

Wednesday Mar 8, 2017

Last time:

- Capacitors
- Capacitance as a geometric quantity
- General Capacitors, relating Q to ΔV
- Setting up a process to find Capacitance

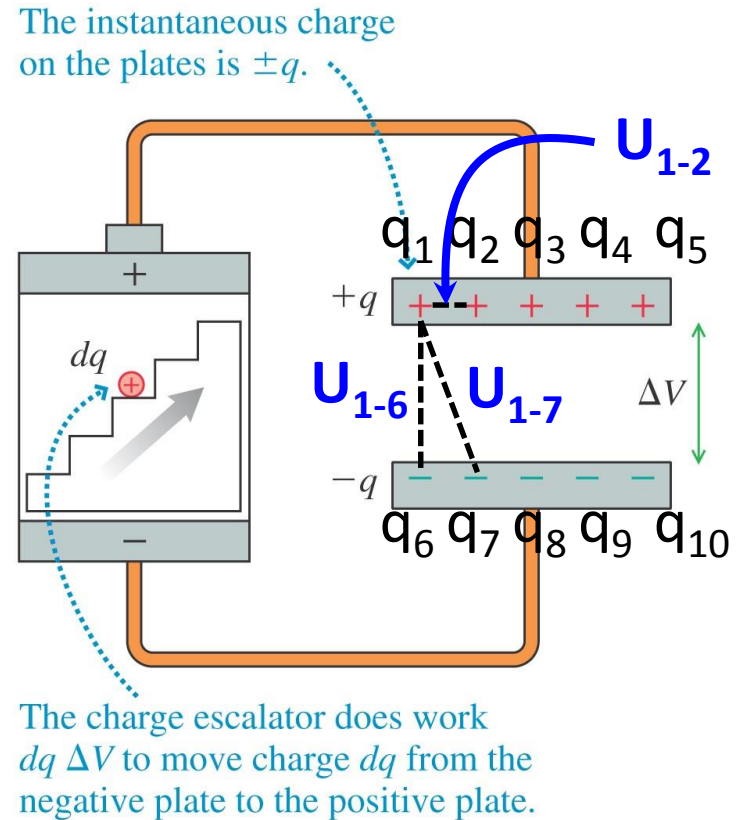
Today:

- Energy stored in parallel plate, spherical, and cylindrical capacitors
- Potential energy stored in the electric field itself.
- Capacitors in electric circuits: how charges move
- Kirchhoff's loop rule with capacitors
- Capacitors in series and parallel
- More complicated capacitor circuit

Energy Storage in Capacitors

We want to calculate this potential energy stored in the capacitor.

It is **way** too hard to add up all the potential energies of every pair of charges in the capacitor:



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$$U = U_{1-2} + U_{1-3} + \dots + U_{1-10} + U_{2-1} + \mathbf{U_{ij} \text{ of every other pair}}$$

Easier way!

Move a tiny charge, dq , from the negative plate to the positive plate.

It moves through a potential difference ΔV .

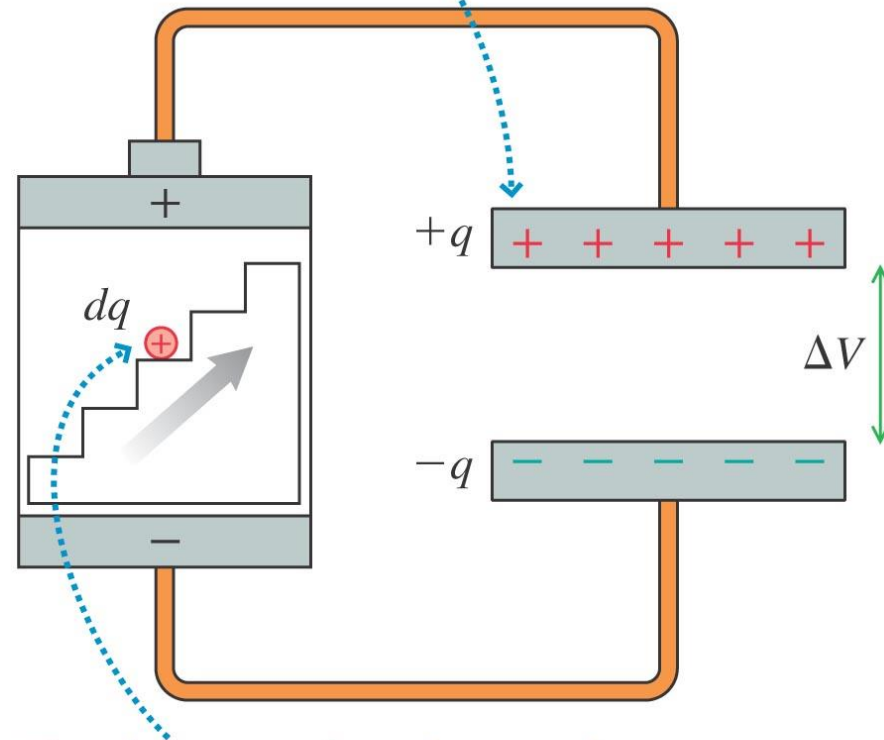
So its potential energy increases by an amount

$$dU = dq DV_C$$

But we also know $DV_C = \frac{q}{C}$

$$dU = dq \frac{q}{C} = \frac{q dq}{C}$$

The instantaneous charge on the plates is $\pm q$.



The charge escalator does work $dq \Delta V$ to move charge dq from the negative plate to the positive plate.

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$$U = \frac{1}{C} \int_0^Q q dq = \frac{1}{2} \frac{Q^2}{C}$$

Potential Energy in a Capacitor

Energy storage in terms of the charge on the plates:

$$U = \frac{1}{2} \frac{Q^2}{C}$$

Use the general relation for a capacitor to swap charge for voltage

$$Q = CDV_c$$

Energy storage in terms of the voltage across the plates:

$$U = \frac{1}{2} \frac{(CDV_c)^2}{C}$$
$$= \frac{1}{2} C (DV_c)^2$$

Where is the Energy Stored?

$$U = \frac{1}{2} C (DV_C)^2$$

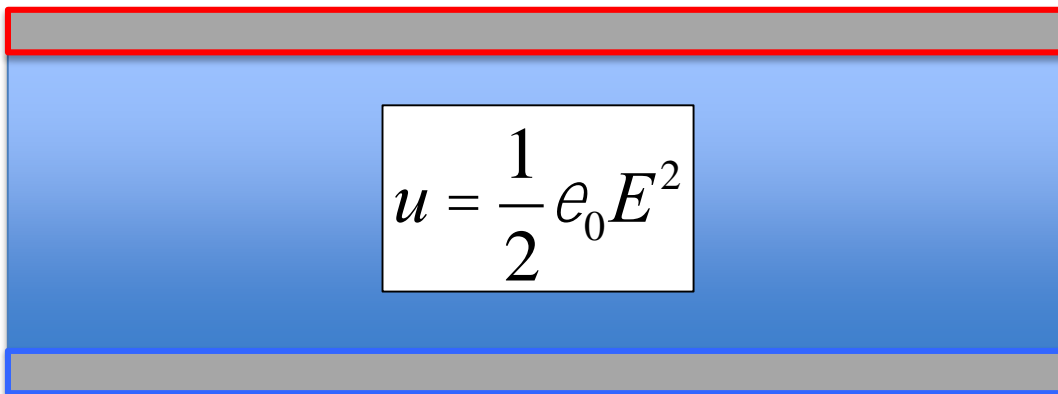
$$= \frac{1}{2} C E^2 d^2$$

$$= \frac{1}{2} \frac{e_0 A}{d} E^2 d^2 = \frac{1}{2} e_0 E^2 (Ad)$$

$$DV = Ed$$

$$C = \frac{e_0 A}{d}$$

$$u = \frac{U}{Ad}$$

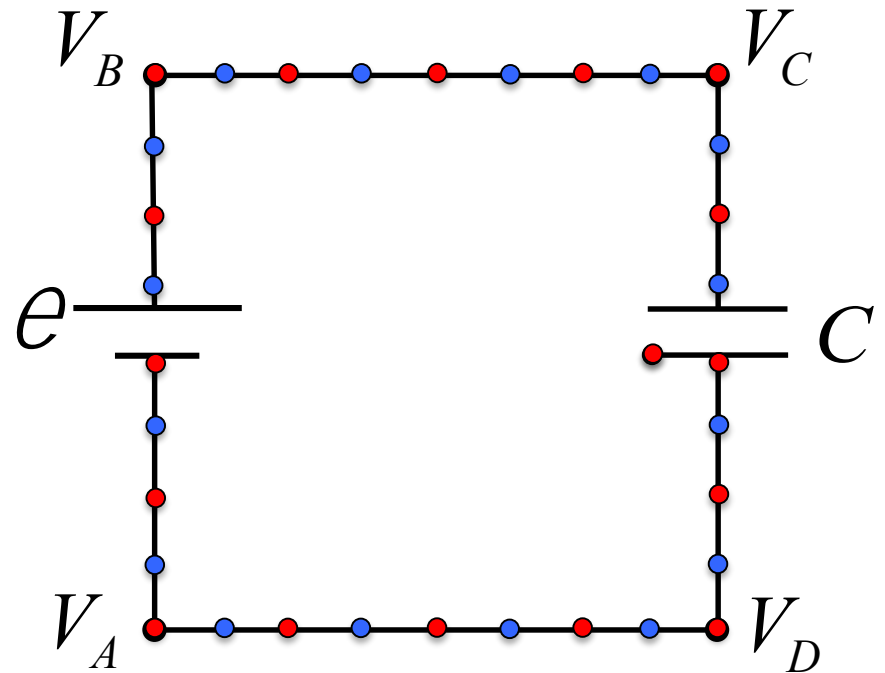


The capacitor's energy is stored in the electric field between the plates!

