# Momentum Distribution Properties of Ultracold Particles

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### Introduction

# Research Summary

- We study strongly-interacting, ultracold particles confined in a one-dimensional harmonic trap
- Strongly-interacting 1D systems often exhibit peculiar quantum behavior

# Two Central Questions:

- How does the momentum distribution decay for large momentum (e.g., exponentially, inverse power law)?
- What happens to momentum distribution when trapping potential is turned off?

#### Introduction

### Relevant Factors

- Particle type (e.g., fermion, boson, anyon)
- Particles with spin degree of freedom
- Systems with non-zero temperature
- Number of particles
- Interaction strength

### Background: Types of Particles

### **Bosons and Fermions**

- In nature, we only observe bosons and fermions
- Distinguishing characteristic is how their wavefunctions behave when two particles are exchanged

$$\psi_B(x_1, x_2) = \psi_B(x_2, x_1)$$
$$\psi_F(x_1, x_2) = -\psi_F(x_2, x_1)$$

- Exchange Statistics: how wave function behaves under particle exchange
- Identical particles  $\rightarrow$  exchanging particles can only change  $\psi$  by complex phase

# Background: Types of Particles

# Anyons

- As a theoretical idea, we can define another type of particle the anyon
- Described by its anyonic statistic  $\kappa \in [0, 2)$
- Defining characteristic:

$$\psi_{\kappa}(x_1, x_2) = e^{i\pi\kappa\epsilon(x_1 - x_2)}\psi_{\kappa}(x_2, x_1)$$

where  $\epsilon(x)$  is the sign function

- $\kappa = 0 \to e^{i\pi\kappa} = 1 \to \text{Bosons}$
- $\kappa = 1 \to e^{i\pi\kappa} = -1 \to \text{Fermions}$

# **Background: Strong Interaction**

### **Hard-Core Particles**

- We consider systems 'hard-core' particles: infinite, repulsive contact interactions
- Means that no two of the particles can occupy same position
- If particles are spinless, we can write the Hamiltonian

$$H = \sum_{j=1}^{N} \left[ \frac{-1}{2} \frac{\partial^2}{\partial x_j^2} + \frac{1}{2} x_j^2 \right] + g \sum_{j < \ell} \delta(x_j - x_\ell)$$

where  $\hbar = m = \omega = 1$  and in the hard-core limit  $g \to \infty$ 

• Our research will consider hard-core bosons and hard-core anyons

### Methods: Mappings

# Bose-Fermi Mapping

- Bosons and fermions are very different, but hard-core bosons (HCB) are similar to fermions (e.g., no two can occupy same position)
- M. Girardeau (1960) showed there was a mapping between these two systems
- ullet Procedure for turning N-particle fermion solution into boson solution

$$\psi_B(x_1, \dots, x_N) = A(x_1, \dots, x_n)\psi_F(x_1, \dots, x_N)$$
$$A(x_1, \dots, x_N) = \prod_{1 \le \ell \le j \le N} \epsilon(x_j - x_\ell)$$

# Anyon-Fermion Mapping

- Girardeau (2006) showed mapping existed for hard-core anyons (HCA)
- Procedure for turning fermionic solution into bosonic solution

$$\psi_{\kappa}(x_1,\ldots,x_N) = A_{\kappa}(x_1,\ldots,x_N)\psi_F(x_1,\ldots,x_N)$$

$$A_{\kappa}(x_1, \dots, x_N) = \exp\left(i\pi(1-\kappa)\sum_{1\leq \ell < j\leq N} \theta(x_j - x_\ell)\right)$$

where  $\theta$  is the heaviside theta function

# Methods: One Body Density Matrix

• After obtaining the wave function for our system using the mappings, we study its properties using the One-Body Density Matrix (OBDM)

$$\rho(x,x') = \int_{-\infty}^{\infty} dx_2 \cdots dx_N \psi^*(x,x_2,\ldots,x_N) \psi(x',x_2,\ldots,x_N)$$

