

Laboratoire de Chimie Théorique

# One dimensional model for Relativistic Quantum Chemistry

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Thesis work with Julien Toulouse

**1. Introduction: Why going into relativistic?**

**2. Relativistic Quantum Chemistry**

**3. The 1D Model**

**4. QED Effects**

**5. Basis set convergence**

**6. Conclusion**

# **1. Introduction: Why going into relativistic?**

## 2. Relativistic Quantum Chemistry

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## 6. Conclusion

**Why are we interested in relativistic quantum chemistry?**

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- Important for fast-moving particles: average electron velocity in hydrogen-like ground-state

$$v = (\alpha Z)c, \text{ with } \alpha \approx 1/137$$

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- Explains some physical properties of heavy elements

### 1 yellow color of gold

P Pykkö, ACIE, (2004)

P Romaniello and P L. de Boeij, JCP (2005)

### 2 lead-acid battery electro-chemical potential

R. Ahuja, A. Blomqvist, P Larsson, P Pykkö and P Zaleski-Ejgierd, PRL, (2011)

### 3 liquid state of mercury at room temperature

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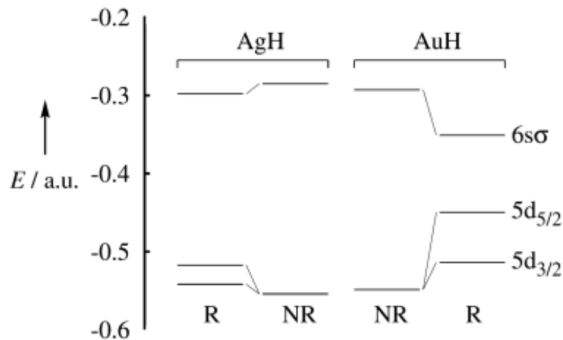
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- Includes the electron spin and spin-orbit coupling

- Fast-moving electrons in *s* and *p* orbitals
- Mass increases due to relativity
- Contraction of those orbitals
- Stronger screening of the nucleus
- Destabilize *d* and *f* orbitals
- Change the gap energy between 5*d* and 6*s*, from UV to yellow



**Figure 1:** Orbitals diagram of Gold and Silver, with and without relativistic effects

<sup>1</sup> P. Pyykkö, *Angew. Chem. Int. Ed.* 43, 4412 (2004)



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## Free particle energy

- Non relativistic:

$$E = \frac{p^2}{2m} \Rightarrow E \in [0, +\infty),$$

- Relativistic:

$$E^2 = m^2 c^4 + c^2 p^2 \Rightarrow E \in (-\infty, -mc^2] \cup [mc^2, +\infty),$$

## Free particle energy

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## Link between both

- Taylor expansion of the square root for the positive part

$$E = \underbrace{mc^2}_{\text{Rest mass}} + \underbrace{\frac{p^2}{2m}}_{\text{Kinetic energy}} - \underbrace{\frac{p^4}{8m^3c^2}}_{\text{1st order}} + \dots$$

## From classical to quantum - Klein-Gordon equation

$$\begin{aligned} E^2 &= p^2 c^2 + m^2 c^4 \\ -\hbar^2 \partial_t^2 &= -\hbar^2 c^2 \nabla^2 + m^2 c^4 \end{aligned}$$

- Negative energy solutions give negative probability...
- Does not include spin

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## From classical to quantum - Dirac equation

$$\begin{aligned} E &= \pm c \sqrt{p^2 + m^2 c^2} \\ i\hbar \partial_t &= c \sqrt{p^2 + m^2 c^2} \end{aligned}$$

- Expand the term under the square root as a perfect square:

$$\begin{aligned} p^2 + m^2 c^2 &= (\vec{\alpha} \cdot \vec{p} + \beta mc)^2 \\ i\hbar \partial_t &= c \vec{\alpha} \cdot \vec{p} + \beta mc^2 \end{aligned}$$

## Dirac operator

$$D_0(x) = c (\vec{\alpha} \cdot \vec{p}) + \beta mc^2, \quad (1)$$

- $\vec{\alpha} = \begin{pmatrix} 0 & \vec{\sigma} \\ \vec{\sigma} & 0 \end{pmatrix}$  where  $\vec{\sigma} = \{\sigma_x, \sigma_y, \sigma_z\}$  are the Pauli matrices
- $\beta = \begin{pmatrix} \mathbb{I}_2 & 0 \\ 0 & -\mathbb{I}_2 \end{pmatrix}$

