Electricity and Magnetism

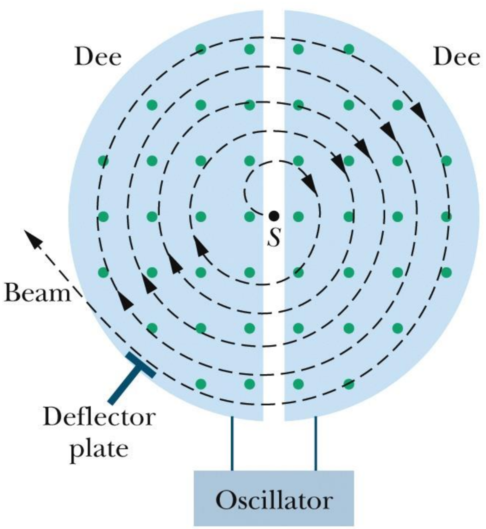
- Physics 259 – L02
- Lecture 39

Chapter 28: Magnetic fields

Last Section:

- ✓ The Cyclotron
- ✓ The proton Synchrotron

The protons spiral outward in a cyclotron, picking up energy in the gap.



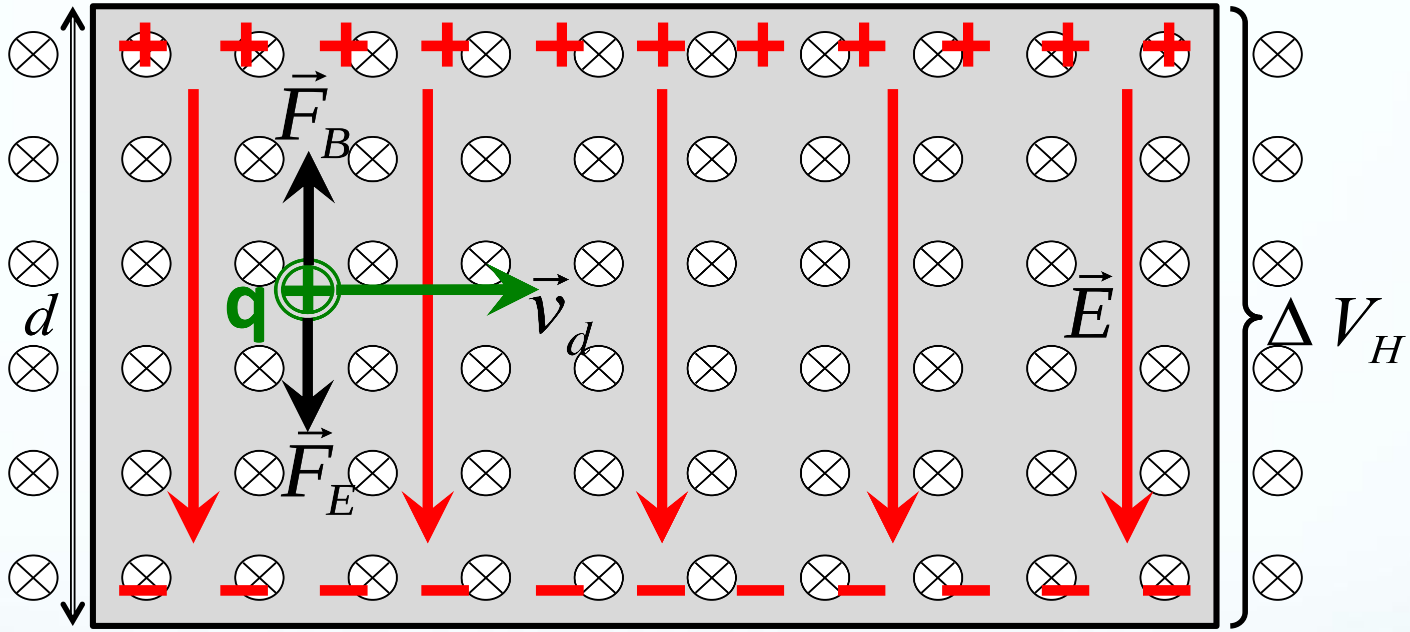
$$\boxed{R = \frac{mv}{|q|B}}$$

$$\boxed{T_{cyc} = \frac{2\pi m}{|q|B}}$$

$$\boxed{f_{cyc} = \frac{|q|B}{2\pi m}}$$

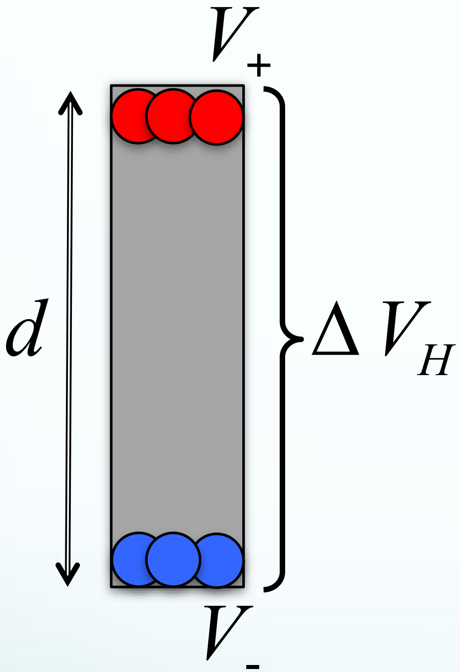
Last Section:

