

# Electricity and Magnetism

- Physics 259 – L02
- Lecture 43



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## Chapter 29: Magnetic field due to current



## Last time:

- Biot-Savart Law (like Coulomb's Law for magnetism)
- B-field of a line of current
- Magnetic force between parallel current-carrying wires

## Today:

- Ampere's law
- Applications of Ampere's law



For a single charge →

$$\vec{F}_B = q \vec{v}_d \times \vec{B}$$

For N charges moving through the wire  
(current carrying wire) →

$$\vec{F}_B = i \vec{\ell} \times \vec{B}$$

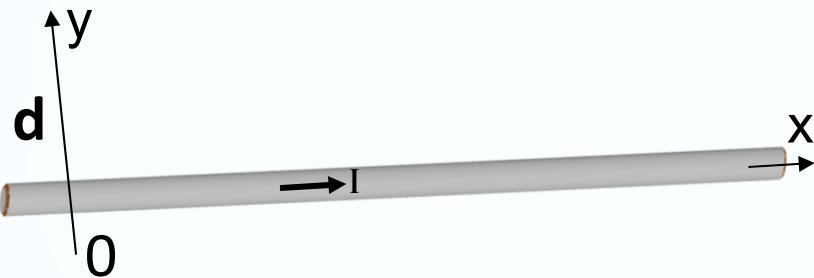
The Biot-Savart Law →

$$\vec{B}_{\text{point charge}} = \frac{\mu_0}{4\pi} \frac{q \vec{v} \times \hat{r}}{r^2}$$

For an electric current →

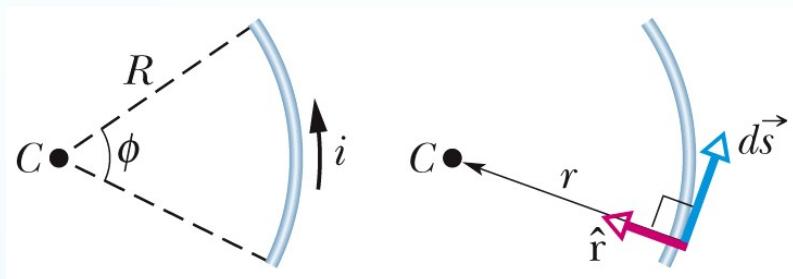
$$\vec{B}_{\text{current segment}} = \frac{\mu_0}{4\pi} \frac{I \Delta \vec{s} \times \hat{r}}{r^2}$$

