

Electricity and Magnetism

- Physics 259 – L02
 - Lecture 38

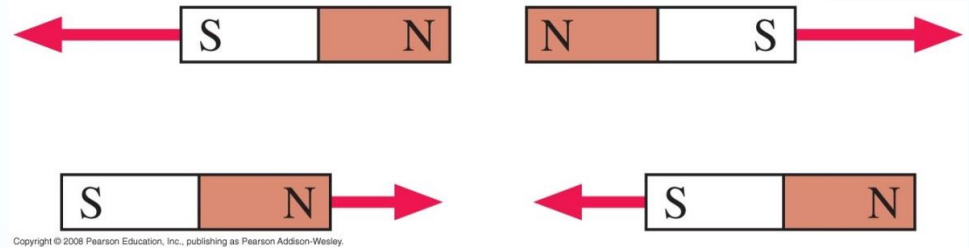
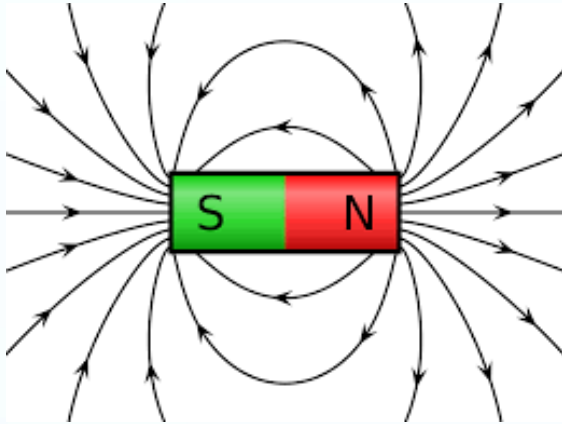


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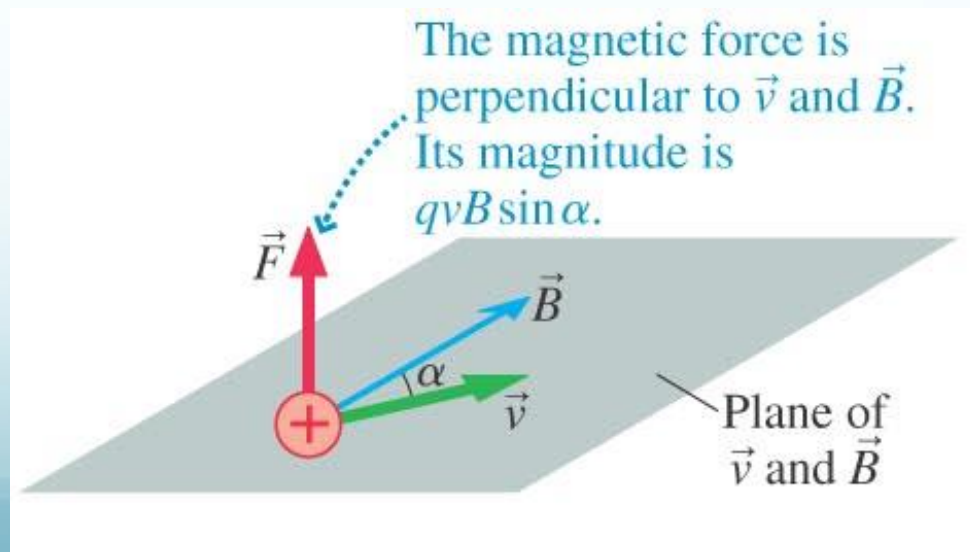
Chapter 28: Magnetic fields



28.1: Magnetic fields



$$\vec{F}_B = q\vec{v} \times \vec{B}$$

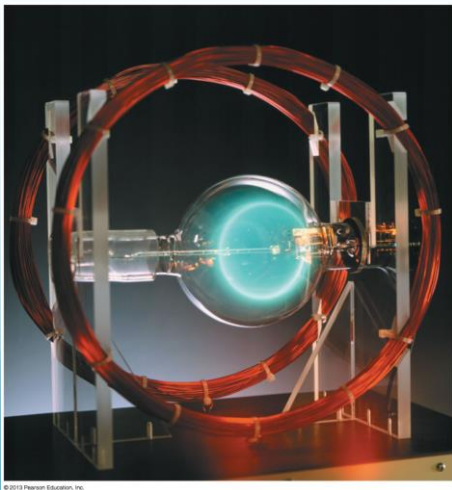


28.4: A circulating charged particle

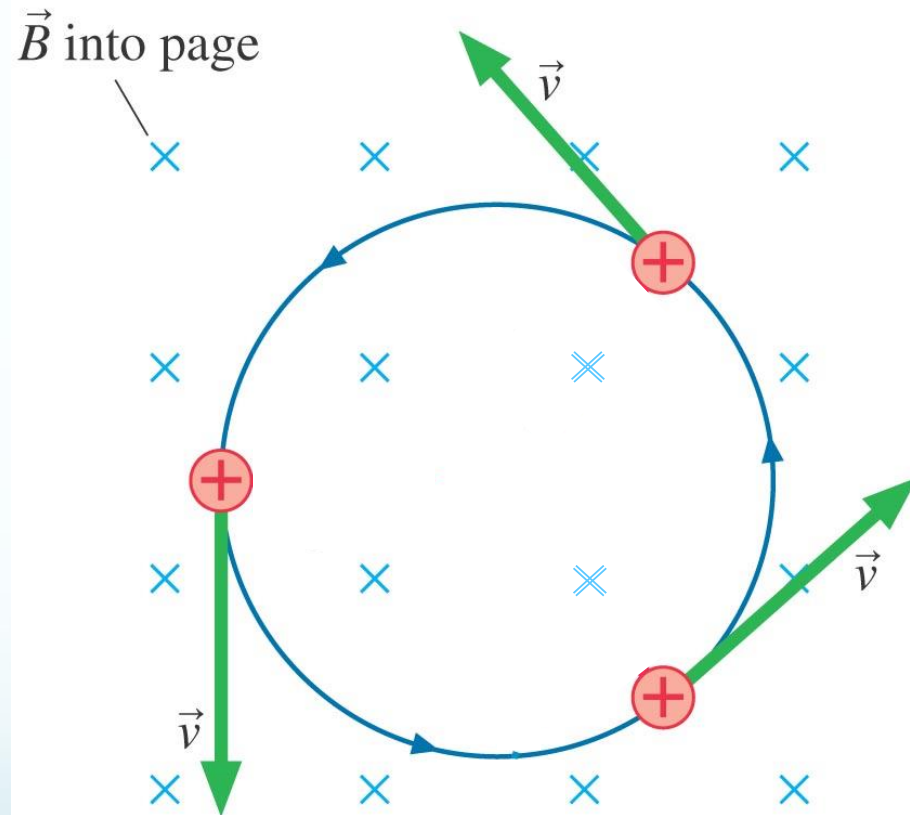
Charged particles in uniform magnetic fields undergo **uniform circular motion**.