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## Dirac eigenvalue equation

$$(D_0(x) + V(x))\psi(x) = \mathcal{E}\psi(x) \quad (2)$$

- $\psi = \begin{pmatrix} \psi^L \\ \psi^S \end{pmatrix}$  where  $\psi^{L/S} = \begin{pmatrix} \psi_{\alpha}^{L/S} \\ \psi_{\beta}^{L/S} \end{pmatrix}$
- $\mathcal{E}$  is the energy of the state  $\psi$

## Non-relativistic spectrum

$$\mathcal{H}(x) = -\frac{\Delta}{2m} + V(x) \quad (3)$$

Spectrum:

- If  $\mathcal{E} > 0$ : continuum
- If  $\mathcal{E} < 0$ : bound states

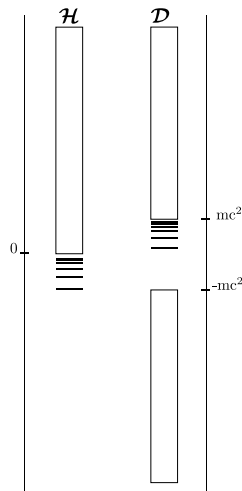
## Relativistic spectrum

$$D(x) = c(\vec{\alpha} \cdot \vec{p}) + \beta mc^2 + V(x) \quad (4)$$

Spectrum:

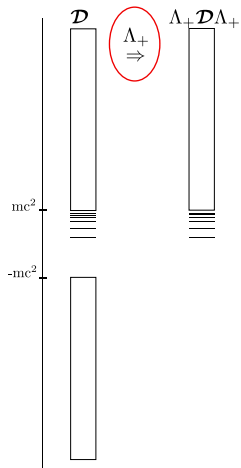
- If  $\mathcal{E} \in (-\infty, -mc^2] \cup [mc^2, +\infty)$ : continuum
- If  $\mathcal{E} \in (-mc^2, mc^2)$ : bound states

⇒ Negative continuum spectrum

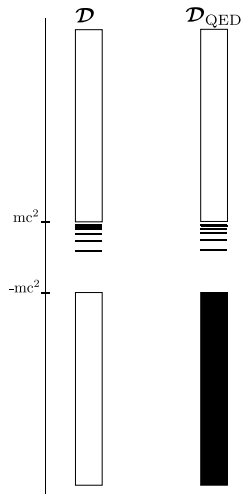




- No-pair approximation



- QED description



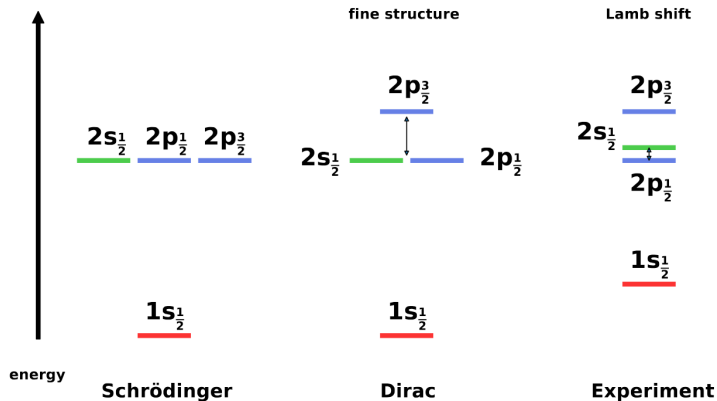


Figure 2: Lamb shift in hydrogen spectra<sup>2</sup>

<sup>2</sup>M. Salman, Ph.D. thesis (2022)

- Different levels of relativistic corrections
- **Zeroth level:** Non-relativistic chemistry

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<sup>3</sup> P Schwerdtfeger *et al.*, PRL, (2017)

- Different levels of relativistic corrections
- **Zeroth level:** Non-relativistic chemistry
- **First level:** Solving Dirac's equation or adding scalar-relativistic and spin-orbit correction to Schrödinger's equation (ZORA, DKHn, X2C)