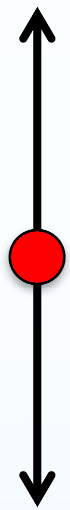
Hall effect



$$\boxed{\Delta V_H = v_d B d}$$

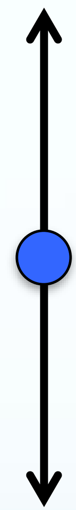
Conductors moving in B-fields





$F_B = qvB$

$F_E = q\dfrac{\Delta V_H}{d}$



$F_E = q\dfrac{\Delta V_H}{d}$

$F_B = qvB$

In equilibrium, forces balance, leading to a constant voltage

$$q\frac{\Delta V_H}{d} = qvB$$

$$\boxed{\Delta V_H = vBd}$$

28-6 Magnetic Force on a Current-Carrying Wire

Conductors in B-fields

Free charges moving in a B-field feel→

$$\vec{F}_B = q\vec{v} \times \vec{B}$$

Conductors are full of charges that are free to move around→

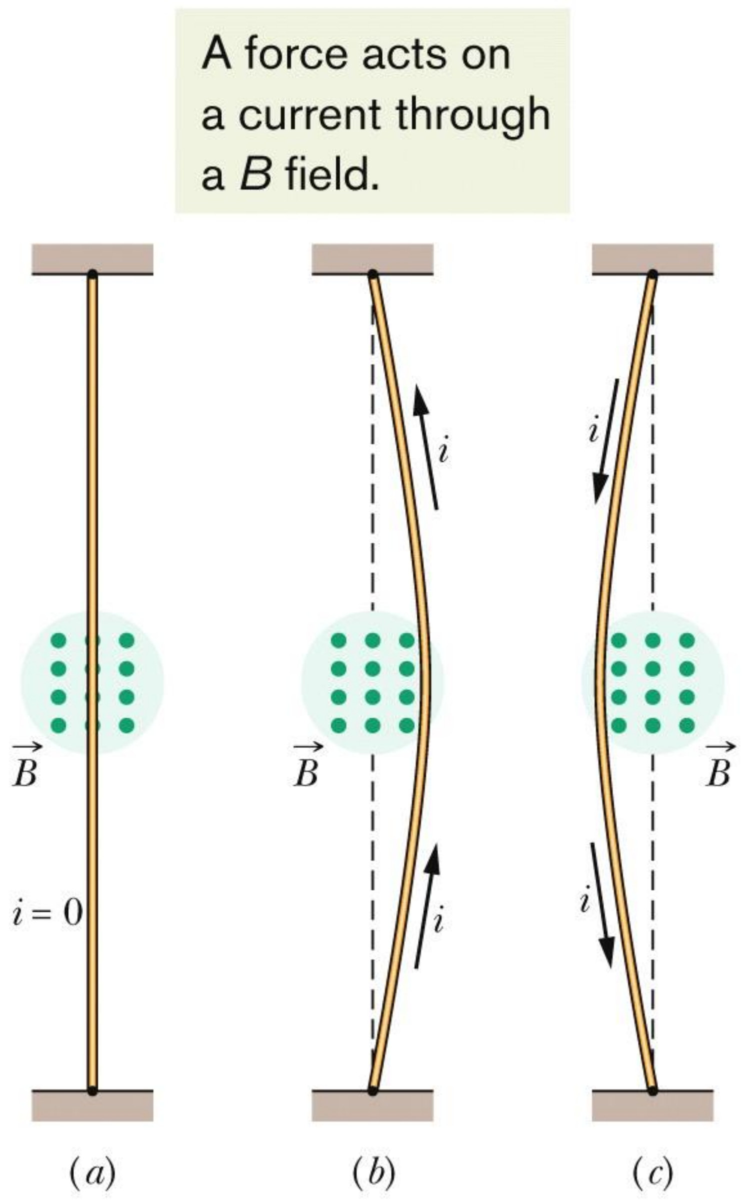
If a conductor moves in a magnetic field, these charges also feel a magnetic force

28-6 Magnetic Force on a Current-Carrying Wire

A straight wire carrying a current *i* in a uniform magnetic field experiences a sideways force

$$\vec{F}_B = i\vec{L} \times \vec{B} \quad \text{(force on a current).}$$

Here ***L*** is a length vector that has magnitude *L* and is directed along the wire segment in the direction of the (conventional) current.



