

Last time

- Review of electrostatic equations
- Obtaining electric field from electric potential
- Equipotential surface

This time

- Electric potential due to a dipole
- Electric field from electric potential for a dipole
- Electric potential of a solid spherical conductor

Moving in the same direction of electrical field

$$dV = -\vec{E} \cdot d\vec{l} = -Edl \cos 0 = -Edl < 0$$

Moving in the opposite direction of electrical field

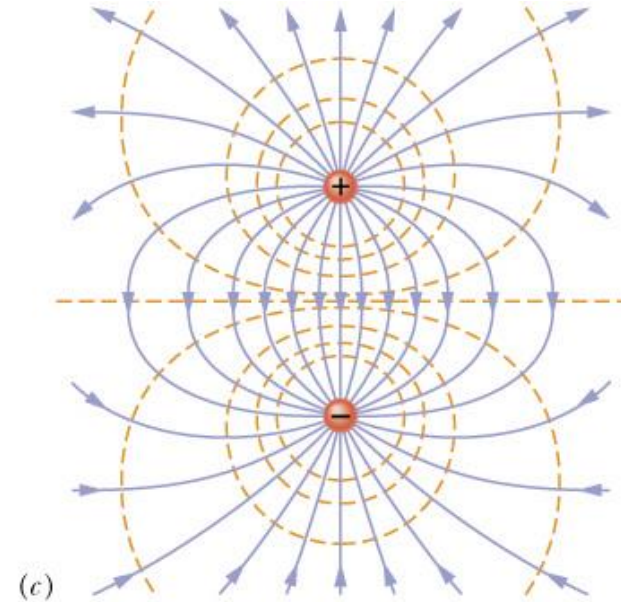
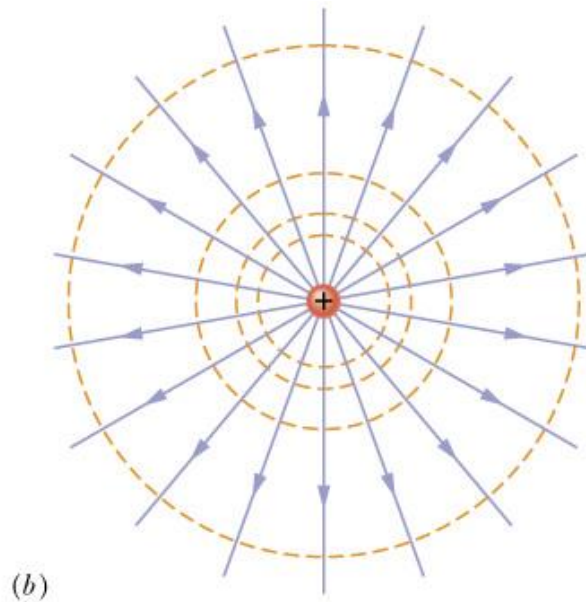
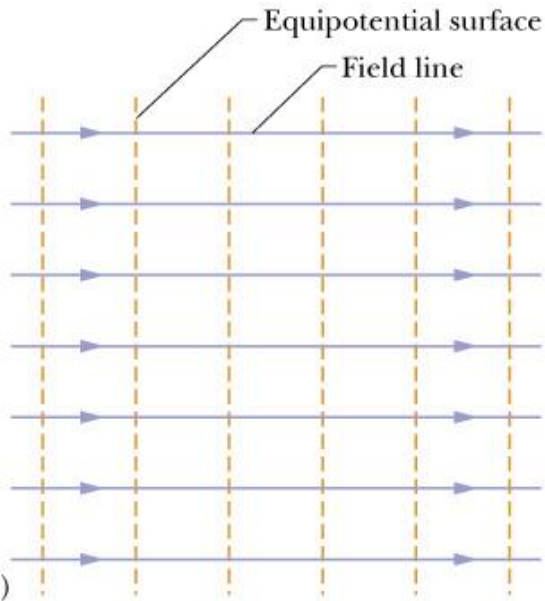
$$dV = -\vec{E} \cdot d\vec{l} = -Edl \cos 180 = Edl > 0$$

Moving perpendicular to the direction of electrical field

$$dV = -\vec{E} \cdot d\vec{l} = -Edl \cos 90 = 0$$

Moving with the electrical field decreases the electrical potential. Moving against the field increases it. Moving perpendicular to it does not change the electric potential.

$$V_f - V_i = -\int_i^f \vec{E} \cdot d\vec{l} = -\int_i^f E dl \cos 90 = 0$$



Electric potential doesn't change. A surface with this property is called an **equipotential surface**.