The radius of the circle depends on how fast the particle is moving:

$$R = \frac{mv}{|q|B}$$



$$\vec{F}_B = q \vec{v} \times \vec{B}$$



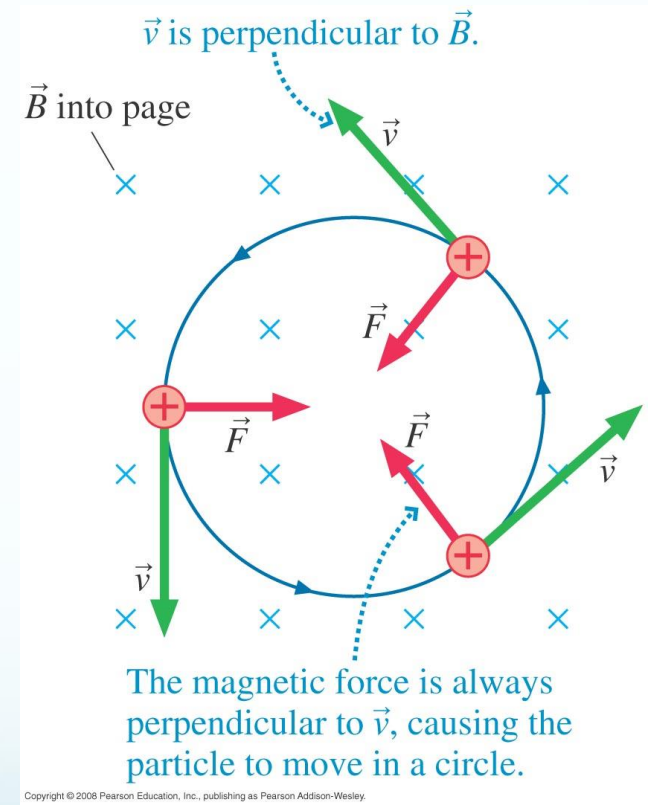
$$v = \frac{2\rho R}{T_{cyc}}$$

T_{cyc} is period (time it takes to make one cycle)

$$R = \frac{mv}{|q|B}$$

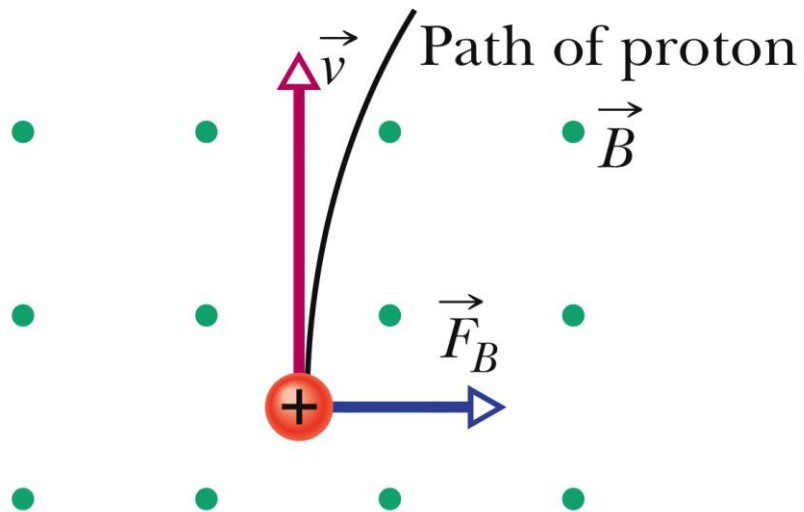
$$T_{cyc} = \frac{2\pi m}{|q|B}$$

$$f_{cyc} = \frac{|q|B}{2\pi m}$$



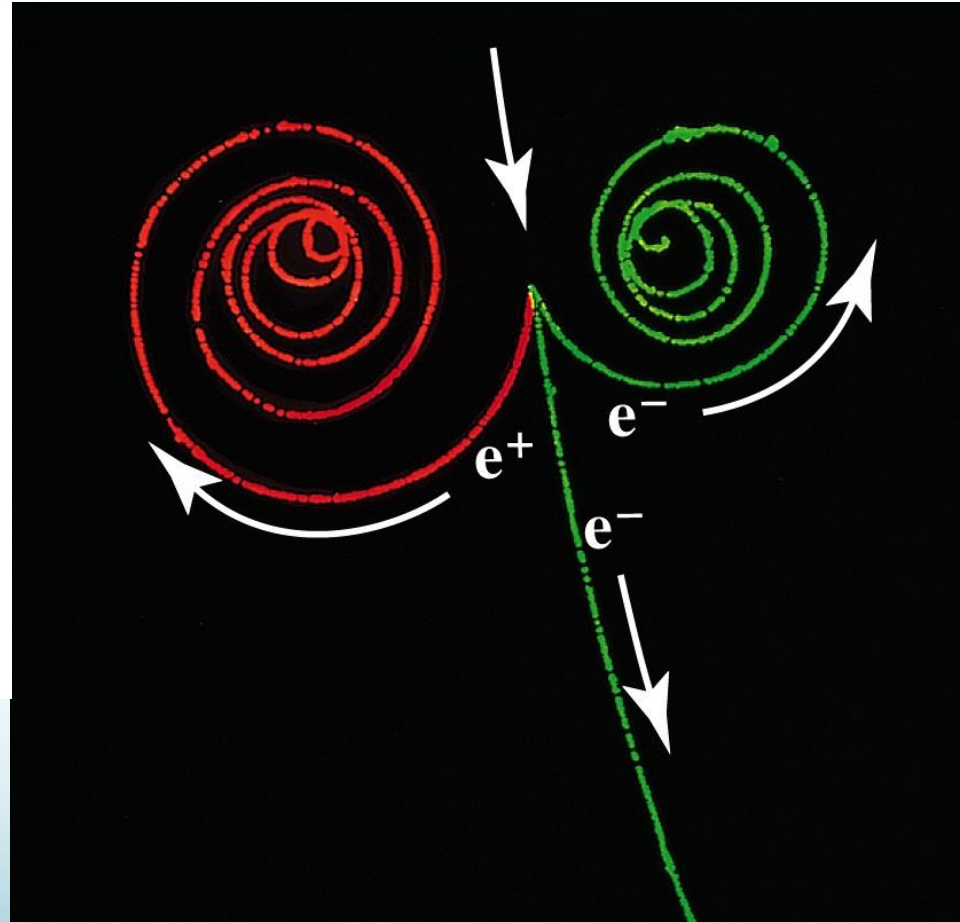
The period (and also the frequency of the circular motion) depend on the B-field strength and the charge-to-mass ratio q/m