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- Different levels of relativistic corrections
- **Zeroth level:** Non-relativistic chemistry
- **First level:** Solving Dirac's equation or adding scalar-relativistic and spin-orbit correction to Schrödinger's equation (ZORA, DKHn, X2C)
- **Second level:** Quantum electrodynamics (QED) effects (e.g. IP and EA of gold, Lamb shift in hydrogen)

	IP	Error	EA	Error
DC-HF	7.6892	-1.5363	0.6690	-1.6396
DC-CCSD	9.1164	-0.1092	2.1070	-0.2017
DC-CCSD(T)	9.2938	0.0683	2.3457	0.0371
DC-CCSDTQP	9.2701	0.0446	2.3278	0.0192
+Breit	9.2546	0.0290	2.3188	0.0102
+QED	9.2288	0.0032	2.3072	-0.0014
Experiment [31,32]	9.2256		2.3086	

**Figure 3:** IP and EA of gold with different levels of approximation<sup>3</sup>

<sup>3</sup> P Schwerdtfeger *et al.*, PRL, (2017)

- We need accurate QED calculations for high precision spectroscopy
- Also studying core properties need to go into QED details
- So far, state of the art QED for many-electrons systems is based on no-pair approximation
- **Next challenge in QC:** can we go beyond the NPA?

## Goal

Use a model system to understand the issue of developing new methods based on QED.

### To do list:

- ☐ Study the 1D hydrogen and compute the QED effects.
- ☐ Develop these quantities in a finite basis set.
- ☐ Use our insight to solve the 3D problem.

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## Free Dirac operator

$$\mathcal{D}_{0,x} = c \boldsymbol{\alpha}_x p_x + \boldsymbol{\beta} mc^2,$$

- From this 1D operator we can make an unitary transformation

$$\mathbf{U} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

- Define a new operator  $\mathcal{D}'_{0,x} = \mathbf{U} \mathcal{D}_{0,x} \mathbf{U}^{-1} = \begin{pmatrix} \mathbf{D}_0 & \mathbf{0}_2 \\ \mathbf{0}_2 & \mathbf{D}_0 \end{pmatrix}$

- Where

$$\mathbf{D}_0 = c \boldsymbol{\sigma}_x p_x + \boldsymbol{\sigma}_z mc^2$$

is a  $2 \times 2$  operator!

<sup>4</sup>T. Audinet, J. Toulouse, J. Chem. Phys. 158, 244108 (2023)



## Free 1D Dirac operator

$$\mathbf{D}_0(x) = -ic\sigma_x \frac{d}{dx} + \sigma_z mc^2 \quad (5)$$

Lapidus, AJP (1983) ; T. Audinet, J. Toulouse, JCP (2023)

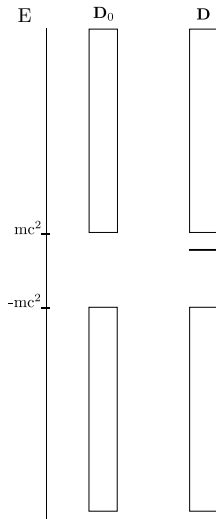
## Dirac eigenvalue equations

$$\mathbf{D}_0(x)\psi_p^0(x) = \varepsilon_p \psi_p^0(x) \quad (6)$$

- Continuum:  $\varepsilon_k = \pm \sqrt{m^2 c^4 + k^2 c^2}$

$$(\mathbf{D}_0(x) - Z\delta(x))\psi_p^Z(x) = \varepsilon_p \psi_p^Z(x) \quad (7)$$

- Bound State:  $\varepsilon_b < mc^2$
- Continuum:  $\varepsilon_k = \pm \sqrt{m^2 c^4 + k^2 c^2}$



## 1D Schrödinger equation

$$\hat{H} = -\frac{1}{2} \frac{d^2}{dx^2} - Z\delta(x) \quad (8)$$

- follows from the generalization of  $\hat{H}$  to arbitrary dimensions<sup>5</sup>
- recovers the energy, cusp and radial part of the 3D ground state of hydrogen<sup>6</sup>
- easily solvable model

## 1D Dirac equation

$$\hat{H} = -ic\sigma_x \frac{d}{dx} + \sigma_z mc^2 - Z\delta(x) \quad (9)$$

- gives analytical QED results
- recovers the non-relativistic hydrogen 1D properties

<sup>5</sup>D. R. Herschbach, J. Chem. Phys., 84 (1986)