Electric field and electric potential

$$V_f - V_i = -\int_i^f \vec{E} \cdot d\vec{l}$$

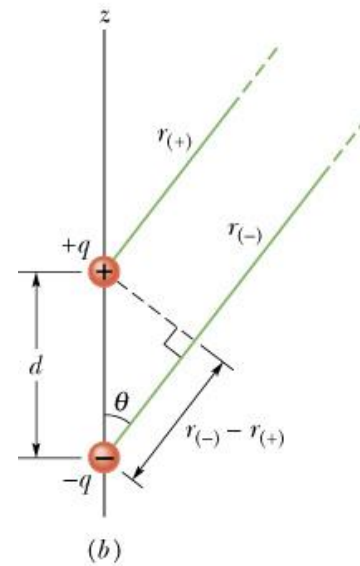
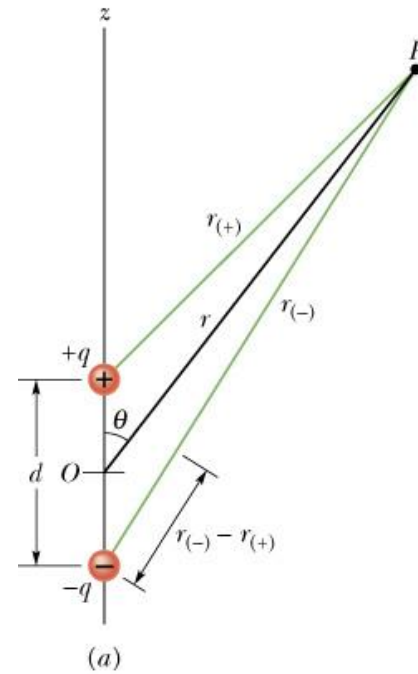
$$\frac{\partial V}{\partial x} = -E_x$$

$$\frac{\partial V}{\partial y} = -E_y$$

$$\frac{\partial V}{\partial z} = -E_z$$

$$\vec{E} = -\vec{\nabla} V = -\frac{\partial V}{\partial x} \hat{i} - \frac{\partial V}{\partial y} \hat{j} - \frac{\partial V}{\partial z} \hat{k}$$