Motion of charges in B-field



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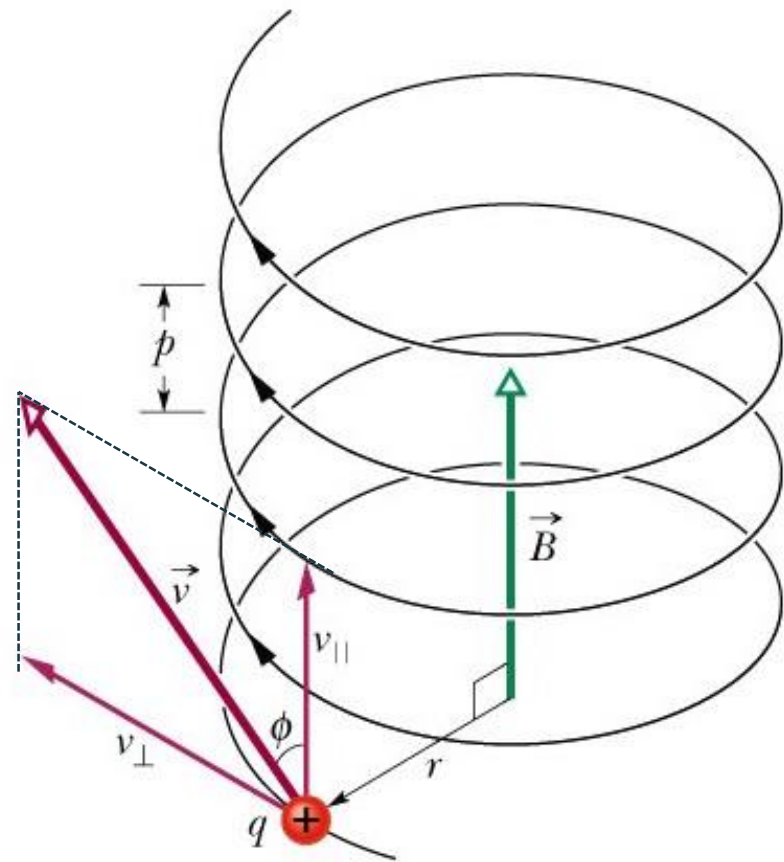
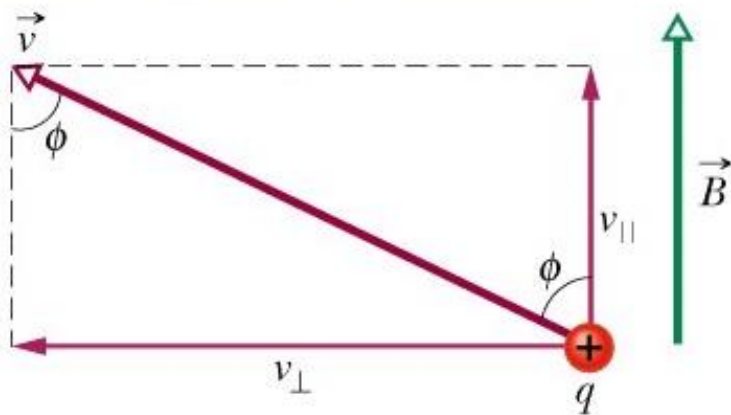


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Helical Paths Through a B-field

Splitting up the velocity into a component parallel to B-field and a component perpendicular to B-field immediately leads to helical motion

The velocity component perpendicular to the field causes circling, which is stretched upward by the parallel component.



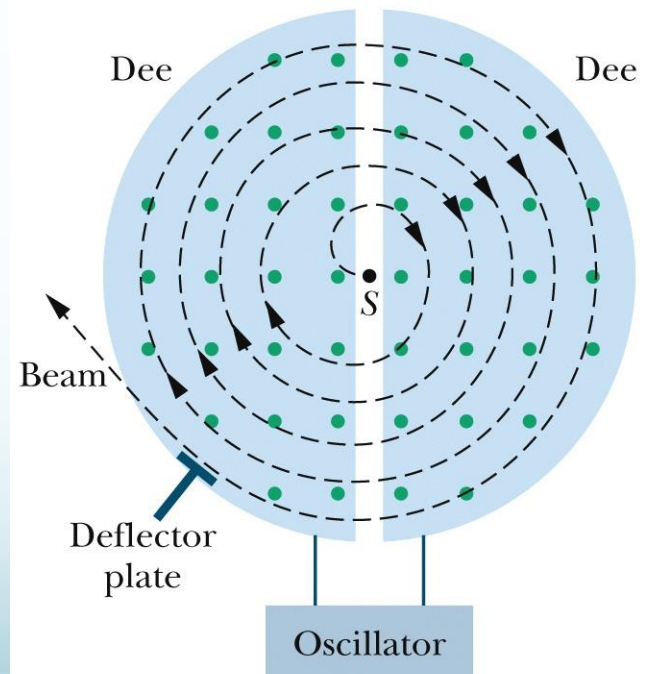
28.5: Cyclotrons and Synchrotrons

We need beams of high-energy particles→

Two accelerators that employ a magnetic field to repeatedly bring particles back to an accelerating region are:

- ✓ The Cyclotron
- ✓ The proton Synchrotron

The protons spiral outward in a cyclotron, picking up energy in the gap.