Proof

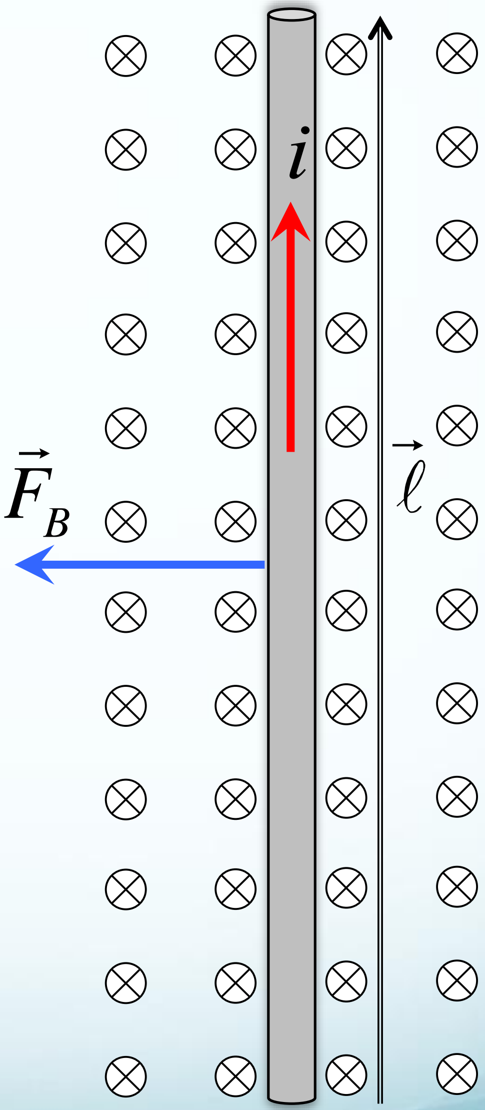
Current in wires is nothing more than charges in motion.

It doesn't matter if we consider $-q$ moving opposite i or $+q$ moving in the same direction as i

For a single charge→

$$\vec{F}_B = q\,\vec{v}_d \times \vec{B}$$

For N charges moving through the wire:



$$\boxed{\vec{F}_B = i\vec{\ell} \times \vec{B}}$$

Crooked Wire

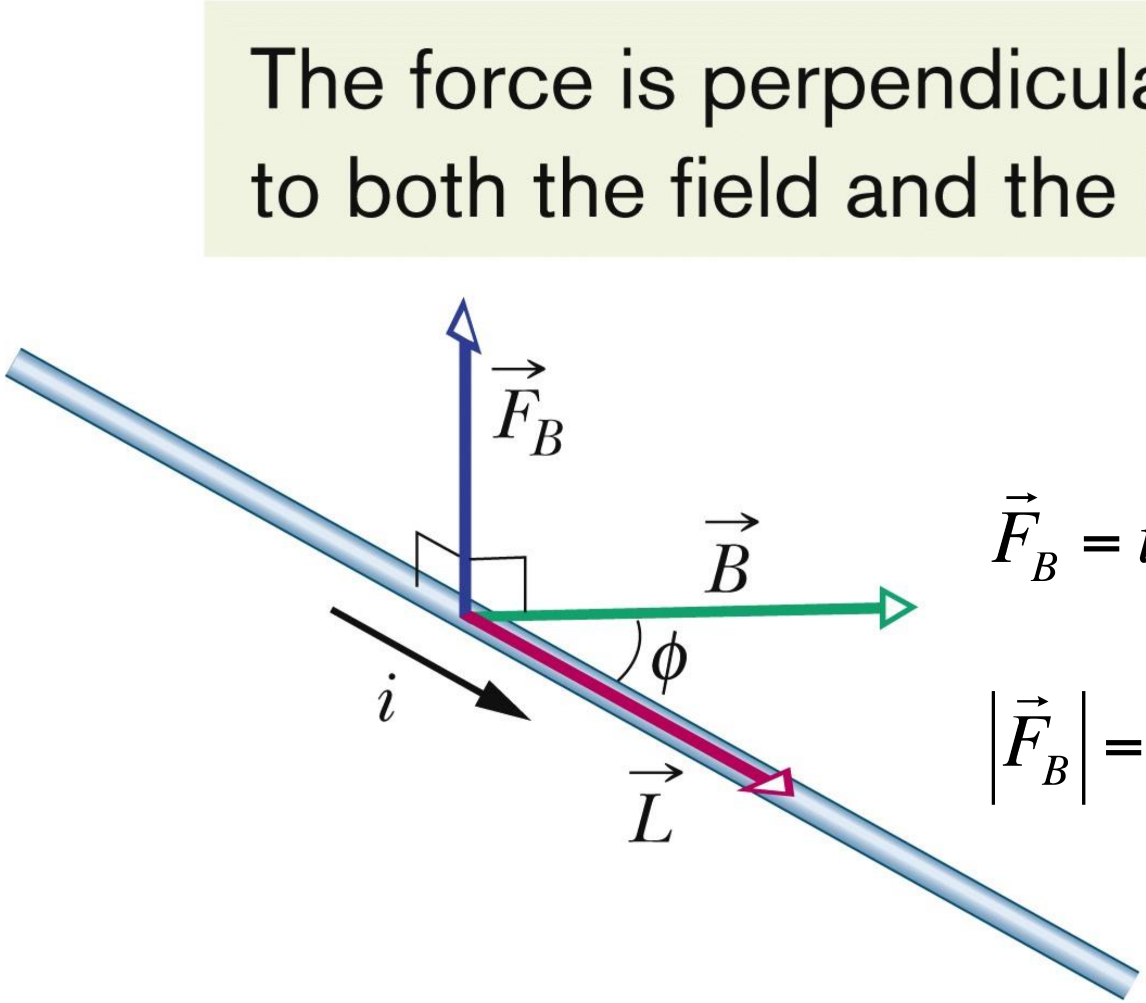
If a wire **is not straight** or the **field is not uniform**→
we can imagine the wire broken up into small straight segments.

