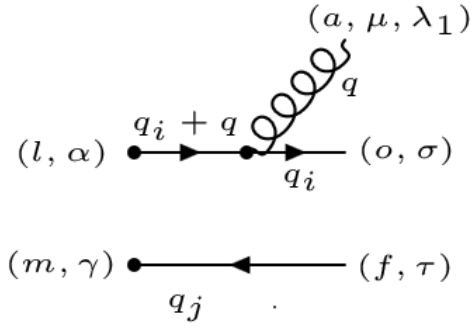
$$\sum_a T_{ij}^a T_{kl}^a = \frac{1}{2} (\delta_{il} \delta_{kj} - \frac{1}{N} \delta_{ij} \delta_{kl}) \quad (18)$$

# Kapitel 1

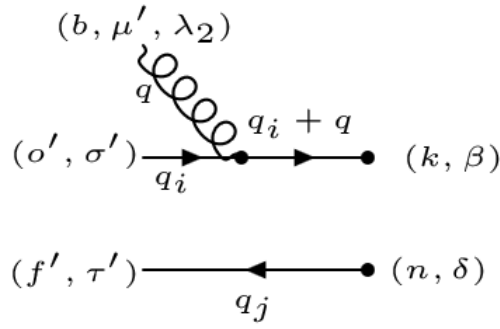
## Quark antiquark gluon emission kernel



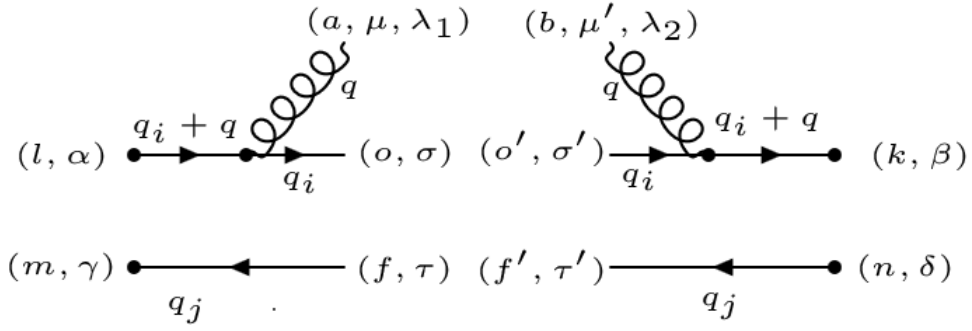


1.1  $qg\text{-}\bar{q}$ 

$$M_1 = [\bar{u}_\sigma(q_i)(-ig_s\gamma^\mu \times [T^a]_{o^l}^l) \frac{i(\not{q}_i + \not{q})}{(q_i + q)^2} \varepsilon^{\lambda_1}_\mu(q)] [v_\tau(q_j)] \quad (1.1)$$



$$M_1^\dagger = [\frac{-i(\not{q}_i + \not{q})}{(q_i + q)^2} (ig_s\gamma^{\mu'} \times [T^b]_{o'^k}^k) u_{\sigma'}(q_i) \varepsilon^{\lambda_2}_{\mu'}(q)] [\bar{v}_{\tau'}(q_j)] \quad (1.2)$$



$$|M_1|^2 = M_1 M_1^\dagger = [\bar{u}_\sigma(q_i) (-ig_s \gamma^\mu \times [T^a]_{o'}^l) \frac{i(\not{q}_i + \not{q})}{(q_i + q)^2} \varepsilon^{\lambda_1}_\mu(q) [v_\tau(q_j)]$$

$$[\frac{-i(\not{q}_i + \not{q})}{(q_i + q)^2} (ig_s \gamma^{\mu'} \times [T^b]_{o'}^k) u_{\sigma'}(q_i) \varepsilon^{\lambda_2}_{\mu'}^*(q) [\bar{v}_{\tau'}(q_j)]] \quad (1.3)$$

$$|M_1|^2 = [\frac{-i(\not{q}_i + \not{q})}{(q_i + q)^2} (ig_s \gamma^{\mu'} \times [T^b]_{o'}^k) \bar{u}_\sigma(q_i) u_{\sigma'}(q_i) \varepsilon^{\lambda_2}_{\mu'}^*(q) \varepsilon^{\lambda_1}_\mu(q)$$

$$\times (-ig_s \gamma^\mu \times [T^a]_{o'}^l) \frac{i(\not{q}_i + \not{q})}{(q_i + q)^2} [\bar{v}_{\tau'}(q_j) v_\tau(q_j)]] \quad (1.4)$$

and after sum over the lorenz index  $(\sigma, \sigma')$  and  $(\tau, \tau')$  and unsing the spin addition relation:

$$\sum_{\sigma, \sigma'} \bar{u}_\sigma(q_i) u_{\sigma'}(q_i) = \not{q}_i \delta^{\sigma\sigma'},$$

$$\sum_{\tau, \tau'} \bar{v}_\tau(q_j) v_{\tau'}(q_j) = \not{q}_j \delta^{\tau\tau'} \quad (1.5)$$

and sum over polarization index  $(\lambda_1, \lambda_2)$  :

$$\sum_{\mu, \mu'} \varepsilon^{\lambda_2}_{\mu'}^*(q) \varepsilon^{\lambda_1}_\mu(q) = -g_{\mu\mu'} \delta^{ab} \quad (1.6)$$

$$|M_1|^2 = \frac{-g_s^2 [T^a]_{o'}^k [T^a]_{o'}^l}{(q_i + q)^2 (q_i + q)^2} [(\not{q}_i + \not{q}) \gamma^{\mu'} \not{q}_i g_{\mu'\mu} \gamma^\mu (\not{q}_i + q)] [\not{q}_j] \quad (1.7)$$

from here and after contraction between all indices we can actually make statements about the last result.

$$|M_1|^2 = \frac{-g_s^2 [T^a]_{o'}^k [T^a]_{o'}^l}{(q_i + q)^2 (q_i + q)^2} [(\not{q}_i + \not{q}) \gamma^{\mu'} \not{q}_i \gamma_{\mu'} (\not{q}_i + q)] [\not{q}_j] \quad (1.8)$$

In other words we expect the tree level diagram from LO and a number: Which means:

$$|M^2| = \left| \begin{array}{c} \text{diagram with two shaded circles and momenta } P_i, P_j \\ \text{contribution from LO} \end{array} \right|^2 \otimes \left| \begin{array}{c} \text{diagram with a loop and momenta } q_i, q, q_i+q \\ \text{a complex number} \end{array} \right|^2$$

$$|M_1|^2 = \frac{-g_s^2 [T^a]_o^k [T^a]_o^l}{(q_i + q)^2 (q_i + q)^2} [P_i][P_j] \otimes (\text{a complex number}) \quad (1.9)$$