A Basic Circuit with Capacitor

The simplest capacitor circuit has an ideal battery, ideal wires, and a single capacitor.

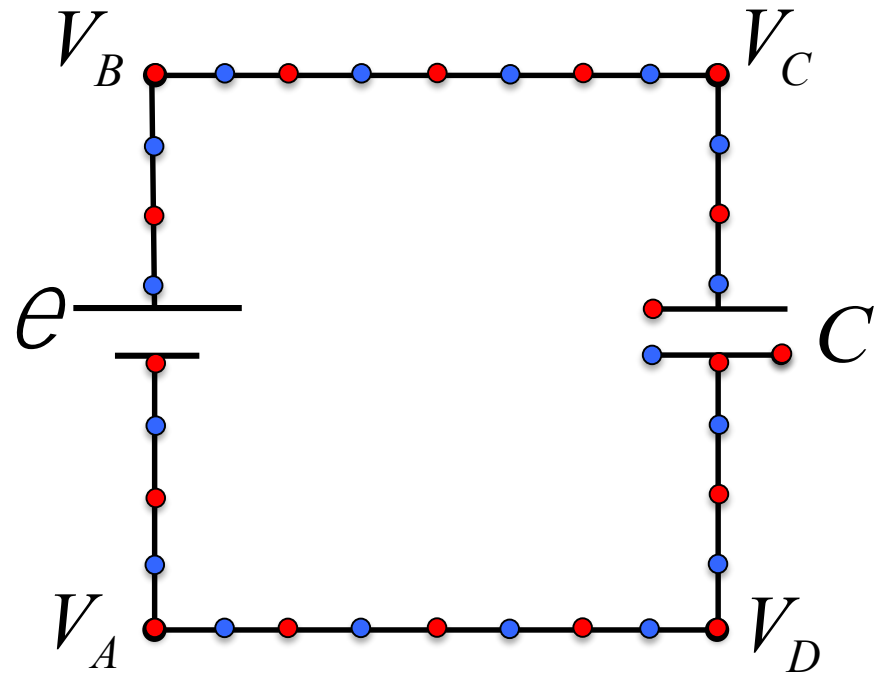
The battery causes charges to flow from the bottom plate to the top plate. This creates a potential ΔV_C between the two plates. Remember charges never “jump the gap” between the two plates of a capacitor.



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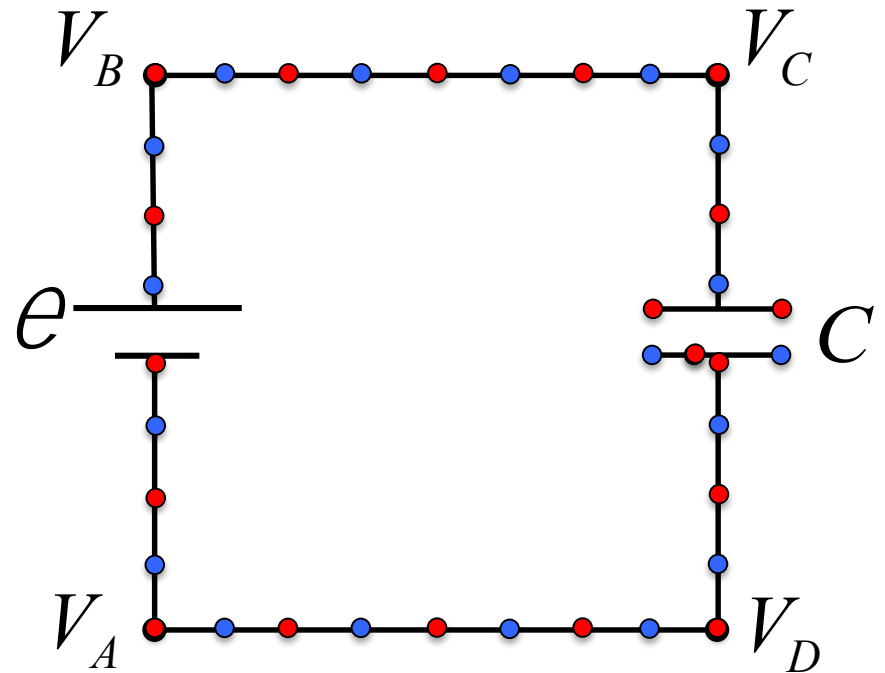
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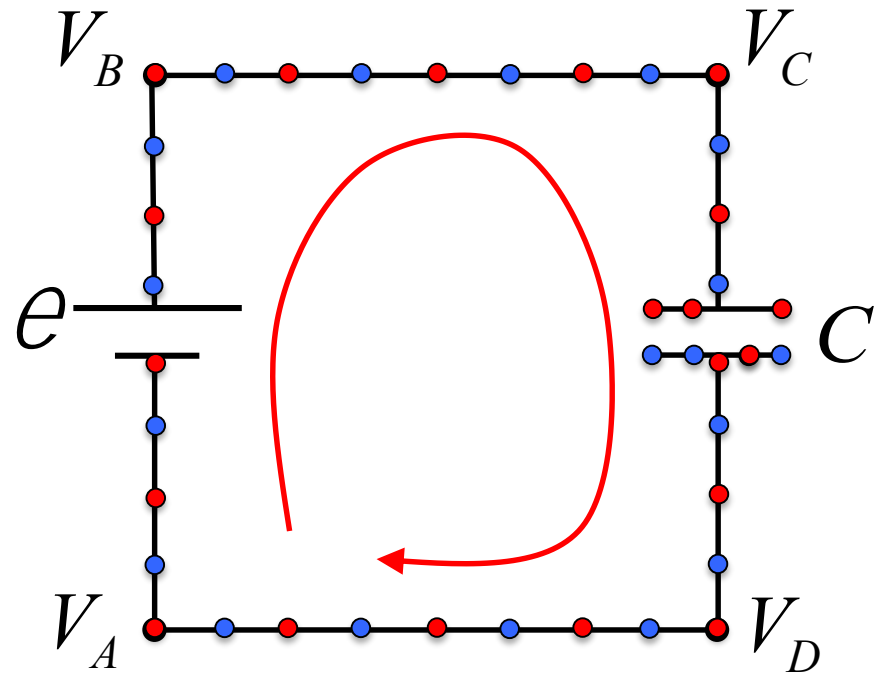
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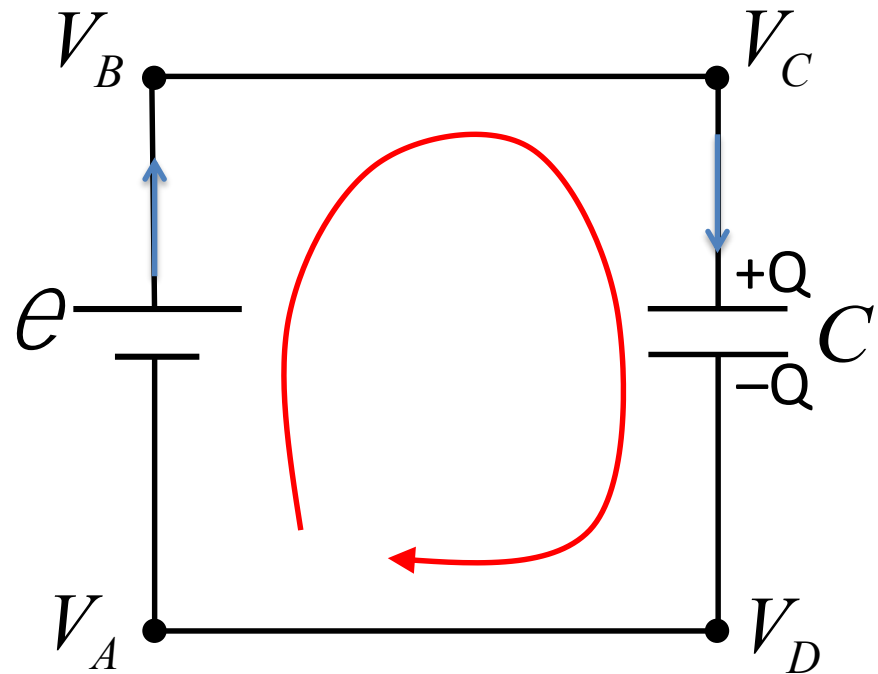
The battery causes charges to flow from the bottom plate to the top plate. This creates a potential ΔV_C between the two plates. Remember charges never “jump the gap” between the two plates of a capacitor.