✓ The Cyclotron

The key to the operation of the cyclotron is that the frequency f at which the proton circulates in the magnetic field (and that does not depend on its speed) must be equal to the fixed frequency f_{osc} of the electrical oscillator, or

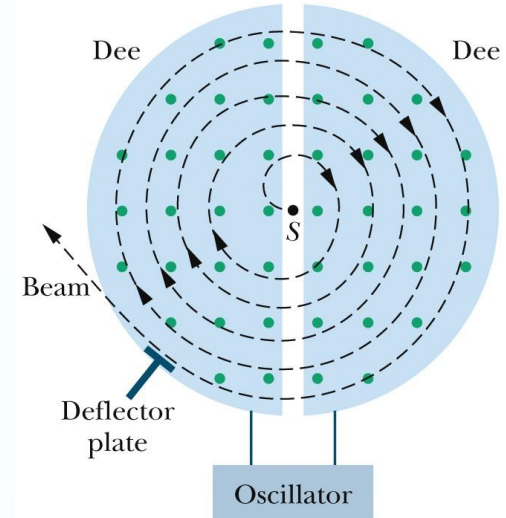
$$f = f_{osc} \quad (\text{resonance condition}).$$

✓ The proton Synchrotron

Magnetic field B and the oscillator frequency f_{osc} , instead of having fixed values as in the conventional cyclotron, are made to vary with time during the accelerating cycle →

- (1) the frequency of the circulating protons remains in step with the oscillator at all times
- (2) the protons follow a circular — not a spiral — path. Thus, the magnet need extend only along that circular path, not over some $4 \times 10^6 \text{ m}^2$.

The protons spiral outward in a cyclotron, picking up energy in the gap.

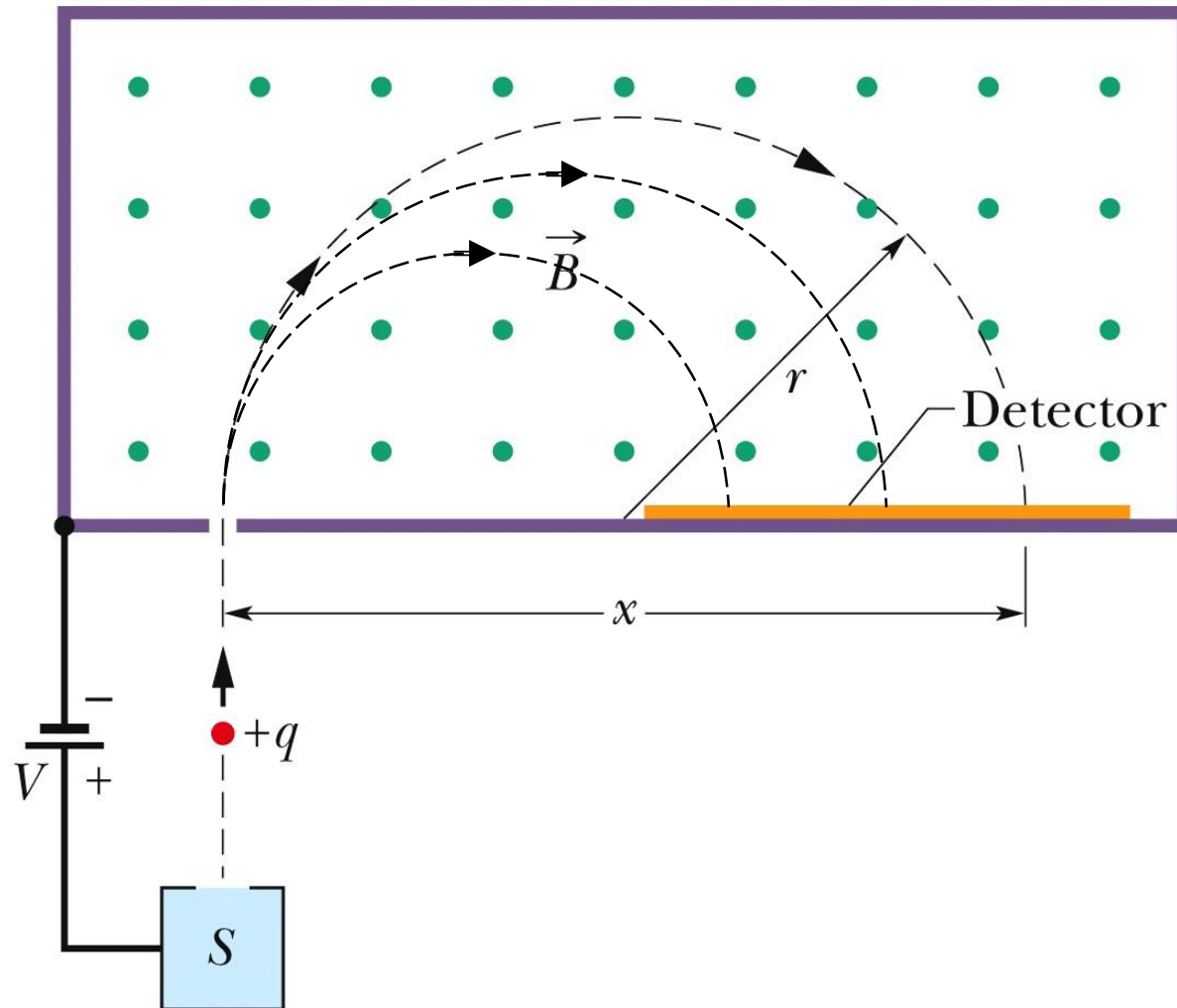


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$$R = \frac{mv}{|q|B} \quad T_{cyc} = \frac{2\rho m}{|q|B} \quad f_{cyc} = \frac{|q|B}{2\rho m}$$