Let's calculate the contribution and compare the final result with this expectation:

$$\begin{aligned} N &=: \gamma^{\mu'} \not{q}_i \gamma_{\mu'} = q_{i\sigma} \gamma^{\mu'} \gamma^\sigma \gamma_{\mu'} \\ &= q_{i\sigma} (\{\gamma^{\mu'}, \gamma^\sigma\} - \gamma^\sigma \gamma^{\mu'}) \gamma_{\mu'} \\ &= q_{i\sigma} 2g^{\mu'\sigma} \gamma_{\mu'} - d \gamma^\sigma \\ &= (2 - d) \not{q}_i \end{aligned} \quad (1.10)$$

$$|M_1|^2 = -(2 - d) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{(q_i + q)^2 (q_i + q)^2} [(\not{q}_i + \not{q}) \not{q}_i (\not{q}_i + q)] [\not{q}_j] \quad (1.11)$$

$$|M_1|^2 = -(2 - d) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{(q_i + q)^2 (q_i + q)^2} [\not{q}_i \not{q}_i \not{q}_i + \not{q}_i \not{q}_i \not{q} + \not{q} \not{q}_i \not{q}_i + \not{q} \not{q}_i \not{q}] [\not{q}_j] \quad (1.12)$$

For the momenta are on-shell which means:

$$\begin{aligned} \not{q}_i \not{q}_i &= q_i^2 = m_i^2 \\ \not{q} \not{q} &= q^2 = m^2 \\ \not{q}_j \not{q}_j &= q_j^2 = m_j^2 \end{aligned} \quad (1.13)$$

we can first neglect the mass of patrons and ignore each term with  $\not{q}_i \not{q}_i$  and  $\not{q} \not{q}$  as well.

$$|M_1|^2 = -(2 - d) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{(2q_i q)(2q_i q)} [\not{q} \not{q}_i \not{q}] [\not{q}_j] \quad (1.14)$$

$$\begin{aligned}
L &= \not{q}_i \not{q}_\mu (\{\gamma^\mu, \gamma^\sigma\} - \gamma^\sigma \gamma^\mu) \\
&= \not{q}_i [2q_i^\mu q_\mu - q_{i\sigma} q_\mu \gamma^\mu \gamma^\sigma] \\
&= \not{q}_i (2q_i q) - q_\mu q_{i\sigma} q_\mu [\gamma^\mu \gamma^\mu \gamma^\sigma] \\
&= \not{q}_i (2q_i q) - q_\mu q_{i\sigma} q_\mu \left[ \frac{\gamma^\mu \gamma^\mu}{2} + \frac{\gamma^\mu \gamma^\mu}{2} \right] \gamma^\sigma \\
&= \not{q}_i (2q_i q) - q_\mu q_{i\sigma} q_\mu [g^{\mu\mu}] \gamma^\sigma \\
&= \not{q}_i (2q_i q) - q_\mu q_{i\sigma} q^\mu \gamma^\sigma \\
&= \not{q}_i (2q_i q) - q^2 \not{q}_i \\
&= \not{q}_i (2q_i q)
\end{aligned} \tag{1.15}$$

After inserting the last result of  $L$  and simplify the term  $(2q_i q)$  from the denominator and nominator because , we get:

$$|M_1|^2 = -(2-d) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{q}_i] [\not{q}_j] \tag{1.16}$$

Now we are going to use the parametrisation from equation (1) to reduce the 3-member matrix element to 2-member and take out the singularity term from the amplitude.

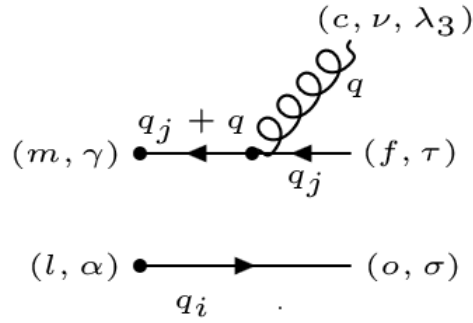
$$|M_1|^2 = (d-2) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{2y(1-2z+2z^2)(p_i \cdot p_j)} [(1-z) \not{p}_i + zy \not{p}_j - \sqrt{zy(1-z)} \not{m}_\perp] [(1-y) \not{p}_j] \tag{1.17}$$

Multiplying the both sides

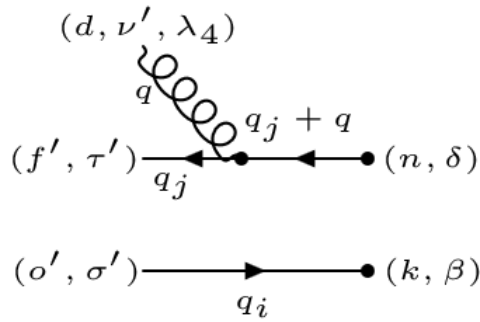
$$\begin{aligned}
|M_1|^2 &= (d-2) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{2y(1-2z+2z^2)(p_i \cdot p_j)} [(1-z)(1-y) \not{p}_i \not{p}_j \\
&\quad + zy(1-y) \not{p}_j \not{p}_j + (1-y) \sqrt{zy(1-z)} \not{m}_\perp \not{p}_j]
\end{aligned} \tag{1.18}$$

and under consideration of the fact that  $p_i$  and  $p_j$  are the on-shell momenta of the emitter and spectator partons, we can ignore the terms with  $\not{p}_i \not{p}_i$  and  $\not{p}_j \not{p}_j$ . The  $p_i \cdot m_\perp$  and  $p_j \cdot m_\perp$  are always 0 because the  $p_i$  and  $p_j$  are lightlike, i.e. zero transverse component. So those terms can be neglected.

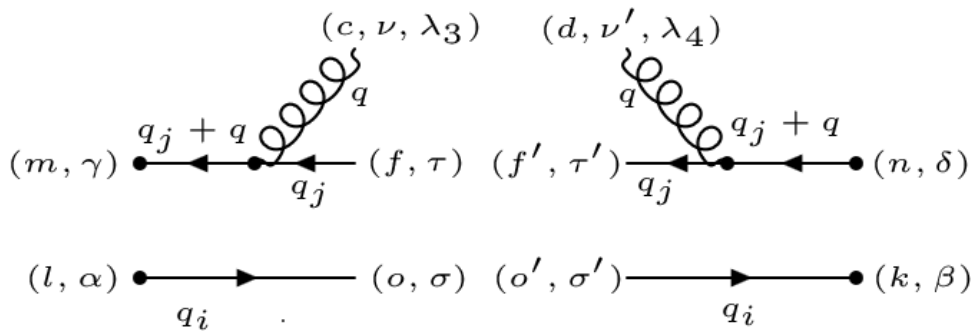
$$|M_1|^2 = (d-2)(1-z)(1-y) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i] [\not{p}_j] \tag{1.19}$$