$$\Delta V_{AB} + \Delta V_{BC} + \Delta V_{CD} + \Delta V_{DA} = 0$$

A Basic Circuit

The voltage across a capacitor is **negative** if you are going around the loop in the direction **from the + plate to the – plate**. Current flows **from the negative terminal to the positive terminal**



ideal wires

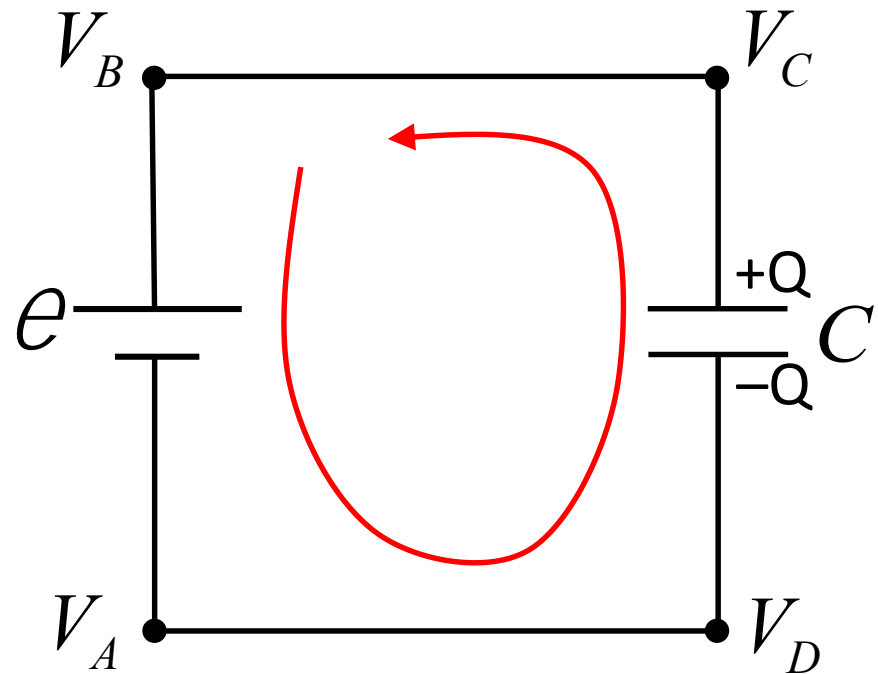
$$\Delta V_{AB} + \cancel{\Delta V_{BC}} + \Delta V_{CD} + \cancel{\Delta V_{DA}} = 0$$

$$e - \frac{Q}{C} = 0$$

A Basic Circuit

The voltage across a capacitor is **positive** if you are going around the loop in the direction **from – plate to + plate**.

Voltage across a battery is **negative** going **from positive to negative**



$$\Delta V_{BA} + \cancel{\Delta V_{AD}} + \Delta V_{DC} + \cancel{\Delta V_{CB}} = 0$$

ideal wires

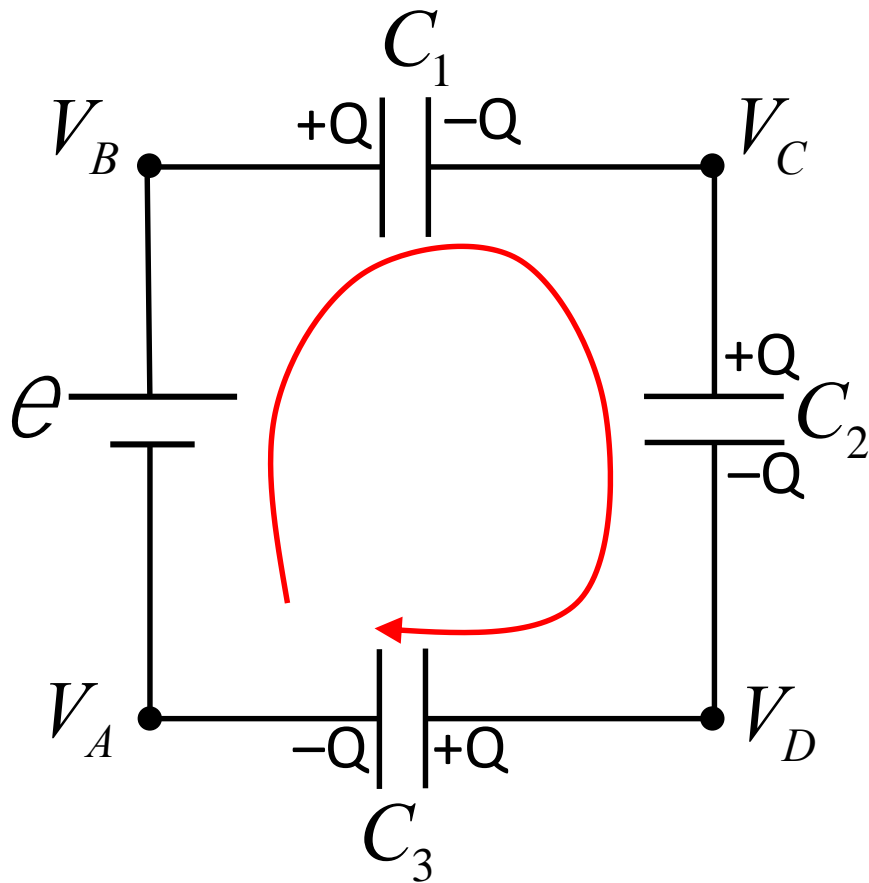
$$-e + \frac{Q}{C} = 0$$

Same as before

TopHat Question

Capacitors in Series

A slightly more complicated circuit has multiple capacitors in series



Kirchhoff's Loop Rule:

$$\Delta V_{AB} + \Delta V_{BC} + \Delta V_{CD} + \Delta V_{DA} = 0$$

Charge on each plate is the same

$$e - \frac{Q}{C_1} - \frac{Q}{C_2} - \frac{Q}{C_3} = 0$$

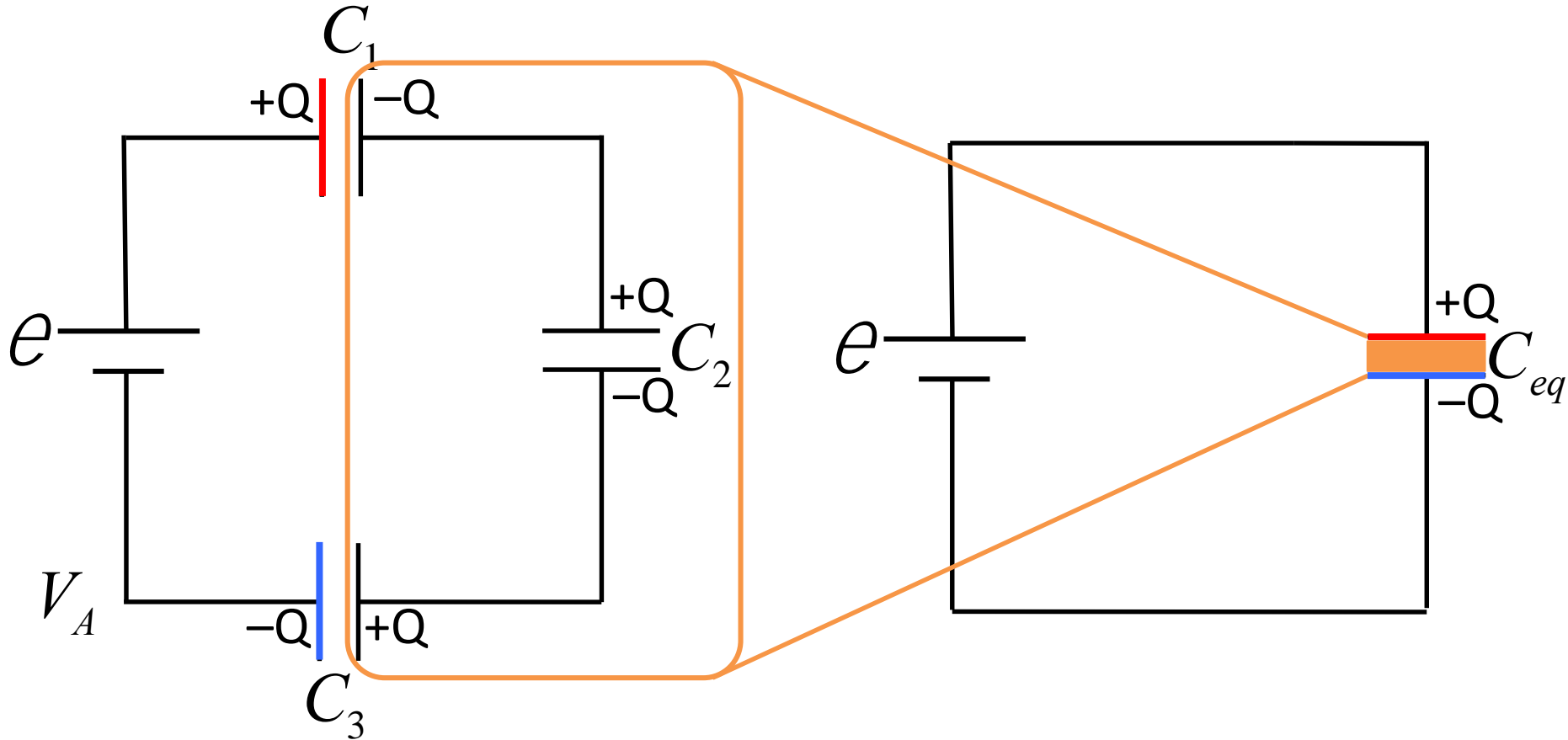
Rewrite this as

$$e - Q \left(\frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} \right) = 0$$