The force on the wire as a whole is then the vector sum of all the forces on the segments that make it up. In the differential limit, we can write

$$\boxed{d\vec{F}_B = i\, d\vec{L} \times \vec{B}.}$$

and the direction of length vector ***L*** or *d**L*** is in the direction of *i.*

Forces on Current-Carrying Wires: B and L not perpendicular



$$\vec{F}_B = i\vec{L} \times \vec{B}$$

$$\left|\vec{F}_B\right| = \left|i\vec{L} \times \vec{B}\right| = iLB\sin\phi$$

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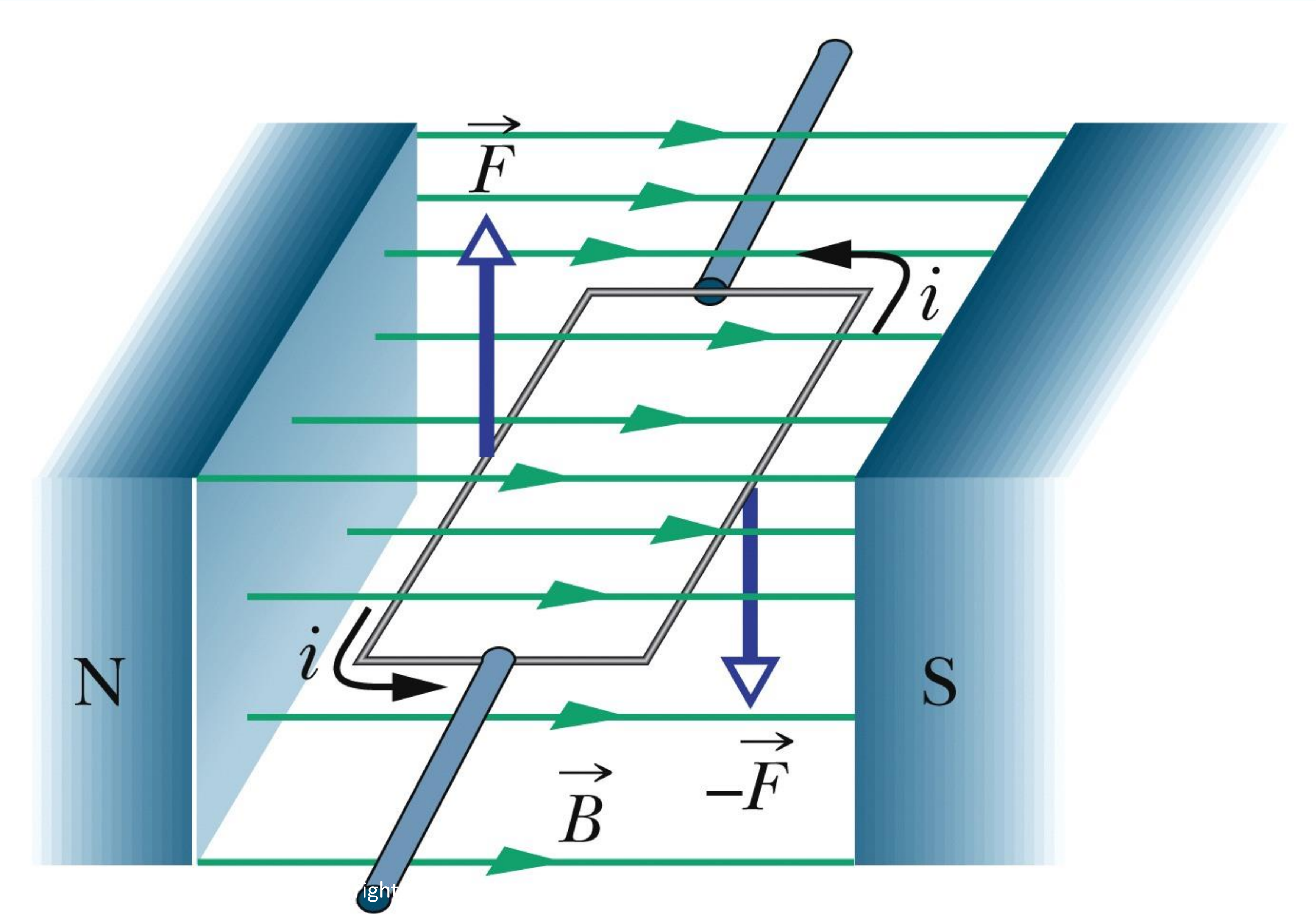
28-7 Torque on a Current Loop