1.2  $\bar{q}g$ -q

$$M_2 = \left[ \frac{i(\not{q}_j + \not{q})}{(q_j + q)^2} (-ig_s \gamma^\nu \times [T^c]_f^m) v_\tau(q_j) \varepsilon^{\lambda_3}_\nu(q) [u_\sigma(q_i)] \right] \quad (1.20)$$



$$M_2^\dagger = [\bar{v}_{\tau'}(q_j) (ig_s \gamma^{\nu'} \times [T^d]_{f'}^n) \frac{-i(\not{q}_j + \not{q})}{(q_j + q)^2} \varepsilon^{\lambda_4}_{\nu'}(q) [\bar{u}_{\sigma'}(q_i)]] \quad (1.21)$$



$$|M_2|^2 = M_2 M_2^\dagger = \left[ \frac{i(\not{q}_j + \not{q})}{(q_j + q)^2} (-ig_s \gamma^\nu \times [T^c]_f^m) v_\tau(q_j) \varepsilon^{\lambda_3}_\nu(q) [u_\sigma(q_i)] \right] \quad (1.22)$$

$$\left[ \bar{v}_{\tau'}(q_j) (ig_s \gamma^{\nu'} \times [T^d]_{f'}^n) \frac{-i(\not{q}_j + \not{q})}{(q_j + q)^2} \varepsilon^{\lambda_4}_{\nu'}(q) [\bar{u}_{\sigma'}(q_i)] \right]$$

$$|M_2|^2 = \frac{g_s^2 [T^c]_f^m [T^d]_{f'}^n}{(q_j + q)^2 (q_j + q)^2} [(\not{q}_j + \not{q}) \gamma^\nu v_\tau(q_j) \bar{v}_{\tau'}(q_j) \varepsilon^{\lambda_3}_\nu(q) \varepsilon^{\lambda_4}_{\nu'}(q) \gamma^{\nu'} (\not{q}_j + \not{q})] \quad (1.23)$$

$$[u_\sigma(q_i)] [\bar{u}_{\sigma'}(q_i)]$$

and after sum over the lorenz index  $(\sigma, \sigma')$  and  $(\tau, \tau')$  and unsing the spin addition relation:

$$\sum_{\sigma, \sigma'} \bar{u}_\sigma(q_i) u_{\sigma'}(q_i) = \not{q}_i \delta^{\sigma\sigma'}, \quad (1.24)$$

$$\sum_{\tau, \tau'} \bar{v}_\tau(q_j) v_{\tau'}(q_j) = \not{q}_j \delta^{\tau\tau'}$$

and sum over polarization index  $(\lambda_3, \lambda_4)$  :

$$\sum_{\nu, \nu'} \varepsilon^{\lambda_4*}_{\nu'}(q) \varepsilon^{\lambda_3}_\nu(q) = -g_{\nu\nu'} \delta^{cd} \quad (1.25)$$

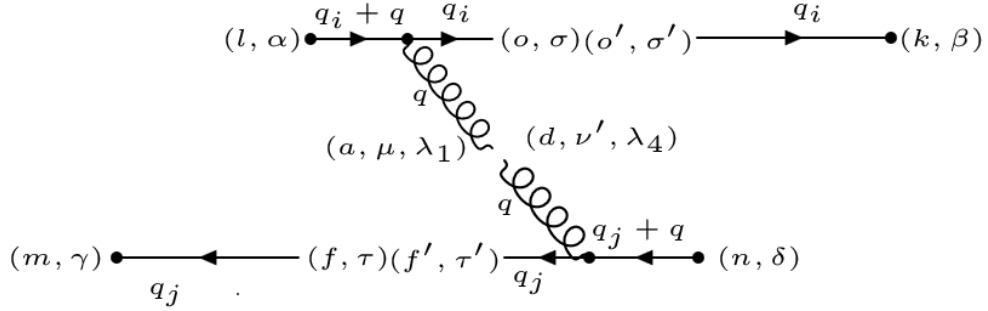
$$|M_2|^2 = \frac{g_s^2 [T^c]_f^m [T^c]_{f'}^n}{(q_j + q)^2 (q_j + q)^2} [(\not{q}_j + \not{q}) \gamma^\nu \not{q}_j (-g_{\nu\nu'}) \gamma^{\nu'} (\not{q}_j + \not{q})] [\not{q}_i] \quad (1.26)$$

After the same calculation from the last part, we'll get:

$$|M_2|^2 = (d - 2) \frac{g_s^2 [T^c]_f^m [T^c]_{f'}^n}{(2qq_j)} [\not{q}] [\not{q}_i] \quad (1.27)$$

finally:

$$|M_2|^2 = -(d - 2) y z^2 \frac{g_s^2 [T^c]_f^m [T^c]_{f'}^n}{2(1 - z)(1 - y)(p_i \cdot p_j)} [\not{p}_i] [\not{p}_j] \quad (1.28)$$

1.3  $M_1 M_2^\dagger$ 

$$M_1 M_2^\dagger = [\bar{u}_\sigma(q_i) (-ig_s \gamma^\mu \times [T^a]_o^l) \frac{i(\not{q}_i + \not{q})}{(q_i + q)^2} \varepsilon^{\lambda_1}_\mu(q)] [v_\tau(q_j)]$$

$$[\bar{v}_{\tau'}(q_j) (ig_s \gamma^{\nu'} \times [T^d]_{f'}^n) \frac{-i(\not{q}_j + \not{q})}{(q_j + q)^2} \varepsilon^{\lambda_4}_{\nu'}(q)] [u_{\sigma'}(q_i)] \quad (1.29)$$

$$M_1 M_2^\dagger = \frac{g_s^2 [T^a]_o^l [T^d]_{f'}^n}{(2q_i q)(2q_j q)} [\not{q}_i \gamma^\mu (\not{q}_i + \not{q})] \varepsilon^{\lambda_1}_\mu(q) \varepsilon^{\lambda_4}_{\nu'}(q)$$

$$[\not{q}_j \gamma^{\nu'} (\not{q}_j + \not{q})] \quad (1.30)$$

$$M_1 M_2^\dagger = \frac{g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(2q_i q)(2q_j q)} [\not{q}_i \gamma^\mu (\not{q}_i + \not{q})] - g_{\mu\nu'}$$

$$[\not{q}_j \gamma^{\nu'} (\not{q}_j + \not{q})] \quad (1.31)$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(2q_i q)(2q_j q)} [\not{q}_i \gamma^\mu (\not{q}_i + \not{q})] [\not{q}_j \gamma_\mu (\not{q}_j + \not{q})] \quad (1.32)$$

Expectation:

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(2q_i q)(2q_j q)} [(\not{q}_i + \not{q}) \gamma^\mu \not{q}_i] [(\not{q}_j + \not{q}) \gamma_\mu \not{q}_j] \quad (1.33)$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(2q_i q)(2q_j q)} [-(\not{q}_i + \not{q}) \not{q}_i \gamma^\mu + 2(\not{q}_i + \not{q}) q_i^\mu]$$

$$[-(\not{q}_j + \not{q}) \not{q}_j \gamma_\mu + 2(\not{q}_j + \not{q}) q_{j\mu}] \quad (1.34)$$

$$|M^2| = \left| \begin{array}{c} \text{diagram with two shaded circles connected by two horizontal lines labeled } P_i \text{ and } P_j \\ \text{contribution from LO} \end{array} \right|^2 \otimes \left| \begin{array}{c} \text{diagram with a shaded circle and a wavy line} \\ \text{a complex number} \end{array} \right|^2$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(2q_i q)(2q_j q)} \begin{aligned} & [(\not{q}_i + \not{q}) \not{q}_i \gamma^\mu][(\not{q}_j + \not{q}) \not{q}_j \gamma_\mu] \\ & - 2[(\not{q}_i + \not{q}) \not{q}_i \gamma^\mu][(\not{q}_j + \not{q}) q_{j\mu}] \\ & - 2[(\not{q}_i + \not{q}) q_i^\mu][(\not{q}_j + \not{q}) \not{q}_j \gamma_\mu] \\ & + 4[(\not{q}_i + \not{q}) q_i^\mu][(\not{q}_j + \not{q}) q_{j\mu}] \end{aligned} \quad (1.35)$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(2q_i q)(2q_j q)} \begin{aligned} & [\not{q} \not{q}_i \gamma^\mu][\not{q} \not{q}_j \gamma_\mu] \\ & - 2[\not{q} \not{q}_i \gamma^\mu][(\not{q} + \not{q}_j) q_{j\mu}] \\ & - 2[(\not{q}_i + \not{q}) q_i^\mu][\not{q} \not{q}_j \gamma_\mu] \\ & + 4[(\not{q}_i + \not{q}) q_i^\mu][(\not{q}_j + \not{q}) q_{j\mu}] \end{aligned} \quad (1.36)$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{4(1-z)(1-y)y(1-2z+2z^2)(p_i \cdot p_j)(p_i \cdot p_j)} \begin{aligned} & [y(1-2z+2z^2) \not{p}_i \not{p}_j \gamma^\mu][(1-z)(1-y) \not{p}_i \not{p}_j \gamma_\mu] \\ & + 4(q_i^\mu \cdot q_{j\mu})[(\not{q}_i + \not{q})][(\not{q}_j + \not{q})] \end{aligned} \quad (1.37)$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{4(1-z)(1-y)y(1-2z+2z^2)(p_i \cdot p_j)(p_i \cdot p_j)} \begin{aligned} & [y(1-2z+2z^2) \not{p}_i \not{p}_j \gamma^\mu][(1-z)(1-y) \not{p}_i \not{p}_j \gamma_\mu] \\ & + 4(p_i \cdot p_j)[(\not{p}_i + y \not{p}_j)][(1-z) \not{p}_i + (1+yz-y) \not{p}_j - \sqrt{zy(1-z)} \not{m}] \end{aligned} \quad (1.38)$$

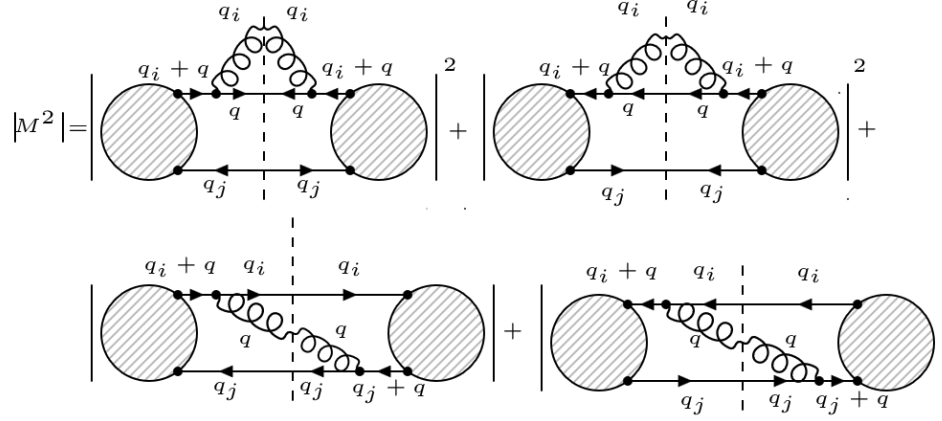
$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(1-z)(1-y)y(1-2z+2z^2)(p_i \cdot p_j)} z(1-y)[\not{p}_i][\not{p}_j] \quad (1.39)$$

$$M_1 M_2^\dagger = \frac{-g_s^2 [T^a]_o^l [T^a]_{f'}^n}{(1-z)y(1-2z+2z^2)(p_i \cdot p_j)} z[\not{p}_i][\not{p}_j] \quad (1.40)$$

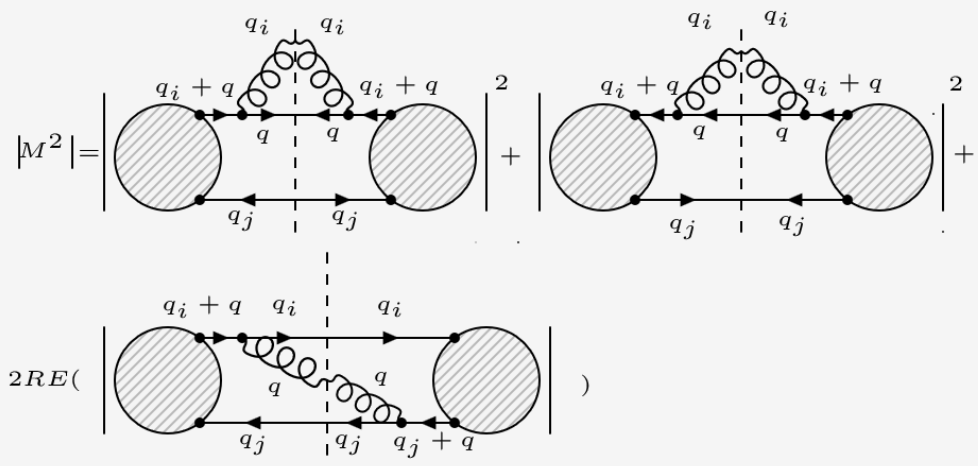


1.4  $|M^2|$ 

$$|M|^2 = |M_1|^2 + |M_2|^2 + M_1 M_2^\dagger + M_1^\dagger M_2 \quad (1.41)$$



$$|M|^2 = |M_1|^2 + |M_2|^2 + 2RE(M_1 M_2^\dagger) \quad (1.42)$$



$$\begin{aligned}
 |M|^2 = & (d-2)(1-z)(1-y) \frac{g_s^2 [T^a]_o^k [T^a]_o^l}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\
 & - (d-2)yz^2 \frac{g_s^2 [T^c]_f^m [T^c]_f^n}{2(1-z)(1-y)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\
 & + 2RE\left(\left(\frac{-2z}{z-1}\right) \frac{g_s^2 [T^a]_o^l [T^a]_f^n}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j]\right)
 \end{aligned} \quad (1.43)$$

$$T^a_{ok} T^a_{lo} = \frac{1}{2}(\delta_{oo}\delta_{lk} - \frac{1}{N}\delta_{ok}\delta_{lo}) = \frac{1}{2}(N\delta_{lk} - \frac{1}{N}\delta_{lk}) = C_F\delta_{lk} \quad (1.44)$$