Define an equivalent capacitance

$$e - \frac{Q}{C_{eq}} = 0$$

Capacitors in Series

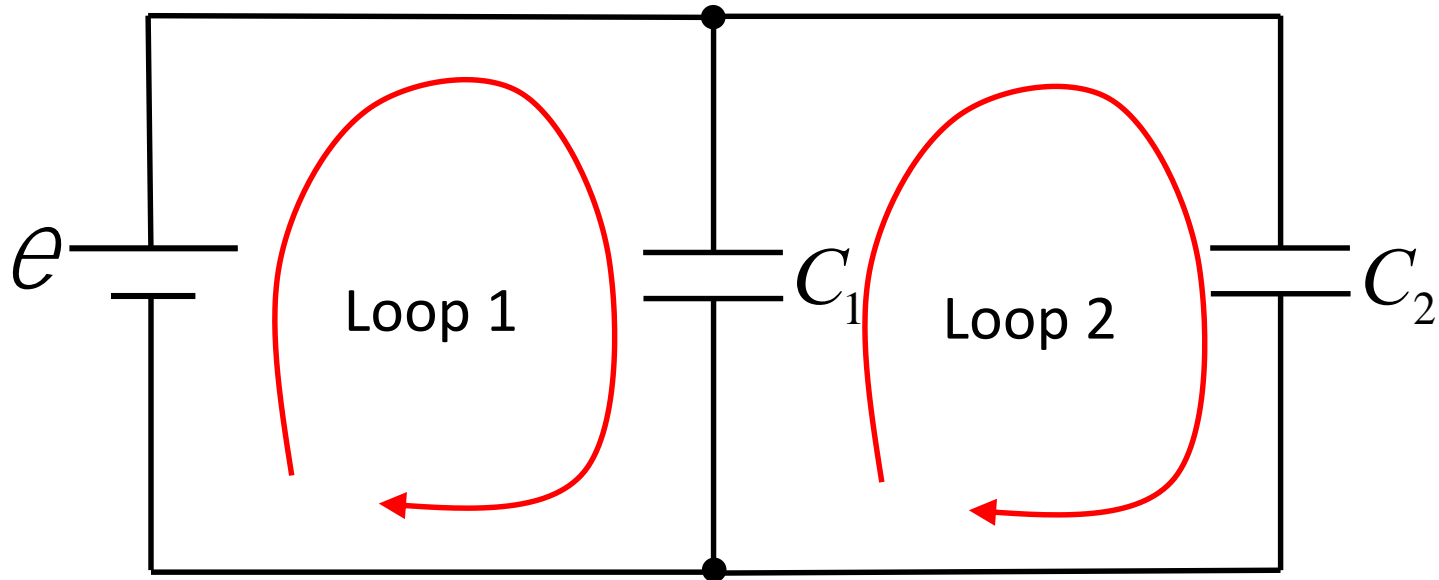


Capacitors in series act like a single equivalent capacitor:

$$C_{eq} = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3}}$$

Capacitors in Parallel

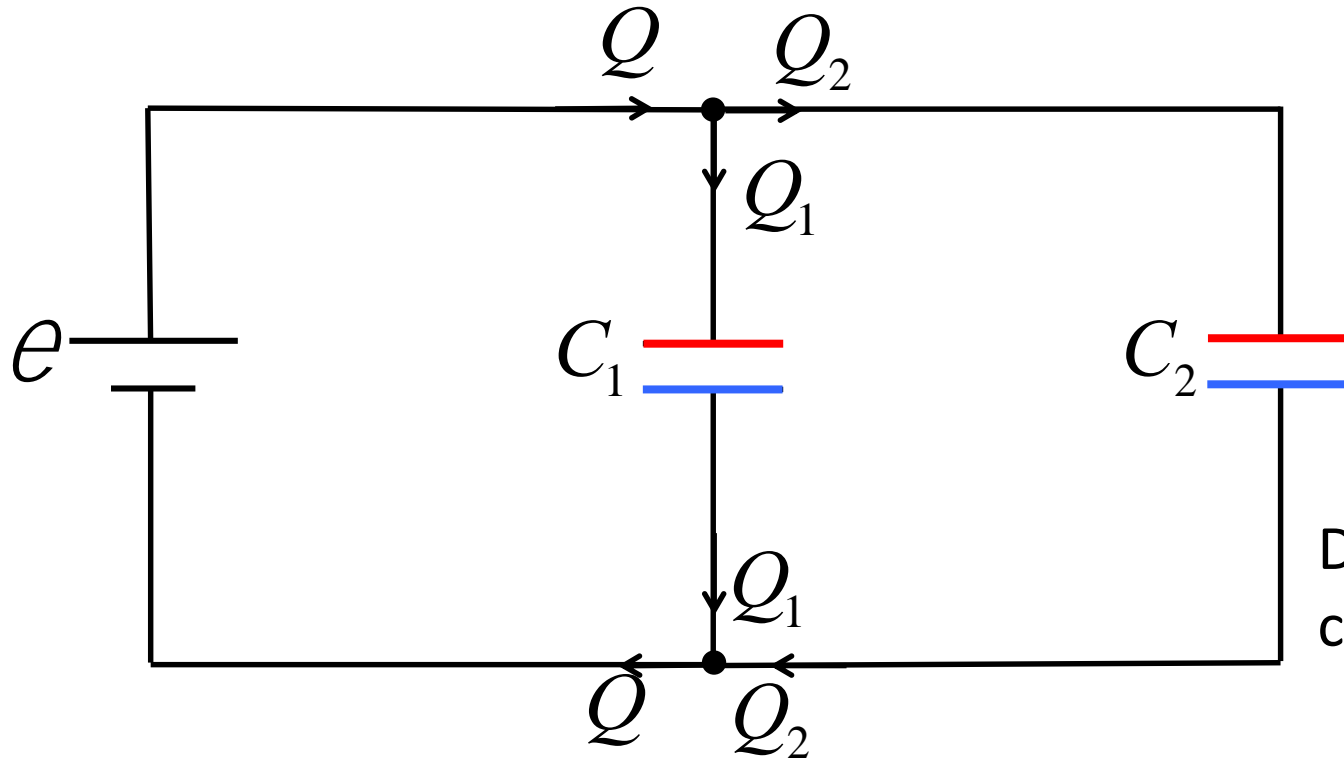
A slightly more complicated circuit has multiple branches with capacitors in parallel



Capacitors in parallel have the same voltage across their plates

$$\text{Loop 1: } e - \Delta V_{C_1} = 0 \quad \text{Loop 2: } \Delta V_{C_1} - \Delta V_{C_2} = 0$$

Capacitors in Parallel



$$Q = Q_1 + Q_2$$

$$Q_1 = \Delta V_{C_1} C_1$$

$$Q_2 = \Delta V_{C_2} C_2$$

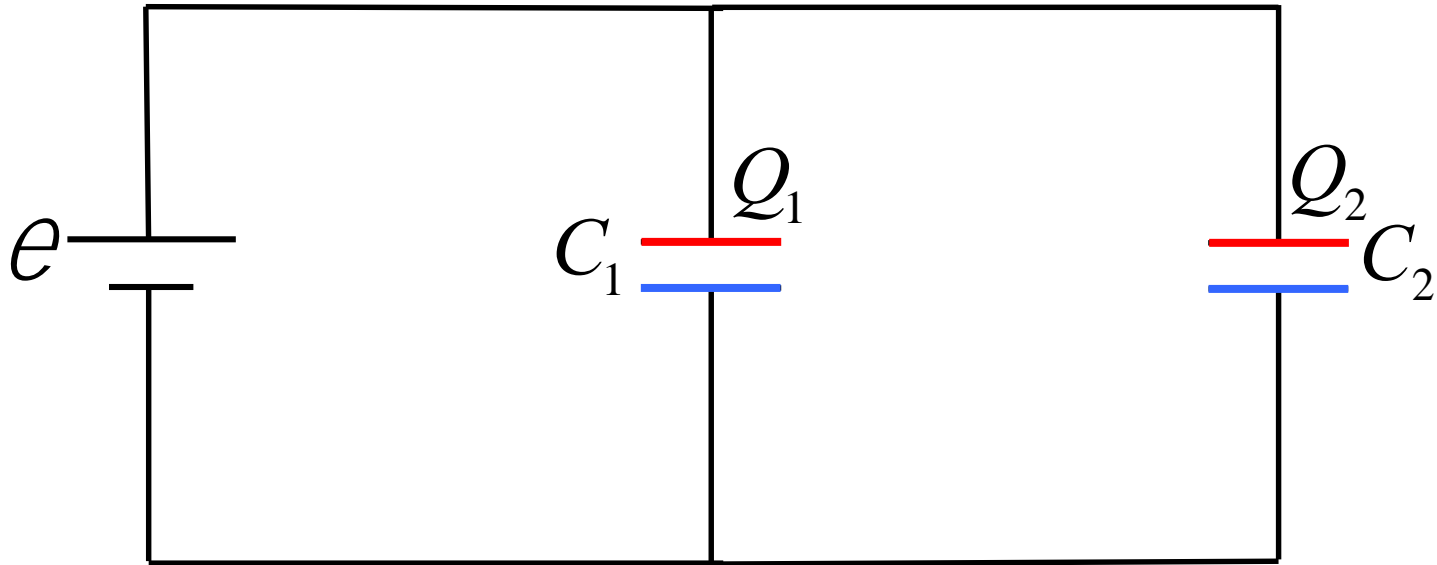
Define an equivalent capacitance:

$$Q = eC_{eq}$$

From conservation of charge: $eC_{eq} = eC_1 + eC_2$

For capacitors in parallel: $C_{eq} = C_1 + C_2$

Capacitors in Parallel



$$Q = Q_1 + Q_2$$

$$C_{eq} = C_1 + C_2$$

Summary of Capacitors

Relation between charge and voltage across plates

$$\Delta V_C = \frac{Q}{C}$$

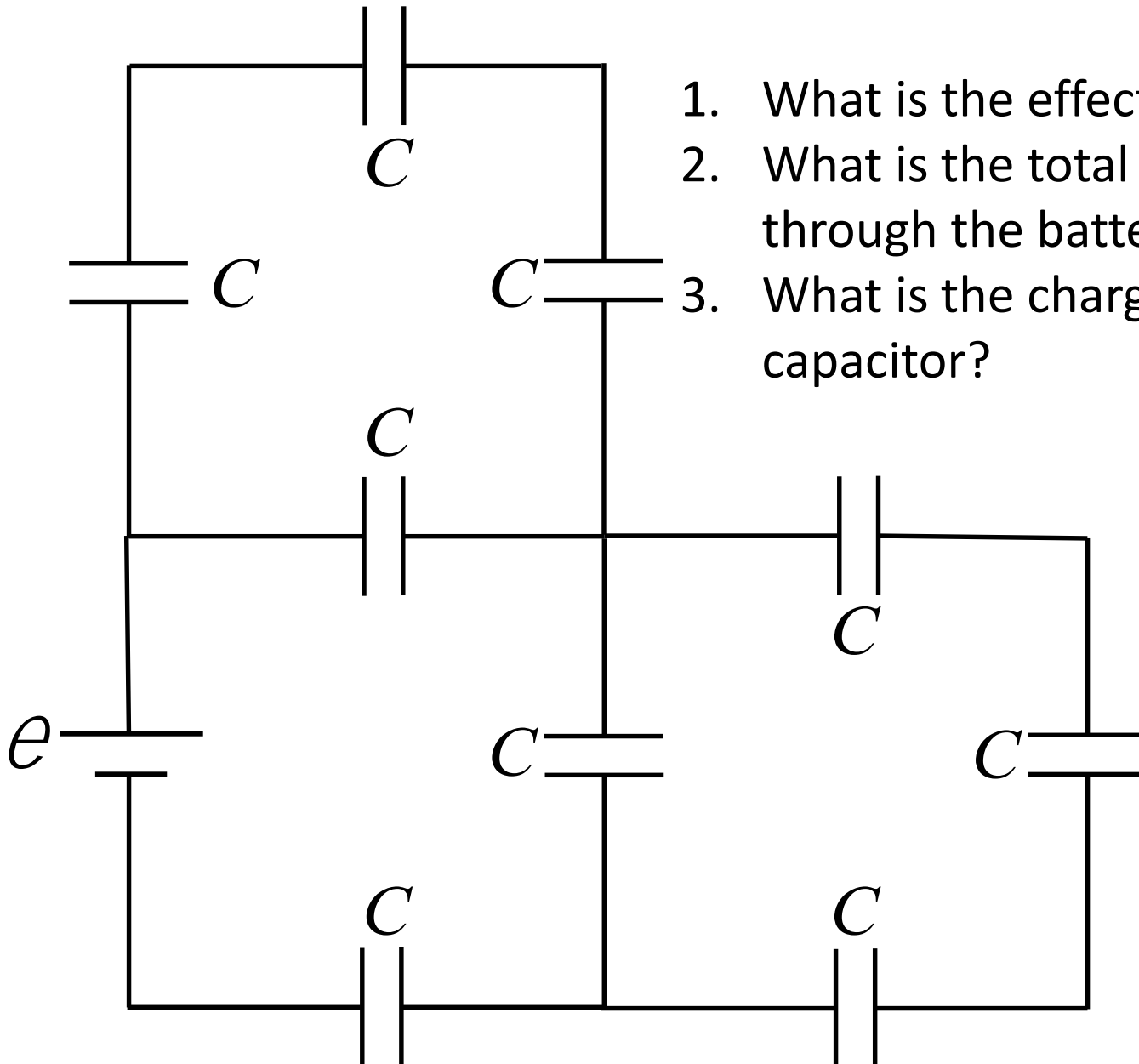
Capacitors in Series: store the same amount of charge

$$C_{eq} = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2} + \dots + \frac{1}{C_N}}$$

Capacitors in Parallel: have the same voltage across them

$$C_{eq} = C_1 + C_2 + \dots + C_N$$

On document camera



1. What is the effective capacitance?
2. What is the total charge moved through the battery?
3. What is the charge on each capacitor?

Wednesday March 8, 2017 class 2

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Today:

- Subtlety with capacitors: series or parallel?
- Linear dielectric materials: an atomic perspective
- Effect of dielectrics on capacitance
- Applications of dielectrics and capacitors

Summary of Capacitors

Relation between charge and voltage across plates

$$\Delta V_C = \frac{Q}{C}$$

Capacitors in Series: store the same amount of charge

$$C_{eq} = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2} + \dots + \frac{1}{C_N}}$$

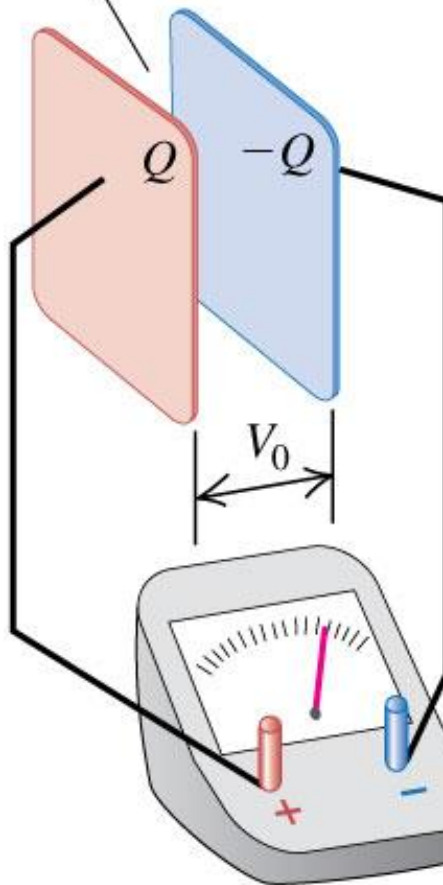
Capacitors in Parallel: have the same voltage across them

$$C_{eq} = C_1 + C_2 + \dots + C_N$$

$$Q = C_0 V_0$$

(a) $V_0 = E_0 d$

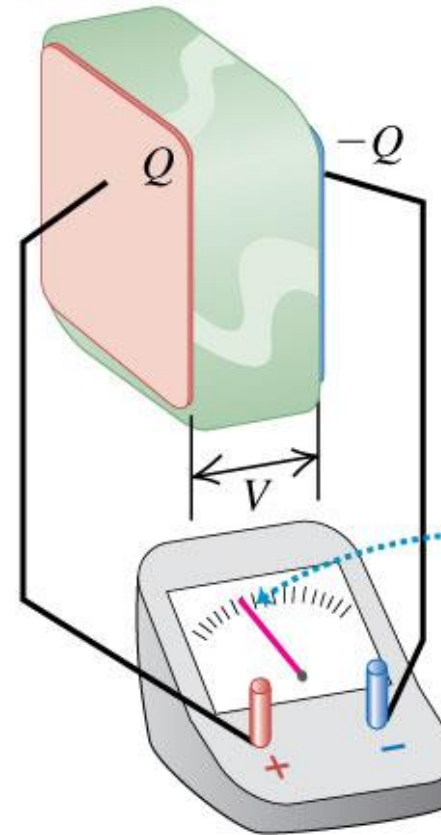
Vacuum



Electrometer
(measures potential
difference across
plates)