Potential due to an electric dipole

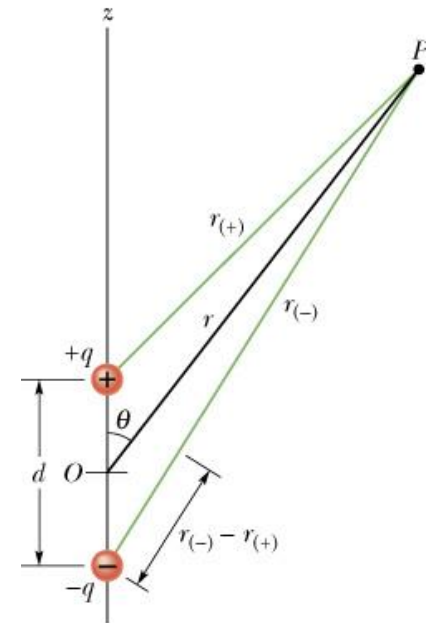


Remember V is a scalar

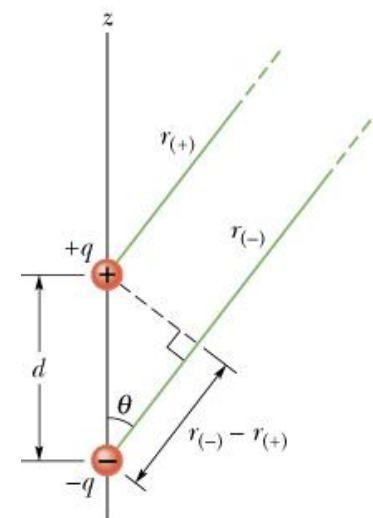
$$V(r) = V_+(r) + V_-(r)$$

$$V(r) = \frac{q}{4\pi\epsilon_0 r_+} - \frac{q}{4\pi\epsilon_0 r_-}$$

$$V(r) = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r_+} - \frac{1}{r_-} \right) = \frac{q}{4\pi\epsilon_0} \left(\frac{r_- - r_+}{r_+ r_-} \right)$$



(a)



(b)

$$V(r) = \frac{q}{4\pi\epsilon_0} \left(\frac{r_- - r_+}{r_+ r_-} \right)$$

$$\text{For } r \gg d \Rightarrow r_+ \approx r_- \approx r \Rightarrow r_+ r_- \approx r^2$$

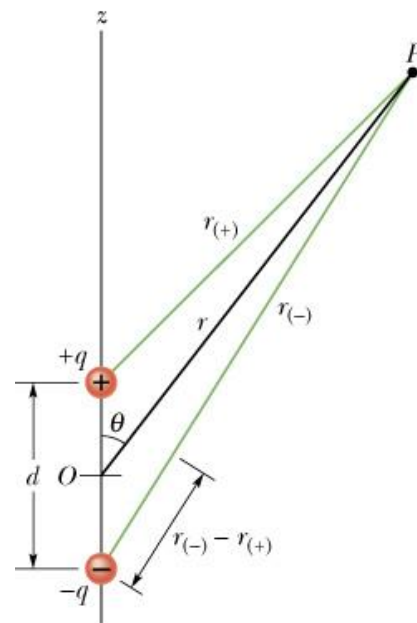
$$r_- - r_+ \approx d \cos \theta$$

$$V(r) = \frac{qd \cos \theta}{4\pi\epsilon_0 r^2}$$

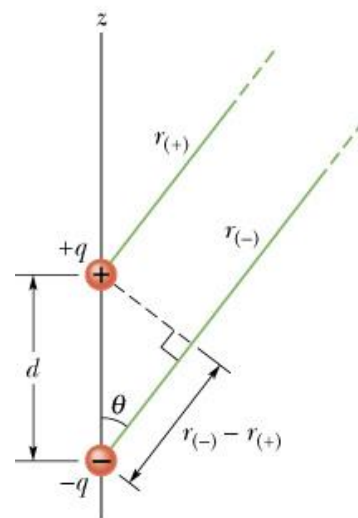
$$\cos \theta = \hat{k} \cdot \hat{r}$$

$$\vec{p} = qd\hat{k}$$

$$V(r) = \frac{\vec{p} \cdot \hat{r}}{4\pi\epsilon_0 r^2}$$



(a)



(b)

Now we can calculate the electric field for an arbitrary point P at a distance r from the center of the dipole.

$$\begin{aligned} V(r) &= \frac{p \cos \theta}{4\pi\epsilon_0 r^2} \\ &= \frac{pz/r}{4\pi\epsilon_0 r^2} \\ &= \frac{pz}{4\pi\epsilon_0 r^3} \\ &= \frac{pz}{4\pi\epsilon_0 (x^2 + y^2 + z^2)^{3/2}} \\ &= \frac{pz}{4\pi\epsilon_0} (x^2 + y^2 + z^2)^{-3/2} \end{aligned}$$

Note that $z = r \cos \theta$.

$$\begin{aligned} E_x &= -\frac{\partial V}{\partial x} = -\frac{\partial}{\partial x} \left[\frac{pz}{4\pi\epsilon_0} (x^2 + y^2 + z^2)^{-3/2} \right] \\ &= -\frac{pz}{4\pi\epsilon_0} \left(-\frac{3}{2} 2x \right) (x^2 + y^2 + z^2)^{-5/2} \\ &= \frac{3pxz}{4\pi\epsilon_0 r^5} \end{aligned}$$