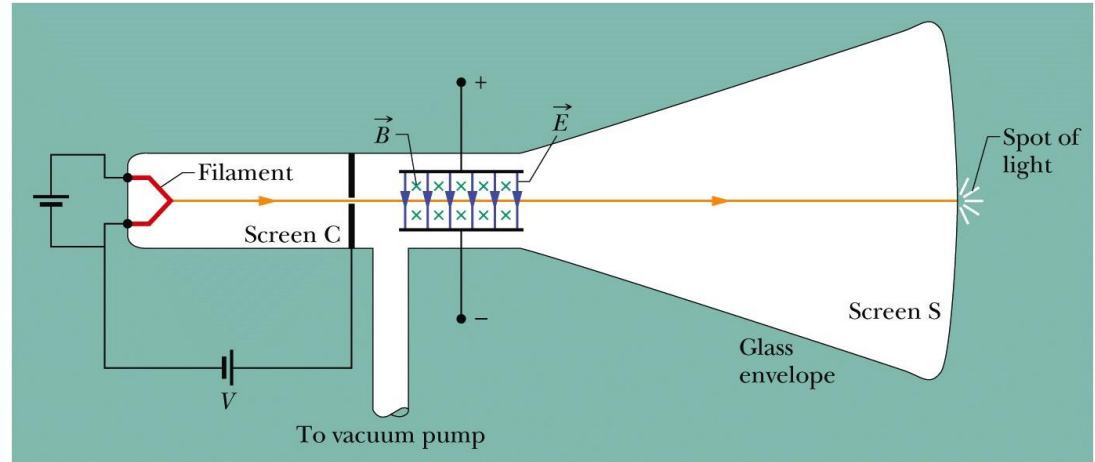
Application: Mass Spectrometer

$$R = \frac{mv}{|q|B}$$



28-2 Crossed Fields: Discovery of The Electron

A modern version of J.J. Thomson's apparatus for measuring the ratio of mass to charge for the electron



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If a charged particle moves through a region containing both an electric field and a magnetic field, it can be affected by both an electric force and a magnetic force.

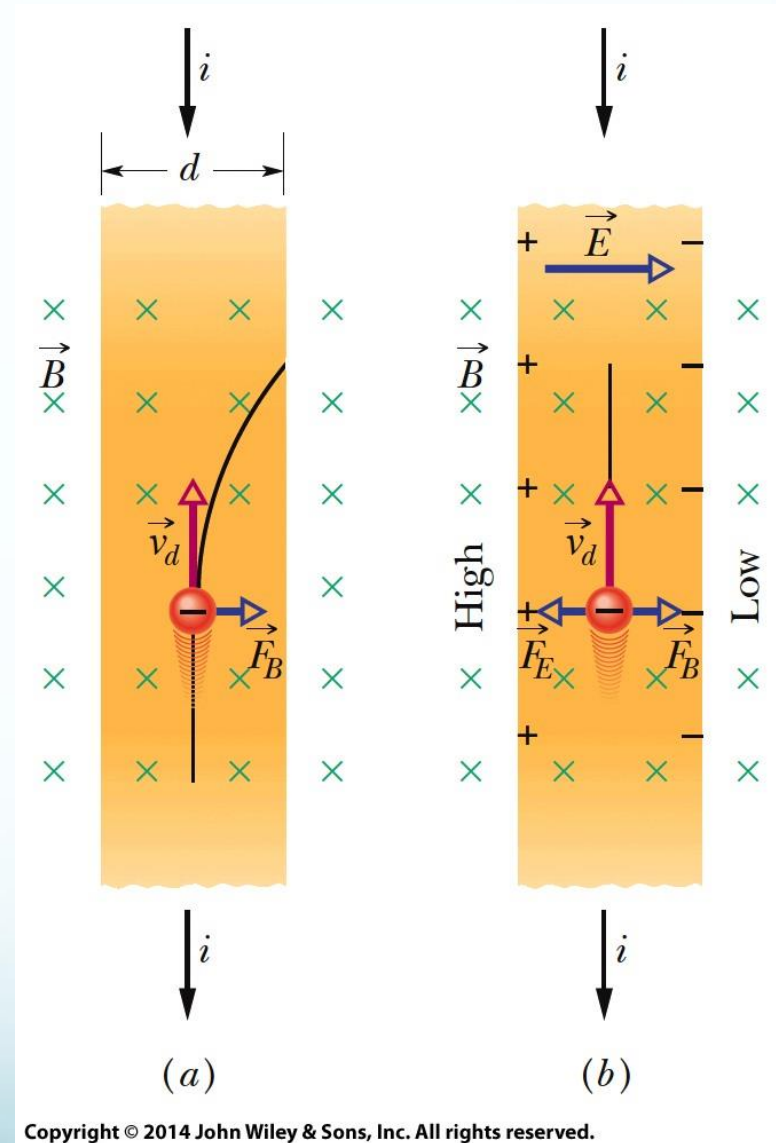
When the two fields are perpendicular to each other, they are said to be **crossed fields**.

28-3 Crossed Fields: The Hall Effect

A beam of electrons in a vacuum can be deflected by a magnetic field.