After summation over the final colour states and averaging over initial colour states we get:

$$T^a_{ok} T^a_{lo} = C_F \delta_{lk} = \frac{1}{N} \sum_{l=1}^N \delta_{lk} C_F = C_F \quad (1.45)$$

The same calculation for  $T^c_{mf} T^c_{fn}$  and  $T^a_{ol} T^a_{fn}$  turns  $C_F$  out as the colour factor. Now we are going to compute the splitting function in the case of the colinearity, which means, if:

$$y \longrightarrow 0 \quad (1.46)$$

$$\begin{aligned} |M|^2 = & (d-2)(1-z)(1-y) \frac{g_s^2 C_F}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\ & - (d-2)yz^2 \frac{g_s^2 C_F}{2(1-z)(1-y)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\ & + 2RE\left(\left(\frac{-2z}{z-1}\right) \frac{g_s^2 C_F}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j]\right) \end{aligned} \quad (1.47)$$

$$|M|^2 = C_F \left( (d-2)(1-z) - \frac{4z}{z-1} \right) \frac{g_s^2}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \quad (1.48)$$

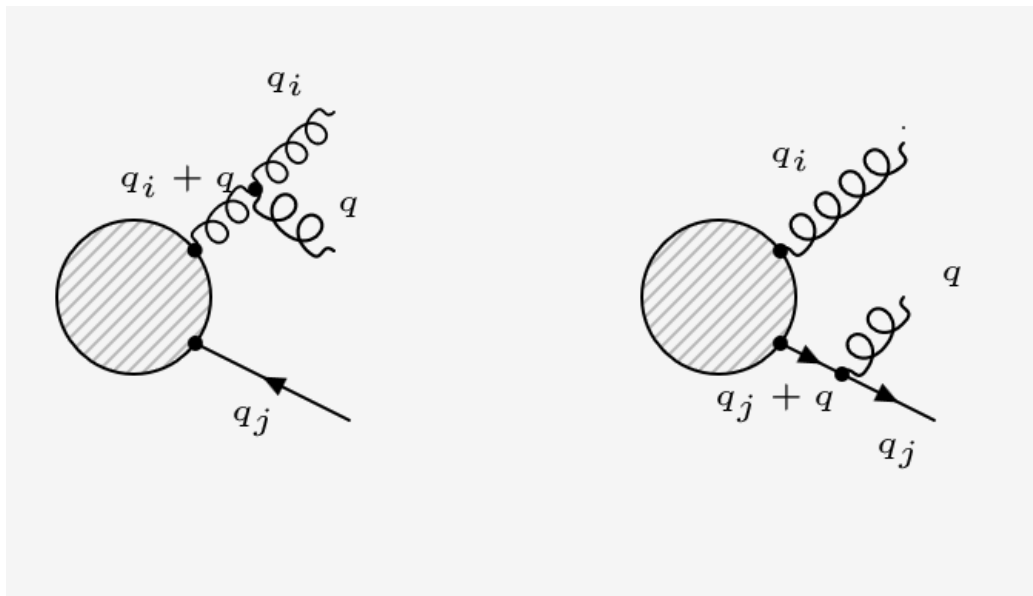
for

$$d = 4 - 2\epsilon \quad (1.49)$$

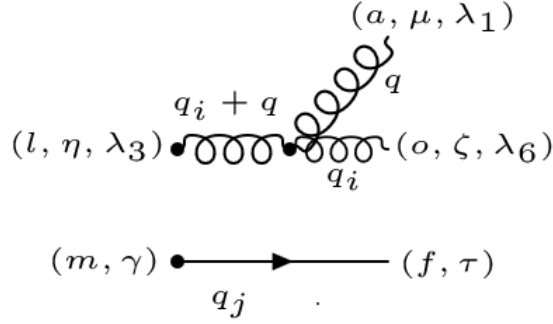
$$\begin{aligned} |M|^2 = & C_F \left( (4-2\epsilon-2)(1-z) + \frac{4z}{1-z} \right) \frac{g_s^2}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\ = & C_F \left( \frac{2(1-\epsilon)(1-z)^2 + 4z}{1-z} \right) \frac{g_s^2}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\ = & C_F \left( \frac{2-4z+2z^2-\epsilon(1-z)^2+4z}{1-z} \right) \frac{g_s^2}{2y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\ = & C_F \left( \frac{(1+z^2)}{1-z} - \epsilon(1-z) \right) \frac{g_s^2}{y(1-2z+2z^2)(p_i \cdot p_j)} [\not{p}_i][\not{p}_j] \\ = & \langle \hat{P}_{qq} \rangle \frac{g_s^2}{q_i \cdot q} [\not{p}_i][\not{p}_j] \end{aligned} \quad (1.50)$$