# Magnetic field due to current in long straight wire

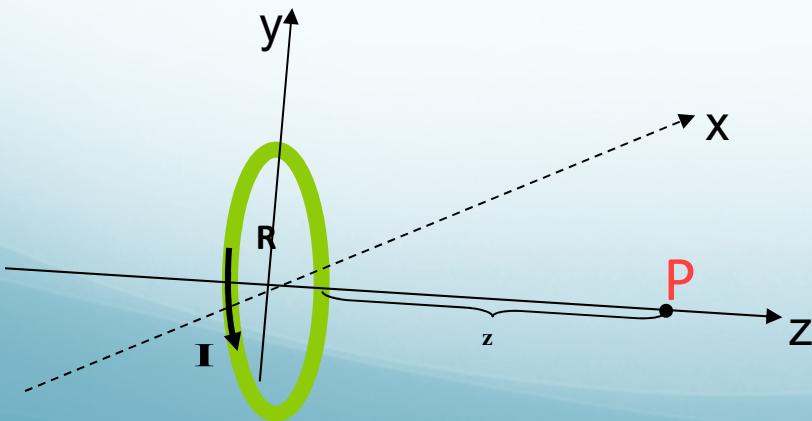


$$B_z = \frac{\mu_0}{2\pi} \frac{I}{d}$$

Non-infinite straight wire → Appendix 1-chapter 22



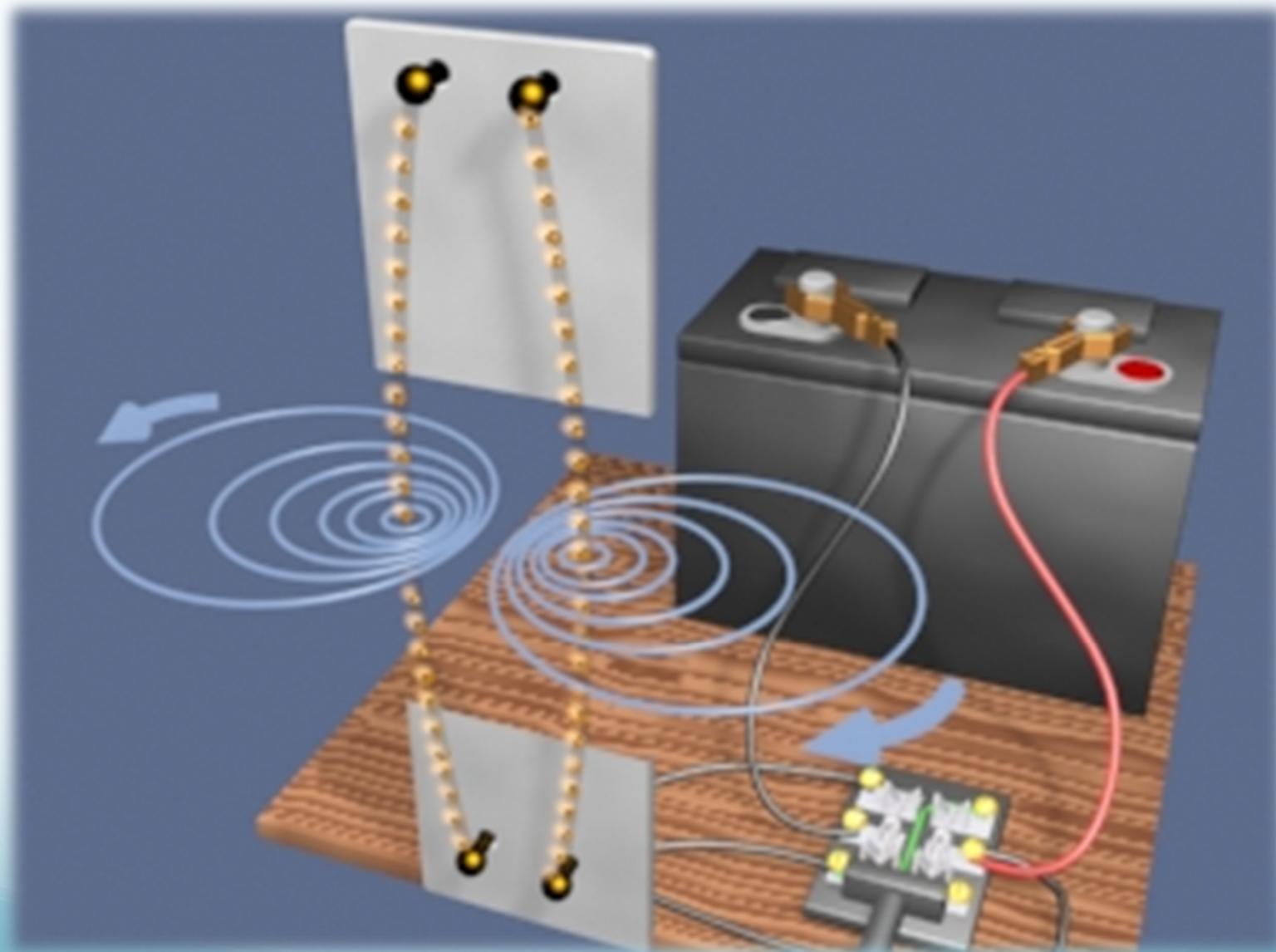
$$B = \frac{\mu_0 i \phi}{4\pi R}$$



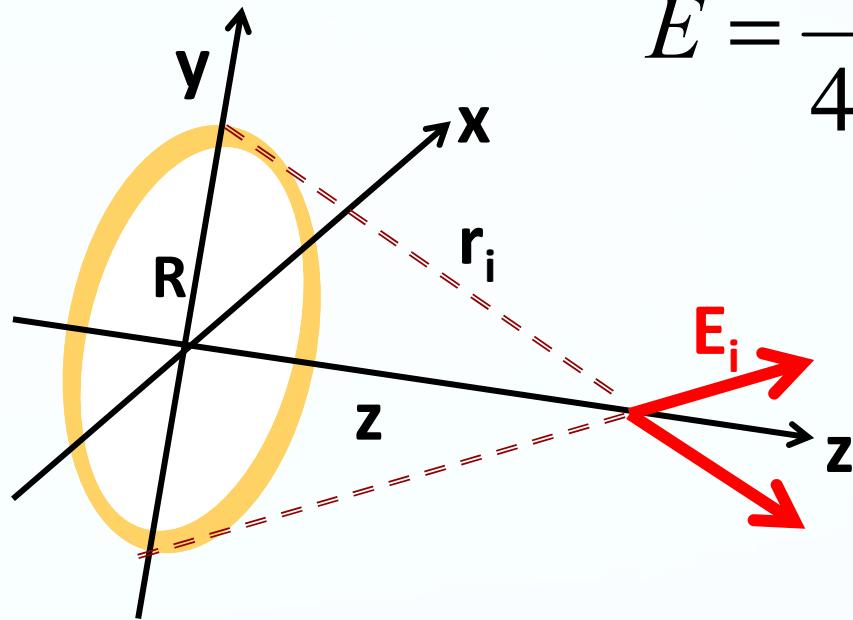
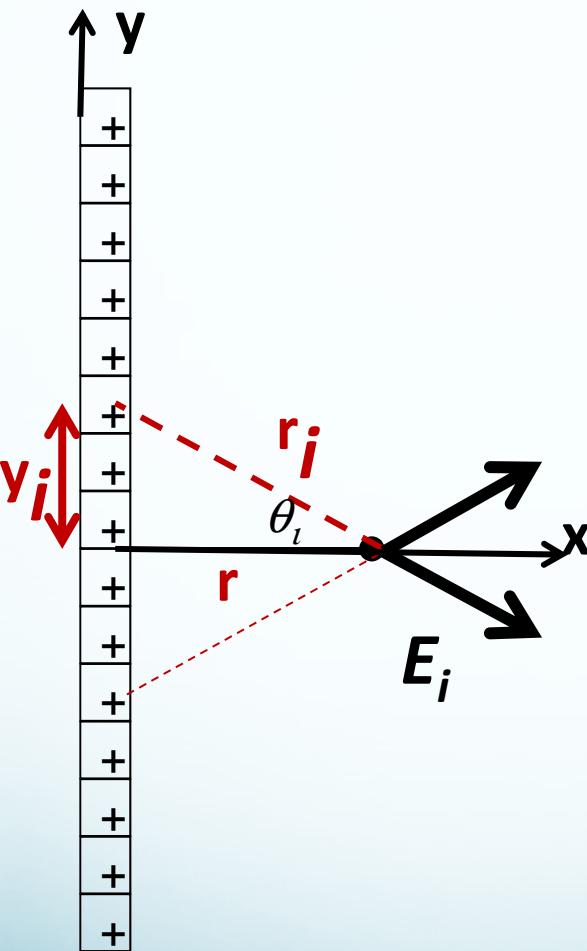
$$\vec{B} = \frac{\mu_0}{2} \frac{IR^2}{(z^2 + R^2)^{3/2}} \hat{k}$$

if  $z = 0$        $\vec{B}_{\text{center}} = \frac{\mu_0}{2} \frac{I}{R} \hat{k}$

## 29.2: Force between two antiparallel currents



# Electrostatics



$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r}$$

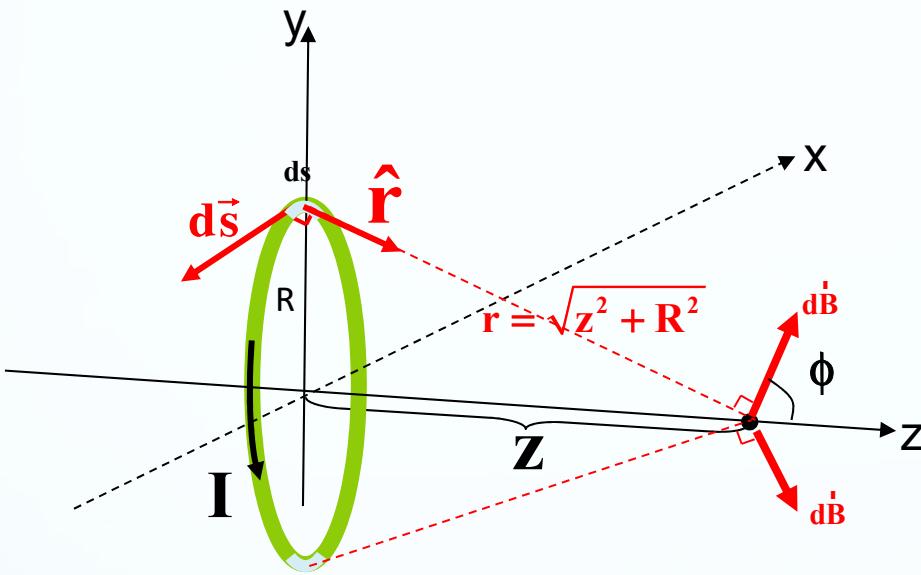
Savior:  
**Gausses' law**

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{in}}{\epsilon_0}$$



# Magnetostatics

$$\vec{B}_{\text{current segment}} = \frac{\mu_o}{4\pi} \frac{I \Delta \vec{s} \times \hat{r}}{r^2}$$