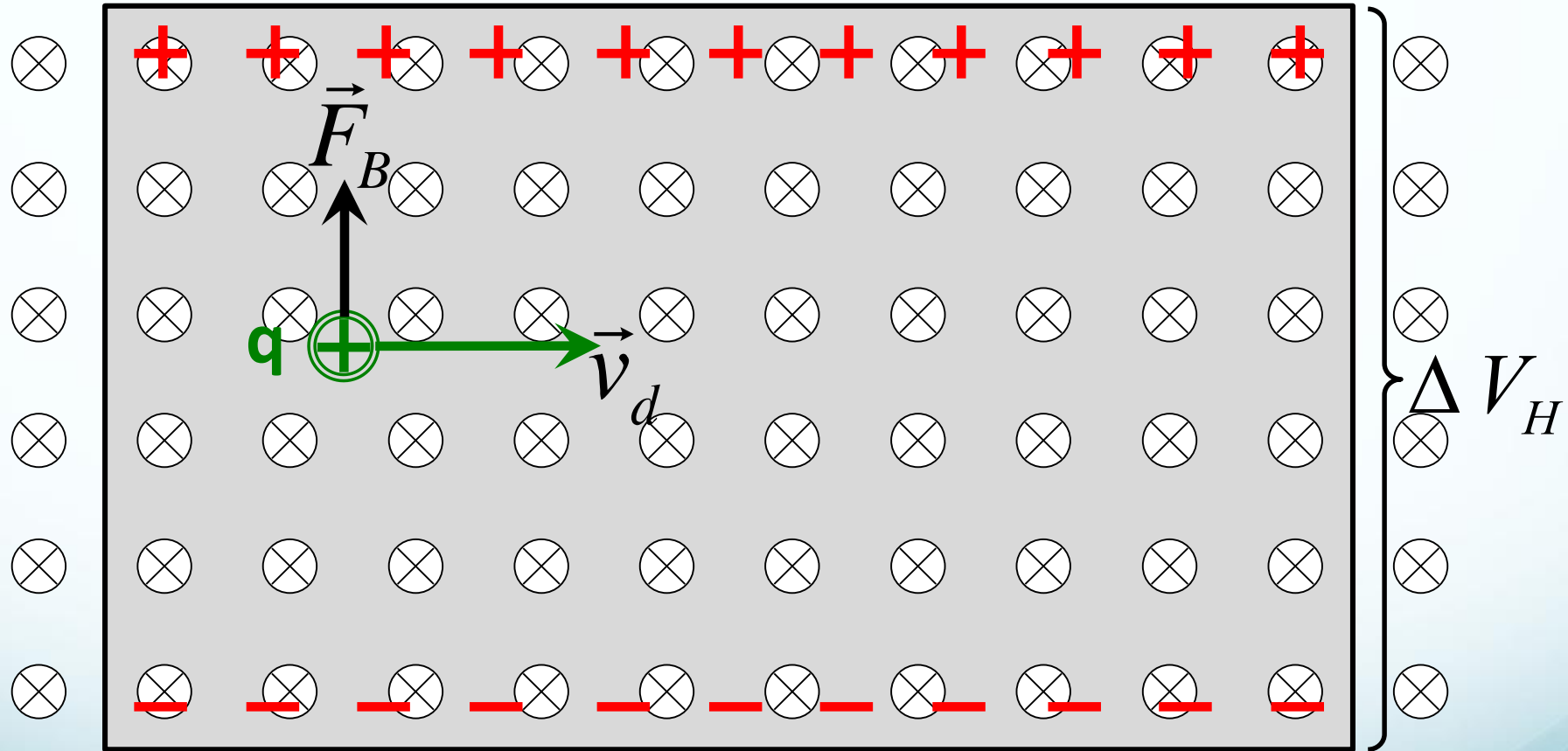
Can the drifting conduction electrons in a copper wire also be deflected by a magnetic field?

In 1879, Edwin H. Hall, then a 24-year-old graduate student at the Johns Hopkins University, showed that they can.