Similarly

$$\begin{aligned}E_y &= -\frac{\partial V}{\partial y} = -\frac{\partial}{\partial y} \left[\frac{pz}{4\pi\epsilon_0} (x^2 + y^2 + z^2)^{-3/2} \right] \\&= -\frac{pz}{4\pi\epsilon_0} \left(-\frac{3}{2} 2y \right) (x^2 + y^2 + z^2)^{-5/2} \\&= \frac{3pzy}{4\pi\epsilon_0 r^5}\end{aligned}$$

and

$$\begin{aligned}E_z &= -\frac{\partial V}{\partial z} = -\frac{\partial}{\partial z} \left[\frac{pz}{4\pi\epsilon_0} (x^2 + y^2 + z^2)^{-3/2} \right] \\&= -\frac{p}{4\pi\epsilon_0} \left[\frac{1}{r^3} + z \left(-\frac{3}{2} 2z \right) (x^2 + y^2 + z^2)^{-5/2} \right] \\&= -\frac{p}{4\pi\epsilon_0} \left[\frac{1}{r^3} - \frac{3z^2}{r^5} \right] \\&= -\frac{p}{4\pi\epsilon_0 r^5} [r^2 - 3z^2] \\&= -\frac{p}{4\pi\epsilon_0 r^5} [x^2 + y^2 - 2z^2]\end{aligned}$$

For a point $P(0, 0, z)$, we have

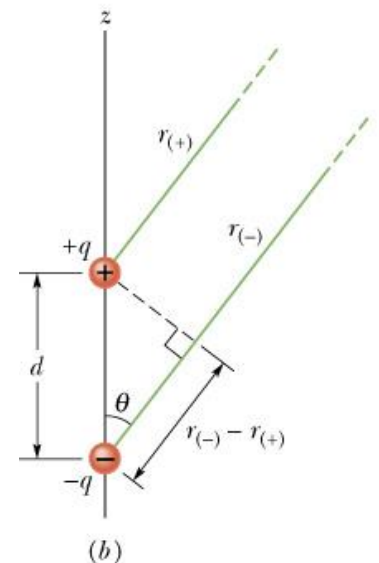
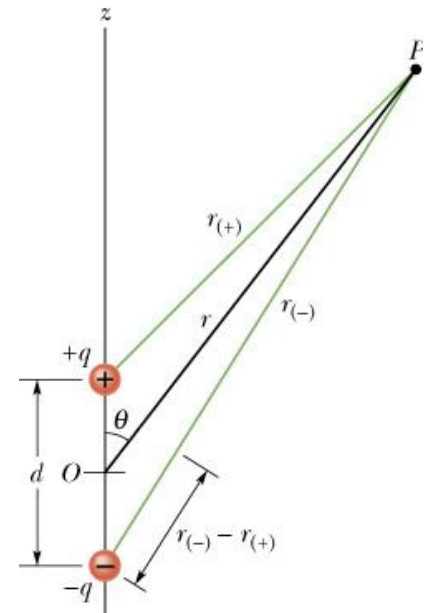
$$\begin{aligned} E_x &= \frac{3pzx}{4\pi\epsilon_0 r^5} \\ &= 0 \end{aligned}$$

$$\begin{aligned} E_y &= \frac{3pzy}{4\pi\epsilon_0 r^5} \\ &= 0 \end{aligned}$$

$$\begin{aligned} E_z &= -\frac{p}{4\pi\epsilon_0 r^5} [x^2 + y^2 - 2z^2] \\ &= \frac{p}{2\pi\epsilon_0 z^3} \end{aligned}$$

or

$$\begin{aligned} \vec{E} &= \frac{p\hat{z}}{2\pi\epsilon_0 z^3} \\ &= \frac{\vec{p}}{2\pi\epsilon_0 z^3} \end{aligned}$$