Savior:

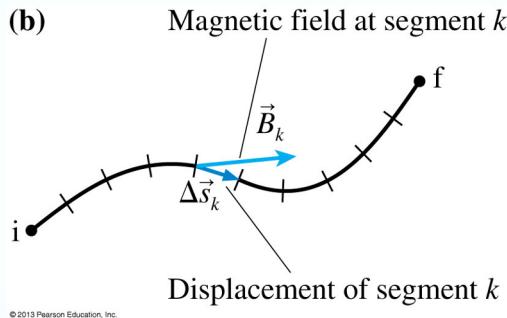
**Ampere's law**

Expression?



# Ampère's law

The line integral of  $\mathbf{B}$  along the path:



$$\int_i^f \vec{B} \cdot d\vec{l}$$

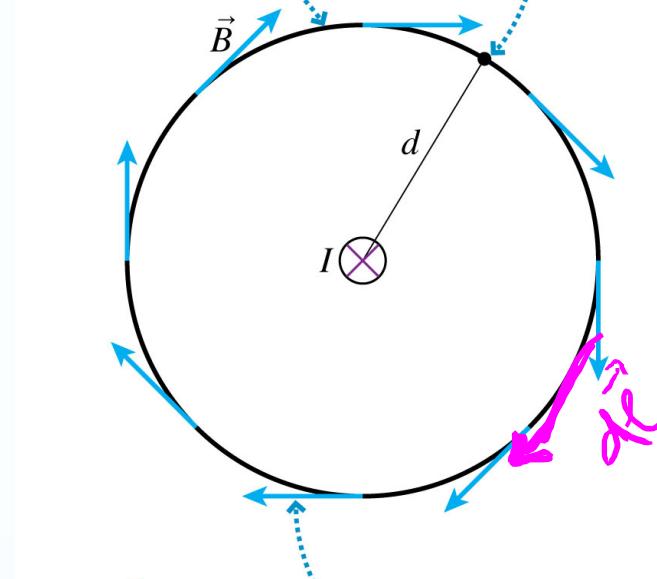
$$\oint \vec{B} \cdot d\vec{l} = (2\pi r) \left( \frac{\mu_0 I}{2\pi r} \right)$$

$\Rightarrow$  i.e.  $\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{encl}$

Ampère's Law is true for any shape of path and any current distribution

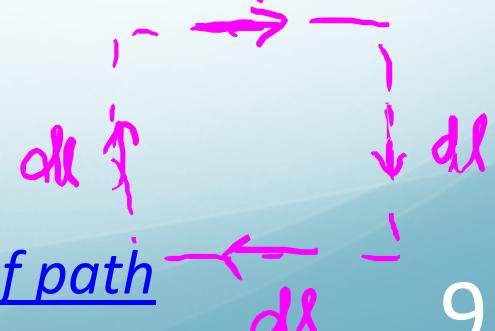
The integration path is a circle of radius  $d$ .

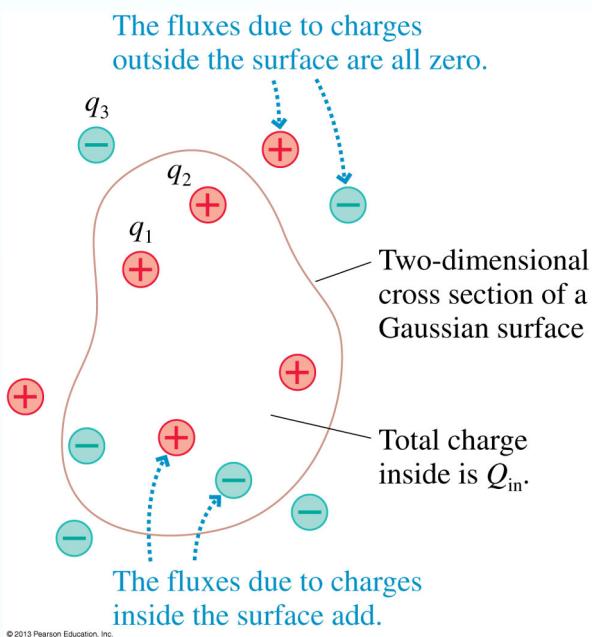
The integration starts and stops at the same point.



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Infinite wire  $\rightarrow B = \frac{\mu_0 I}{2\pi r}$

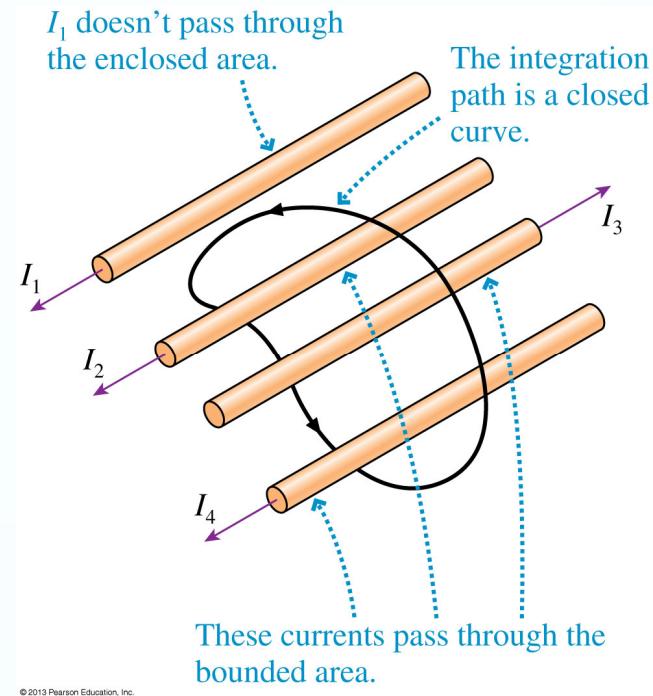




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$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{in}}{\epsilon_0}$$