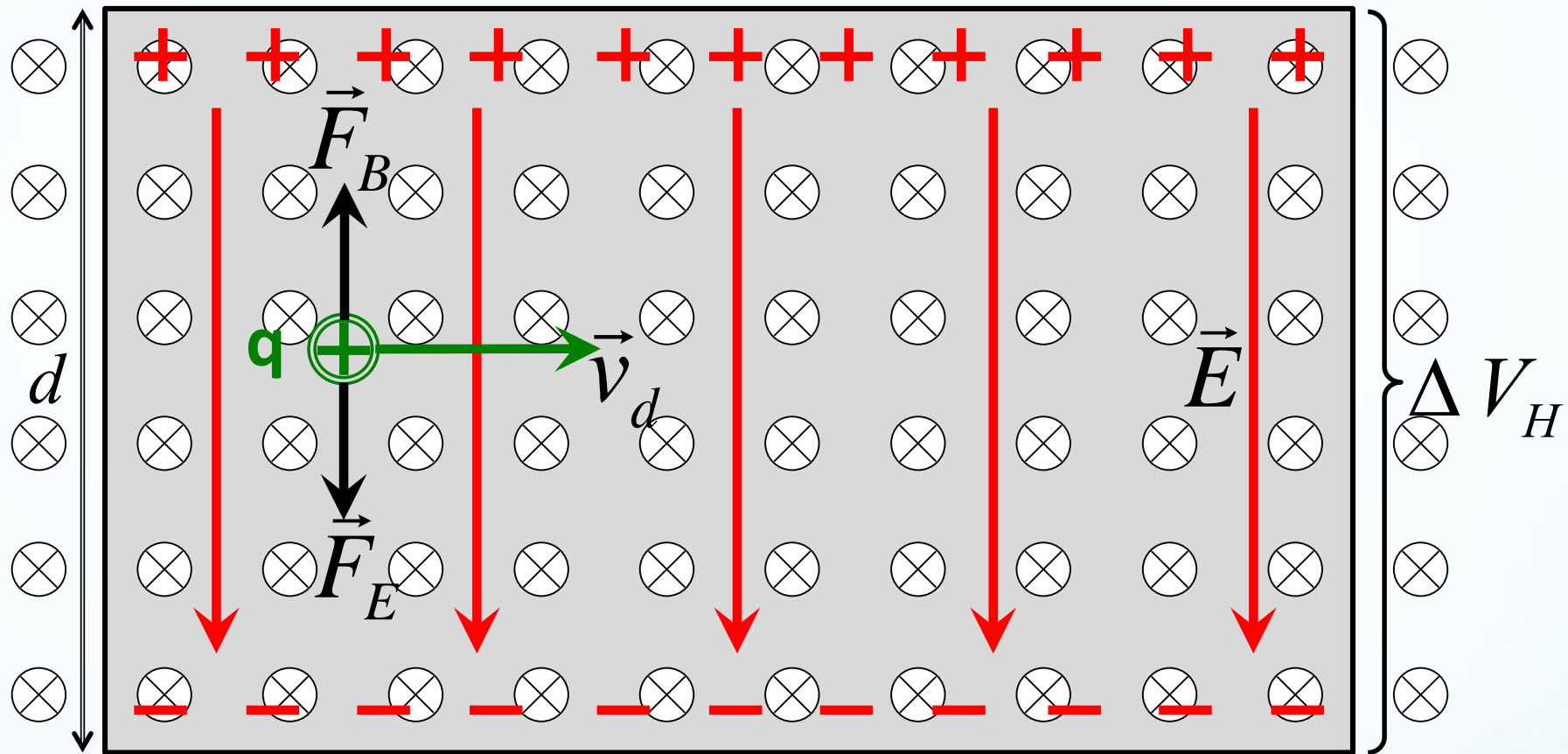


Explanation: Let's talk about positive charge

Due to the B-field, net charge build up on the edges.



In equilibrium, current still flows. Need to balance the magnetic and electric forces on the charge carriers.



$$F_B = q v_d B$$

$$F_E = q \frac{\Delta V_H}{d}$$

$$q \frac{\Delta V_H}{d} = q v_d B$$

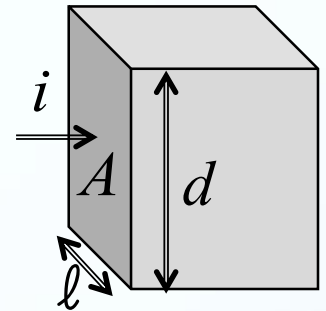
$$\Delta V_H = v_d B d$$

→ Voltage established across a conductor carrying a current in a magnetic field →

$$\Delta V_H = v_d B d$$

We previously related the drift speed to the current via

$$v_d = \frac{i}{neA} \quad \text{where } A = \ell d \text{ and } n \text{ is a material property}$$



We can then relate the Hall voltage to known quantities:

$$\Delta V_H = \frac{i}{ne\cancel{\ell d}} B \cancel{d} = \frac{iB}{ne\ell}$$

In practical applications, you measure ΔV_H to find B :

$$B = \frac{ne\ell}{i} \Delta V_H$$

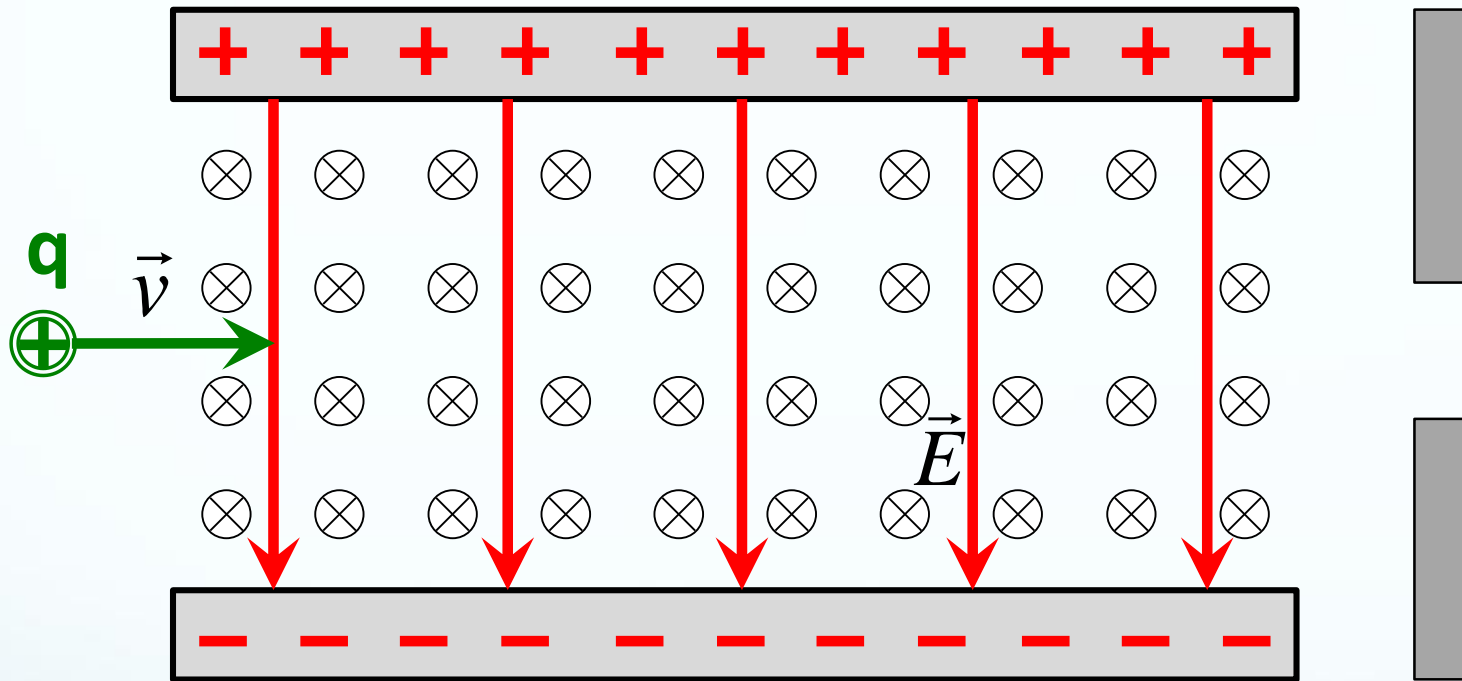
How the B-field probe used in the next lab works