(b)

Dielectric



$$Q = CV$$

$$V < V_0$$

$$C > C_0$$

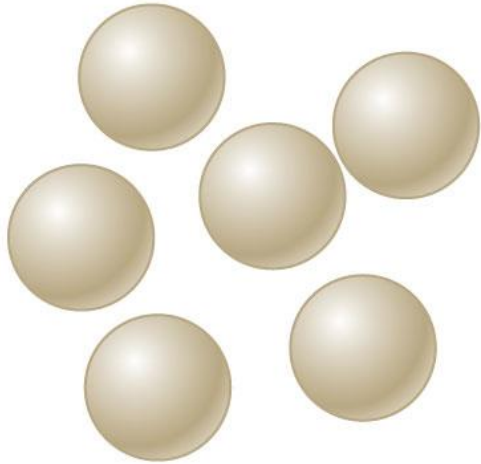
Adding the dielectric
reduces the potential
difference across the
capacitor.

$$V = Ed$$

$$E < E_0$$

non-polar molecules

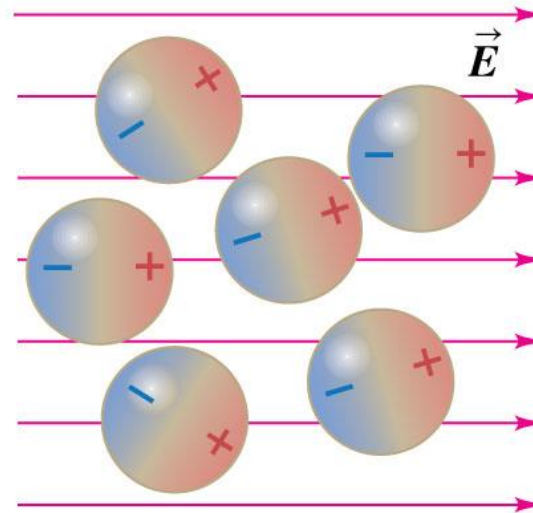
(a)



In the absence of an electric field, nonpolar molecules are not electric dipoles.

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(b)

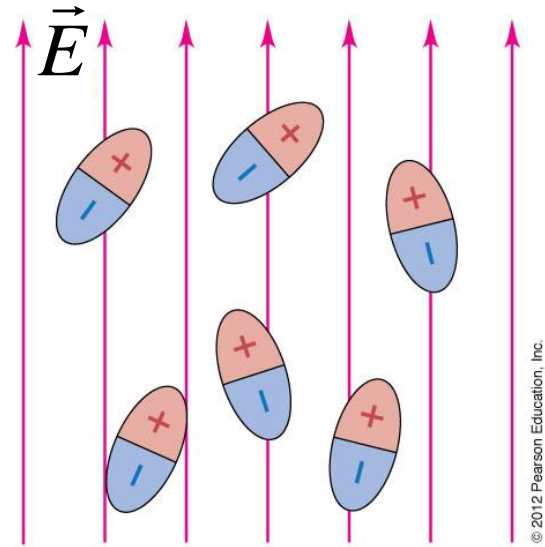
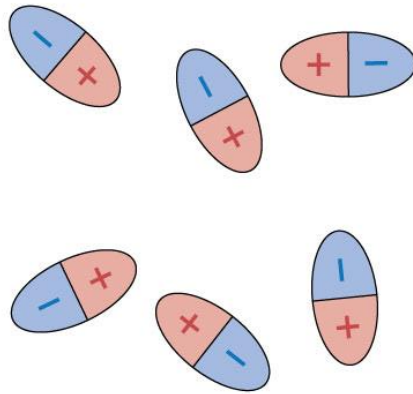


An electric field causes the molecules' positive and negative charges to separate slightly, making the molecule effectively polar.

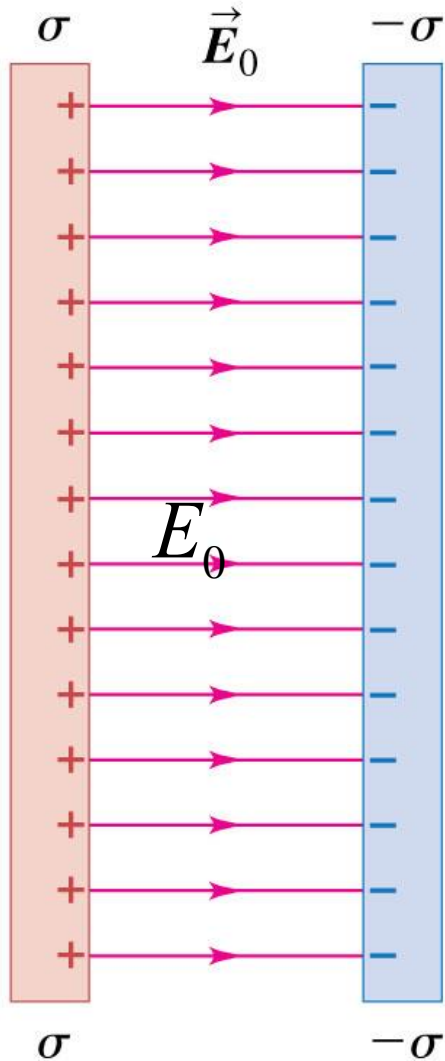
Recall the balloon on the wall example from week 1



polar molecules

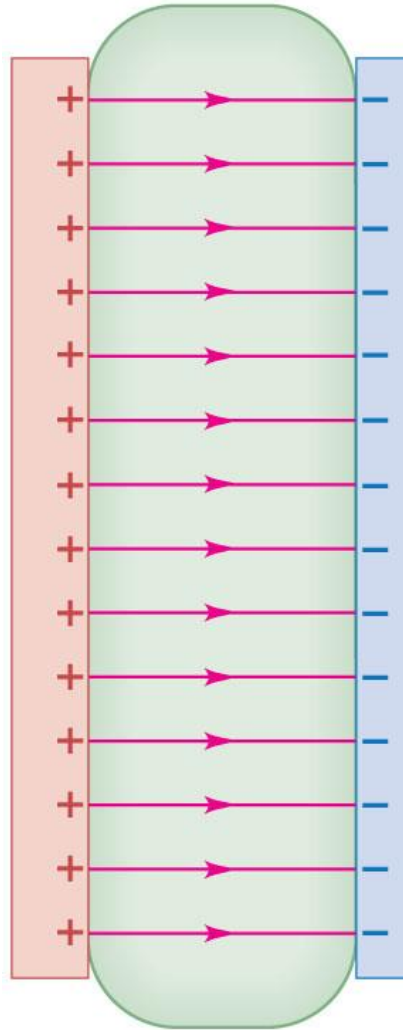


(a) No dielectric



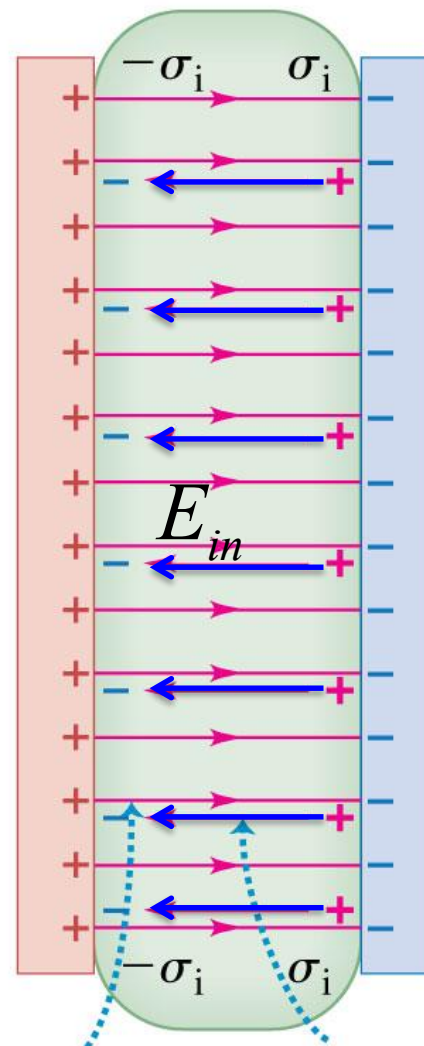
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(b) Dielectric just inserted



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(c) Induced charges create electric field