For a closed surface enclosing  
total Charge Q

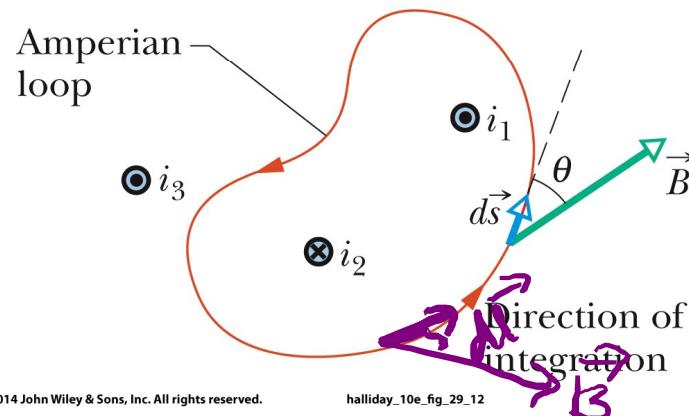


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$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enclosing}$$

Current I passes through an area bounded by a closed curve

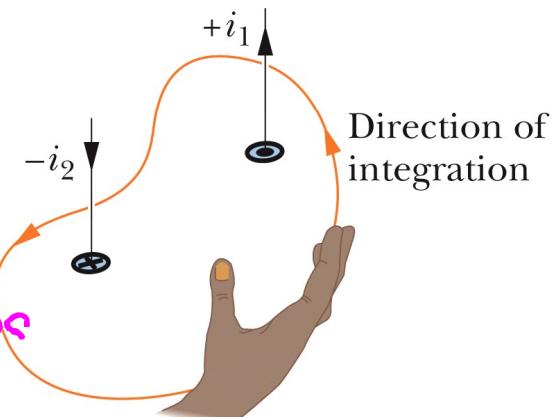
Only the currents encircled by the loop are used in Ampere's law.



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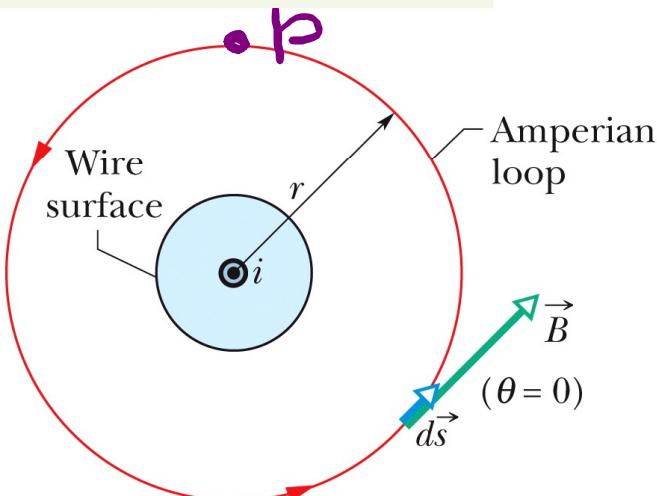
This is how to assign a sign to a current used in Ampere's law.



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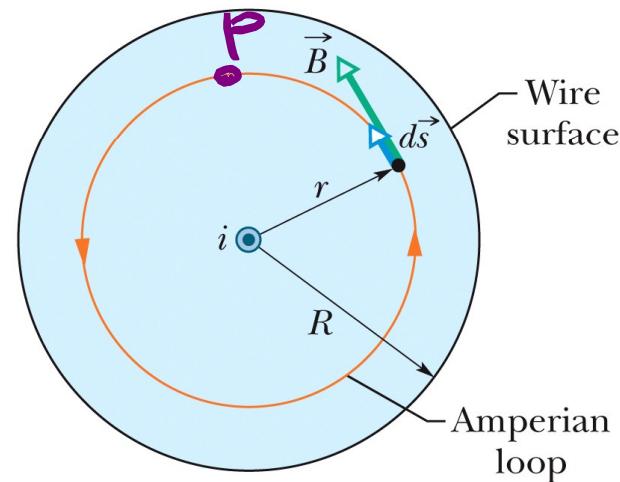
All of the current is encircled and thus all is used in Ampere's law.



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Only the current encircled by the loop is used in Ampere's law.



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# TopHat Question

What is  $I_{\text{encl}}$  here, where all three wires have 5 A?

A) -5 A

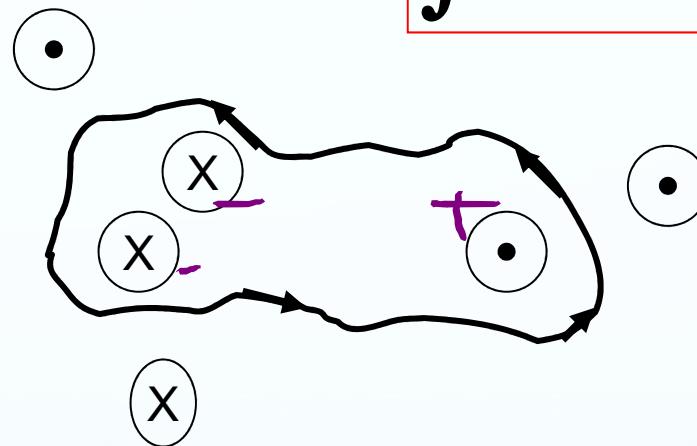
B) 5A

C) 15 A

D) -15 A

E) other

$$\oint \vec{B} \cdot d\vec{L} = \mu_0 I_{\text{through}}$$



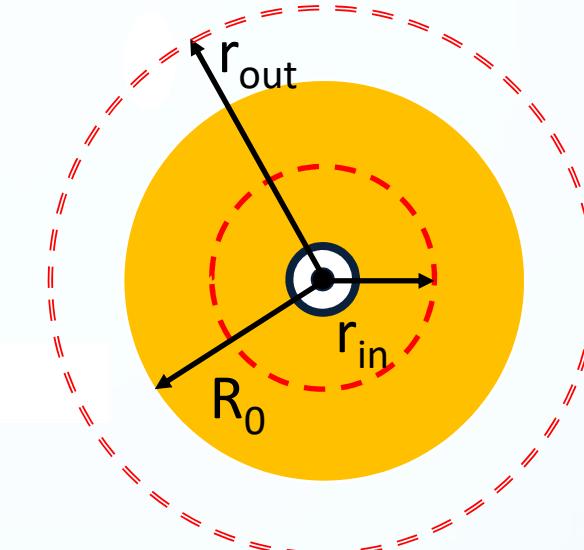
# Ampère's law: application

- (a) Using Ampère's law, calculate the magnetic field **outside** a solid current carrying wire a distance  $r_{\text{out}}$  from its axis