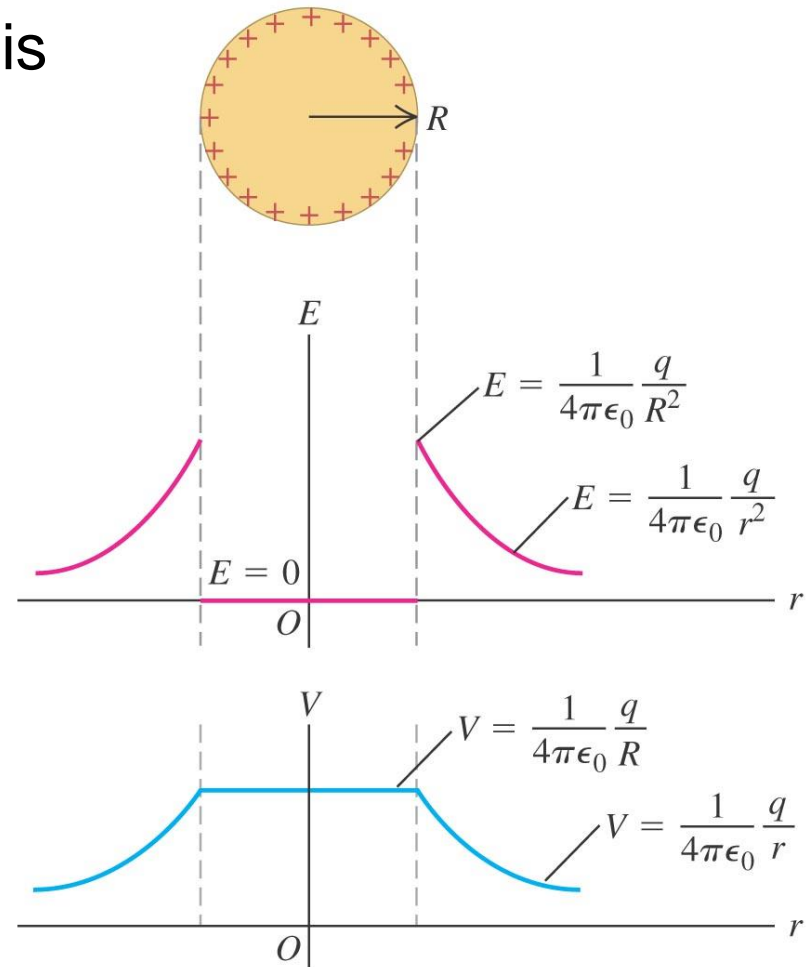
Calculation of electrical potential

The electric field everywhere inside the conductor and on the surface is zero.

$$\vec{E} = 0 \quad \text{For } r \leq R$$

$$V(r \leq R) = -\int \vec{E} \cdot d\vec{l} = C_1$$

We therefore conclude the all points inside and on the surface of the solid conductor are at the same potential, **the conductor is an equipotential object.**