Original
electric field

Weaker field in dielectric
due to induced (bound) charges

$$E_{in} = E_0 - E_{diel}$$

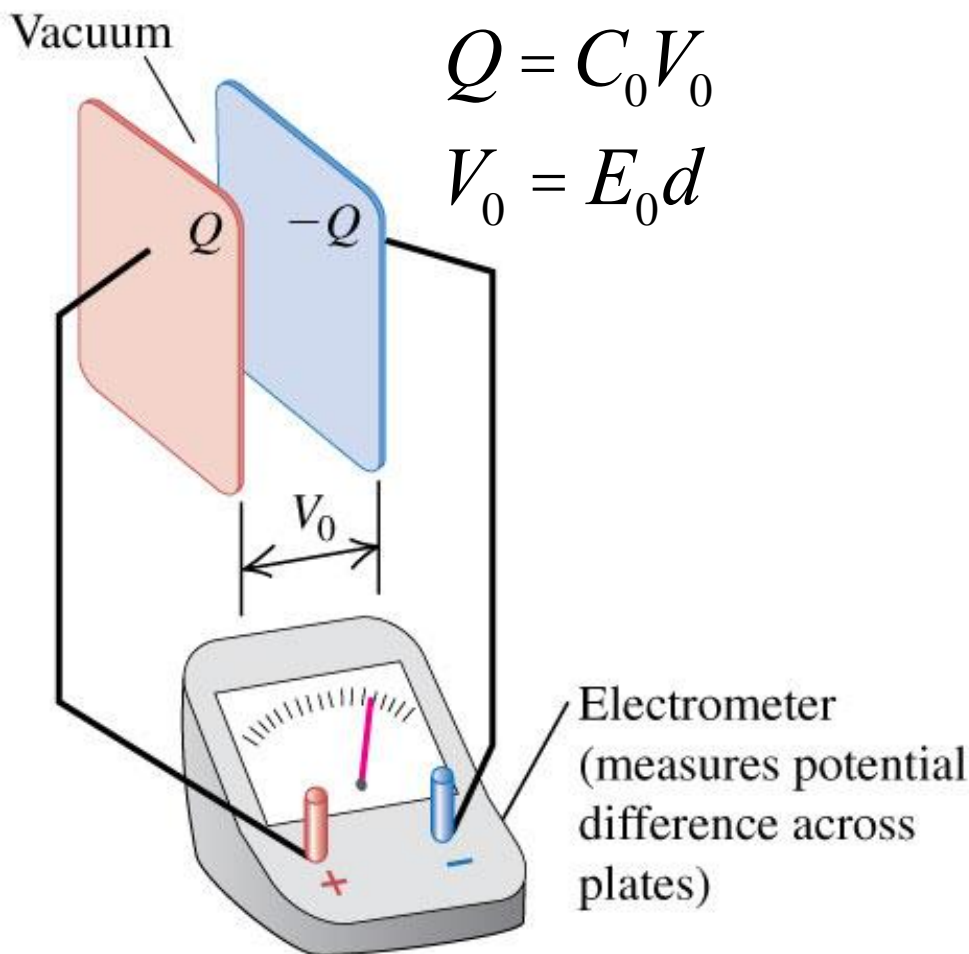
$$E_{in} < E_0$$

$$E_{in} = \frac{E_0}{K}$$

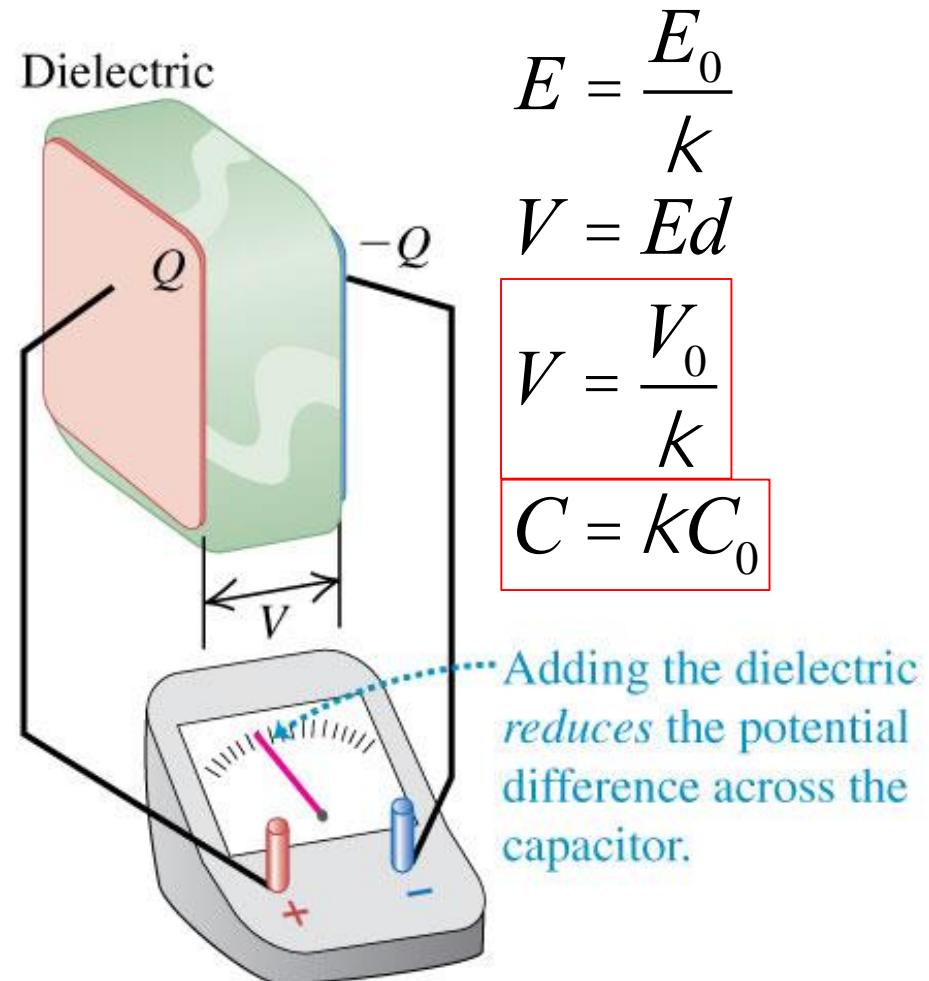
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Linear Dielectrics

(a)



b)



Parallel Plate Capacitors With and Without Dielectrics

$$E_0 = \frac{S}{e_0}$$

$$E = \frac{E_0}{k} = \frac{S}{ke_0}$$

$$V_0 = E_0 d = \frac{Sd}{e_0}$$

$$V = \frac{V_0}{k} = \frac{Sd}{ke_0}$$

$$C_0 = \frac{Q}{V_0} = \frac{SA}{Sd} e_0 = \frac{Ae_0}{d}$$

$$C = kC_0 = \frac{Ake_0}{d}$$

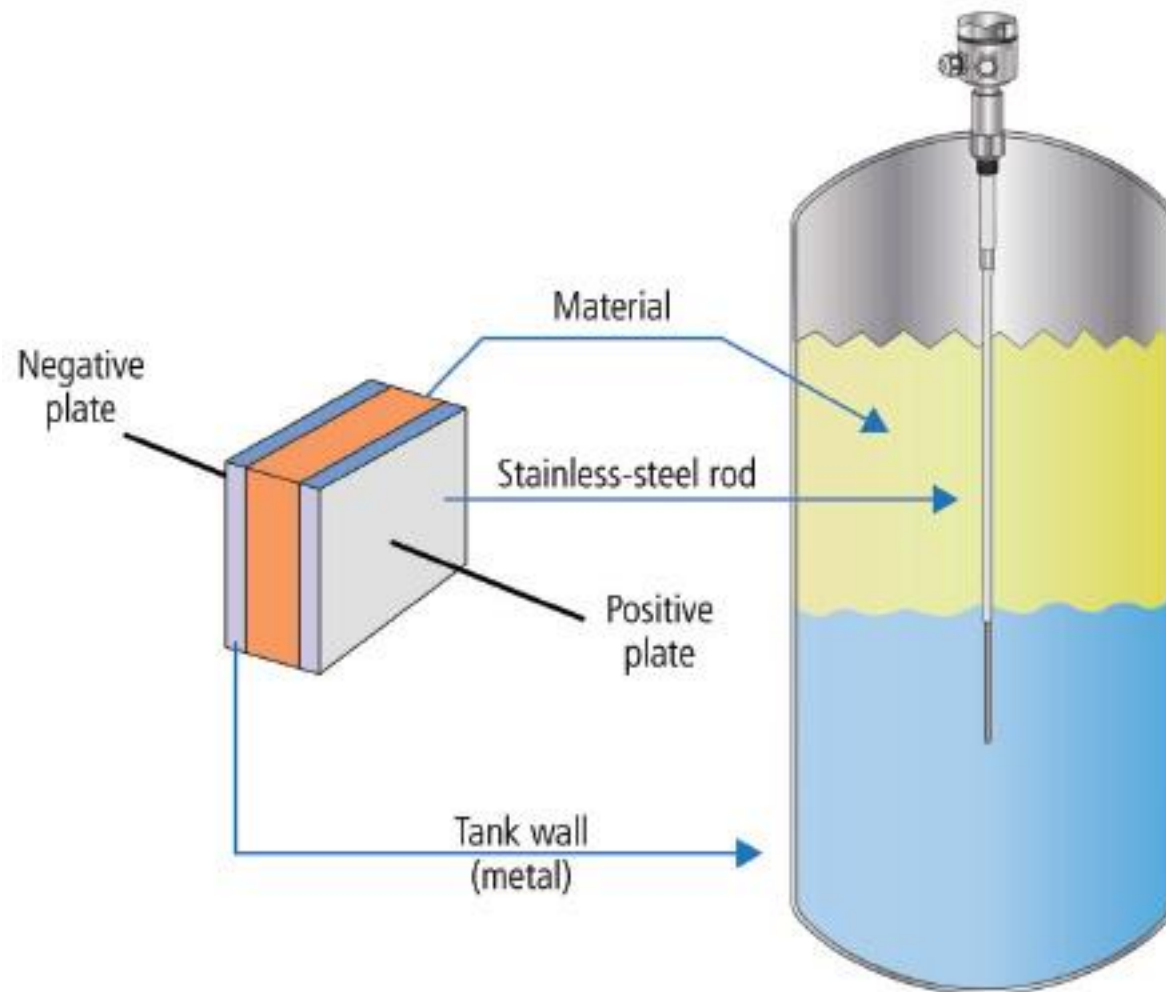
Conclusion: for a linear dielectric, all the regular electrostatic equations hold if $e_0 \rightarrow e \equiv ke_0$