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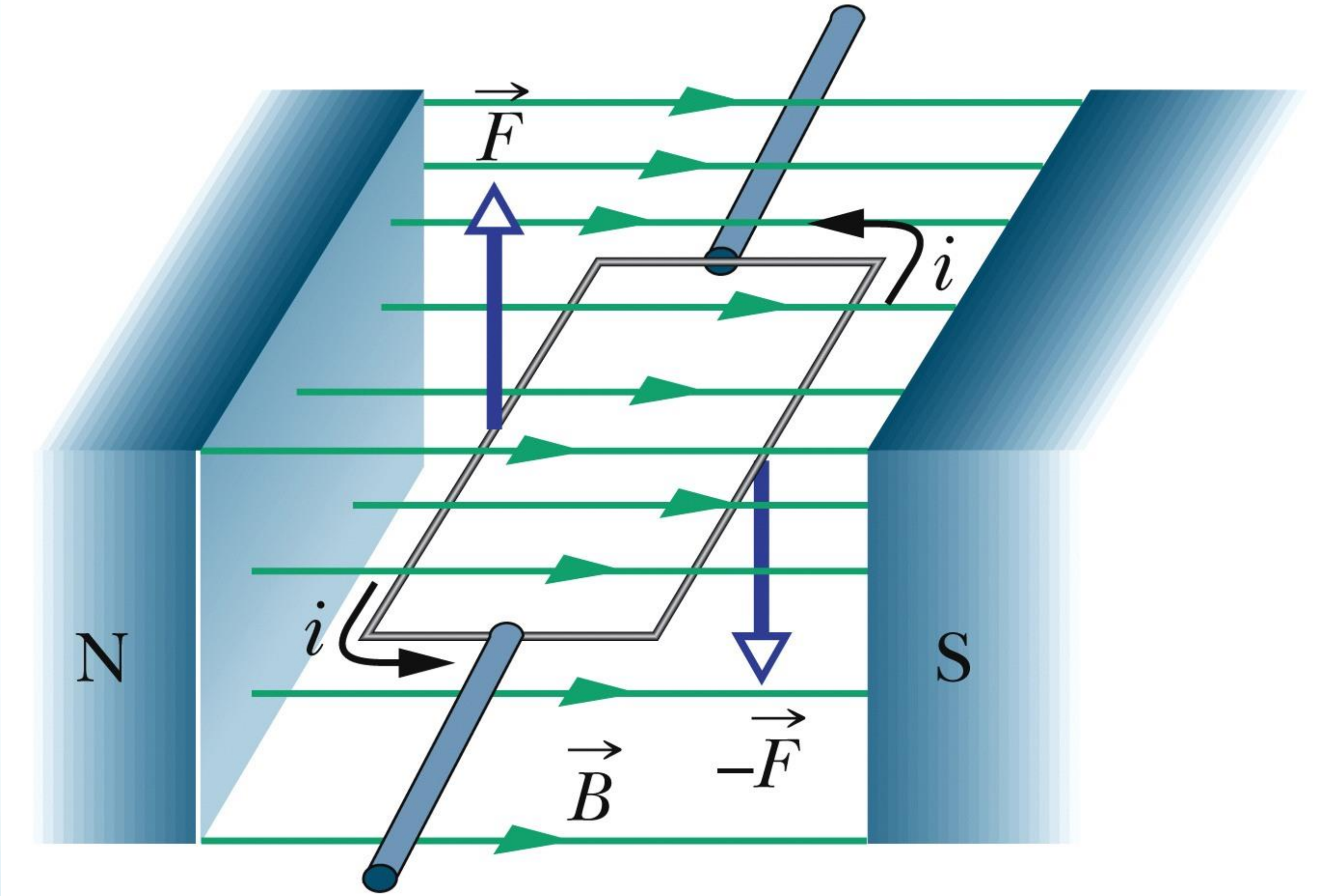
$$\vec{F}_B = i\vec{L} \times \vec{B} \quad \text{(force on a current).}$$