The electric field outside the conductor is given by

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} \quad \text{For } r > R$$

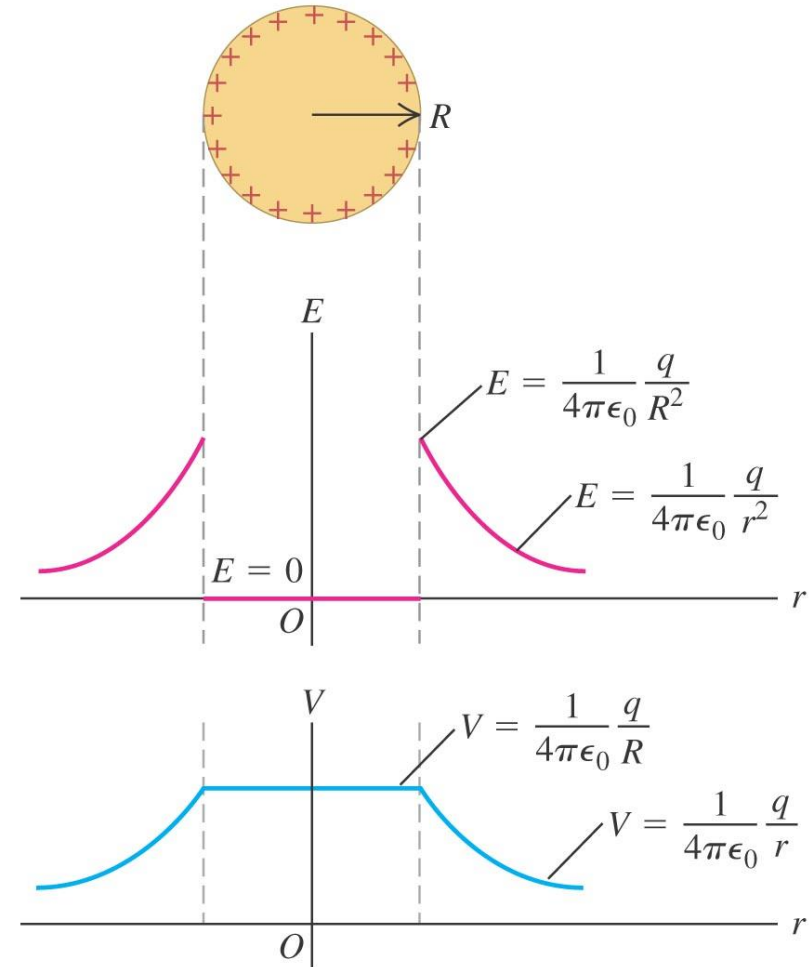
$$V(r > R) = -\frac{q}{4\pi\epsilon_0} \int \frac{\hat{r}}{r^2} \cdot d\vec{l}$$

Because the electrostatic field is a conservative field, we can choose

$$d\vec{l} = dr\hat{r}$$

$$V(r > R) = -\frac{q}{4\pi\epsilon_0} \int \frac{\hat{r}}{r^2} \cdot (dr\hat{r})$$

$$V(r > R) = -\frac{q}{4\pi\epsilon_0} \int \frac{dr}{r^2} = \frac{q}{4\pi\epsilon_0} \frac{1}{r} + C_2$$



Boundary condition I:

$$V(r \rightarrow \infty) = 0$$

$$0 = \frac{q}{4\pi\epsilon_0} \frac{1}{r} \Big|_{r=\infty} + C_2 \Rightarrow C_2 = 0$$

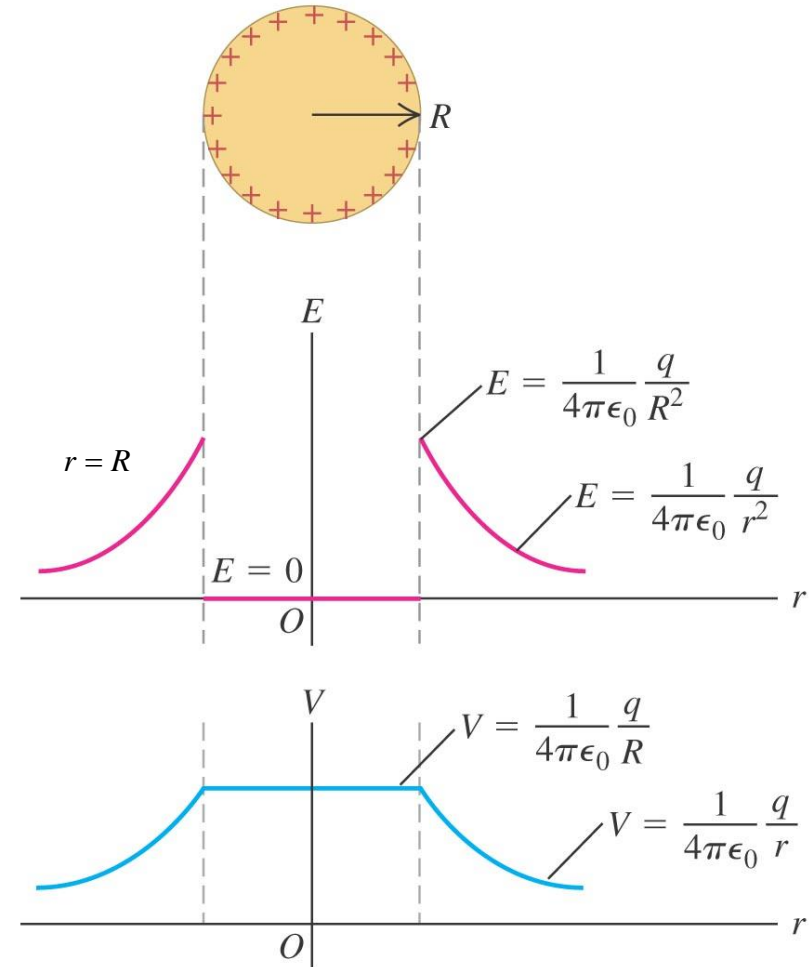
$$V(r > R) = \frac{q}{4\pi\epsilon_0} \frac{1}{r}$$