<sup>6</sup>D. Traore, E. Giner, J. Toulouse, J. Chem. Phys., 156 (2022)

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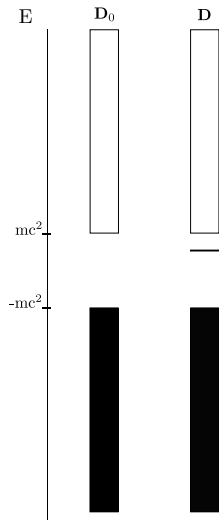
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- Spontaneous creation of electron positron pairs due to the external potential
- Creates a charge density that will interact through the two-electron interaction

## Vacuum Polarization Density

$$n^{\text{vp}}(x) = \sum_{\epsilon_p < 0} \psi_p^{Z\dagger}(x) \psi_p^Z(x) - \sum_{\epsilon_p < 0} \psi_p^{0\dagger}(x) \psi_p^0(x)$$

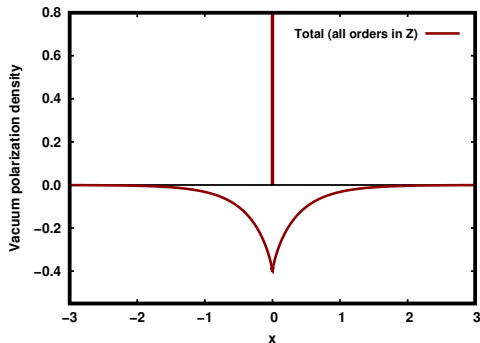
- In 3D this quantity diverges, needs to be renormalized
- We want to have a better understanding of this quantity and its influence on the energy spectrum



## Vacuum Polarization Density

$$n^{\text{vp}}(x) = \sum_{\varepsilon_p < 0} \psi_p^{Z\dagger}(x) \psi_p^Z(x) - \sum_{\varepsilon_p < 0} \psi_p^{0\dagger}(x) \psi_p^0(x) = \mathcal{N}_0^{\text{vp}} \delta(x) + n_{\text{reg}}^{\text{vp}}(x) \quad (10)$$

Audinet, Morellini, Levitt, Toulouse, in preparation

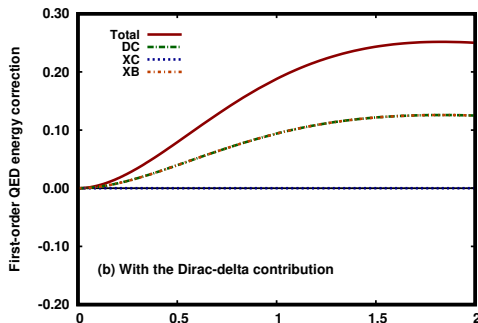


## Lamb shift

Correction to the 1s orbital energy of the hydrogen spectrum due to the interaction between the electrons and the vacuum polarization density

$$\mathcal{E}_{\text{Lamb}}^{\text{D}} = \int n_{1s}(x) n^{\text{vp}}(x) dx$$

$$\mathcal{E}_{\text{Lamb}}^{\text{X}} = - \int \text{Tr} [\mathbf{n}_{1s}(x) \mathbf{n}^{\text{vp}}(x)] dx$$



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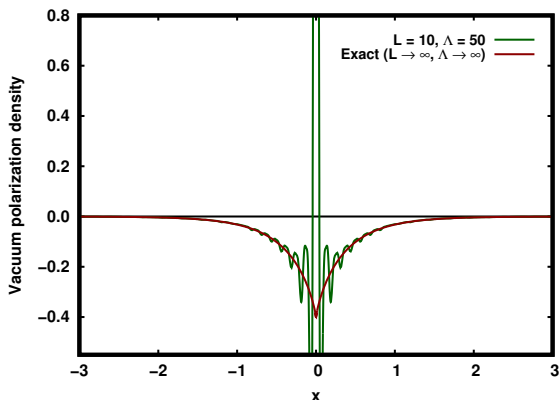
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## Plane waves basis