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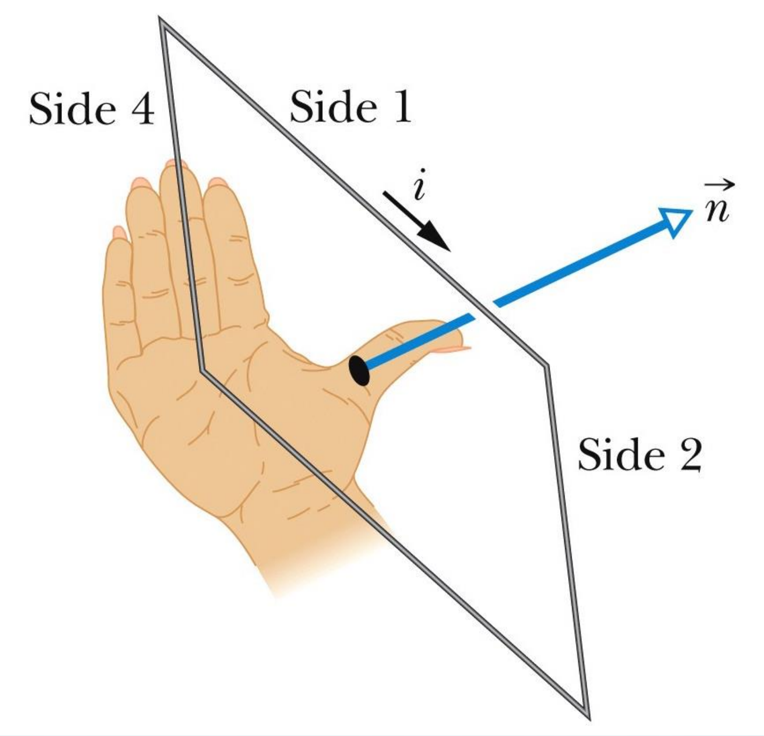
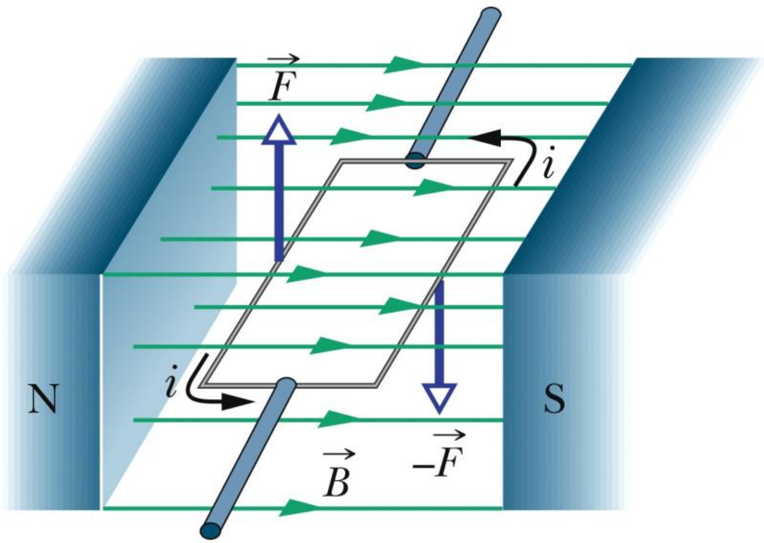
Pick the normal vector to the loop area by RHR: curl your fingers in the direction of i, thumb points in direction of n