(The length of the solid wire is infinite and the current  $I$  is uniformly distributed throughout the solid wire)

- b) Calculate the magnetic field **inside** a solid current carrying wire a distance  $r_{\text{in}}$  from its axis.

$$\oint \vec{E} \cdot d\vec{A} = \frac{\rho_{\text{enc}}}{\epsilon_0}$$



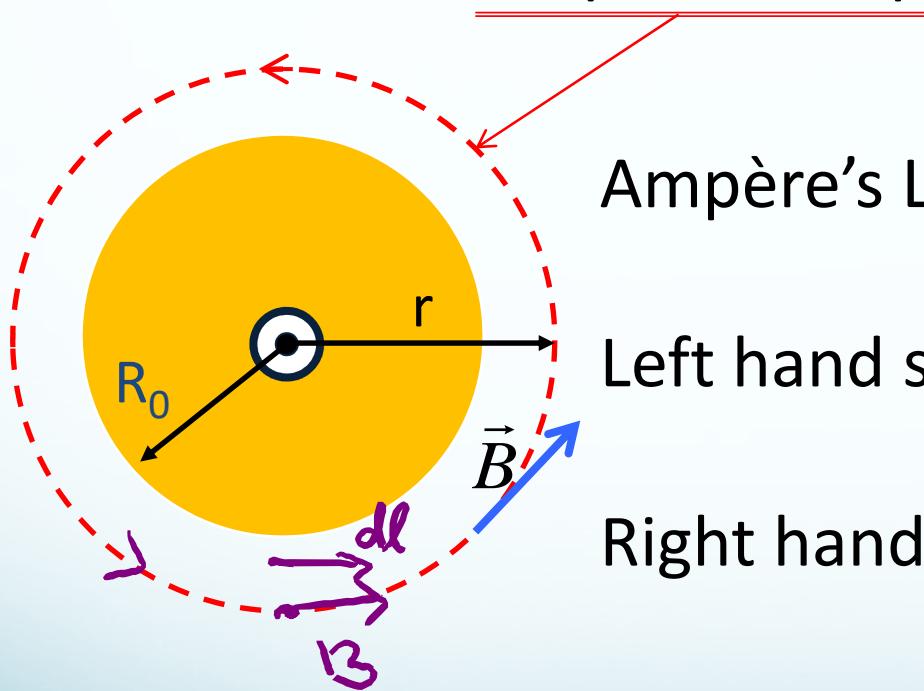
End view:  
Wire with radius  $R$   
and current  $I$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enc}}$$
$$J = \frac{I}{A}$$

# Ampère's law: application(1)

## (a) B-field outside

We want to know the B-field a distance  $r$ , so we choose an Ampèrian loop with radius  $r > R_0$ .



Ampère's Law:

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enc}$$

Left hand side:

$$\oint \vec{B} \cdot d\vec{l} = B L = B 2\pi r$$

Right hand side:

$$\mu_0 I_{enc} = \mu_0 I$$

Combine together:

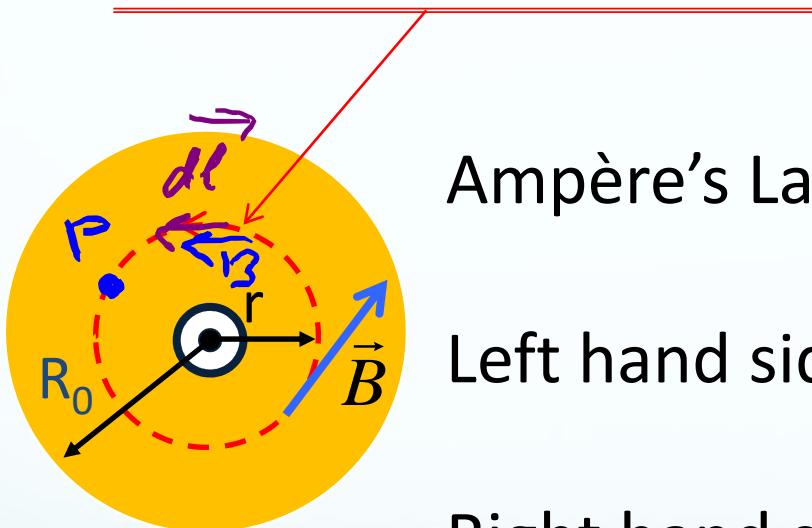
$$B 2\pi r = \mu_0 I$$

$$B = \frac{\mu_0 I}{2\pi r}$$

# Ampère's law: application(1)

## (b) B-field inside

We want to know the B-field at a distance  $r$ , so we choose an Amperian circular loop with radius  $r < R_0$ .



Ampère's Law:

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enc}$$

Left hand side:

$$\oint \vec{B} \cdot d\vec{l} = BL = \boxed{B2\pi r}$$

Right hand side:  $\mu_0 I_{enc} = \mu_0 JA = \mu_0 \frac{I}{\pi R^2} \pi r^2$

$$\rightarrow \mu_0 I_{enc} = \mu_0 JA = \mu_0 \frac{I}{\pi R^2} \pi r^2 = \mu_0 \frac{I}{A} \pi r^2 = \mu_0 \frac{I}{\pi R_0^2} \pi r^2$$

Combine together:

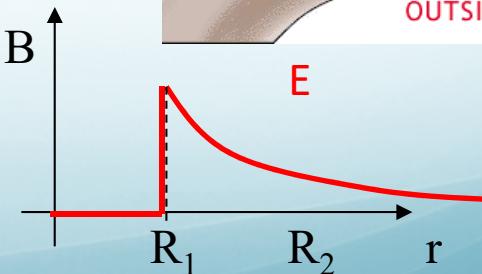
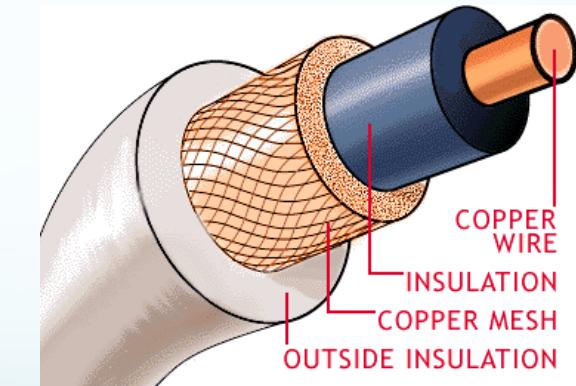
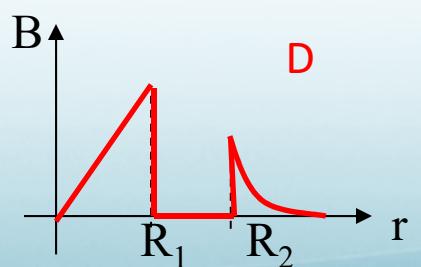
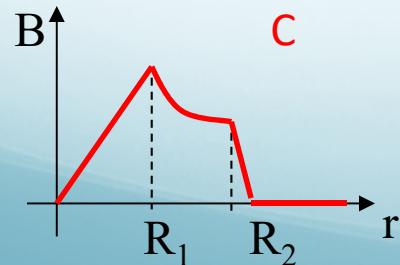
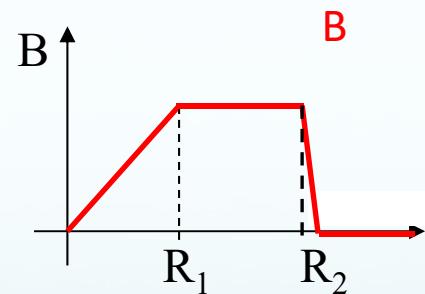
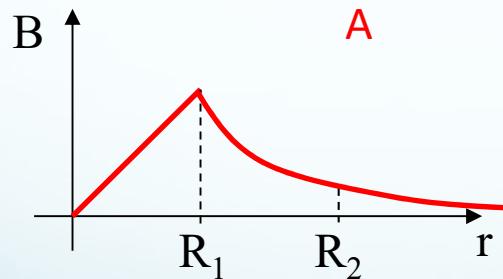
$$B2\pi r = \mu_0 \frac{I}{R_0^2} r^2$$

$$\boxed{B = \frac{\mu_0 I r}{2\pi R_0^2}}$$

A coaxial cable consists of a wire (radius  $R_1$ ) surrounded by an insulating sleeve and another cylindrical conducting shell (inner radius  $R_2$ ) and finally another insulating sleeve. **The wire and the shell carry the same current I but in opposite directions.**

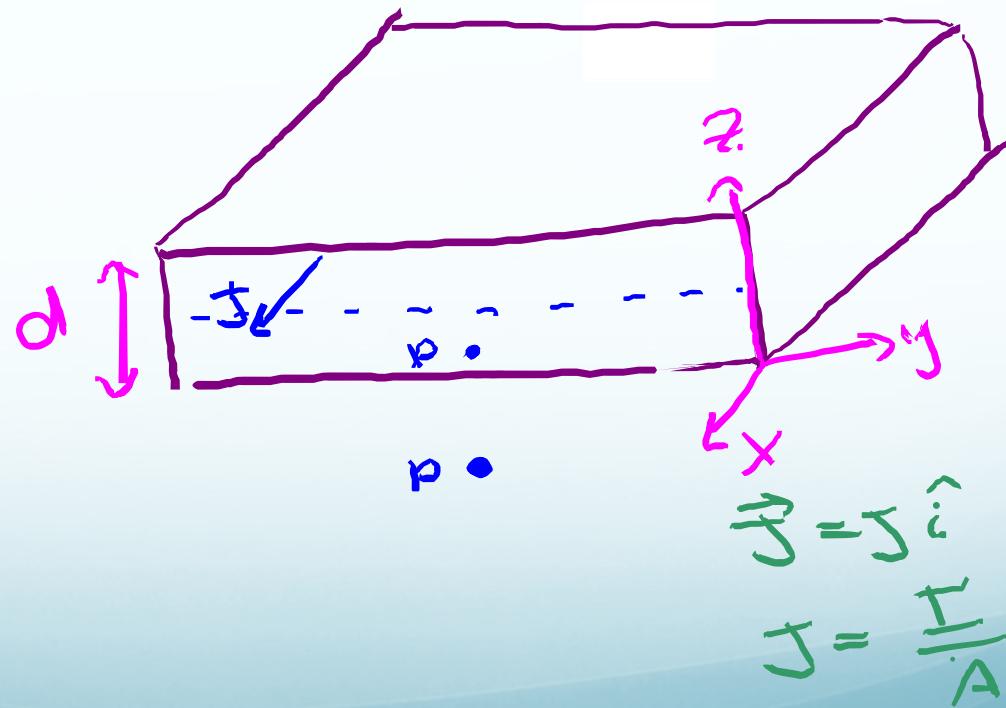


Which diagram best represents the **magnetic field** as a function of radial distance from the cable's axis?



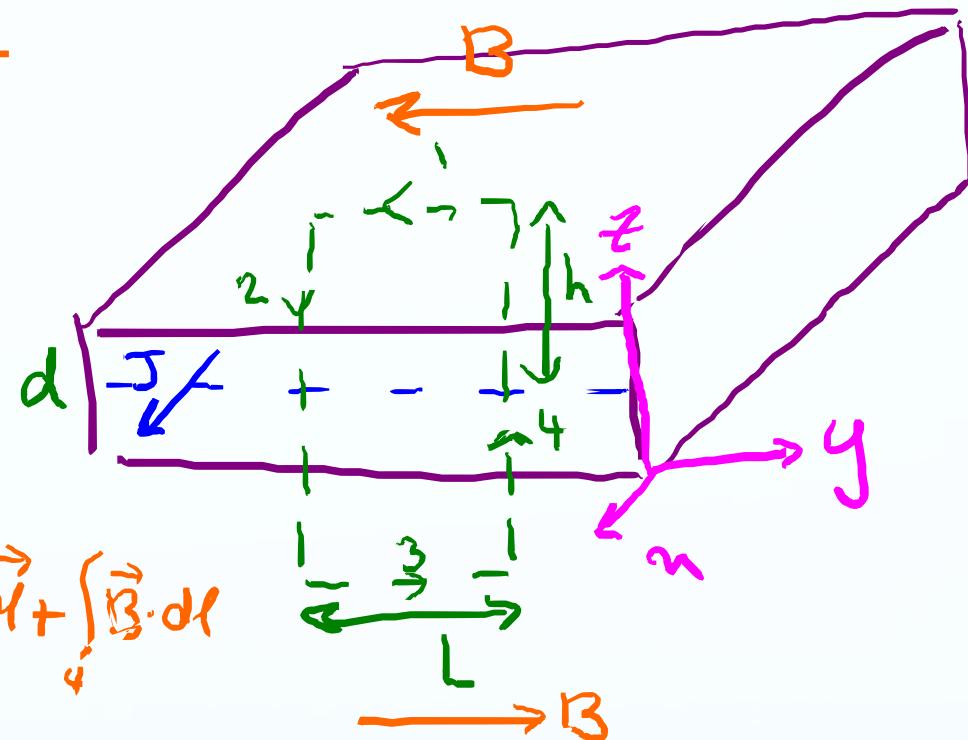
# Ampère's law: application (3)