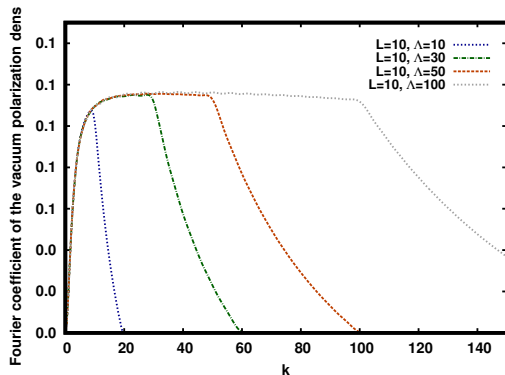
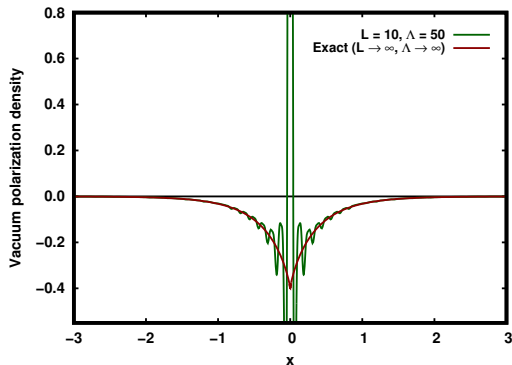
$$\forall x \in \mathbb{R}, k \in \frac{2\mathbb{Z}\pi}{L}, |k| \leq \Lambda, \chi_k(x) = \frac{1}{\sqrt{L}} e^{ikx} \quad (11)$$





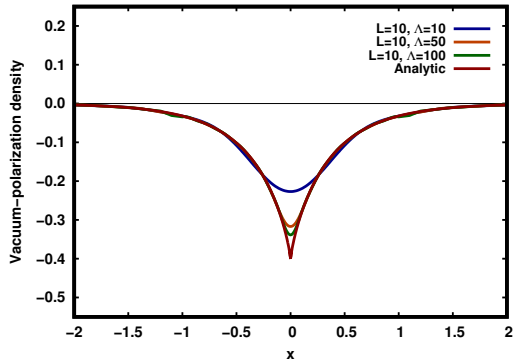
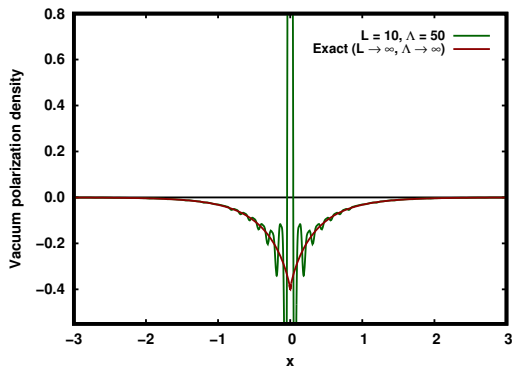
## Regularization with Fourier Transform

$$n_{\text{reg}}^{\text{vp}}(x) = \mathcal{F}^{-1} \left[ \mathcal{F}[n^{\text{vp}}](k) \theta(k_{\text{max}} - k) - \mathcal{F}[n^{\text{vp}}](k_{\text{max}}) \right] \quad (12)$$



## Regularization with Fourier Transform

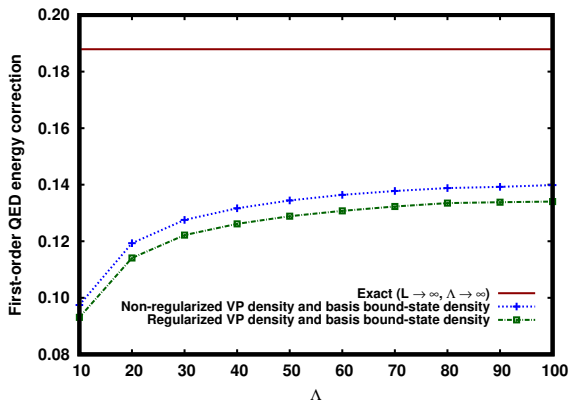
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## Computing the Lamb shift