## Kapitel 2

### Gluon quark gluon emission kernel

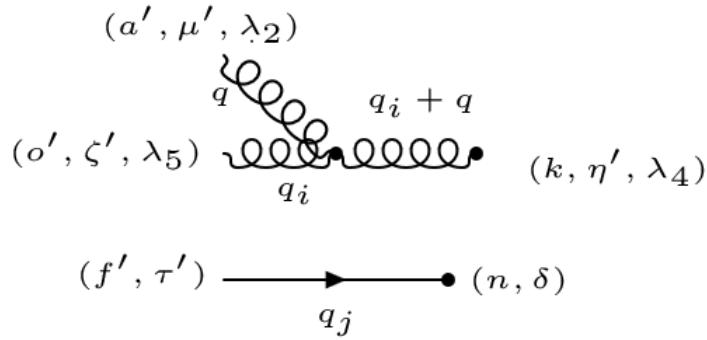


## 2.1 gg-q

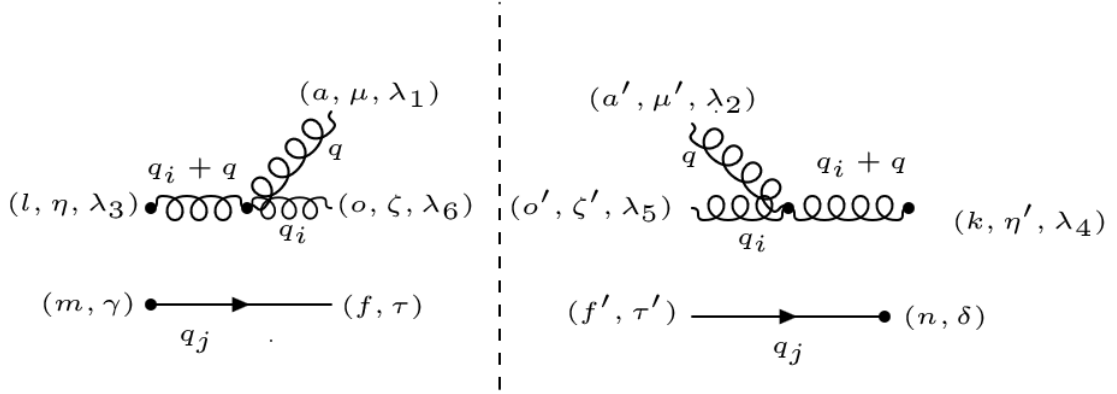


$$M_1 = \left[ \frac{-i}{(q_i + q)^2} (-g_s f^{a o l} (g^{\mu \zeta} (-q_i + q)^\eta + g^{\zeta \eta} (-q - (q_i + q))^\mu + g^{\eta \mu} (q_i + q + q_i)^\zeta) \right. \\ \left. \varepsilon^{\lambda_1}_{\mu}(q) \varepsilon^{\lambda_6}_{\zeta}(q_i) \right] [\bar{u}_\tau(q_j)] \quad (2.1)$$

$$M_1 = \left[ \frac{-i}{(q_i + q)^2} (-g_s f^{a o l} (g^{\mu \zeta} (q - q_i)^\eta - g^{\zeta \eta} (2q_i + q)^\mu + g^{\eta \mu} (2q_i + q)^\zeta) \right. \\ \left. \varepsilon^{\lambda_1}_{\mu}(q) \varepsilon^{\lambda_6}_{\zeta}(q_i) \right] [\bar{u}_\tau(q_j)] \quad (2.2)$$



$$M_1^\dagger = \left[ \frac{i}{(q_i + q)^2} (-g_s f^{a' o' k} (g^{\mu' \zeta'} (q - q_i)^{\eta'} - g^{\zeta' \eta'} (2q_i + q)^{\mu'} + g^{\eta' \mu'} (2q_i + q)^{\zeta'}) \right. \\ \left. \varepsilon^{\lambda_2}_{\mu'}(q) \varepsilon^{\lambda_5}_{\zeta'}(q_i) \right] [u_{\tau'}(q_j)] \quad (2.3)$$



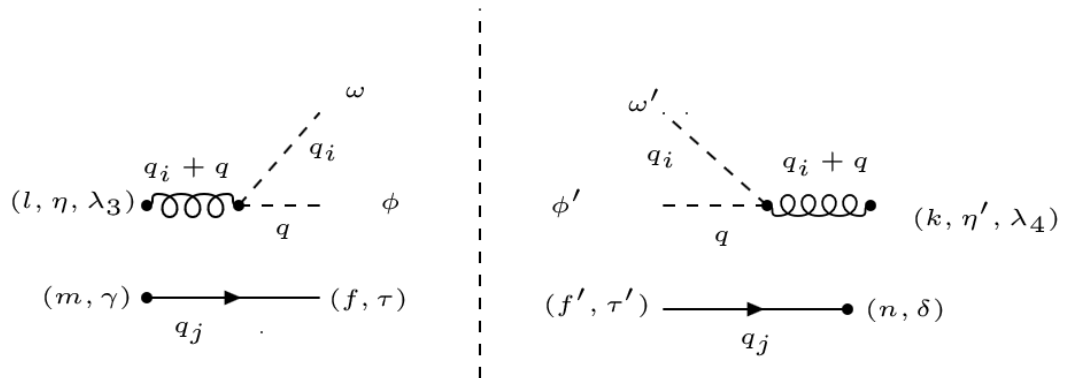
$$\begin{aligned}
 |M_1|^2 = & \left[ \frac{-i}{(q_i + q)^2} (-g_s f^{a o l} (g^{\mu \zeta} (q - q_i)^\eta - g^{\zeta \eta} (2q_i + q)^\mu + g^{\eta \mu} (2q_i + q)^\zeta) \right. \\
 & \left. \varepsilon^{\lambda_1}_{\mu}(q) \varepsilon^{\lambda_2}_{\mu'}(q) \varepsilon^{\lambda_6}_{\zeta}(q_i) \varepsilon^{\lambda_5}_{\zeta'}(q_i) \right. \\
 & \left. (-g_s f^{a' o' k} (g^{\mu' \zeta'} (q - q_i)^{\eta'} - g^{\zeta' \eta'} (2q_i + q)^{\mu'} + g^{\eta' \mu'} (2q_i + q)^{\zeta'}) \frac{i}{(q_i + q)^2} [\bar{u}_\tau(q_j) u_{\tau'}(q_j)] \right] \quad (2.4)
 \end{aligned}$$

$$\begin{aligned}
 |M_1|^2 = & \frac{g_s^2 f^{a o l} f^{a' o' k}}{(q_i + q)^2 (q_i + q)^2} \\
 & [(g^{\mu \zeta} (q - q_i)^\eta - g^{\zeta \eta} (2q_i + q)^\mu + g^{\eta \mu} (2q_i + q)^\zeta) \\
 & g_{\mu \mu'} g^{\zeta \zeta'} (g_{\mu' \zeta'} (q - q_i)_{\eta'} - g_{\zeta' \eta'} (2q_i + q)_{\mu'} + g_{\eta' \mu'} (2q_i + q)_{\zeta'})] [\not{A}_j] \quad (2.5)
 \end{aligned}$$

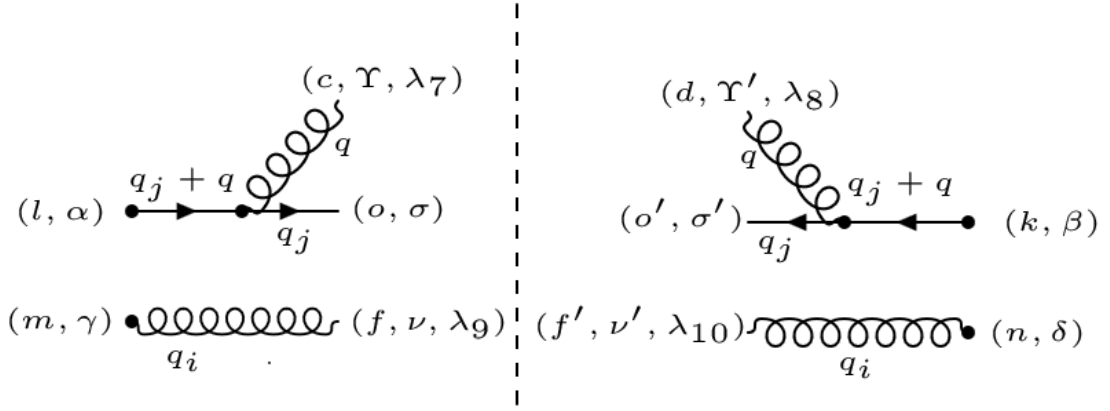
$$\begin{aligned}
 |M_1|^2 = & \frac{g_s^2 f^{a o l} f^{a' o' k}}{(q_i + q)^2 (q_i + q)^2} \\
 & [(g^{\mu' \zeta} (q - q_i)^\eta - g^{\zeta \eta} (2q_i + q)_{\mu'} + g^{\eta}_{\mu'} (2q_i + q)^\zeta) \\
 & (g^{\eta}_{\mu'} (q - q_i)_{\eta'} - g^{\zeta}_{\eta'} (2q_i + q)_{\mu'} + g_{\eta' \mu'} (2q_i + q)^\zeta)] [\not{A}_j] \quad (2.6)
 \end{aligned}$$