- (a) Using Ampère's law, calculate the magnetic field **above** the current carrying slab
- b) Calculate the magnetic field **inside** the current carrying slab



$\vec{B}$  above or below the slab  $\Rightarrow$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enc}}$$



$$\oint \vec{B} \cdot d\vec{l} \Rightarrow$$

$$\rightarrow \int_1 \vec{B} \cdot d\vec{l} + \int_2 \vec{B} \cdot d\vec{l} + \int_3 \vec{B} \cdot d\vec{l} + \int_4 \vec{B} \cdot d\vec{l}$$

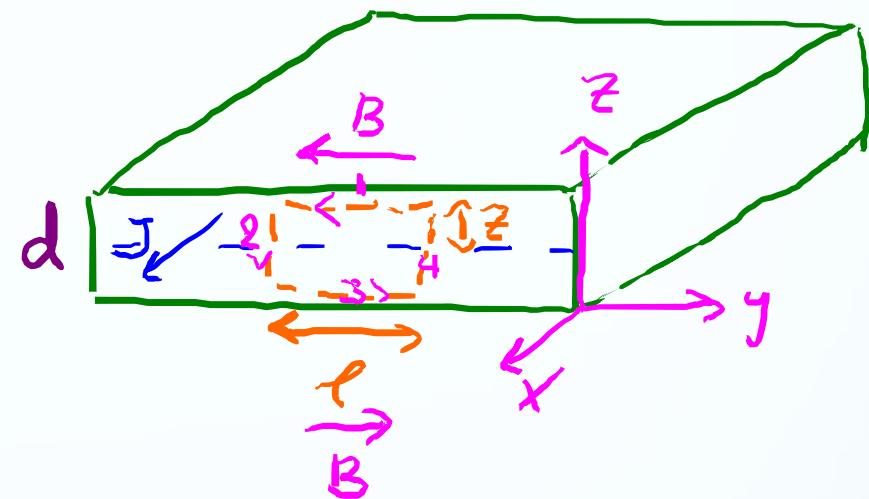
$$\rightarrow 2B \int d\ell = \mu_0 I_{\text{enc}} = \mu_0 J A = \mu_0 J L d$$

$$\rightarrow B = \frac{\mu_0 J d}{2}$$

$\vec{B}$  inside the slab  $\Rightarrow$

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{enc}}$$

$$\rightarrow \left\{ \int_2 \vec{B} \cdot d\vec{l} + \int_3 \vec{B} \cdot d\vec{l} + \right\} + \int_4$$

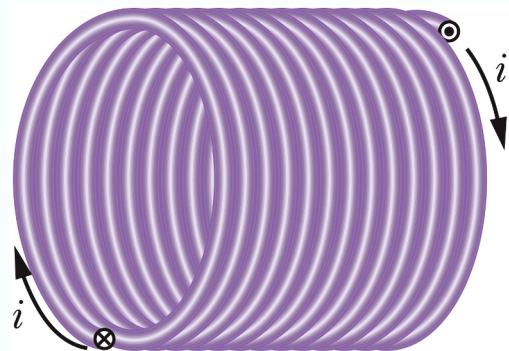


$$\rightarrow 2Bl = \mu_0 I_{\text{enc}} = \mu_0 JA_{\text{in}} = \mu_0 J 2z\ell$$

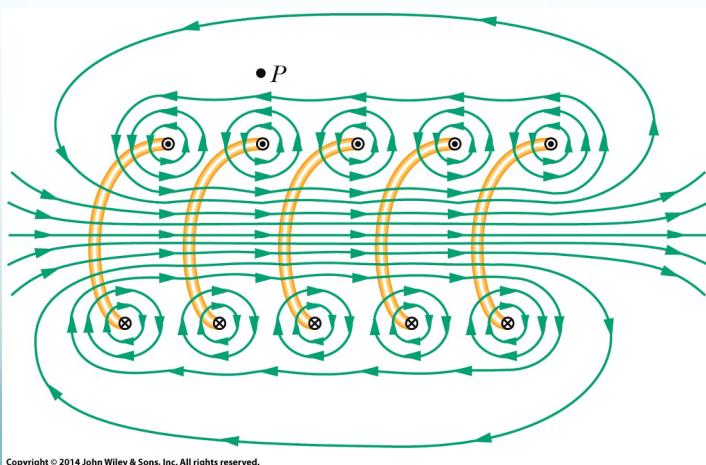
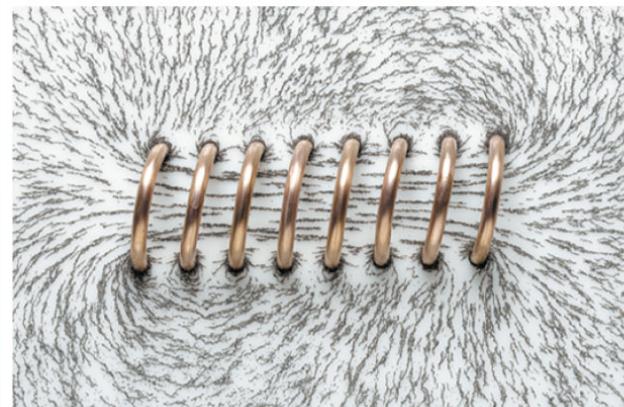
$$\rightarrow B = \mu_0 J z \quad \& \quad \text{if } z = \frac{d}{2} \rightarrow B = \frac{\mu_0 J d}{2}$$