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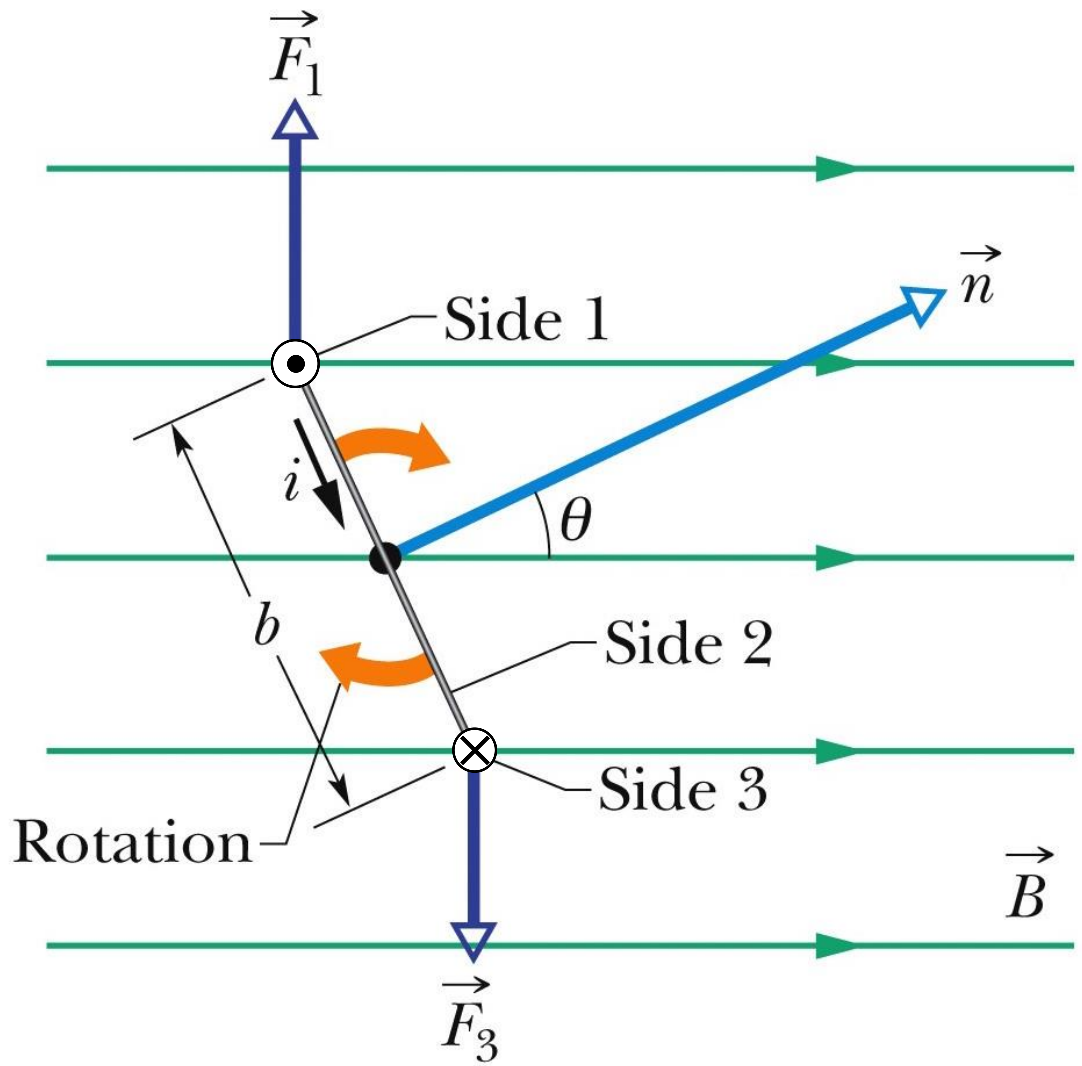
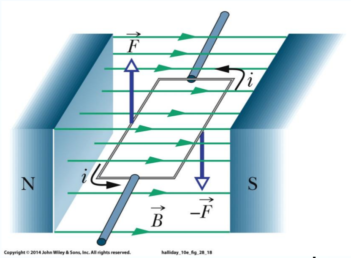
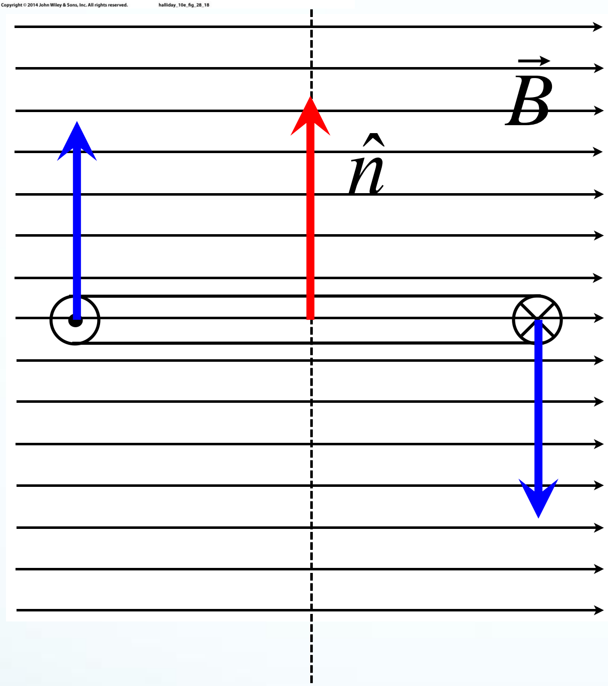
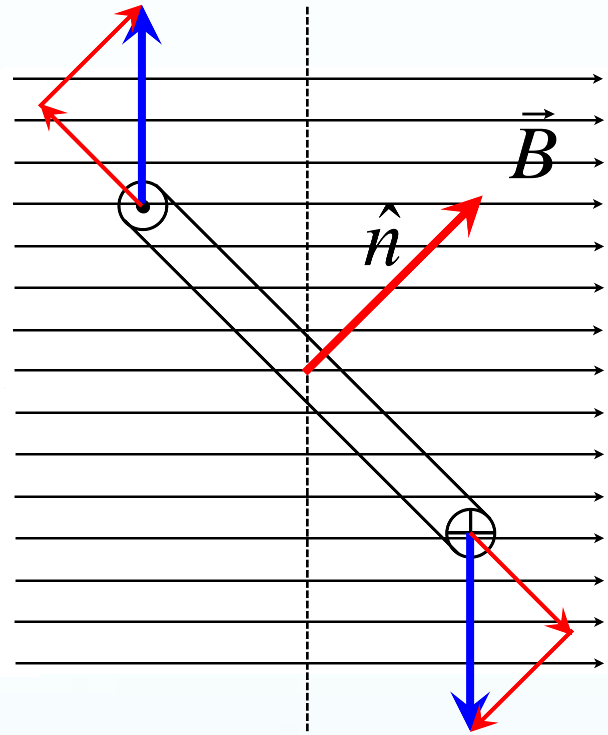