$$\begin{aligned}
 |M_1|^2 = & \frac{g_s^2 f^{a o l} f^{a' o' k}}{(q_i + q)^2 (q_i + q)^2} \\
 & [g^{\mu' \zeta} (q - q_i)^\eta g^{\eta}_{\mu'} (q - q_i)_{\eta'} - g^{\mu' \zeta} (q - q_i)^\eta g^{\zeta}_{\eta'} (2q_i + q)_{\mu'} \\
 & + g^{\mu' \zeta} (q - q_i)^\eta g_{\eta' \mu'} (2q_i + q)^\zeta - g^{\zeta \eta} (2q_i + q)_{\mu'} g^{\eta}_{\mu'} (q - q_i)_{\eta'} \\
 & + g^{\zeta \eta} (2q_i + q)_{\mu'} g^{\zeta}_{\eta'} (2q_i + q)_{\mu'} - g^{\zeta \eta} (2q_i + q)_{\mu'} g_{\eta' \mu'} (2q_i + q)^\zeta \\
 & + g^{\eta}_{\mu'} (2q_i + q)^\zeta g^{\eta}_{\mu'} (q - q_i)_{\eta'} - g^{\eta}_{\mu'} (2q_i + q)^\zeta g^{\zeta}_{\eta'} (2q_i + q)_{\mu'} \\
 & + g^{\eta}_{\mu'} (2q_i + q)^\zeta g_{\eta' \mu'} (2q_i + q)^\zeta] [\not{A}_j] \quad (2.7)
 \end{aligned}$$

## 2.1.1 One-loop corrections to the gluon self-energy diagram



## 2.2 qg-g



$$|M_2|^2 = M_2 M_2^\dagger = [\bar{u}_\sigma(q_j) (-ig_s \gamma^\Upsilon \times [T^c]_o^l) \frac{i(\not{q}_j + \not{q})}{(q_j + q)^2} \varepsilon^{\lambda_7}_\Upsilon(q)] [\varepsilon^{\lambda_9}_\nu(q_i)]$$

$$[\frac{-i(\not{q}_j + \not{q})}{(q_j + q)^2} (ig_s \gamma^{\Upsilon'} \times [T^d]_{o'}^k) u_{\sigma'}(q_j) \varepsilon^{\lambda_8}_{\Upsilon'}(q)] [\varepsilon^{\lambda_{10}}_{\nu'}(q_i)] \quad (2.8)$$

$$|M_2|^2 = [\frac{-i(\not{q}_j + \not{q})}{(q_j + q)^2} (-ig_s \gamma^\Upsilon \times [T^c]_o^l) \bar{u}_\sigma(q_j) u_{\sigma'}(q_j) \varepsilon^{\lambda_7}_\Upsilon(q) \varepsilon^{\lambda_8}_{\Upsilon'}(q)]$$

$$\times (ig_s \gamma^{\Upsilon'} \times [T^d]_{o'}^k) \frac{i(\not{q}_j + \not{q})}{(q_j + q)^2} [\varepsilon^{\lambda_{10}}_{\nu'}(q_i) \varepsilon^{\lambda_9}_\nu(q_i)] \quad (2.9)$$

$$|M_2|^2 = \frac{g_s^2 [T^c]_o^l [T^d]_{o'}^k}{(q_j + q)^2 (q_j + q)^2} [(\not{q}_j + \not{q}) \gamma^\Upsilon \not{q}_j (-g_{\Upsilon\Upsilon'})$$

$$\gamma^{\Upsilon'} (\not{q}_j + \not{q})] [-g_{\nu'\nu}] \quad (2.10)$$

Expect:

$$|M_2|^2 = \left| \text{diagram} \right|^2 \otimes \left| \text{diagram} \right|^2$$

contribution from LO                      a complex number

The diagram on the left shows two shaded circles representing hard subprocesses. A quark line with momentum \$p\_i\$ connects the top of the two circles. A gluon line with momentum \$q\_i\$ connects the bottom of the two circles. The gluon line is labeled \$g\_{\nu\nu'}\$. The diagram on the right shows a quark line with momentum \$q\_i\$ entering a vertex, from which a gluon with momentum \$q\$ is emitted. The quark line continues with momentum \$q\_i + q\$ to a second vertex, from which a gluon with momentum \$q\_i\$ is emitted. The quark line then continues with momentum \$q\_i + q\$ to a third vertex, from which a gluon with momentum \$q\$ is emitted.

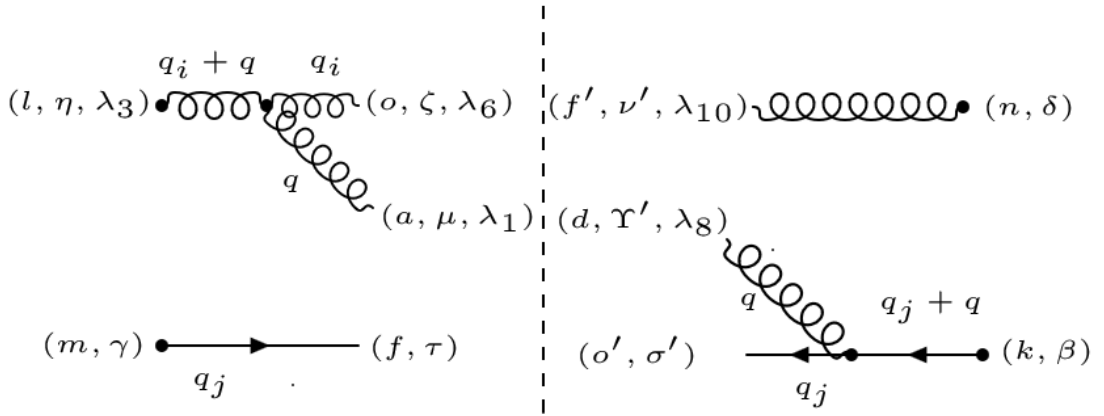
$$|M_2|^2 = -\frac{g_s^2 [T^c]_o^l [T^d]_{o'}^k}{(q_j + q)^2 (q_j + q)^2} [(\not{q}_j + \not{q}) \gamma^\Upsilon \not{q}_j \gamma^\Upsilon (\not{q}_j + \not{q})] [-g_{\nu'\nu}] \quad (2.11)$$

$$|M_2|^2 = (d-2) \frac{g_s^2 [T^c]_o^l [T^d]_{o'}^k}{(2qq_j)} [\not{q}] [-g_{\nu'\nu}] \quad (2.12)$$

$$|M_2|^2 = (d-2) \frac{g_s^2 [T^c]_o^l [T^d]_{o'}^k}{2(1-z)(1-y)(p_i \cdot p_j)} [(1-z) \not{p}_i + zy \not{p}_j - \sqrt{zy(1-z)} \not{m}_\perp] [-g_{\nu'\nu}] \quad (2.13)$$

$$|M_2|^2 = (d-2)zy \frac{g_s^2 [T^c]_o^l [T^d]_{o'}^k}{2(1-z)(1-y)(p_i \cdot p_j)} [\not{p}_j] [-g_{\nu'\nu}] \quad (2.14)$$