# Ampère's law: application(2)

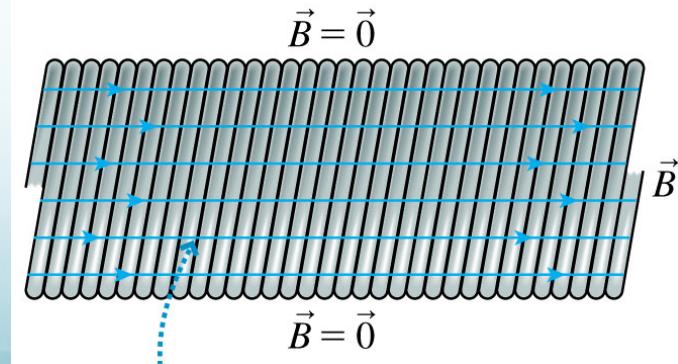
## 29.3: Solenoids and Toroids



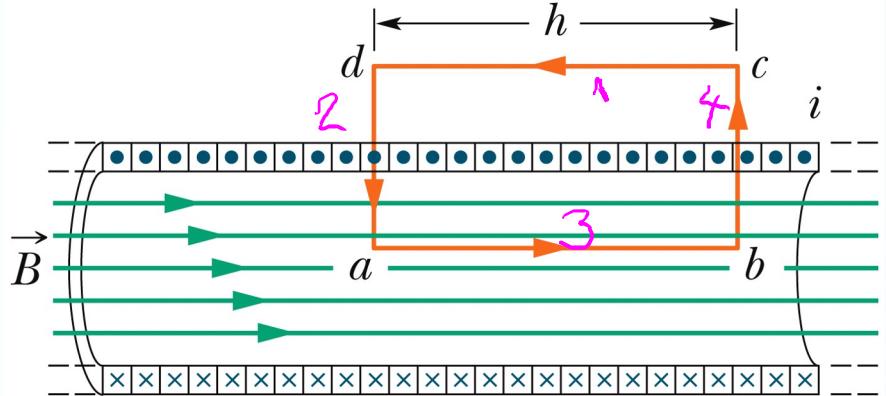
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$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i_{enc}$$

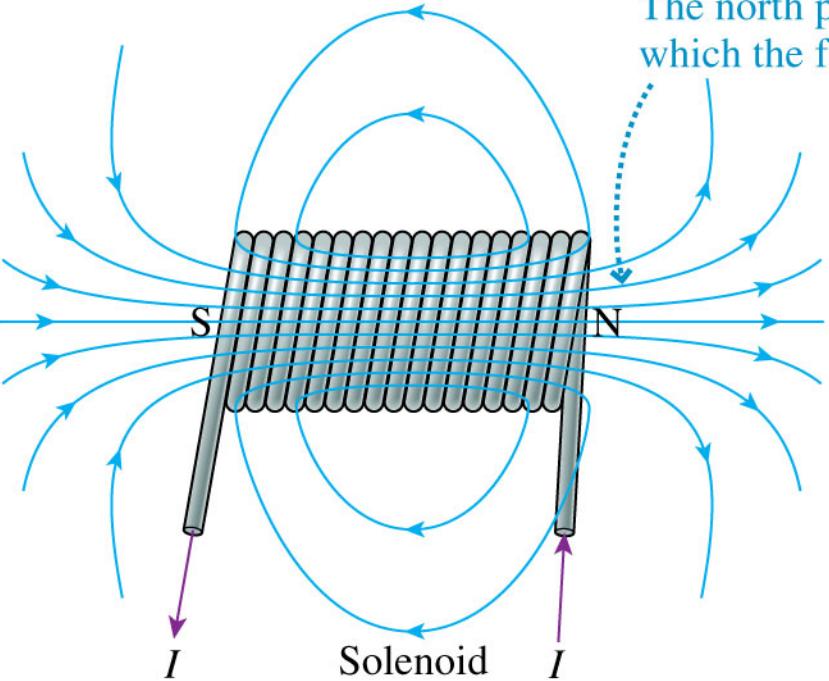


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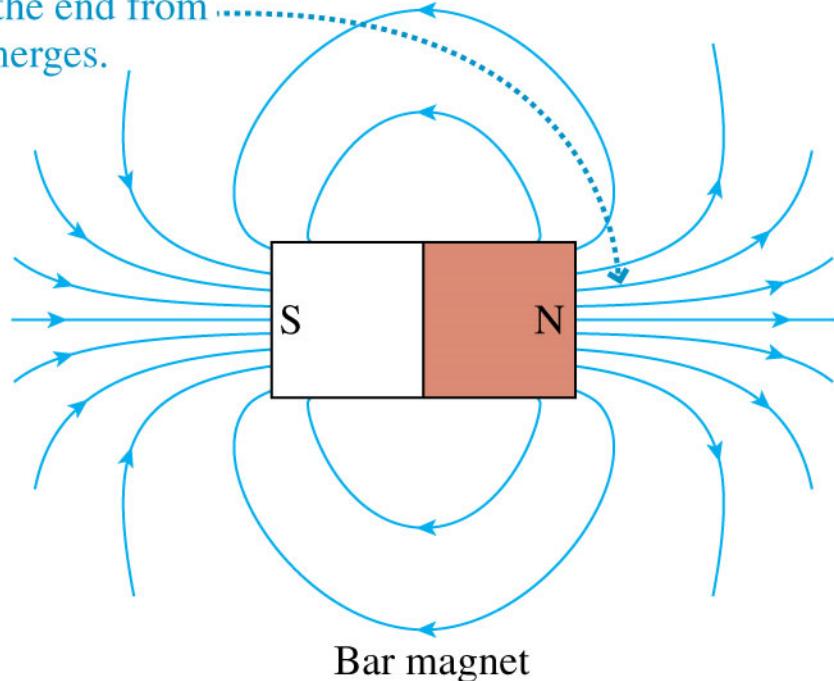
$$n = \frac{N}{L} \text{ number of turns per unit length}$$

$$i_{enc} = n i h$$

$$B_{Solenoid} = \mu_0 n i$$



The north pole is the end from which the field emerges.

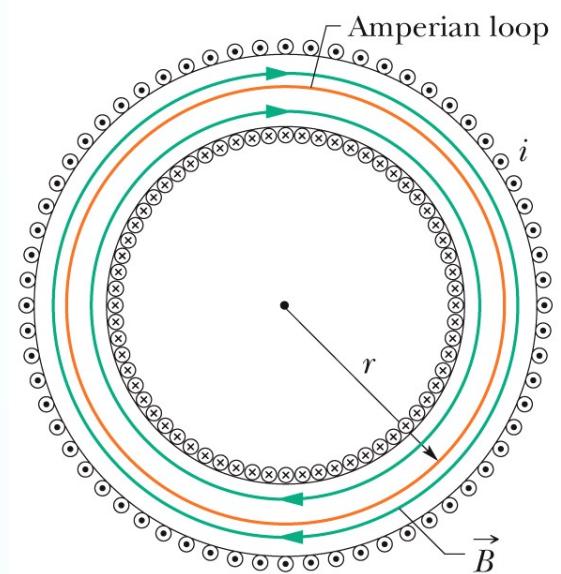


$$B_{\text{Solenoid}} = \mu_0 n i$$

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(a)



(b)

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This section we talked about:

Chapter 29

*See you on Friday*