This section we talked about:
Chapter 28.2, 28.4, 28.3, 28.5

See you on Thursday



Similar concept: velocity selector

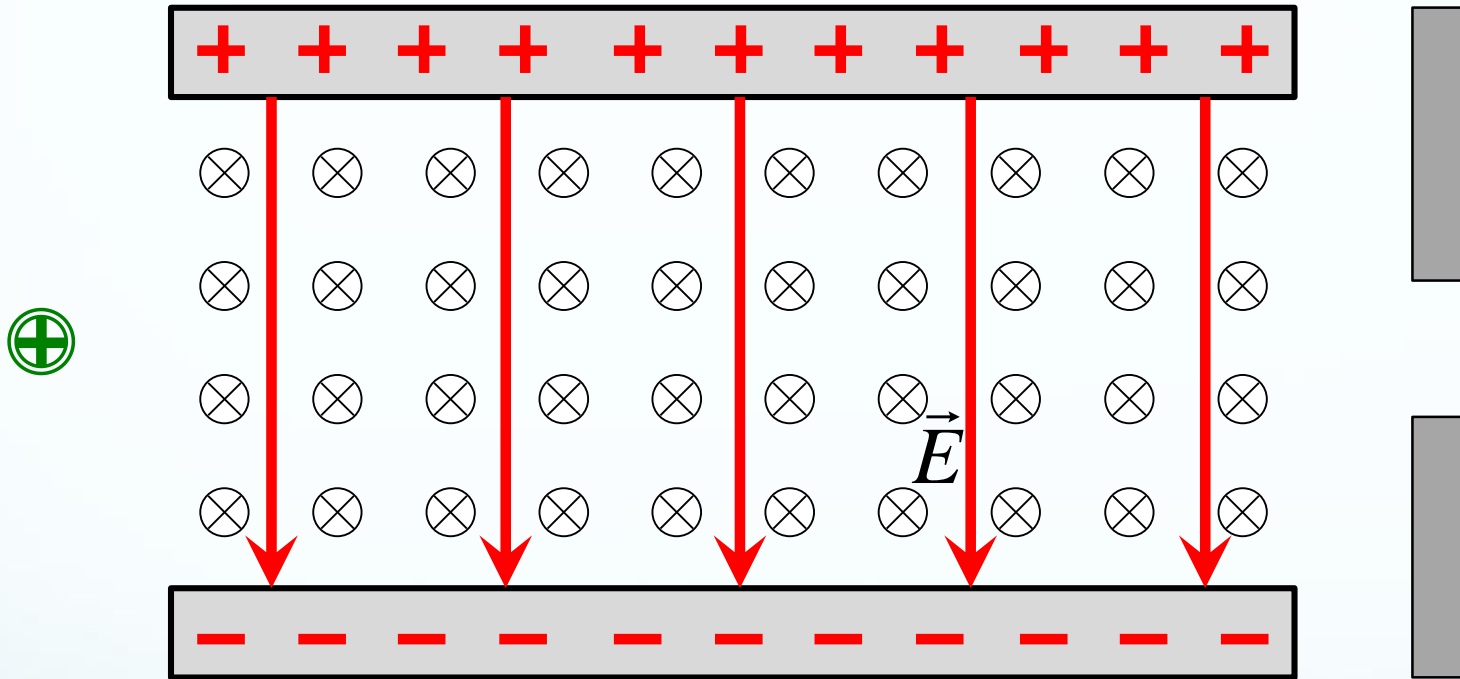


In a velocity selector, you send a charge through a region with crossed E and B fields, which leads to electric and magnetic forces:

$$\vec{F}_e = q\vec{E} \quad \vec{F}_B = q\vec{v} \times \vec{B} \quad qE = qvB \quad v = \frac{E}{B}$$

If the forces balance ($F_{\text{net}} = 0$) the charge makes it through the slit

Similar concept: velocity selector



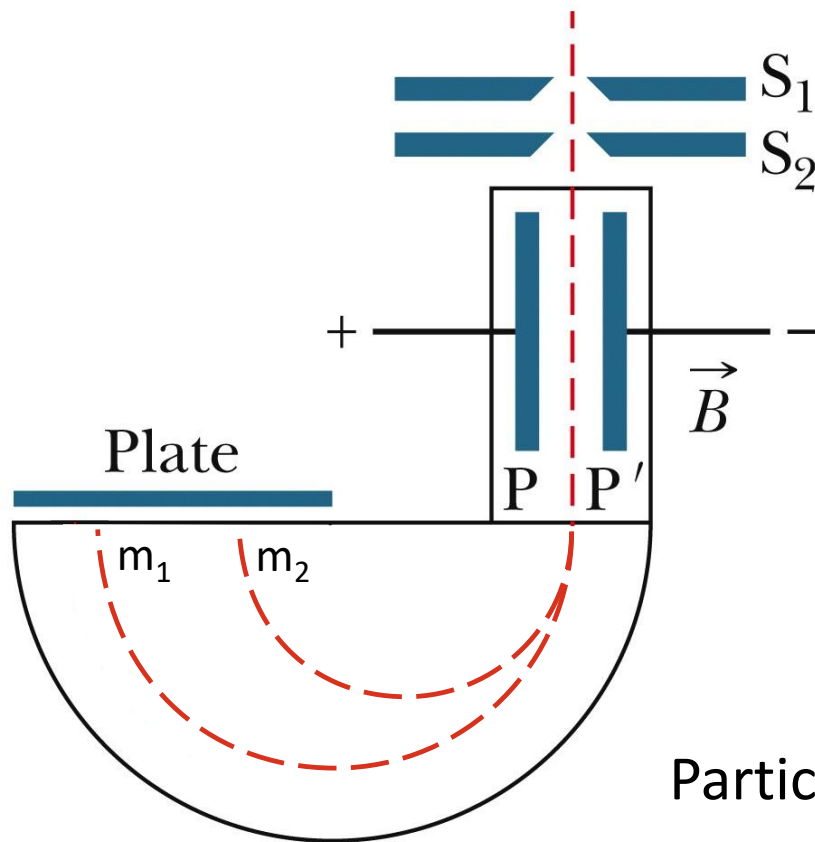
If the forces don't balance the charge hits the wall

$$qE - qvB = ma$$

We pick the E and B magnitudes to select the speeds we want

Bainbridge Mass Spectrometer

Accelerate charges through ΔV so they all have same Kinetic Energy



The slits S_1 and S_2 ensure the beam of particles is collimated.

The beam enters a region of crossed E and B-fields

A narrow slit ensures only particles with a specific speed enter

Particles with same KE but different masses and charges will have different radius in B field