### 2.3 $M_1 M_2^\dagger$



$$M_1 M_2^\dagger = \left[ \frac{-ig_\eta \zeta}{(q_i + q)^2} (-g_s f^{aol} (g^{\mu\zeta} (q - q_i)^\eta - g^{\zeta\eta} (2q_i + q)^\mu + g^{\eta\mu} (2q_i + q)^\zeta) \right. \\ \left. \varepsilon^{\lambda_1}_\mu(q_i) \varepsilon^{\lambda_6}_\zeta(q) [\bar{u}_\tau(q_j)] \left[ \frac{-i(\not{q}_j + \not{q})}{(q_j + q)^2} (ig_s \gamma^{\mu'} \times [T^d]_{o'}^k) u_{\sigma'}(q_j) \varepsilon^{\lambda_8}_{\Upsilon'}(q) [\varepsilon^{\lambda_{10}}_{\nu'}(q_i)] \right] \right] \quad (2.15)$$

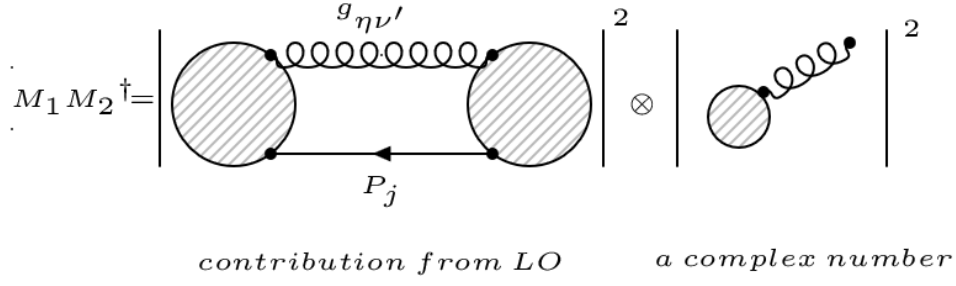
$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(qq_i)(qq_j)} [(g^{\mu\zeta} (q - q_i)^\eta - g^{\zeta\eta} (2q_i + q)^\mu + g^{\eta\mu} (2q_i + q)^\zeta)] \\ (-g_{\zeta\nu'}) (-g_{\mu\Upsilon'}) [(\not{q}_j + \not{q}) \gamma^{\Upsilon'} \not{q}_j] \quad (2.16)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(qq_i)(qq_j)} [(g^{\mu\nu'} (q - q_i)^\eta - g^{\eta\nu'} (2q_i + q)^\mu + g^{\eta\mu} (2q_i + q)_{\nu'})] \\ [(\not{q}_j + \not{q}) \gamma_\mu \not{q}_j] \quad (2.17)$$

Expect:







$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(qq_i)(qq_j)} [-g^{\eta}_{\nu'}(2q_i + q)^\mu] [-(\not{q}_j + \not{q}) \not{q}_j \gamma_\mu + 2(\not{q}_j + \not{q}) q_{j\mu}] \quad (2.18)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(qq_i)(qq_j)} [-g^{\eta}_{\nu'}(2q_i + q)^\mu] [-(qq_j) \gamma_\mu + 2(\not{q}_j + \not{q}) q_{j\mu}] \quad (2.19)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(qq_i)(qq_j)} [(qq_j) g^{\eta}_{\nu'}(2 \not{q}_i + \not{q}) - 2((2q_i + q) \cdot q_j) g^{\eta}_{\nu'}(\not{q}_j + \not{q})] \quad (2.20)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)(1-y)y(1-2z+z^2)(p_i \cdot p_j)(p_i \cdot p_j)} [g^{\eta}_{\nu'}][(qq_j)(2 \not{q}_i + \not{q}) - 2((2q_i + q) \cdot q_j) (\not{q}_j + \not{q})] \quad (2.21)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)(1-y)y(1-2z+z^2)(p_i \cdot p_j)(p_i \cdot p_j)} [g^{\eta}_{\nu'}][(1-z)(1-y)(p_i \cdot p_j)(2 \not{q}_i + \not{q}) - (4q_i \cdot q_j + 2q \cdot q_j) (\not{q}_j + \not{q})] \quad (2.22)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)(1-y)y(1-2z+z^2)(p_i \cdot p_j)(p_i \cdot p_j)} [g^{\eta}_{\nu'}][(1-z)(1-y)(p_i \cdot p_j)(2 \not{q}_i + \not{q}) - (4(z(1-y)(p_i \cdot p_j)) + 2((1-z)(1-y)(p_i \cdot p_j))) (\not{q}_j + \not{q})] \quad (2.23)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)(1-y)y(1-2z+z^2)(p_i \cdot p_j)} [g^{\eta}_{\nu'}][(1-z)(1-y)(2 \not{q}_i + \not{q}) - (4(z(1-y)) + 2((1-z)(1-y))) (\not{q}_j + \not{q})] \quad (2.24)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)(1-y)y(1-2z+z^2)(p_i \cdot p_j)} [g^\eta_{\nu'}] [(1-z)(1-y)(2y(1-z) \not{p}_j + yz \not{p}_j) - 2(1+z)(1-y)((1+yz-z) \not{p}_j)] \quad (2.25)$$

$$M_1 M_2^\dagger = \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)(1-y)y(1-2z+z^2)(p_i \cdot p_j)} [g^\eta_{\nu'}] [-2(1+z)(1-y)(1+yz-z) \not{p}_j] \quad (2.26)$$

$$M_1 M_2^\dagger = -2(1+z)(1+yz-z) \frac{ig_s^2 f^{aol} [T^d]_{o'}^k}{4(1-z)y(1-2z+z^2)(p_i \cdot p_j)} [g^\eta_{\nu'}] [\not{p}_j] \quad (2.27)$$

2.4  $|M^2|$ 

$$|M^2| = \left| \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \\ \text{Diagram 3} \end{array} \right|^2 + 2 \operatorname{Re} E \left( \begin{array}{c} \text{Diagram 4} \end{array} \right)$$