Now you see why e_0 is called “permittivity of free space” (i.e. in vacuum)

Gauss' law for dielectrics

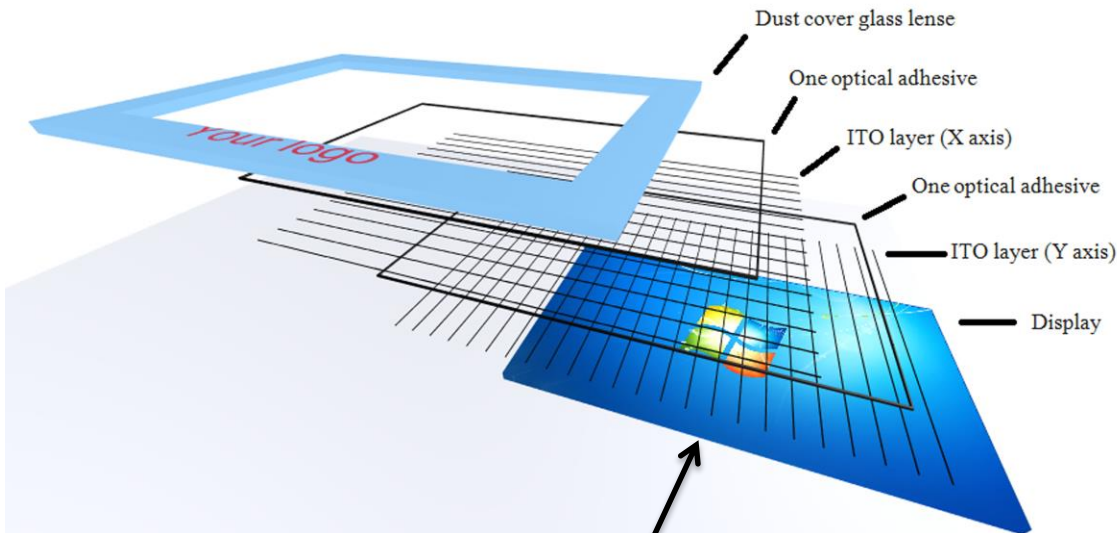
- $\oiint \vec{E} \cdot d\vec{A} = \frac{q_{f, enc}}{\kappa \epsilon_0}$

Application: Capacitive Fuel Gauge

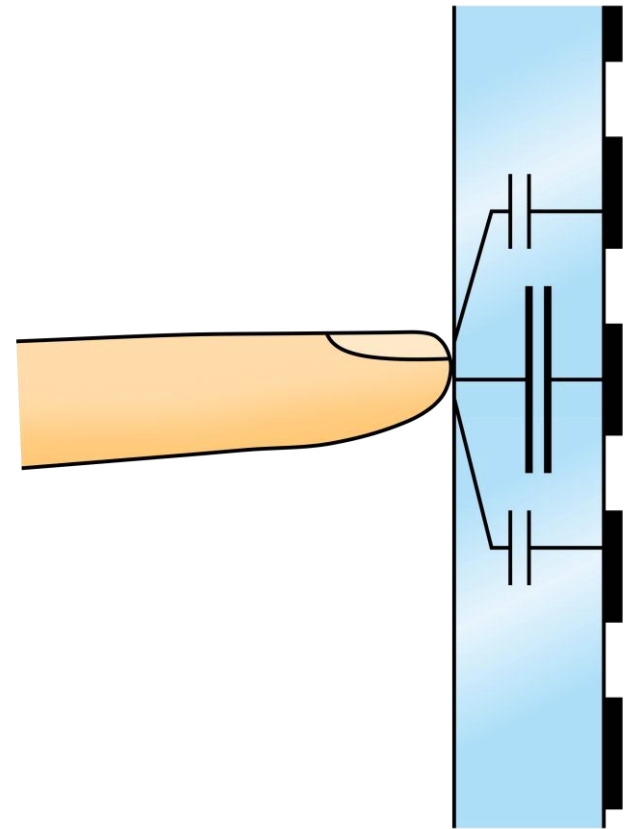


Application: Capacitive Touch Screen

Row and column stack up layers



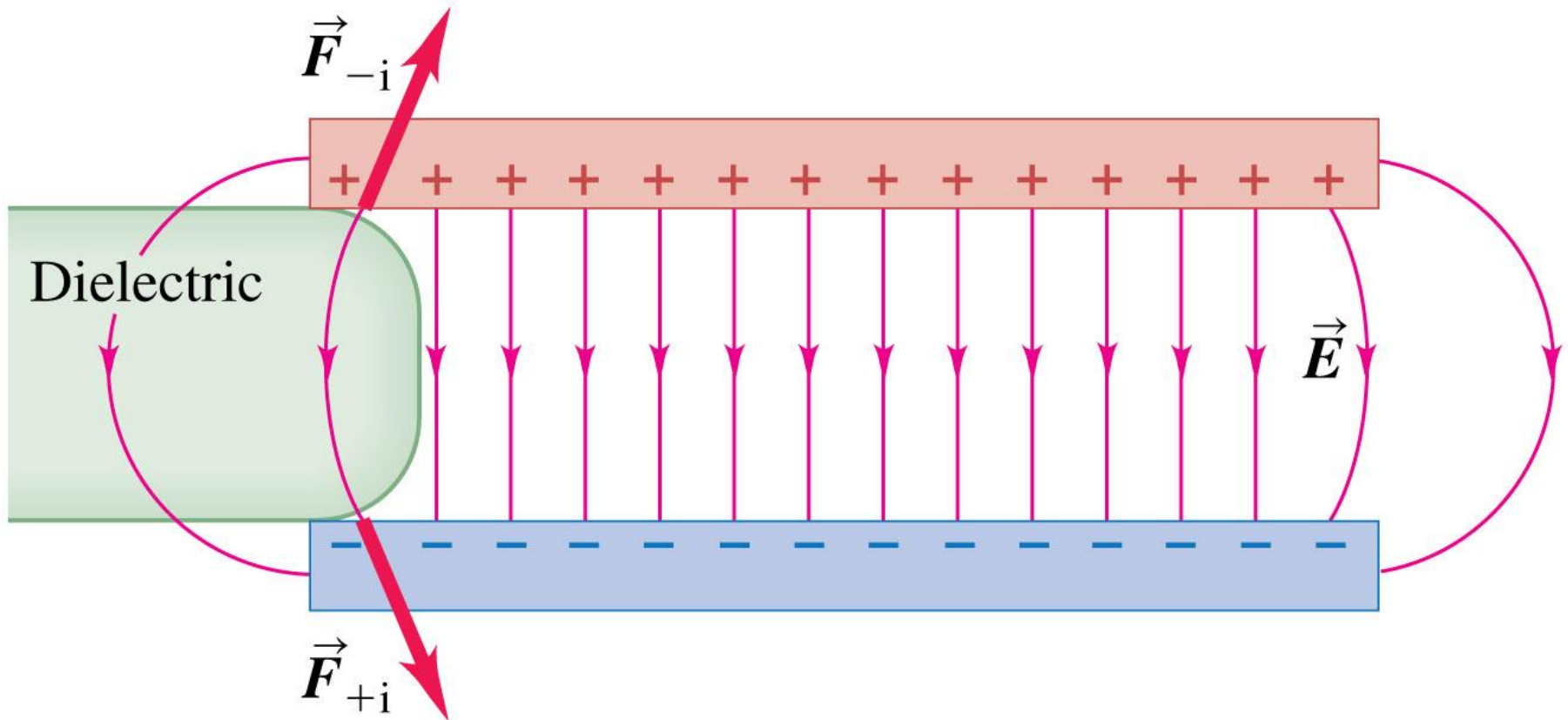
Should be 



<https://www.youtube.com/watch?v=qBbxSEp3-6o>

Can find the force using $F_x = -dU/dx$ and

$$W = \int dq \left(\frac{q}{C} \right) = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} C V^2$$



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Fringe Electric Field pulls the dielectric into the gap