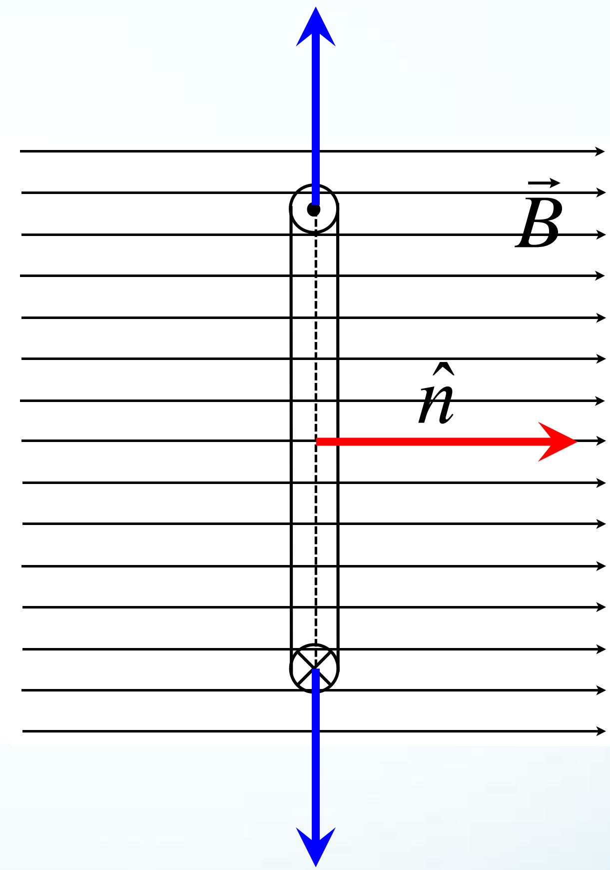




The normal vector is at right angles to the B-field: all magnetic force causes rotation of the loop



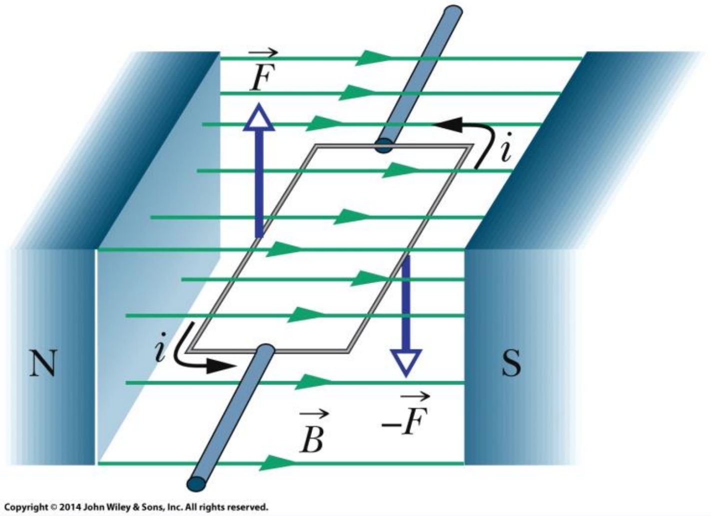
The normal vector is at some angle to the B-field: some of the magnetic force causes rotation of the loop



The normal vector is parallel to the B-field: none of the magnetic force causes rotation of the loop

Conclusion: components of magnetic force (anti)parallel to normal vector cause torque

Conclusion: Torque on a Current Loop



Net force on the loop is the vector sum of the forces acting on its four sides and comes out to be zero. The net torque acting on the coil has a magnitude given by

$$\tau = NiAB \sin \theta,$$

where N is the number of turns in the coil, A is the area of each turn, i is the current, B is the field magnitude, and θ is the angle between the magnetic field $\boldsymbol{B}$ and the normal vector to the coil $\boldsymbol{n}$.

This section we talked about:

Chapter 28

See you on Friday