Boundary condition II:

The electric potential must be continuous at $r = R$. **Why?**

$$V(r < R) \Big|_{r=R} = V(r > R) \Big|_{r=R}$$

$$C_1 = \frac{q}{4\pi\epsilon_0} \frac{1}{R}$$



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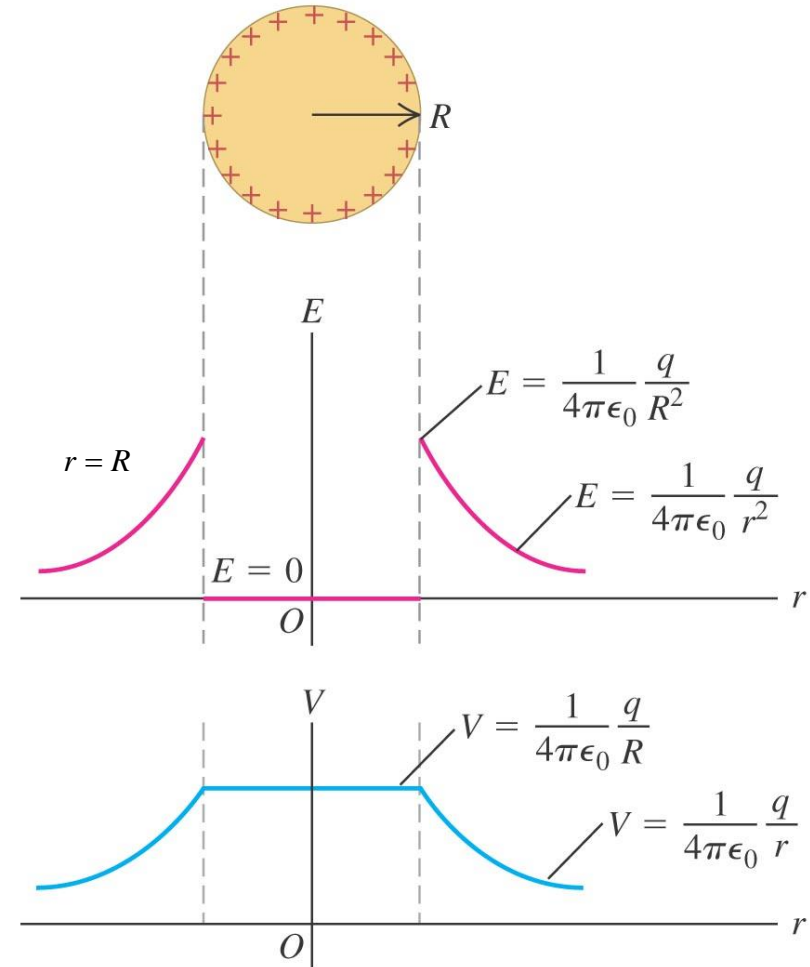
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The electric potential must be continuous at $r = R$. **Why?**

$$V(r < R) \Big|_{r=R} = V(r > R) \Big|_{r=R}$$

$$C_1 = \frac{q}{4\pi\epsilon_0} \frac{1}{R}$$



Calculation of electrical potential

$$V(r) = \frac{q}{4\pi\epsilon_0} \frac{1}{R} \quad \text{For } r \leq R$$

$$V(r) = \frac{q}{4\pi\epsilon_0} \frac{1}{r} \quad \text{For } r \geq R$$