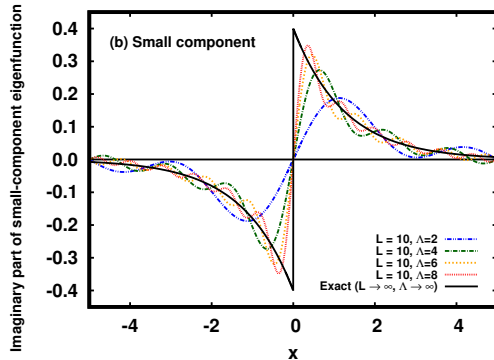
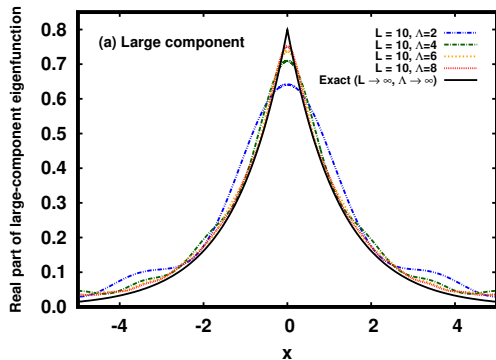
• Non-regularized:  $\mathcal{E}_{\text{Lamb}} = \int n_{1s}(x)n^{\text{VP}}(x)dx$

• Regularized:  $\mathcal{E}_{\text{Lamb}}^{\text{reg}} = \int n_{1s}(x)n_{\text{reg}}^{\text{VP}}(x)dx + \mathcal{N}_0^{\text{VP}}n_{1s}(0)$



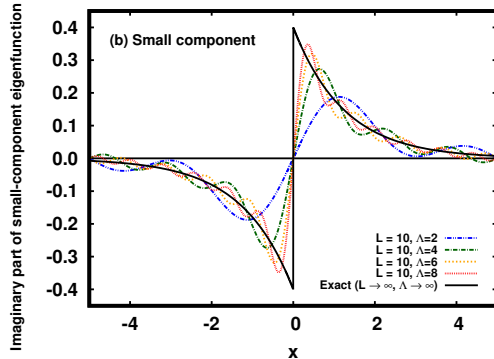
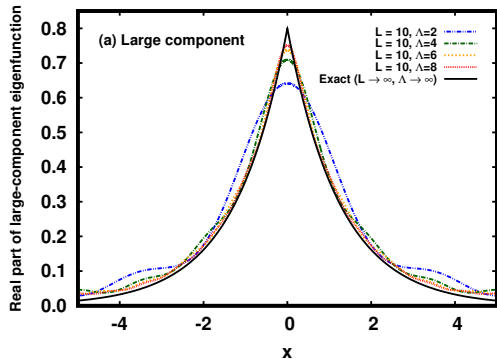
Value at  $x = 0$

$$n_{1s}(x=0) = |\psi_{1s}^L(0)|^2 + |\psi_{1s}^S(0)|^2 \quad (14)$$



Value at  $x = 0$

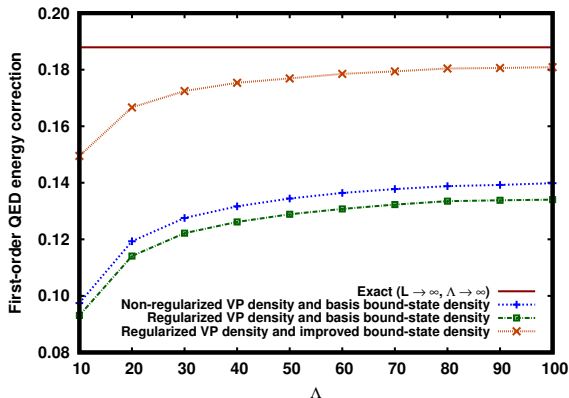
$$n_{1s}^{\text{improved}}(x=0) = |\psi_{1s}^L(0)|^2 + \max_{x \in \mathbb{R}} (|\psi_{1s}^S(x)|^2) \quad (15)$$



## Computing the Lamb shift

• Non-regularized:  $\mathcal{E}_{\text{Lamb}} = \int_{\mathbb{R}} n_{1s}(x) n^{\text{vp}}(x) dx$

• Regularized:  $\mathcal{E}_{\text{Lamb}}^{\text{reg}} = \int n_{1s}(x) n_{\text{reg}}^{\text{vp}}(x) dx + \mathcal{N}_0^{\text{vp}} n_{1s}^{\text{improved}}(0)$



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- ✓ Develop these quantities in a finite basis set.
- Use our insight to solve the 3D problem.

- We developed an effective QED theory including electron-positron pairs.
- Some singularities appears in a finite basis.
- This 1D model help us to understand this problem.