• Then, we use this to find an expression for the momentum space density profiles

$$n(p) = \frac{N}{2\pi} \int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dx' e^{ip(x-x')} \rho(x, x')$$

• Difficult to solve analytically, even for two hard-core bosons

# Methods: Momentum Space Density Profiles

# Background: Two Hard-Core Bosons

### Tan Contact

- Momentum distribution for spinless fermions in a harmonic trap will decay exponentially
- $\bullet\,$  Minguzzi et al. (2002) showed for 2 spinless hard-core bosons

$$\lim_{p \to \infty} n_B(p) = \frac{C_B}{p^4} = \left(\frac{2}{\pi}\right)^{3/2} \frac{1}{p^4}$$

where  $C_B$  is the Tan Contact (see Tan (2008)) or momentum tail coefficient

### Background: Two Hard-Core Bosons

# Dynamical Fermionization

- Hard-core bosons and fermions have different momentum distributions
- When trapping potential is turned off for hard-core bosons, Minguzzi et al. (2005) showed that

$$\lim_{t \to \infty} n_B(p;t) = n_F(p)$$

using a scaling transformation

• This process is known as dynamical fermionization

# Progress: Two Hard-Core Anyons

### Tan Contact

• Calculate Tan Contact for two hard-core anyons

$$\lim_{p \to \infty} n_{\kappa}(p) = \frac{C_{\kappa}}{p^4} = \cos^2\left(\frac{\pi\kappa}{2}\right) \frac{C_B}{p^4}$$

• And numerically showed dynamical fermionization occurs

### Progress: Finite Temperature Tan Contact

- Hao et al. (2017) calculate the Tan Contact for N Hard-core Anyons at finite temperature in grand canonical ensemble
- Their calculation more or less follows Vignolo et al. (2013)

$$\rho_B(x,y) = \sum_{N,\alpha} P_{N,\alpha} N \int_{-\infty}^{\infty} dx_2 \cdots dx_N \psi_{N,\alpha}(x, x_2, \dots, x_N)$$
$$\times \psi_{N,\alpha}^*(y, x_2, \dots, x_N)$$

- N is number of particles and  $\alpha = \{\nu_1, \nu_2, \dots, \nu_N\}$  gives the quantum number for each particle
- $P_{N,\alpha}$  gives the probability that the system has N particles with state  $\alpha$

$$P_{N,\alpha} = \frac{\exp(-\beta(E_{N,\alpha} - \mu N))}{Z}$$

### Progress: Finite Temperature Tan Contact

- Density matrix  $\rho_B$  can expressed as sum of special functions
- Only first term contributes at high momentum

$$\rho_B^{(j=1)}(x, x') \sim \frac{|x - x'|^3}{3} F(R)$$

where  $R = \frac{x+x'}{2}$  can be ignored in Fourier transform

• Use asymptotics of Fourier transforms to evaluate the momentum distribution

$$\lim_{p \to \infty} n_B(p) = \frac{C}{p^4}$$

$$C = \frac{2}{\pi} \int_{-\infty}^{\infty} dR F(R)$$

### Plans: Two Spinor Case

- The results so far have been with spinless particles
- For two hard-core bosons with spin degree of freedom (called spinors) we ask:
  - What is the total Tan Contact?
  - Will dynamical fermionization occur?
- We will use tools from past research in Pu group:
  - Yang et al. (2015) simplifies OBDM for strongly-interacting spinor particles by separating into spatial part and spin part
  - Yang et al. (2017) exploits connection between spinless HCA and spinors to calculate momentum distribution
  - Other unpublished work, Li Yang has shown that the Tan Contact does exist for spinor system

### Conclusion

### Past Research:

- Calculated Tan Contact for two, spinless hard-core anyons
- Showed dynamical fermionization would occur
- $\bullet$  Studied N-particle, strongly-intereacting systems with finite temperature

### **Future Work:**

- Use current tools to calculate total Tan Contact for two spinors
- Determine whether dynamical fermionization will occur

### References